

Lec #23: Tuesday Nov 10th

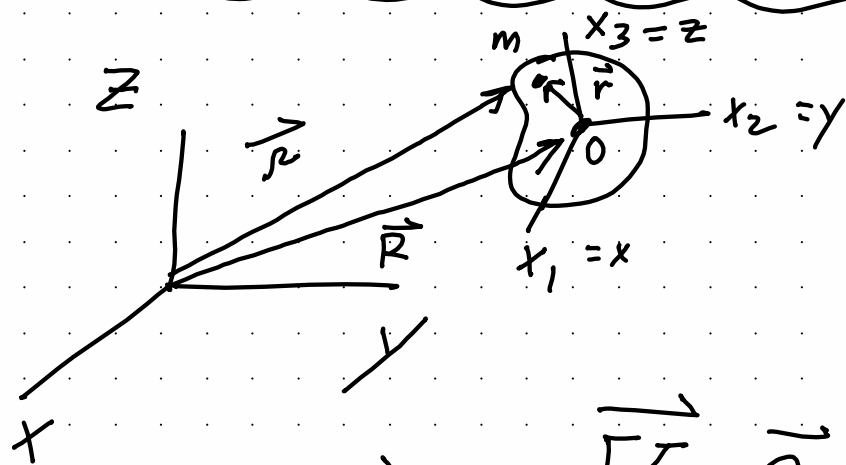
— Quiz #5: Thursday

— Midterm #2: Next Thursday 11/19 (Scattering, small oscillations, some rigid body)

— Today: Rigid body motion

(Sec 31-36, 38, 39)

└ non-inertial
└ static equilibrium



O: at COM (usually)
 x_1, x_2, x_3 : Fixed in RB

$$\vec{v} = \frac{d\vec{r}}{dt} = \vec{V} + \vec{\Omega} \times \vec{r}$$

$\vec{\Omega}$: angular-velocity vector

$$\left. \begin{aligned} \vec{V} &= \frac{d\vec{R}}{dt} \\ \vec{\Omega} &= \frac{d\vec{\phi}}{dt} \end{aligned} \right\} \dot{q}_i$$

$(\vec{R}, \vec{\phi})$: 6 DOF = \dot{q}_i

$$T = \frac{1}{2} m v^2$$

$$\vec{p} = m \vec{v}$$

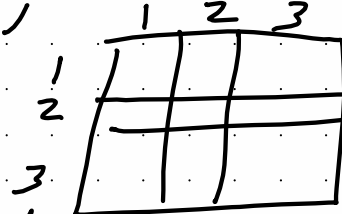
rotational quantities

$$T_{\text{rot}} = \frac{1}{2} I \Omega^2$$

$$\vec{M} = I \vec{\Omega} \rightarrow M_i = \sum_j I_{ij} \Omega_j$$

I : moment of inertia

$\rightarrow I_{ij}$: inertia tensor



$$M_i = I_{ij} \Omega_j$$

KE:

$$T = \frac{1}{2} \sum_a m_a |\vec{v}_a|^2$$

$$\vec{v}_a = \vec{V} + \vec{\Omega} \times \vec{r}_a$$

$$= \frac{1}{2} \sum_a m_a \left| \vec{V} + \vec{\Omega} \times \vec{r}_a \right|^2$$

$$\left| \vec{A} + \vec{B} \right|^2 = A^2 + B^2 + 2\vec{A} \cdot \vec{B}$$

$$= \frac{1}{2} \sum_a m_a \left(|\vec{V}|^2 + |\vec{\Omega} \times \vec{r}_a|^2 + 2\vec{V} \cdot (\vec{\Omega} \times \vec{r}_a) \right)$$

$$\textcircled{1} = \frac{1}{2} \sum_a m_a |\vec{V}|^2 = \boxed{\frac{1}{2} \mu V^2} \quad \text{total mass}$$

$$\begin{aligned} \textcircled{3} &= \sum_a m_a \vec{V} \cdot (\vec{\Omega} \times \vec{r}_a) \\ &= \left(\sum_a m_a \vec{r}_a \right) \cdot (\vec{V} \times \vec{\Omega}) \\ &= \mu \vec{R}_{\text{com}} \cdot (\vec{V} \times \vec{\Omega}) \end{aligned}$$

$$\begin{aligned} \vec{A} \cdot (\vec{B} \times \vec{C}) \\ &= \vec{B} \cdot (\vec{C} \times \vec{A}) \\ &= \vec{C} \cdot (\vec{A} \times \vec{B}) \end{aligned}$$

$\underbrace{\quad}_{=0 \text{ for } \vec{O} \text{ at com}}$

$$\begin{aligned}
(2) &= \frac{1}{2} \sum_a m_a |\vec{\Omega} \times \vec{r}_a|^2 \\
&= \frac{1}{2} \sum_a m_a (\vec{\Omega} \times \vec{r}_a) \cdot (\vec{\Omega} \times \vec{r}_a) \\
&= \frac{1}{2} \sum_a m_a \vec{\Omega} \cdot (\vec{r}_a \times (\vec{\Omega} \times \vec{r}_a)) \\
&= \frac{1}{2} \sum_a m_a \vec{\Omega} \cdot (\vec{\Omega} r_a^2 - \vec{r}_a (\vec{r}_a \cdot \vec{\Omega})) \\
&= \frac{1}{2} \sum_a m_a (\Omega^2 r_a^2 - (\vec{\Omega} \cdot \vec{r}_a)^2) \\
&= \frac{1}{2} \sum_a m_a \left(\sum_{i,j} \Omega_i \Omega_j \delta_{ij} r_a^2 - \sum_{i,j} \Omega_i r_{ai} \Omega_j r_{aj} \right) \\
&= \frac{1}{2} \sum_{i,j} \left(\sum_a m_a (r_a^2 \delta_{ij} - r_{ai} r_{aj}) \right) \Omega_i \Omega_j \\
&= \boxed{\frac{1}{2} \sum_{i,j} I_{ij} \Omega_i \Omega_j} \quad \left(= \frac{1}{2} \begin{array}{|c|c|c|} \hline \square & \square & \square \\ \hline \square & \square & \square \\ \hline \square & \square & \square \\ \hline \end{array} \right)
\end{aligned}$$

$$\vec{A} \times (\vec{B} \times \vec{C}) = \vec{B} (\vec{A} \cdot \vec{C}) - \vec{C} (\vec{A} \cdot \vec{B})$$

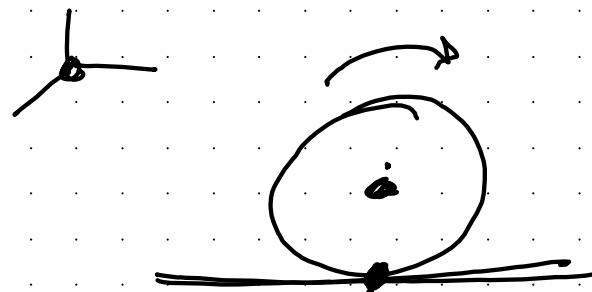
$$T = \underbrace{\frac{1}{2} M V^2}_{\text{trans}} + \underbrace{\frac{1}{2} \sum_{i,j} I_{ij} \Omega_i \Omega_j}_{\text{rotational}} + \textcircled{Q}$$

$\left(\frac{1}{2} I \Omega^2 \right)$
 Freshman physics

for COM at origin
OF RB Frame

\vec{M} : wrt COM of body

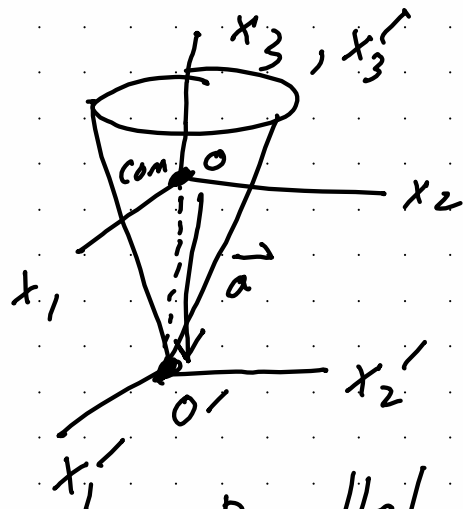
$$\begin{aligned}
 \vec{M} &= \sum_a \vec{r}_a \times \vec{p}_a \\
 &= \sum_a m_a \vec{r}_a \times \vec{v}_a \\
 &= \sum_a m_a \vec{r}_a \times (\vec{V} + \vec{\Omega} \times \vec{r}_a) \\
 &= \sum_a m_a \vec{r}_a \times (\vec{\Omega} \times \vec{r}_a)
 \end{aligned}$$



$$= \sum_a m_a (\Omega^2 r_a^2 - \vec{r}_a (\vec{r}_a \cdot \vec{\Omega}))$$

$$M_i = \sum_j I_{ij} \Omega_j$$

$$\vec{M} = I \vec{\Omega} \quad (\text{Fresh. physics})$$



parallel-axis
theorem

I_{ij}'
wrt O'

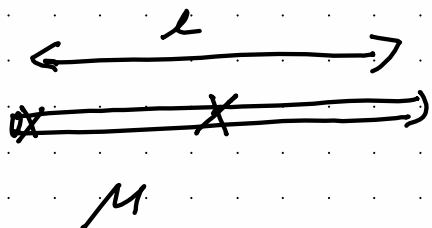
I_{ij}
wrt O (com)

$$I_{ij}' = I_{ij} + \mu(a^2 \delta_{ij} - a_i a_j)$$

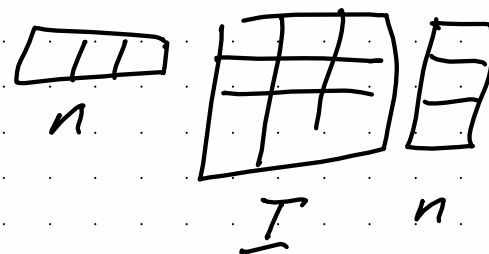
\vec{a} : Vector from O to O'

$$I(\hat{n}) = \sum_{i,j} I_{ij} n_i n_j$$

moment
of inertia



\hat{n} : axis of
rotation



$$I_{com} = \frac{1}{12} \mu l^2$$

$$I_{end} = \frac{1}{3} \mu l^2$$

$$\mathbf{I}_{ij} = \sum_a m_a (r_a^2 \delta_{ij} - r_{ai} r_{aj})$$

$\underbrace{\quad}_{\substack{\text{3x3 real} \\ \text{symmetric}}}$
 $\underbrace{\quad}_{dm}$
 $= \int \rho dV (r^2 \delta_{ij} - r_i r_j)$

can always be diagonalized

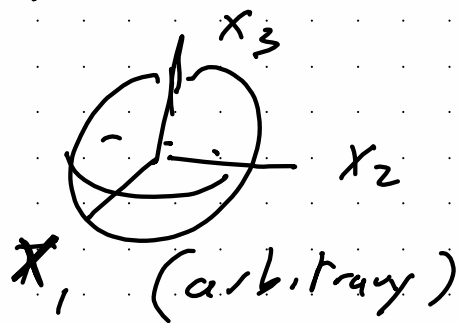
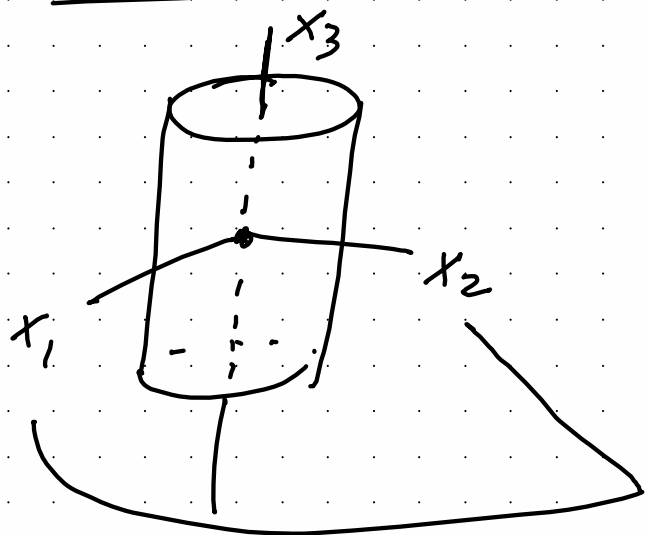
I_1	0	0
0	I_2	0
0	0	I_3

$$\mathbf{I}_{ij} = I_i \delta_{ij}$$

$$T_{\text{rot}} = \frac{1}{2} \sum_{i,j} \mathbf{I}_{ij} \Omega_i \Omega_j$$

principle axes: (x_1, x_2, x_3)

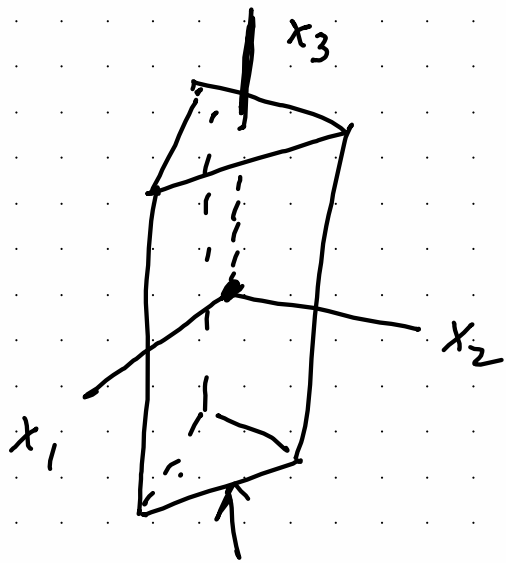
$$= \frac{1}{2} (I_1 \Omega_1^2 + I_2 \Omega_2^2 + I_3 \Omega_3^2)$$



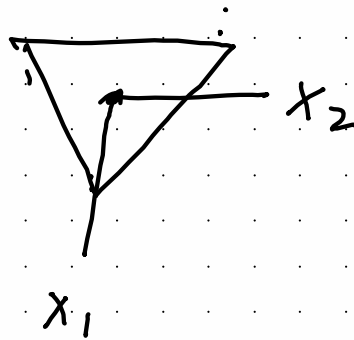
arbitrary

$$M_i = \sum_j \mathbf{I}_{ij} \Omega_j$$

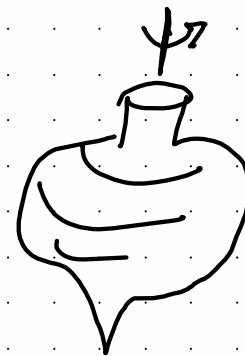
$$M_1 = I_1 \Omega_1, M_2 = I_2 \Omega_2, M_3 = I_3 \Omega_3$$

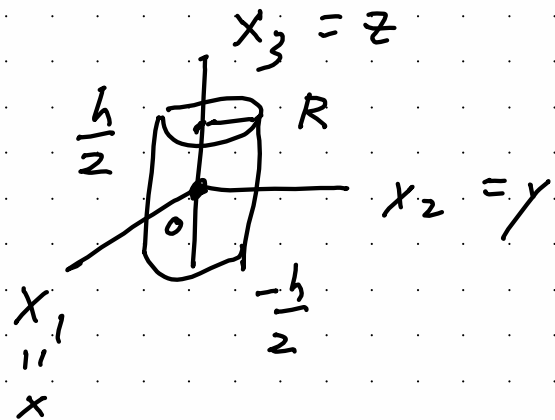


equilateral
triang



$I_1 = I_2 \neq I_3$
symmetrical
top





total mass M

$$\rho = \frac{M}{\text{volume}} = \frac{M}{\pi R^2 h}$$

$$I_3 = I_{33} = \int \rho dV \left(\underbrace{r^2}_{=1} \underbrace{\delta_{33}}_{=2} - \underbrace{\underbrace{r_3}_{=z} \underbrace{r_3}_{=z}}_{=z^2} \right)$$

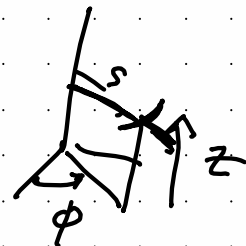
$$= \int \rho dV (r^2 - z^2)$$

$$= \int \rho dV (x^2 + y^2) = \int \rho dV s^2$$

cylindrical:

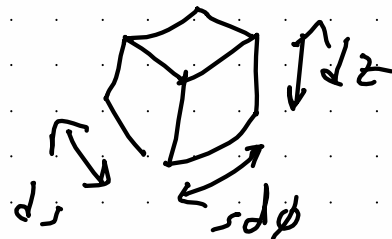
s, ϕ, z

$$s^2 = x^2 + y^2$$



$$dV = ds \, s d\phi \, dz$$

$$= s ds d\phi dz$$

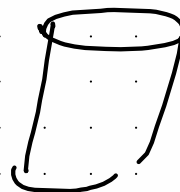
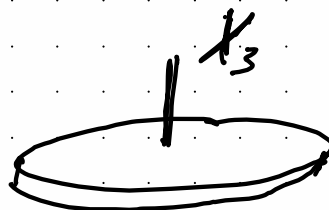


$$\begin{aligned}
 I_3 &= \int \rho dV s^2 \\
 &= \frac{M}{\pi R^2 h} \int_0^{2\pi} d\phi \int_{-h/2}^{h/2} dz \int_0^R s^3 ds
 \end{aligned}$$

$\underbrace{\quad}_{2\pi} \quad \underbrace{\quad}_h \quad \underbrace{\quad}_{\frac{R^4}{4}}$

$$= \frac{M}{\pi R^2 h} \cdot 2\pi h \frac{R^4}{4}$$

$$= \boxed{\frac{1}{2} M R^2}$$



$$I_1 = I_2 \equiv I$$

$$\begin{aligned}
 I_1 &= \int \rho dV (r^2 - x^2) \\
 + I_2 &= \int \rho dV (r^2 - y^2)
 \end{aligned}$$

$$2I = \int p dV (2r^2 - x^2 - y^2)$$

$$\left(\begin{array}{l} r^2 = s^2 + z^2 \\ x^2 + y^2 = s^2 \end{array} \right)$$

$$2I = \int p dV (s^2 + 2z^2)$$

$$\boxed{I = \frac{1}{2} \underbrace{\int p dV s^2}_{I_3} + \int p dV z^2}$$

$$\boxed{= \frac{1}{2} I_3 + \int p dV z^2}$$

easy-to-evaluate

$$\int \rho dV z^2 = \frac{M}{\pi R^2 h} \int_0^{2\pi} d\phi \int_{-h/2}^{h/2} dz \cdot z^2 \int_0^R s ds$$

$\underbrace{\quad}_{2\pi} \quad \underbrace{\quad}_{\frac{z^3}{3} \Big|_{-h/2}^{h/2}} \quad \underbrace{\quad}_{\frac{R^2}{2}}$

$$= \frac{2}{3} \left(\frac{h}{2} \right)^3$$

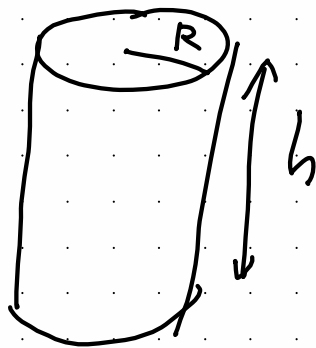
$$= \frac{h^3}{12}$$

$$= \frac{M}{\cancel{\pi R^2 h}} \cdot \cancel{2\pi} \cdot \frac{h^3}{12} \cdot \frac{\cancel{R^2}}{\cancel{2}}$$

$$= \boxed{\frac{M h^2}{12}}$$

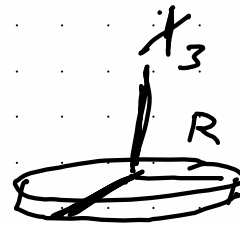
$$\begin{aligned}
 I &= \frac{1}{2} \left(\frac{1}{2} M R^2 \right) + \frac{M h^2}{12} \\
 &= \frac{1}{4} M R^2 + \frac{1}{12} M h^2 \\
 &= \frac{1}{4} M \left(R^2 + \frac{1}{3} h^2 \right) = I_1, I_2
 \end{aligned}$$

$$I_3 = \frac{1}{2} M R^2$$



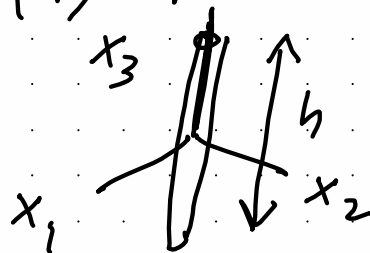
Limiting cases

(i) Disk ($h \rightarrow 0$)



$$\begin{aligned}
 I_3 &= \frac{1}{2} M R^2 \\
 I_1, I_2 &= \frac{1}{4} M R^2
 \end{aligned}$$

(i) thin rod ($R \rightarrow 0$)



$$I_3 = 0$$

$$I_1 = I_2 = \frac{1}{12} M h^2$$

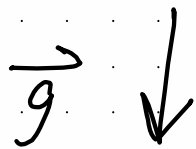
$$L = T - U$$

$$= \frac{1}{2} \mu V^2 + \frac{1}{2} \sum_{i,j} I_{ij} \Omega_i \Omega_j - U$$

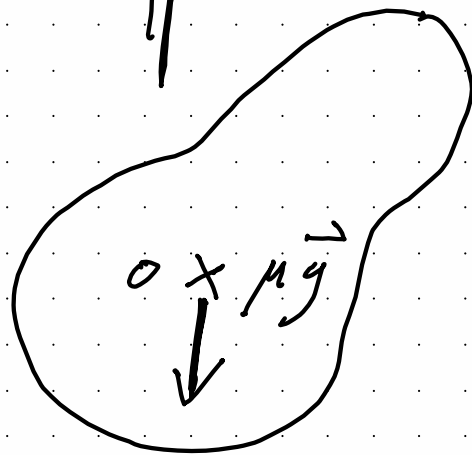
$$\vec{V} = \dot{\vec{R}}$$

$$\vec{\Omega} = \dot{\vec{\phi}}$$

$$L(\vec{R}, \vec{\phi}, \dot{\vec{R}}, \dot{\vec{\phi}})$$



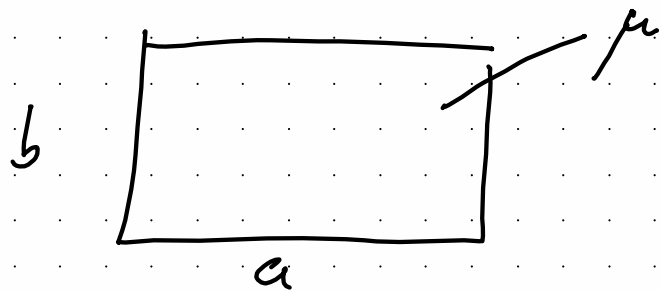
uniform
field



Lecture #24 : Thursday 11/12

- Quiz #5 (today)
- Midterm #2 (next Thursday) (scattering, small oscillations, RB motion)
- Today's topics:
 - (1) RB EOMs
 - (2) Euler's equation,
 - (3) Euler angles

Q5: Calculate the principal moments of inertia for a 2-d rectangle with sidelengths a, b . uniform

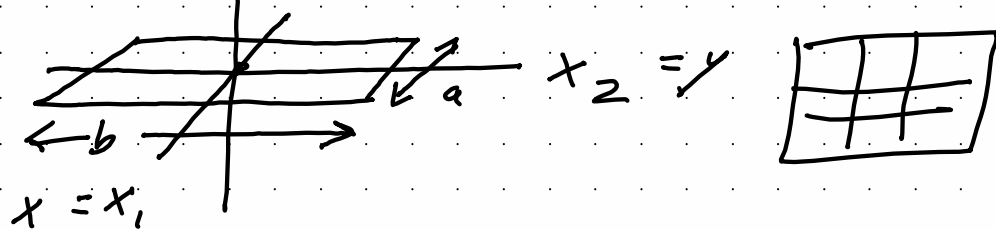


name-25.pdf

$$z = x_3$$

$$I_{ij} = \int \rho dV (r^2 \delta_{ij} - r_i r_j)$$

$$r^2 = x^2 + y^2 + z^2$$



$$I_{11} = \int \rho dV (r^2 - x^2)$$

$$= \int \rho dV y^2$$

$$= \frac{M}{ab} \int_{-a/2}^{a/2} dx \int_{-b/2}^{b/2} dy y^2$$

$$= \frac{M}{ab} a \left[\frac{y^3}{3} \right]_{-b/2}^{b/2}$$

$$= \frac{M}{b} \frac{2}{3} \left(\frac{b}{2} \right)^3$$

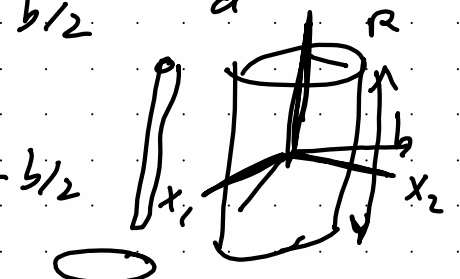
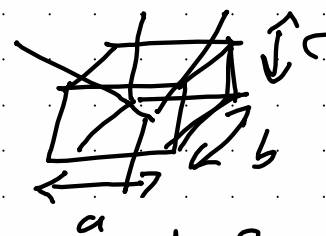
$$= \frac{M}{b} \frac{2}{3} \frac{b^3}{8}$$

$$= \left[\frac{1}{12} M b^2 \right]$$

$$dm = \rho dV$$

$$= \sigma dx dy$$

$$= \frac{M}{ab} dx dy$$



$$I_2 = \frac{1}{12} M a^2$$

$$I_3 = \int \sigma dx dy$$

$$(x^2 + y^2)$$

$$= I_1 + I_2$$

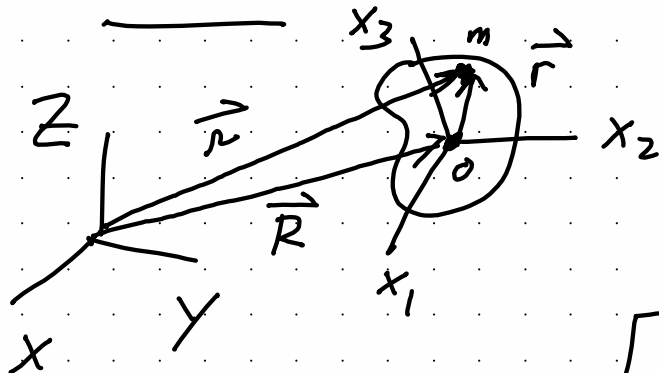
$$I_3 = \frac{1}{12} M (a^2 + b^2)$$

EOMs:

Final result:

$$\frac{d\vec{P}}{dt} = \vec{F} = \sum \vec{f}$$

$$\frac{d\vec{M}}{dt} = \vec{K} = \sum \vec{r} \times \vec{f}$$



$$\vec{r} = \vec{R} + \vec{r}$$

$$\delta \vec{r} = \delta \vec{R} + \delta \phi \times \vec{r}$$



$$L = T - U$$

$$= \frac{1}{2} \mu \vec{V}^2 + \frac{1}{2} \sum_{i,j} I_{ij} \Omega_i \Omega_j - U(\vec{r})$$

$$\delta L = \mu \vec{V} \cdot \delta \vec{V} + \sum_{i,j} I_{ij} \Omega_i \delta \Omega_j - \sum \frac{\partial U}{\partial \vec{r}} \cdot \delta \vec{r}$$

$$= \sum \frac{\partial U}{\partial \vec{r}} \cdot (\delta \vec{R} + \delta \phi \times \vec{r})$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \vec{V}} \right) = \frac{\partial L}{\partial \vec{R}}$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \vec{\Omega}} \right) = \frac{\partial L}{\partial \phi}$$

$$\frac{1}{2} \sum_{i,j} I_{i,j} \cdot \underline{\underline{\delta \Omega_i}} \cdot \Omega_j + \frac{1}{2} \sum_{i,j} I_{i,j} \cdot \Omega_i \cdot \delta \Omega_j$$

$$= \quad \quad \quad // \quad + \frac{1}{2} \sum_{i,i} I_{i,i} \cdot \Omega_i \cdot \underline{\underline{\delta \Omega_i}}$$

$$= \quad \quad \quad // \quad + \frac{1}{2} \sum_{\substack{j,i \\ \text{swap}}} I_{i,j} \cdot \Omega_j \cdot \delta \Omega_i$$

$$= \frac{1}{2} \sum_{i,j} I_{i,j} \cdot \Omega_j \cdot \delta \Omega_i + \frac{1}{2} \sum_{i,j} I_{i,j} \cdot \Omega_j \cdot \delta \Omega_i$$

$$= \boxed{\sum_{i,j} I_{i,j} \cdot \Omega_j \cdot \delta \Omega_i}$$

$$\delta L = \underbrace{m \vec{V}}_{\vec{p}} \cdot \delta \vec{V} + \sum_{i,j} I_{ij} \cdot \delta \Omega_i \quad \left/ \begin{array}{l} \vec{A} \cdot (\vec{B} \times \vec{C}) \\ \vec{B} \cdot (\vec{C} \times \vec{A}) \end{array} \right.$$

$$- \left(\sum \frac{\partial U}{\partial \vec{r}} \right) \cdot \delta \vec{R} - \sum \frac{\partial U}{\partial \vec{r}} \cdot (\delta \vec{\phi} \times \vec{r})$$

$$\left. \begin{array}{l} \sum_i m_i \delta \Omega_i \\ = \vec{M} \cdot \delta \vec{\Omega} \end{array} \right\} \quad \begin{array}{l} = - \sum \delta \vec{\phi} \cdot \left(\vec{r} \times \frac{\partial U}{\partial \vec{r}} \right) \\ = - \delta \vec{\phi} \cdot \left(\sum \vec{r} \times \frac{\partial U}{\partial \vec{r}} \right) \end{array}$$

$$dF = \cancel{\otimes} dx + \cancel{\otimes} dy$$

$$\delta L = \vec{p} \cdot \delta \vec{V} + \vec{M} \cdot \delta \vec{\Omega} + \left(\sum \vec{F} \right) \cdot \delta \vec{R} + \delta \vec{\phi} \cdot \left(\sum \vec{r} \times \vec{F} \right)$$

$$\vec{p} = \frac{\partial L}{\partial \vec{V}}, \quad \vec{M} = \frac{\partial L}{\partial \vec{\Omega}}, \quad \sum \vec{F} = \frac{\partial L}{\partial \vec{R}}, \quad \sum \vec{r} \times \vec{F} = \frac{\partial L}{\partial \vec{\phi}}$$

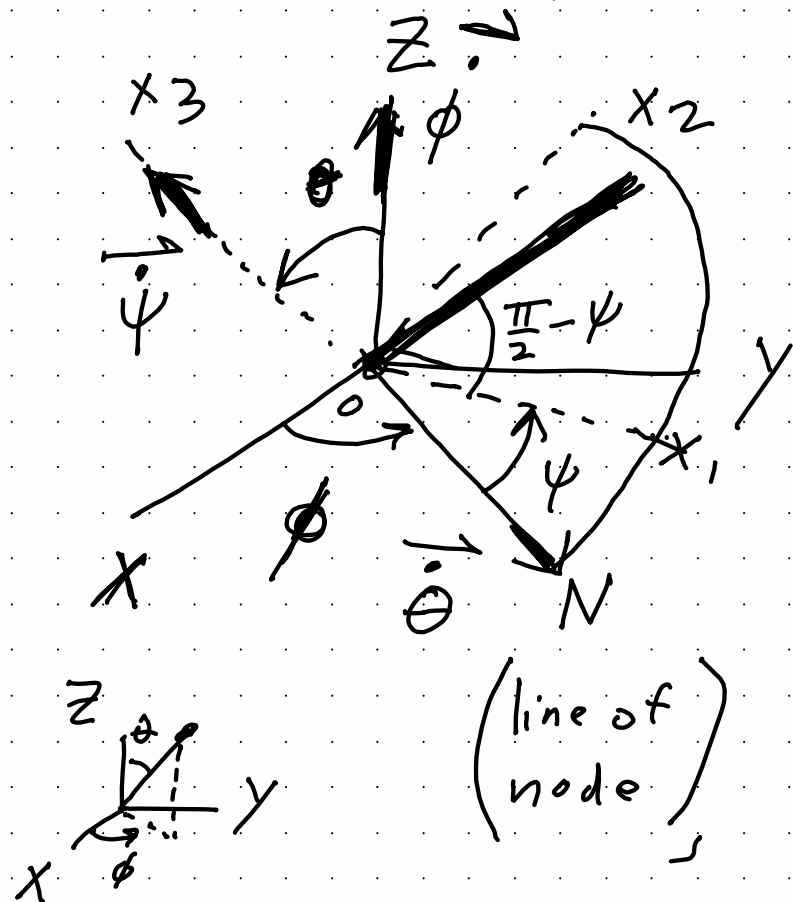
$$\frac{d}{dt} \left(\frac{\partial L}{\partial \vec{V}} \right) = \frac{\partial L}{\partial \vec{R}} \rightarrow \left[\frac{d\vec{p}}{dt} = \sum \vec{F} \equiv \vec{F} \right] \quad \left| \quad \frac{d}{dt} \left(\frac{\partial L}{\partial \vec{\Omega}} \right) = \frac{\partial L}{\partial \vec{\phi}} \rightarrow \left[\frac{d\vec{M}}{dt} = \sum \vec{r} \times \vec{F} \equiv \vec{K} \right]$$

Euler's equations / Euler angle, ★

$\vec{R}, \vec{\phi}$

Ω_i : i th component of $\vec{\Omega}$

\angle wrt \hat{x}_i



$$\vec{\Omega} = \dot{\phi} \vec{e}_1 + \dot{\theta} \vec{e}_2 + \dot{\psi} \vec{e}_3$$

$$\dot{\psi} = \dot{\psi} \hat{x}_3$$

$$\dot{\theta} = \dot{\theta} \cos \psi \hat{x}_1 - \dot{\theta} \sin \psi \hat{x}_2$$

$$\dot{\phi} = \dot{\phi} \cos \theta \hat{x}_3$$

$$+ \dot{\phi} \sin \theta \left(\sin \psi \hat{x}_1 + \cos \psi \hat{x}_2 \right)$$

$\cos(\frac{\pi}{2} - \psi) \quad \sin(\frac{\pi}{2} - \psi)$

Announcements

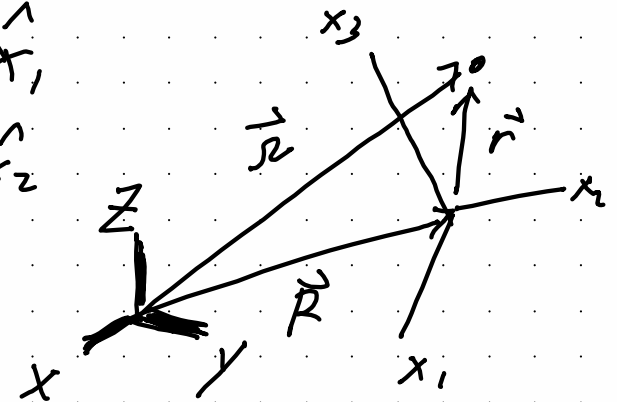
— Midterm II is this Thursday

— Today:

- i) Euler angles
- ii) Euler's equation for RB motion
- iii) Free rotation with $\vec{L} = \text{const}$
- iv) " of a symmetric top ($I_1 = I_2$)
- v) Heavy symmetrical top with one point fixed [prob 35.1]



$$\vec{\Omega} = (\dot{\phi} \sin \theta \sin \psi + \dot{\theta} \cos \psi) \hat{x}_1 + (\dot{\phi} \sin \theta \cos \psi - \dot{\theta} \sin \psi) \hat{x}_2 + (\dot{\psi} + \dot{\phi} \cos \theta) \hat{x}_3$$



$$\begin{aligned} \Omega_1 &= \dot{\phi} \sin \theta \sin \psi + \dot{\theta} \cos \psi \\ \Omega_2 &= \dot{\phi} \sin \theta \cos \psi - \dot{\theta} \sin \psi \\ \Omega_3 &= \dot{\psi} + \dot{\phi} \cos \theta \end{aligned}$$

$$\vec{v} = \frac{d\vec{r}}{dt} = \vec{V} + \vec{\Omega} \times \vec{r}$$

~~$\frac{d\vec{r}}{dt} = \frac{d\vec{r}}{dt} + \vec{\Omega} \times \vec{r}$~~

$\frac{d\vec{r}}{dt}$ wrt RB

Euler's equations: (wrt RB axes) \vec{A} : any vector

$$\begin{aligned} \frac{d\vec{P}}{dt} &= \sum \vec{F} = \vec{F} \\ \frac{d\vec{M}}{dt} &= \sum \vec{r} \times \vec{F} = \vec{K} \end{aligned}$$

$$\frac{d\vec{A}}{dt} = \frac{d'\vec{A}}{dt} + \vec{\Omega} \times \vec{A}$$

wrt inertial frame wrt rotating frame angular velocity of the rotating frame

$$\left(\frac{d' \vec{A}}{dt} \right)_i = \frac{dA_i}{dt} = \dot{A}_i$$

Cartesian
components
wrt rotating
Frame

$$\vec{a} = \frac{d^2 \vec{r}}{dt^2} = \ddot{x} \hat{x} + \ddot{y} \hat{y} + \ddot{z} \hat{z}$$
~~$$= \ddot{x} \hat{x} + \ddot{y} \hat{y} + \ddot{z} \hat{z}$$~~

$$\vec{A} = \sum_i \underline{A_i} \underline{\hat{x}_i}$$

$$\frac{d\vec{A}}{dt} = \underbrace{\sum_i \left(\frac{dA_i}{dt} \right) \hat{x}_i}_{\frac{d'\vec{A}}{dt}} + \underbrace{\sum_i A_i \frac{d\hat{x}_i}{dt}}_{\vec{\Omega} \times \vec{A}}$$

$$\frac{d\vec{A}}{dt} = \frac{d'\vec{A}}{dt} + \vec{\Omega} \times \vec{A}, \quad \left(\frac{d'\vec{A}}{dt}\right)_i = \dot{A}_i$$

$$\vec{F} = \frac{d\vec{P}}{dt} = \frac{d'\vec{P}}{dt} + \vec{\Omega} \times \vec{P}$$

$$\begin{aligned} \rightarrow \boxed{F_1} &= \dot{P}_1 + (\vec{\Omega} \times \vec{P})_1 \\ &= \dot{P}_1 + \Omega_2 P_3 - \Omega_3 P_2 \\ &= \mu(\dot{V}_1 + \Omega_2 V_3 - \Omega_3 V_2) \end{aligned}$$

$$\vec{P} = \mu \vec{V}$$

similar equations
for F_2, F_3

$$\boxed{\vec{K} = \frac{d\vec{M}}{dt}} = \frac{d'\vec{M}}{dt} + \vec{\Omega} \times \vec{M}$$

$$M_i = I_i \Omega_i$$

$$\boxed{K_x = \dot{M}_x}$$

$$\begin{aligned} \boxed{K_1} &= \dot{M}_1 + \Omega_2 M_3 - \Omega_3 M_2 \\ &= I_1 \dot{\Omega}_1 + \Omega_2 I_3 \Omega_3 - \Omega_3 I_2 \Omega_2 \\ &= I_1 \dot{\Omega}_1 + \Omega_2 \Omega_3 (I_3 - I_2) \end{aligned}$$

* (similar
equations
for K_2, K_3)

Free rotation: $\dot{H}_i = 0, \dot{F}_i = 0$

$$0 = I_1 \dot{\Omega}_1 + \Omega_2 \Omega_3 (I_3 - I_2) \quad \text{etc.}$$

$$0 = I_2 \dot{\Omega}_2 + \Omega_3 \Omega_1 (I_1 - I_3)$$

$$0 = I_3 \dot{\Omega}_3 + \Omega_1 \Omega_2 (I_2 - I_1)$$

Free rotation with $\vec{\Omega} = \text{const}$:

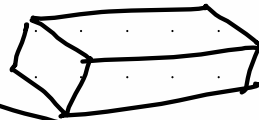
$$0 = \Omega_2 \Omega_3 (I_3 - I_2)$$

$$0 = \Omega_3 \Omega_1 (I_1 - I_3)$$

$$0 = \Omega_1 \Omega_2 (I_2 - I_1)$$

$$\boxed{\Omega_1 = \text{const}, \quad I_1 < \textcircled{I_2} < I_3}$$

$\Omega_2 = 0$
 $\Omega_3 = 0$ ← stable
 unstable ↓



$$\boxed{\Omega_2 = \text{const}, \quad \Omega_1 = 0, \quad \Omega_3 = 0}$$

$$\boxed{\Omega_3 = \text{const}, \quad \Omega_1 = 0, \quad \Omega_2 = 0}$$

$$\begin{cases} \frac{d\vec{\Omega}}{dt} = 0 \\ \frac{d'\vec{\Omega}}{dt} = 0 \end{cases}$$

$$\boxed{\frac{d\vec{\Omega}}{dt} = \frac{d'\vec{\Omega}}{dt} + \vec{\Omega} \times \vec{\Omega}}$$

$$\boxed{\frac{d\vec{M}}{dt} \neq \frac{d'\vec{M}}{dt}}$$

$$\Omega_1 = \text{const}, \quad \Omega_2 = 0, \quad \Omega_3 = 0 \quad : \text{exact}$$

$$\boxed{\begin{aligned} \Omega_1 &= \text{const} + \epsilon_1 \\ \Omega_2 &= \epsilon_2 \\ \Omega_3 &= \epsilon_3 \end{aligned}}$$

$\epsilon_{1,2,3}$: small time dependent perturbation

Keep 0th and 1st order terms.
Ignore 2nd order, e.g. $\epsilon_2 \epsilon_3$

$$0 = I_1 \frac{d}{dt} (\underbrace{\text{const}}_{\Omega_1} + \epsilon_1) + \underbrace{\epsilon_2 \epsilon_3 (\Gamma_3 - \Gamma_2)}_{\text{2nd order} \rightarrow \text{ignore}}$$

$$\approx I_1 \dot{\epsilon}_1$$

$$\rightarrow \epsilon_1 = \text{const} \rightarrow \boxed{\Omega_1 = \text{const}}$$

$$0 = I_2 \dot{\epsilon}_2 + \epsilon_3 \Omega_1 (\Gamma_1 - \Gamma_3)$$

$$0 = I_3 \dot{\epsilon}_3 + \Omega_1 \epsilon_2 (\Gamma_2 - \Gamma_1)$$

} coupled
1st order
diff. equations

Differentiate ..

$$\begin{aligned} 0 &= I_2 \ddot{\epsilon}_2 + \dot{\epsilon}_3 \Omega_1 (\Gamma_1 - \Gamma_3) \\ &= I_2 \ddot{\epsilon}_2 - \underbrace{\Omega_1 \epsilon_2 (\Gamma_2 - \Gamma_1)}_{\Gamma_3} \Omega_1 (\Gamma_1 - \Gamma_3) \end{aligned}$$

$$I_1 < I_2 < I_3$$

$$\begin{aligned} 0 &= I_2 \ddot{\epsilon}_2 + \dot{\epsilon}_3 \Omega_1 (I_1 - I_3) \\ &= I_2 \ddot{\epsilon}_2 - \underbrace{\Omega_1 \epsilon_2 (I_2 - I_1) \Omega_1 (I_1 - I_3)}_{I_3} \end{aligned}$$

$$= \ddot{\epsilon}_2 + \underbrace{\Omega_1^2 (I_2 - I_1) (I_3 - I_1)}_{I_2 I_3} \epsilon_2$$

$$\underbrace{I_2 I_3}_{\omega^2}$$

$$\begin{pmatrix} I_3 - I_2 \\ I_1 - I_2 \end{pmatrix}$$

$$= \ddot{\epsilon}_2 + \omega^2 \epsilon_2$$

$$\begin{aligned} \ddot{\epsilon}_2 &= -\omega^2 \epsilon_2 \rightarrow \text{SHM} & \epsilon_2 &= A \cos \omega t + B \sin \omega t \\ \text{Similarly,} & & &= C \cos(\omega t + \alpha) \\ \ddot{\epsilon}_3 &= -\omega^2 \epsilon_3 \rightarrow \epsilon_3 = D \cos(\omega t + \beta) \end{aligned}$$

(ϵ_2, ϵ_3 are bound by their initial deviation away from 0)

For $\Omega_2 = \text{const}$ solution

perturbation,

$$\ddot{\epsilon}_1 = +\omega^2 \epsilon_1$$

$$\ddot{\epsilon}_3 = +\omega^2 \epsilon_3$$

$$\rightarrow \epsilon_1(t) = A \underbrace{e^{\omega t}}_{\text{grows exponentially}} + B \underbrace{e^{-\omega t}}_{\text{damped exponentially}}$$

~~_____~~
 \rightarrow instability