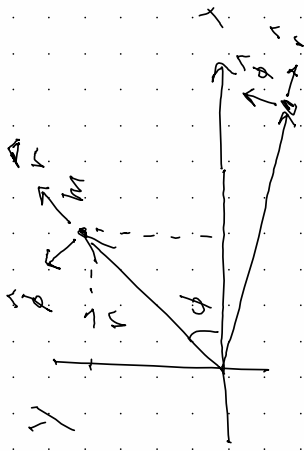


PHY 5306



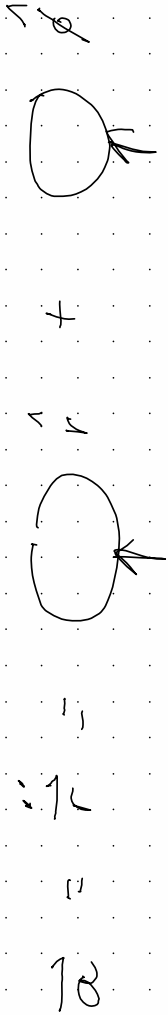
2-d motion:

$$\begin{aligned} \vec{r} &= r \hat{r} \\ \vec{v} &= \dot{r} \hat{r} + r \dot{\phi} \hat{\phi} \\ \vec{a} &= (\ddot{r} - r \dot{\phi}^2) \hat{r} + (2\dot{r}\dot{\phi} + r\ddot{\phi}) \hat{\phi} \end{aligned}$$

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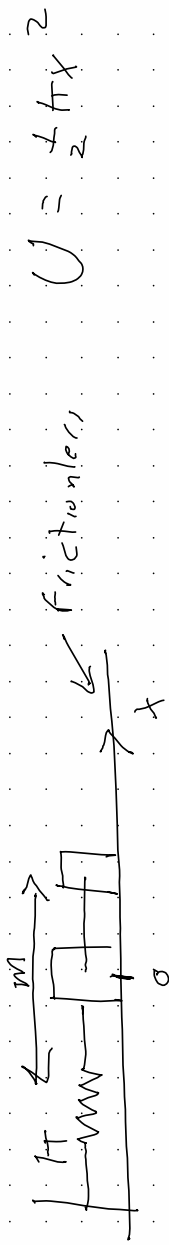
~~kinetic~~ kinetic energy for a single mass m
in Cartesian, sph, polar, and polar coordinates.

$$T = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2)$$

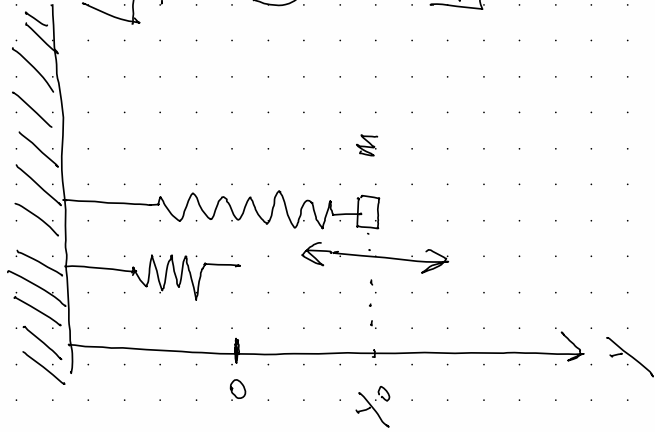
$$= \frac{1}{2} m (\quad) \quad \swarrow \text{spherical } (r, \theta, \phi)$$

$$= \frac{1}{2} m (\quad) \quad \swarrow \text{cylindrical } (\rho, \phi, z)$$

$$x = r \sin \theta \cos \phi, \text{ etc}$$



$$U = \frac{1}{2} kx^2$$



$$L = T - U$$

$$T = \frac{1}{2} m \dot{y}^2$$

$$U = \underbrace{\frac{1}{2} k y^2}_{U_s} - \underbrace{mgy}_{U_g}$$

$$E = T + U$$

$$= \frac{1}{2} m \dot{y}^2 + \underbrace{\frac{1}{2} k y^2 - mgy}_{U_{\text{eff}}(y)}$$

↑
plot this

$$\dot{x}^2 = \left(\dot{x} - l \dot{\phi} \cos \phi \right)^2$$

$$= \dot{x}^2 + l^2 \dot{\phi}^2 \cos^2 \phi - 2l \dot{x} \dot{\phi} \cos \phi$$

$$\dot{y}^2 = l^2 \dot{\phi}^2 \sin^2 \phi$$

$$\rightarrow \dot{x}^2 + \dot{y}^2 = \dot{x}^2 + l^2 \dot{\phi}^2 - 2l \dot{x} \dot{\phi} \cos \phi$$

$$T = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2)$$

$$= \frac{1}{2} m \dot{x}^2 + \frac{1}{2} m l^2 \dot{\phi}^2 - m l \dot{x} \dot{\phi} \cos \phi$$

$$U = -mgl \cos \phi$$

$$L = \frac{1}{2} m \dot{x}^2 + \frac{1}{2} m l^2 \dot{\phi}^2 - m l \dot{x} \dot{\phi} \cos \phi + mgl \cos \phi$$

$$L = L(\phi, \dot{\phi}, t) \rightarrow L' = L + \frac{d}{dt} f(\phi, t)$$

$$L = \cancel{\frac{1}{2} m \dot{\phi}^2} + \frac{1}{2} m l^2 \dot{\phi}^2 - m l \dot{\phi} \cos \phi + m g l \cos \phi$$

~~~~~

prescribed function of time  $\rightarrow$  ignore

$$\frac{\partial L}{\partial \dot{\phi}} = \frac{\partial L}{\partial y}$$

$$\dot{\phi} \cos \phi = \frac{d}{dt} (\dot{\phi} \sin \phi) - \dot{\phi} \sin \phi$$

$$\dot{\phi} \sin \phi + \dot{\phi} \cos \phi$$

$$- m l \dot{\phi} \cos \phi = \underbrace{\frac{d}{dt} (-m l \dot{\phi} \sin \phi)} + m l \dot{\phi} \sin \phi$$

ignore this in Lagrangian

$$L = \frac{1}{2} m l^2 \dot{\phi}^2 + m g l \sin \phi + m g l \cos \phi$$

i) show that both Lagrangians give the same EOMs

$$\frac{\partial \mathcal{L}}{\partial \dot{\phi}} = \left( \frac{\partial \mathcal{L}}{\partial \dot{\phi}} \right)_{\text{p}} \Leftrightarrow \frac{\partial \mathcal{L}}{\partial \phi} = \left( \frac{\partial \mathcal{L}}{\partial \phi} \right)_{\text{p}}$$

$$(ii) \quad \mathcal{L}' = \frac{1}{2} m l^2 \dot{\phi}^2 + m a l \sin \phi + m g l \cos \phi$$

does not depend explicitly on time

$$\rightarrow E = h = H = \underbrace{\dot{\phi} \frac{\partial \mathcal{L}}{\partial \dot{\phi}} - \mathcal{L}}_{\text{more generally}} = \text{const}$$

$$= \sum_i \dot{q}_i \frac{\partial \mathcal{L}}{\partial \dot{q}_i} - \mathcal{L} \quad (\text{more generally})$$

now

$p_i = \text{momentum conjugate to } q_i$

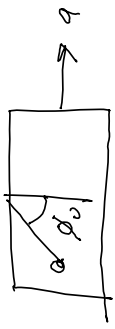
in general  $\rightarrow$

$$\neq T + U \leftarrow \begin{matrix} \text{total} \\ \text{mechanical} \\ \text{energy} \end{matrix}$$



iii) Equil solution:  $\dot{\phi} = 0$

$$\tan \phi_0 = \frac{a}{g}$$



Use EOMJ from Lagrangian to show +411

$$E = \frac{1}{2} m l^2 \dot{\phi}^2 - m l a \sin \phi - m g l \cos \phi$$

$$= \dot{\phi} \frac{\partial L}{\partial \dot{\phi}} - L$$

$$E = \frac{1}{2} m l^2 \dot{\phi}^2 + U_{\text{eff}}(\phi)$$

$$U_{\text{eff}}(\phi) = -m l (a \sin \phi + g \cos \phi)$$

★ graph it



$$0 = \left. \frac{dU_{\text{eff}}}{d\phi} \right|_{\phi=\phi_0}$$

$$U_{\text{eff}}(\phi) = U_{\text{eff}}(\phi_0) + \cancel{\frac{dU_{\text{eff}}}{d\phi} \bigg|_{\phi_0}} (\phi - \phi_0) + \frac{1}{2} \frac{d^2 U_{\text{eff}}}{d\phi^2} \bigg|_{\phi_0} (\phi - \phi_0)^2 + \dots$$

$$f(x) = f(a) + f'(a)(x-a) + \frac{1}{2} f''(a)(x-a)^2$$

$$+ \frac{1}{3!} f'''(a)(x-a)^3 + \dots$$

~~df~~

$$E = T + U_{\text{eff}}(\phi)$$

$$= T + U_{\text{eff}}(\phi_0) + \frac{1}{2} \mathcal{H}(\phi - \phi_0)^2$$

$$= \frac{1}{2} m \dot{\phi}^2 + \underbrace{U_{\text{eff}}(\phi_0)}_{\text{const}} + \frac{1}{2} \mathcal{H}(\phi - \phi_0)^2$$

$$x \equiv \phi - \phi_0 \quad |x| \ll 1 \quad \text{const} \rightarrow \text{ignore}$$

$$x = \phi$$

$$E = \underbrace{\frac{1}{2} m \dot{x}^2}_M + \frac{1}{2} \mathcal{H} x^2, \quad \omega = \sqrt{\frac{\mathcal{H}}{m}} = \sqrt{\frac{\mathcal{H}}{M}}$$

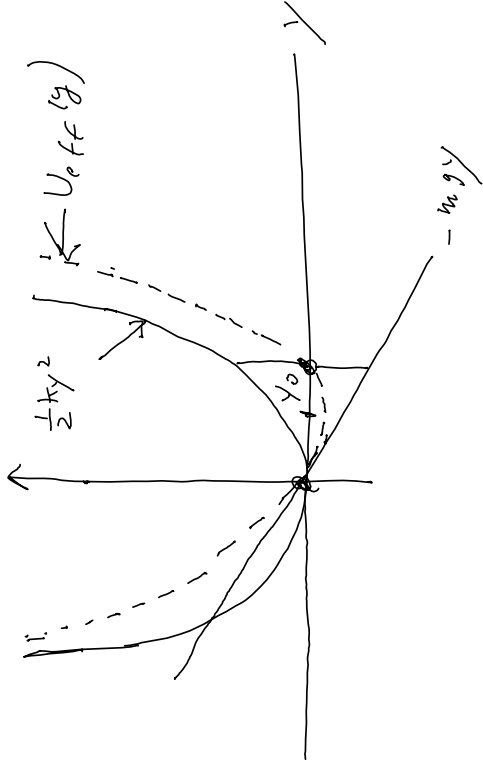
Answer, to problems posed at the end of the last class

$$(1) \vec{a} = (\underbrace{\ddot{r} - r\dot{\phi}^2}_{\text{centrifugal acceleration}}) \hat{r} + (\underbrace{2\dot{r}\dot{\phi} + r\ddot{\phi}}_{\text{tangential acceleration}}) \hat{\phi}$$

$$(2) T = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\phi}^2 + \dot{z}^2) + r^2 \sin^2 \theta \dot{\phi}^2$$

$$= \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\phi}^2 + \dot{z}^2)$$

$$(3) \text{ plot } U_{\text{eff}}(y) = \frac{1}{2} k y^2 - m g y$$



$$U_{\text{eff}}'(y) = U_{\text{eff}}'(y_0) + \frac{1}{2} k (y - y_0)^2$$

where  $y_0$  is the solution to

$$0 = \left. \frac{dU_{\text{eff}}}{dy} \right|_{y_0}$$

and  $k$  is given by

$$k \equiv \left. \frac{d^2 U_{\text{eff}}}{dy^2} \right|_{y_0}$$

$$0 = \frac{dU_{\text{eff}}}{dy} \bigg|_{y_0} = ky_0 - mg \rightarrow y_0 = \frac{mg}{k}$$

$\mathcal{H} = \frac{d^2 U_{\text{eff}}}{dy^2} \bigg|_{y_0} = k$  (so it is the same as  $k$  for this problem)

$$U_{\text{eff}}(y_0) = \frac{1}{2}ky_0^2 - mgy_0 = \frac{1}{2}k\left(\frac{mg}{k}\right)^2 - mg\left(\frac{mg}{k}\right) = -\frac{1}{2}\frac{m^2g^2}{k}$$

$$\rightarrow U_{\text{eff}}(y) = \boxed{-\frac{1}{2}\frac{m^2g^2}{k} + \frac{1}{2}k\left(y - \frac{mg}{k}\right)^2}$$

You can also obtain the same expression for  $U_{\text{eff}}(y)$  by completing the square:

$$\begin{aligned} U_{\text{eff}}(y) &= \frac{1}{2}ky^2 - mgy \\ &= \frac{1}{2}k\left(y^2 - 2\frac{mg}{k}y\right) \\ &= \frac{1}{2}k\left[\left(y - \frac{mg}{k}\right)^2 - \frac{m^2g^2}{k^2}\right] \\ &= \frac{1}{2}k\left(y - \frac{mg}{k}\right)^2 - \frac{1}{2}\frac{m^2g^2}{k} \end{aligned}$$