

to leading order in Ω .

linear accidents,
of origin

[neglecting $\vec{r} \times (\vec{r} \times \vec{v})$
relative to $\vec{r} \times (\vec{v} \times \vec{r})$]

$$O(\Omega^2)$$

put

$$\frac{d\bar{v}}{dt}$$

[1st order OD L for \vec{v}]

↑
oth order

where $\frac{d\vec{v}_1}{dt} = \vec{g}$

ignore side
multiplied
by Ω

$$\rightarrow \vec{V}_2 = -(\vec{\omega} \times \vec{y})t^2 - 2(\vec{\omega} \times \vec{v}_0)t$$

$$\begin{aligned} \rightarrow \vec{v} &= \vec{g}t + \vec{v}_0 - (\vec{\Omega} \times \vec{g})t^2 - 2(\vec{\Omega} \times \vec{v}_0)t \\ &= \vec{v}_0 + \vec{g}t - 2(\vec{\Omega} \times \vec{v}_0)t - (\vec{\Omega} \times \vec{g})t^2 \end{aligned}$$

Integrate:

$$\boxed{\vec{r} = \vec{r}_0 + \vec{v}_0 t + \frac{1}{2}\vec{g}t^2 - (\vec{\Omega} \times \vec{v}_0)t^2 - \frac{1}{3}(\vec{\Omega} \times \vec{g})t^3}$$

standard solution.

Want to find deflection when body hits the ground.

ICS: $\vec{r}_0 = h \hat{z}$
 $\vec{v}_0 = 0$ [dropped]

$$\rightarrow \vec{r}(t) = h \hat{z} - \frac{1}{2}g \hat{z} t^2 - \frac{1}{3}(\vec{\Omega} \times \vec{g})t^3$$

0th order:

$$0 = h - \frac{1}{2}gT^2$$

$$T = \sqrt{\frac{2h}{g}}$$

Thus,
$$\begin{aligned} \vec{r}(T) &= \left(h - \frac{1}{2}gT^2 \right) \hat{z} - \frac{1}{3}(\vec{\Omega} \times \vec{g})T^3 \\ &= -\frac{1}{3} \left[-\Omega \sin \theta \hat{x} + \Omega \cos \theta \hat{z} \right] \times (-g \hat{z}) \left(\frac{2h}{g} \right)^{3/2} \\ &= \underbrace{+\frac{1}{3} \Omega g \sin \theta \left(\frac{2h}{g} \right)^{3/2}}_{\text{positive (deflected to the East)}} \hat{y} \end{aligned}$$

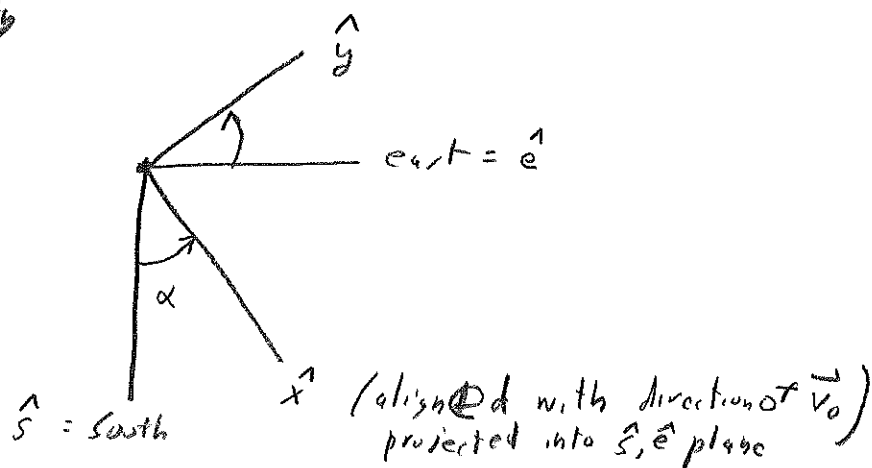
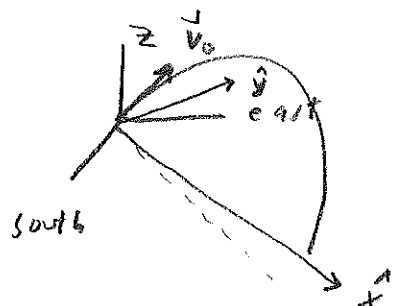
NOTE:

- 1) $h = 100 \text{ m}$, $\theta = \pi/2$ (equator) $\rightarrow \Delta y = \frac{1}{3} \left(\frac{2\pi}{1 \text{ day}} \right) \left(\frac{9.8 \text{ m}}{\text{s}^2} \right) \left(\frac{2 \cdot 100 \text{ m}}{9.8 \text{ m/s}^2} \right)^{3/2} = \frac{0.02 \text{ m}}{2 \text{ cm}}$
- 2) $h = 100 \text{ km} = 10^5 \text{ m}$, $\theta = \pi/2$ $\rightarrow \Delta y = 692 \text{ m} = \boxed{0.69 \text{ km}}$
edge of space

Sec 39, Prob 2:

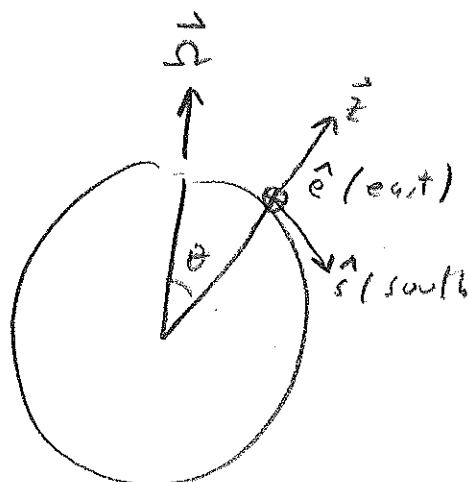
①

Deflection from coplanarity



$$\text{So } \hat{s} = \cos \alpha \hat{x} - \sin \alpha \hat{y}$$

$$\begin{aligned} \boxed{\vec{\Omega}} &= -\Omega \sin \theta \hat{s} + \Omega \cos \theta \hat{z} \\ &= -\Omega \sin \theta \cos \alpha \hat{x} + \Omega \sin \theta \sin \alpha \hat{y} + \Omega \cos \theta \hat{z} \\ \boxed{\vec{\Omega}} &= \Omega_x \hat{x} + \Omega_y \hat{y} + \Omega_z \hat{z} \end{aligned}$$



Using result of Sec 39, Prob 1:

$$\vec{r} = \vec{r}_0 + \vec{v}_0 t + \frac{1}{2} \vec{g} t^2 - (\vec{\Omega} \times \vec{v}_0) t^2 - \frac{1}{2} (\vec{\Omega} \times \vec{g}) t^3$$

I.C.'s:

$$\begin{aligned} \vec{r}_0 &= 0 \\ \vec{v}_0 &= v_{0x} \hat{x} + v_{0z} \hat{z} \end{aligned}$$

Thus,

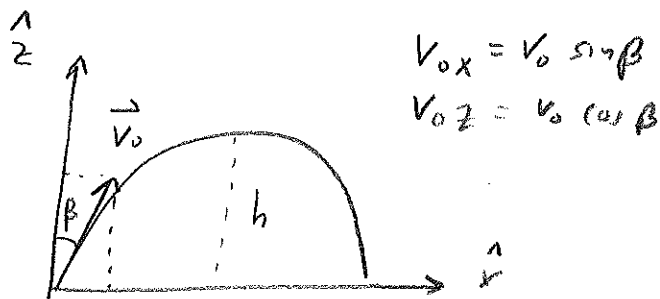
$$\vec{r} = \vec{v}_0 t - \frac{1}{2} g \hat{z} t^2 - (\vec{\Omega} \times \vec{v}_0) t^2 - \frac{1}{2} (\vec{\Omega} \times \vec{g}) t^3$$

0th order:

$$\vec{r} = \vec{v}_0 t - \frac{1}{2} g t^2 \hat{z}$$

$$\text{Max } h = \frac{1}{2} v_{0z}^2 \rightarrow \boxed{h = \frac{v_{0z}^2}{2g}}$$

~~$$h = \frac{v_{0x}^2}{2g}$$~~



Time to height h :

$$V_z = 0 = V_{0z} - g t$$

$$t = \frac{V_{0z}}{g}$$

Time to landing:

$$T = 2t = \boxed{\frac{2V_{0z}}{g}}$$

Range:

$$R = V_{0x} T$$

$$= V_0 \sin \beta \frac{2V_0 \cos \beta}{g}$$

$$= \frac{2V_0^2}{g} \sin \beta \cos \beta$$

$$= \frac{V_0^2}{g} \sin(2\beta)$$

max range for $\beta = \pi/4$

$$\vec{r} = (V_{0x} \hat{x} + V_{0z} \hat{z}) t - \frac{1}{2} g t^2 \hat{z} - (\vec{\Omega} \times \vec{v}_0) t^2 - \frac{1}{3} (\vec{\Omega} \times \vec{g}) t^3$$

$$\begin{aligned} \rightarrow \int \vec{r} &= -(\vec{\Omega} \times \vec{v}_0) T^2 - \frac{1}{3} (\vec{\Omega} \times \vec{g}) T^3 \\ &= -T^2 \left[(\vec{\Omega} \times \vec{v}_0) + \frac{1}{3} (\vec{\Omega} \times \vec{g}) T \right] \\ &= \frac{-4V_{0z}^2}{g^2} \left[(\vec{\Omega} \times \vec{v}_0) + \frac{2V_{0z}}{3g} (\vec{\Omega} \times \vec{g}) \right] \end{aligned}$$

$$\begin{aligned} \text{Now, } \vec{\Omega} \times \vec{v}_0 &= (\Omega_x \hat{x} + \Omega_y \hat{y} + \Omega_z \hat{z}) \times (V_{0x} \hat{x} + V_{0z} \hat{z}) \\ &= -\Omega_x V_{0z} \hat{y} - \Omega_y V_{0x} \hat{z} + \Omega_y V_{0z} \hat{x} + \Omega_z V_{0x} \hat{y} \\ &= (-\Omega_x V_{0z} + \Omega_z V_{0x}) \hat{y} + \underbrace{\Omega_y V_{0z} \hat{x} - \Omega_y V_{0x} \hat{z}}_{\text{ignore}} \end{aligned}$$

$$\begin{aligned} \vec{\Omega} \times \vec{g} &= (\Omega_x \hat{x} + \Omega_y \hat{y} + \Omega_z \hat{z}) \times (-g \hat{z}) \\ &= +\Omega_x g \hat{y} - \underbrace{g \Omega_y \hat{x}}_{\text{ignore}} \end{aligned}$$

$$\text{Thus, } \int \vec{r}_\perp = \frac{-4V_{0z}^2}{g^2} \left[-\Omega_x V_{0z} + \Omega_z V_{0x} + \frac{2}{3} \frac{V_{0z}}{g} \Omega_x g \right] \hat{y}$$

$L \perp$ to plane

$$\vec{r}_+ = \frac{4V_0^2}{g^2} \left[\frac{1}{3} \Omega_x V_0 z - \Omega_z V_0 x \right] \hat{y}$$

In terms of α, β, θ :

($\theta=0$, NP; $\theta=\frac{\pi}{2}$, equator)

$$V_{0z} = V_0 \cos \beta$$

($\beta=0$, vertical)

$$V_{0x} = V_0 \sin \beta$$

$$\Omega_x = -\Omega \sin \theta \cos \alpha$$

($\alpha=0$, south
 $\alpha=\frac{\pi}{2}$, east)

$$\Omega_z = \Omega \cos \theta$$

$$\vec{r}_+ = \frac{4V_0^2 \Omega \cos^2 \beta}{g^2} \left[-\frac{1}{3} \sin \theta \cos \alpha V_0 \cos \beta - \cos \theta V_0 \sin \beta \right] \hat{y}$$

$$= \frac{4V_0^3 \Omega \cos^2 \beta}{g^2} \left[-\frac{1}{3} \cos \alpha \sin \theta \cos \beta - \cos \theta \sin \beta \right] \hat{y}$$

~~Consider $\beta = \pi/4$~~

In terms of range: $R = \frac{2V_0^2}{g} \sin \beta \cos \beta$ (0th order)

$$\rightarrow V_0 = \sqrt{\frac{gR}{2 \sin \beta \cos \beta}}$$

$$V_0 = \sqrt{\frac{gR}{2 \sin \beta \cos \beta}}$$

$$\begin{aligned}
 \vec{r}_1 &= \frac{4\Omega \cos^2 \beta}{g^2} \left(\frac{gR}{2 \sin \beta \cos \beta} \right)^{3/2} \left[-\frac{1}{3} (\cos \alpha \sin \theta \cos \beta - \cos \theta \sin \beta) \right] \hat{y} \\
 &= \frac{4}{2\sqrt{2}} \Omega \sqrt{\frac{\cos \beta}{g}} \left(\frac{R}{\sin \beta} \right)^{3/2} \left[\quad \quad \quad \right] \hat{y} \\
 &= \Omega \sqrt{\frac{2 \cos \beta}{g}} \left(\frac{R}{\sin \beta} \right)^{3/2} \left[-\frac{1}{3} \cos \alpha \sin \theta \cos \beta - \cos \theta \sin \beta \right] \hat{y}
 \end{aligned}$$

Note: $\beta = \pi/4$ (which gives maximum zeroth order range)

$$\sin \frac{\pi}{4} = \cos \frac{\pi}{4} = \frac{\sqrt{2}}{2}$$

$$\begin{aligned}
 \vec{r}_1 &= \Omega \sqrt{\frac{2}{g}} R^{3/2} \frac{1}{(\sqrt{2}/2)} \left[-\frac{1}{3} \cos \alpha \sin \theta - \cos \theta \right] \left(\frac{\sqrt{2}}{2} \right) \hat{y} \\
 &= \Omega \sqrt{\frac{2}{g}} R^{3/2} \left[-\frac{1}{3} \cos \alpha \sin \theta - \cos \theta \right] \hat{y}
 \end{aligned}$$

NOTE: $-\frac{1}{3} \cos \alpha \sin \theta - \cos \theta = 0$

$$-\frac{1}{3} \cos \alpha \sin \theta = \cos \theta$$



~~$$\tan \theta = \frac{3}{\cos \alpha}$$~~

$$\cos \alpha = -\frac{3}{\tan \theta}$$



Thus, for a given θ can choose α so that trajectory stays in a plane.

$$\rightarrow \theta = \frac{\pi}{2} \text{ (equator)} \rightarrow \text{trajectory}$$

$$\sin \frac{\pi}{2} = 1, \cos \frac{\pi}{2} = 0 \rightarrow -\frac{1}{3} \cos \alpha = 0 \rightarrow \alpha = \pm \frac{\pi}{2} \text{ (aim East or West)}$$

Numerical values

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$$\beta = \pi/4 \quad \begin{array}{c} \nearrow \vec{v}_0 \\ \swarrow \end{array} \quad , \quad R = 100 \text{ m} , \quad g = 9.8 \frac{\text{m}}{\text{s}^2} , \quad \Omega = \frac{2\pi}{24 \cdot 3600 \text{ s}}$$

$$\delta \vec{r}_+ = \sqrt{\frac{2\Omega^2 R}{g}} \cdot R \left[-\frac{1}{3} \cos \alpha \sin \theta - \cos \theta \right] \hat{y}$$

$\theta = 0$: (North pole)

$$\delta \vec{r}_+ = -\sqrt{2} R \sqrt{\frac{\Omega^2 R}{g}} \hat{y} \quad \forall \alpha$$

$$= -0.033 \text{ m } \hat{y}$$

$$= \boxed{-3.3 \text{ cm } \hat{y}}$$

$\theta = \pi/4$:

$$\begin{aligned} \delta \vec{r}_+ &= \sqrt{\frac{2\Omega^2 R}{g}} \cdot R \frac{\sqrt{2}}{2} \left[-\frac{1}{3} \cos \alpha - 1 \right] \hat{y} \\ &= \sqrt{\frac{\Omega^2 R}{g}} \cdot R \left(-\frac{1}{3} \right) (3 + \cos \alpha) \hat{y} \leftarrow \text{always } < 0 \end{aligned}$$

$$= \begin{cases} -\frac{4}{3} R \sqrt{\frac{\Omega^2 R}{g}} \hat{y} = \boxed{-3.1 \text{ cm } \hat{y}} & (\alpha = 0) \text{ South} \\ -R \sqrt{\frac{\Omega^2 R}{g}} \hat{y} = \boxed{-2.3 \text{ cm } \hat{y}} & (\alpha = \pi/2) \text{ East} \\ -\frac{2}{3} R \sqrt{\frac{\Omega^2 R}{g}} \hat{y} = \boxed{-1.55 \text{ cm } \hat{y}} & (\alpha = \pi) \text{ North} \end{cases}$$

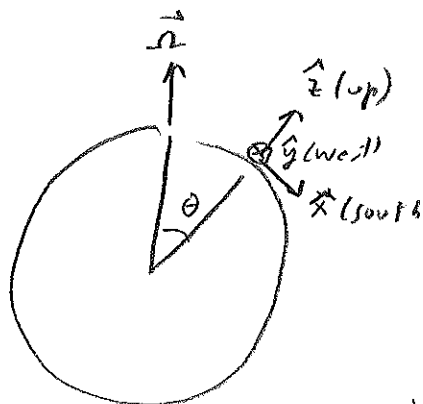
$\theta = \pi/2$: (equator)

$$\delta \vec{r}_+ = -\frac{1}{3} \sqrt{\frac{2\Omega^2 R}{g}} \cdot R \cos \alpha$$

$$= \begin{cases} -\frac{\sqrt{2}}{3} R \sqrt{\frac{\Omega^2 R}{g}} \hat{y} = \boxed{-1.1 \text{ cm}} & (\alpha = 0) \text{ South} \\ \boxed{0} & (\alpha = \pi/2) \text{ East} \\ +\frac{\sqrt{2}}{3} R \sqrt{\frac{\Omega^2 R}{g}} \hat{y} = \boxed{+1.1 \text{ cm}} & (\alpha = \pi) \text{ North} \end{cases}$$

Foucault's pendulum

Sec 39, Prob 3:



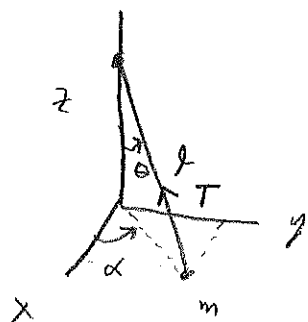
$$m \frac{d\vec{v}}{dt} = -\frac{\partial U}{\partial \vec{r}} = m \vec{W} = m \dot{\vec{R}} \times \vec{r} - 2m \vec{\Omega} \times \vec{v} - m \vec{\Omega} \times (\vec{\Omega} \times \vec{r})$$

$$= m \vec{g}_0 + \vec{T} - m \vec{\Omega} \times (\vec{\Omega} \times \vec{R}) - 2m \vec{\Omega} \times \vec{v} - m \vec{\Omega} \times (\vec{R} \times \vec{r})$$

$$= m \vec{g}_0 - m \vec{\Omega} \times (\vec{\Omega} \times (\vec{R} + \vec{r})) - 2m \vec{\Omega} \times \vec{v} + \vec{T}$$

$\underbrace{\hspace{10em}}_{m \vec{g}}$

Thus,
$$m \frac{d\vec{v}}{dt} = m \vec{g} + \vec{T} - 2m \vec{\Omega} \times \vec{v}$$

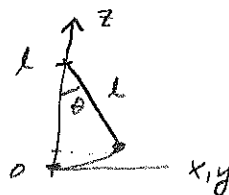


Small oscillations in xy plane:

$$z = l(1 - \cos \theta) \approx 0$$

$$x = l \sin \theta \cos \alpha \approx l \theta \cos \alpha$$

$$y = l \sin \theta \sin \alpha \approx l \theta \sin \alpha$$



Tension: \vec{T} :

$$T_z = T \cos \theta \approx T$$

$$T_x = -T \sin \theta \cos \alpha = -T \frac{x}{l}$$

$$T_y = -T \sin \theta \sin \alpha = -T \frac{y}{l}$$

thus,
$$\vec{T} \approx -T \frac{x}{l} \hat{x} - T \frac{y}{l} \hat{y} + T \hat{z}$$

so
$$m \frac{d\vec{v}}{dt} = -m g \hat{z} + \left(-T \frac{x}{l} \hat{x} - T \frac{y}{l} \hat{y} + T \hat{z} \right) - 2m \vec{\Omega} \times \vec{v}$$

$$\vec{\Omega} = -\Omega \sin \theta \hat{x} + \Omega \cos \theta \hat{z} \approx \Omega_x \hat{x} + \Omega_z \hat{z}$$

$$\vec{v} = v_x \hat{x} + v_y \hat{y} + v_z \hat{z} \approx v_x \hat{x} + v_y \hat{y} \quad (v_z \approx 0)$$

$$\rightarrow \vec{\Omega} \times \vec{v} = (\Omega_x \hat{x} + \Omega_z \hat{z}) \times (v_x \hat{x} + v_y \hat{y}) = \Omega_x v_y \hat{z} + \Omega_z v_x \hat{y} - \Omega_z v_y \hat{x}$$

W.r.to

$$V_x = \dot{x}, V_y = \dot{y}, V_z = \dot{z}$$

$$\frac{d\vec{v}}{dt} = \ddot{x}\hat{x} + \ddot{y}\hat{y} + \ddot{z}\hat{z}$$

then:

$$m\ddot{x} = -\frac{T}{l}x + 2m\Omega_z\dot{y} \quad (1)$$

$$m\ddot{y} = -\frac{T}{l}y - 2m\Omega_z\dot{x} \quad (2)$$

$$m\ddot{z} = -mg + T - 2m\Omega_x\dot{y} \quad (3)$$

For small oscillation in (x, y) plane $\ddot{z} \approx 0$, $T \approx$

$$0 \approx -mg + T - 2m\Omega_x\dot{y} \quad (3)$$

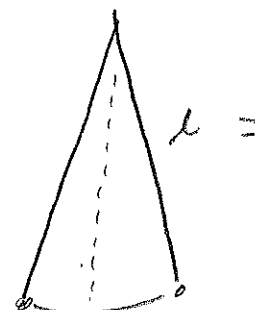
Compare g and $\Omega_x\dot{y}$:

$$\left| \frac{\Omega_x\dot{y}}{g} \right| = \left| \frac{\Omega \sin\theta}{g} \frac{\Delta y}{\Delta t} \right|$$

$$\lesssim \frac{\Omega}{g} \frac{D}{2\pi\sqrt{\frac{l}{g}}}$$

$$\lesssim \Omega \left(\frac{D}{l} \right) \sqrt{\frac{l}{g}}$$

$$\sim \frac{\Omega}{\sqrt{\frac{g}{l}}} \left(\frac{D}{l} \right)$$



period

$$T = \frac{2\pi}{\omega} = 2\pi\sqrt{\frac{l}{g}}$$

$$\Delta y = D \ll l$$

max displacement

Now: $D \ll l$ so $\frac{D}{l} \ll 1$

$$\sqrt{\frac{g}{l}} = \text{angular freq} = \frac{2\pi}{10^5 \text{ of second}}$$

$$\Omega = \frac{2\pi}{24 \times 60} = \frac{2\pi}{24 \times 3600} \approx \frac{2\pi}{100,000 \text{ sec}}$$

$$\begin{array}{r} 3600 \\ 24 \\ \hline 14400 \\ 7200 \\ \hline 86400 \end{array}$$

so

$$\frac{\Omega}{\sqrt{\frac{g}{l}}} \ll 1, \frac{D}{l} \ll 1 \rightarrow \left| \frac{\Omega_x\dot{y}}{g} \right| \ll 1$$

\rightarrow can ignore last term relative to mg

So:

$$0 \approx -mg + T \quad (3)$$

$$\rightarrow \boxed{T = mg}$$

~~→ m~~

Thus, $m \ddot{x} \approx -mg \frac{x}{l} + 2m\Omega_z \dot{y}$

$\ddot{x} \approx -\frac{g}{l}x + 2\Omega_z \dot{y}$	(1)
$\ddot{y} \approx -\frac{g}{l}y - 2\Omega_z \dot{x}$	(2)

Define: $z = x + iy$ [complex-valued]
 $\dot{z} = \dot{x} + i\dot{y}$
 $\ddot{z} = \ddot{x} + i\ddot{y}$

Add: (1) + i(2)

$$\ddot{x} + i\ddot{y} \approx -\frac{g}{l}(x + iy) + 2\Omega_z(\dot{y} - i\dot{x})$$

$$\ddot{z} = -\frac{g}{l}z - i2\Omega_z \dot{z}$$

$$\boxed{\ddot{z} + 2i\Omega_z \dot{z} + \frac{g}{l}z = 0}$$

Ansatz: $z = Ae^{i\lambda t}$
 $\dot{z} = Ae^{i\lambda t} i\lambda$
 $\ddot{z} = -Ae^{i\lambda t} \lambda^2$

$$\rightarrow Ae^{i\lambda t} \left[-\lambda^2 + 2i\Omega_z i\lambda + \frac{g}{l} \right] = 0$$

$$\lambda^2 + 2\lambda\Omega_z - \frac{g}{l} = 0$$

quadratic:

(4)

$$\lambda = \frac{-2\Omega_z \pm \sqrt{4\Omega_z^2 + 4\frac{g}{\lambda}}}{2}$$

$$= -\Omega_z \pm \sqrt{\Omega_z^2 + \frac{g}{\lambda}}$$

$$= -\Omega_z \pm \sqrt{\frac{g}{\lambda}} \sqrt{1 + \frac{\Omega_z^2}{(\frac{g}{\lambda})}}$$

$$\epsilon = \frac{\Omega_z^2}{(\frac{g}{\lambda})} \ll 1$$

$$\approx -\Omega_z \pm \sqrt{\frac{g}{\lambda}} \left(1 + \frac{1}{2} \frac{\Omega_z^2}{(\frac{g}{\lambda})} \right)$$

$$= \pm \sqrt{\frac{g}{\lambda}} - \Omega_z \pm \frac{1}{2} \frac{\Omega_z^2}{\sqrt{\frac{g}{\lambda}}}$$

2nd order in Ω

$$\approx \pm \sqrt{\frac{g}{\lambda}} - \Omega_z$$

Thus, $\zeta(t) = A e^{i(\sqrt{\frac{g}{\lambda}} - \Omega_z)t} + B e^{i(-\sqrt{\frac{g}{\lambda}} - \Omega_z)t}$

$$= e^{-i\Omega_z t} \left[A e^{i\sqrt{\frac{g}{\lambda}} t} + B e^{-i\sqrt{\frac{g}{\lambda}} t} \right]$$

~~$$+ i\gamma = \left[\cos(\Omega_z t) - i \sin(\Omega_z t) \right] \left\{ A \left(\cos\left(\sqrt{\frac{g}{\lambda}} t\right) \pm i \sin\left(\sqrt{\frac{g}{\lambda}} t\right) \right) + B \left(\cos\left(\sqrt{\frac{g}{\lambda}} t\right) - i \sin\left(\sqrt{\frac{g}{\lambda}} t\right) \right) \right\}$$~~

~~$$\zeta(t) = a e^{i(\sqrt{\frac{g}{\lambda}} - \Omega_z)t + \alpha} + b e^{i(-\sqrt{\frac{g}{\lambda}} - \Omega_z)t + \beta}$$

$$= a \left(\cos\left[\left(\sqrt{\frac{g}{\lambda}} - \Omega_z\right)t + \alpha\right] + i \sin\left[\left(\sqrt{\frac{g}{\lambda}} - \Omega_z\right)t + \alpha\right] \right)$$

$$+ b \left(\cos\left[\left(-\sqrt{\frac{g}{\lambda}} - \Omega_z\right)t + \beta\right] + i \sin\left[\left(-\sqrt{\frac{g}{\lambda}} - \Omega_z\right)t + \beta\right] \right)$$~~

$z = x + iy$ where

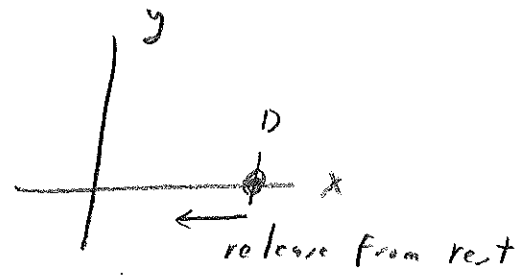
~~$$x(t) = a \cos\left[\left(\sqrt{\frac{g}{L}} \Omega_2\right)t + \alpha\right] + b \cos\left[\left(-\sqrt{\frac{g}{L}} \Omega_2\right)t + \beta\right]$$
$$y(t) = a \sin\left[\left(\sqrt{\frac{g}{L}} \Omega_2\right)t + \alpha\right] + b \sin\left[\left(-\sqrt{\frac{g}{L}} \Omega_2\right)t + \beta\right]$$~~

where a, b, α, β determined by $x(0), y(0), \dot{x}(0), \dot{y}(0)$:

$$z(t) = e^{-i\Omega_2 t} \left[A e^{i\sqrt{\frac{g}{L}} t} + B e^{-i\sqrt{\frac{g}{L}} t} \right]$$

I.C's

$$\begin{aligned} x(0) &= D \\ y(0) &= 0 \\ \dot{x}(0) &= 0 \\ \dot{y}(0) &= 0 \end{aligned}$$



$$z(0) = x(0) + iy(0) = D$$

$$\dot{z}(0) = \dot{x}(0) + i\dot{y}(0) = 0$$

Thus,

$$\begin{aligned} D &= A + B \\ 0 &= -i\Omega_2(A+B) + i\sqrt{\frac{g}{L}}A - i\sqrt{\frac{g}{L}}B \\ &= i \left[A(-\Omega_2 + \sqrt{\frac{g}{L}}) + B(-\Omega_2 - \sqrt{\frac{g}{L}}) \right] \end{aligned}$$

$$\rightarrow B = A \frac{(-\Omega_2 + \sqrt{\frac{g}{L}})}{(\Omega_2 + \sqrt{\frac{g}{L}})}$$

$$\rightarrow D = A \left(1 + \frac{(-\Omega_2 + \sqrt{\frac{g}{L}})}{(\Omega_2 + \sqrt{\frac{g}{L}})} \right) = A \frac{2\sqrt{\frac{g}{L}}}{\Omega_2 + \sqrt{\frac{g}{L}}}$$

Denote ~~$\frac{g}{2}$~~ $\sqrt{\frac{g}{2}} = \omega$

(6)

$$A = D \frac{(\Omega_z + \omega)}{2\omega}$$

$$B = A \frac{(-\Omega_z + \omega)}{(\Omega_z + \omega)} = \frac{D \cancel{(\Omega_z + \omega)}}{2\omega} \frac{(-\Omega_z + \omega)}{\cancel{(\Omega_z + \omega)}} = \frac{D(-\Omega_z + \omega)}{2\omega}$$

for

$$z(t) = e^{-i\Omega_z t} \left(\frac{D}{2\omega} \right) \left[(\Omega_z + \omega) e^{i\omega t} + (-\Omega_z + \omega) e^{-i\omega t} \right]$$

$$= \left(\frac{D}{2\omega} \right) e^{-i\Omega_z t} \left[\underbrace{\Omega_z (e^{i\omega t} - e^{-i\omega t})}_{2i \sin \omega t} + \underbrace{\omega (e^{i\omega t} + e^{-i\omega t})}_{2 \cos \omega t} \right]$$

$$= \left(\frac{D}{\omega} \right) e^{-i\Omega_z t} \left[\omega \cos \omega t + i\Omega_z \sin \omega t \right]$$

$$= \left(\frac{D}{\omega} \right) [\cos(\Omega_z t) - i \sin(\Omega_z t)] [\omega \cos \omega t + i\Omega_z \sin \omega t]$$

$$= \left(\frac{D}{\omega} \right) \left[\omega \cos() \cos() + \Omega_z \sin() \sin() \right]$$

$$- i \omega \sin(\Omega_z t) \cos \omega t + i \Omega_z \cos(\Omega_z t) \sin \omega t$$

$$= x + iy$$

$$x = D \left[\cos() \cos() + \frac{\Omega_z}{\omega} \sin() \sin() \right]$$

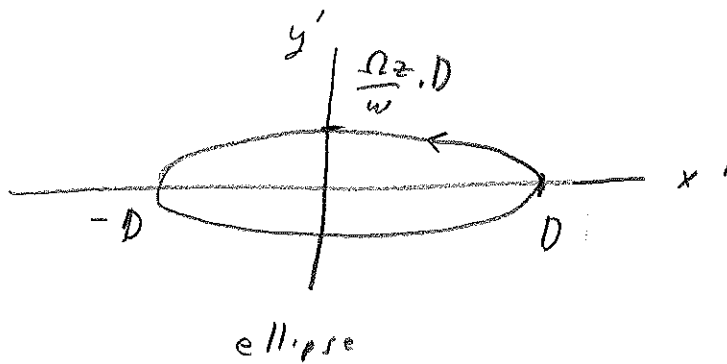
$$y = D \left[-\sin(\Omega_z t) \cos \omega t + \frac{\Omega_z}{\omega} \cos(\Omega_z t) \sin \omega t \right]$$

$$\begin{bmatrix} x \\ y \end{bmatrix} = \underbrace{\begin{bmatrix} \cos(\Omega_z t) & \sin(\Omega_z t) \\ -\sin(\Omega_z t) & \cos(\Omega_z t) \end{bmatrix}}_{\text{Rotation by } \Omega_z t} \begin{bmatrix} D \cos(\omega t) \\ D \frac{\Omega_z}{\omega} \sin(\omega t) \end{bmatrix}$$

In rotating frame:

$$\begin{bmatrix} x' \\ y' \end{bmatrix} = \begin{bmatrix} D \cos(\omega t) \\ D \frac{\Omega_z}{\omega} \sin(\omega t) \end{bmatrix}$$

$$\frac{\Omega_z}{\omega} = \frac{\Omega_{101} \theta}{\sqrt{\frac{g}{L}}} \ll 1$$



Precession period:

$$\begin{aligned} T_{\text{precession}} &= \frac{2\pi}{\Omega_z} = \frac{2\pi}{\Omega_{101} \theta} \\ &= \frac{2\pi}{\left(\frac{2\pi}{1 \text{ day}}\right) \cos \theta} \\ &= \frac{1 \text{ day}}{\cos \theta} \end{aligned}$$

NOTE: $\theta = 0$ (NP) $\rightarrow T_{\text{precession}} = 1 \text{ day}$
 $\theta = \frac{\pi}{2}$ (equator) $\rightarrow T_{\text{precession}} = \infty$