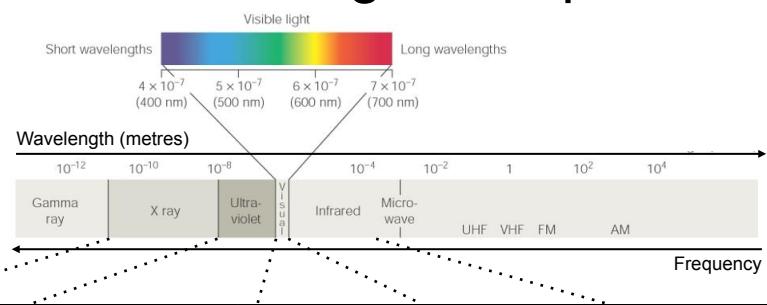
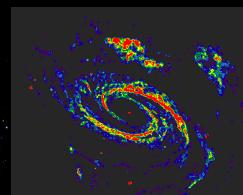


# The Earth's Atmosphere and Astronomical Observations

## The Electromagnetic Spectrum



M81 – “grand design” spiral galaxy



X-ray:  
Black hole  
accretion disks

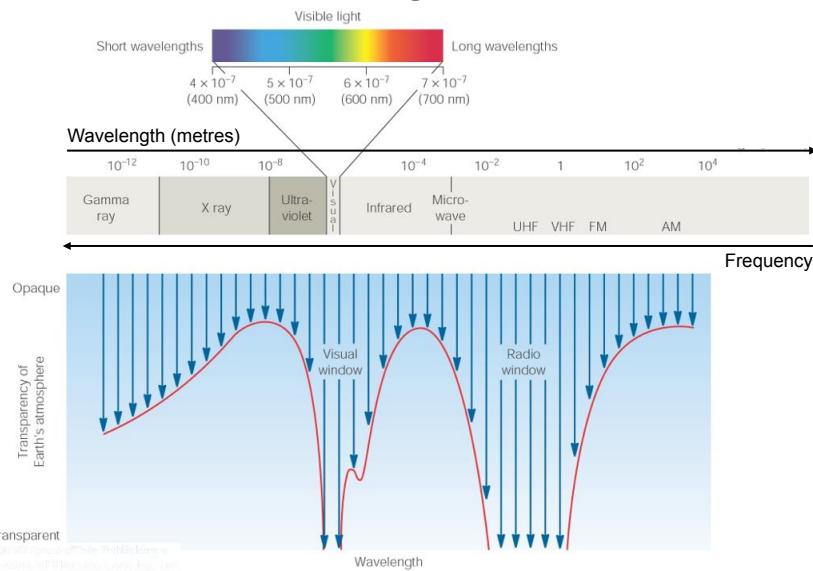
Ultraviolet:  
 $10^6$ - $10^8$ yr-old  
stars

Optical:  
 $>10^9$ yr-old  
stars

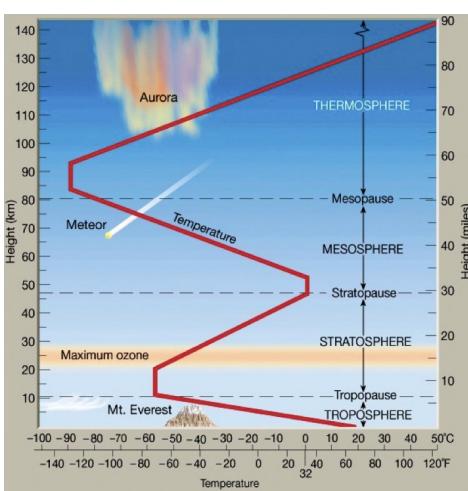
Infrared:  
Cold dust  
 $T < 10^3$ K

Radio:  
Magnetically  
accelerated  
electrons

# The Electromagnetic Spectrum



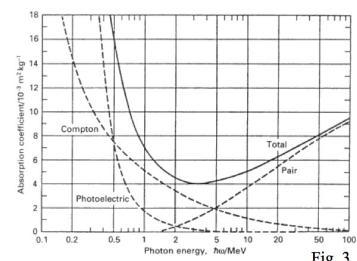
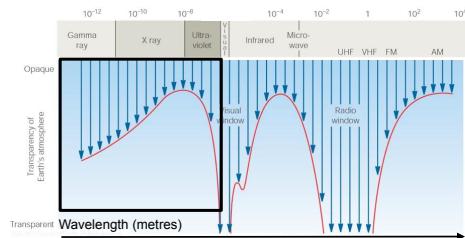
# Structure of the Atmosphere



- Space observatories at  $H>200\text{km}$ 
  - E.g. Hubble, Chandra
- Ionosphere ( $H>\sim 70\text{km}$ ):
  - Reflects radio waves at  $\lambda>10\text{m}$
- Thermosphere:
  - Atomic oxygen and nitrogen absorb UV and X-ray photons
- Stratosphere:
  - Contains Ozone which causes temperature to increase
  - Clouds form at inversion layer ( $dT/dH>0$ )
- Troposphere:
  - Contains 75% of atmospheric gases
  - $\text{O}_2$  and  $\text{CO}_2$  molecules absorb infrared photons
- Highest ground-based observatories at  $H\sim 4\text{km}$ 
  - E.g. Mauna Kea

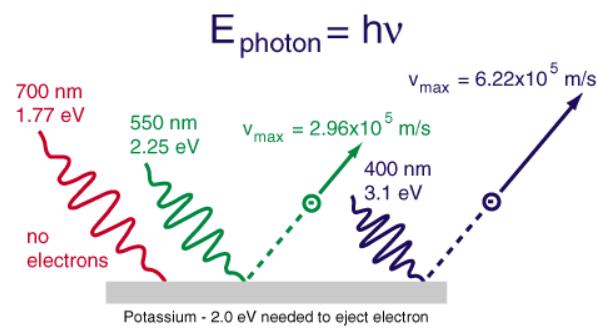
# Atmospheric Attenuation

- Ultraviolet, X-ray, Gamma-ray photons are absorbed very efficiently by atomic oxygen and nitrogen
  - Photo-ionization (photoelectric effect) at:  $\sim 4\text{eV} < E < \sim 0.511\text{MeV}$
  - Compton scattering dominates at:  $\sim 0.511\text{MeV} < E < \sim 5\text{MeV}$
  - Pair production starts at:  $E > 1.022\text{MeV}$



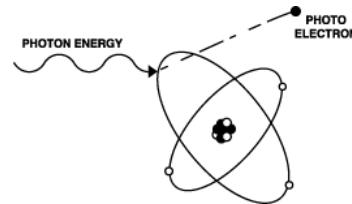
# Photoelectric effect

- Light can make metals emit electrons
- Not all photons can do it – need minimum energy
- But higher-energy light does not produce more electrons
- Instead, they produce faster electrons



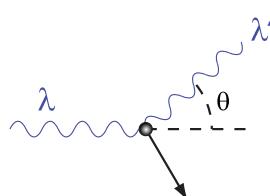
## Photoelectric effect

- Einstein's explanation:
  - Assume light is a stream of photons, each with energy  $h\nu$
  - Assume electrons are bound with some energy  $\phi$  (work function)
  - No electrons emitted if  $h\nu < h\nu_0 = \phi$
- A light packet with energy  $h\nu$  will
  - spend  $\phi$  in releasing the electron and
  - the rest of the energy will go into KE
  - i.e.  $h\nu = KE + \phi$
- Cut-off frequency is that which "ejects" an electron with KE=0:
  - $\nu_0 = \phi / h$



## Compton Scattering

- Photons of energy comparable to the rest mass energy of the electron (0.511MeV) scatter off electrons bound to the atoms, causing electrons to recoil, but not be removed from the atom

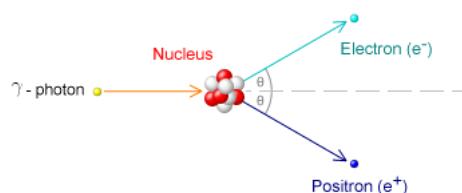


$$\lambda' - \lambda = (h/m_e c) (1 - \cos \theta)$$

h is Planck's constant  
 m<sub>e</sub> is electron mass  
 c is speed of light

## Pair Production

- Photons of energy exceeding twice the rest mass energy of the electron ( $E > 1.022\text{MeV}$ ) can interact with atomic nuclei and produce an electron-positron pair
- To conserve energy and momentum pair production can only take place via interaction of the photon with a nucleus (or other suitable particle)



## Atmospheric Attenuation

- Infrared/millimetre photons ( $10^{-6}\text{m} < \lambda < 10^{-3}\text{m}$ ):
  - Have insufficient energy to ionize atomic Nitrogen and Oxygen
  - Quantum energies that match the energies separating quantum states of molecular vibrations (infrared) and rotations (millimetre)
- Main species responsible for IR attenuation are  $\text{H}_2\text{O}$  and  $\text{CO}_2$

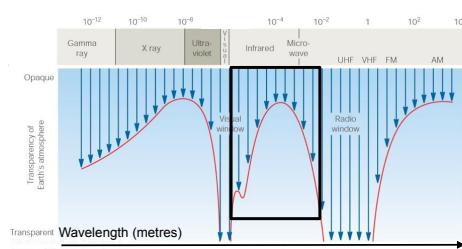


Figure 2.1 from Glass

# Vibrations of simple molecules

## Fundamental or normal modes

$\nu_1$  Symmetric stretch

$\nu_2$  Bend (Scissoring)

$\nu_3$  Asymmetric stretch

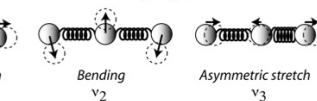
A normal mode is IR-active if the dipole moment changes during mode motion.

Overtones, combinations and differences of fundamental vibrations are also possible (e.g.,  $2\nu_1$ ,  $\nu_1+\nu_3$  etc.)

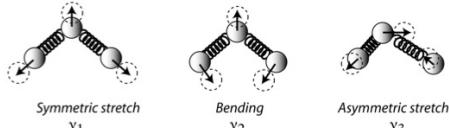
Diatom (N<sub>2</sub>, O<sub>2</sub>, CO)



Linear triatomic (CO<sub>2</sub>, N<sub>2</sub>O)



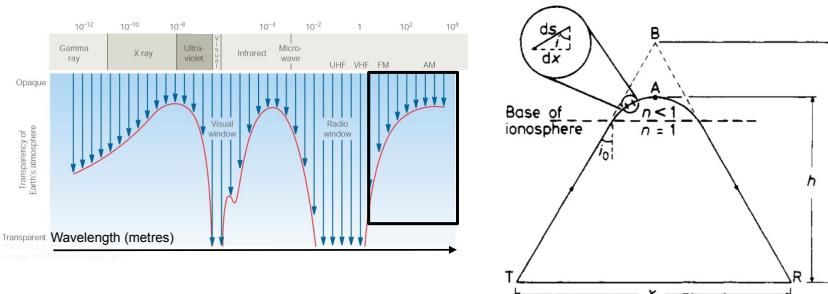
Nonlinear Triatomic (H<sub>2</sub>O, O<sub>3</sub>)



A non-linear molecule of N atoms has  $3N-6$  normal modes of vibration; a linear molecule has  $3N-5$ .

# Atmospheric Attenuation

- Radio waves ( $10^{-6}$ m  $<$   $\lambda$   $<$   $10^{-3}$ m):
  - Transparent window: 3MHz – 0.3THz;  $10^2$ m –  $10^3$ m
  - Upper frequency limit imposed by H<sub>2</sub>O and CO<sub>2</sub> in atmosphere
  - Lower frequency limit imposed by plasma frequency of the ionosphere
  - At lower frequencies the radio waves are reflected by ionosphere



## Physics of the Ionosphere

- UV radiation from the Sun dissociates and ionizes nitrogen and oxygen in the upper atmosphere, generating the ionosphere
- The ionosphere extends from ~60km to 300km above the earth's surface: 1 in  $10^3$  particles is ionized in this region
- Free electrons in the ionosphere refract radio waves and can cause them to be reflected
- Ionospheric reflection of radio waves is responsible for the atmosphere being opaque to low frequency radio waves
- This is helpful for global radio communication and frustrating for radio astronomers!

## Physics of the Ionosphere

- Consider an electromagnetic wave of frequency  $f$  traveling through an ionized medium
- The electric field vector of the wave varies with time as follows:  

$$E = E_0 \sin(2\pi f t)$$
- If this field is in (say) the x-direction then an electron will experience a force  $-eE$  and its equation of motion will be:  

$$m_e d^2x/dt^2 = -eE = -eE_0 \sin(2\pi f t)$$
- Integrate to give:  

$$x = [eE_0 \sin(2\pi f t)] / [4\pi^2 f^2 m_e]$$
- Which then gives:  

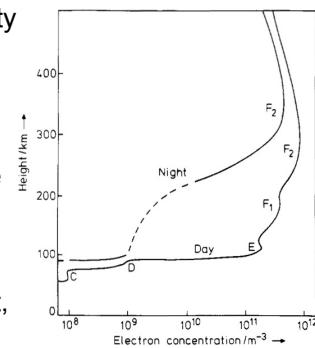
$$d^2x/dt^2 = -(4\pi^2 f^2) x$$
- The free electrons therefore vibrate with simple harmonic motion

## Physics of the Ionosphere

- The oscillation of the electrons can be thought of as inducing a net polarization of the ionosphere
- The induced polarization oscillates in anti-phase with the electric field, and can be written as:  
 $P = -N e x$   
 N is the number density of electrons  
 P is the dipole moment per unit volume
- Then:  $P = -(N e^2 E) / (4 \pi^2 f^2 m_e)$
- From dielectric theory we know that:  $P = \epsilon_0 (\epsilon_r - 1) E$   
 $\epsilon_0$  is the permittivity of free space  
 $\epsilon_r$  is the relative permittivity of the dielectric (i.e. the ionosphere)

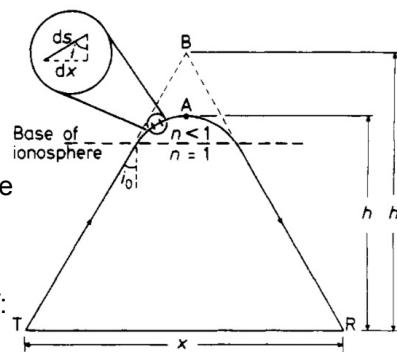
## Physics of the Ionosphere

- We can therefore write:  $\epsilon_r - 1 = -(N e^2) / (4 \pi^2 f^2 m_e \epsilon_0)$
- In vacuo the speed of light is:  $c = (\epsilon_0 \mu_0)^{-0.5}$   
 $\mu_0$  is the free space permeability
- In an isotropic medium of relative permeability  $\mu_r$  the speed of light becomes:  
 $(\mu_0 \mu_r \epsilon_0 \epsilon_r)^{-0.5} = c / (\mu_r \epsilon_r)^{0.5} = c / n$   
 n is the refractive index of the medium
- Taking  $\mu_r=1$  for the ionosphere, the refractive index of the ionosphere is:  
 $n = (\epsilon_r)^{0.5} = [1 - (N e^2) / (4 \pi^2 f^2 m_e \epsilon_0)]^{0.5}$
- N, the electron density, increases with height, so n, refractive index, decreases with height



## Physics of the Ionosphere

- A pulse of radio waves strikes the ionosphere at an angle  $i_0$
- The decreasing refractive index bends the path and at point "A" the angle of refraction becomes  $90^\circ$
- Adopting  $i_0=0$  and applying Snell's law gives:  $\sin(i_0) = n_A = 0$
- The pulse of radio waves is therefore reflected if:  $(N e^2) / (4 \pi^2 f^2 m_e \epsilon_0) = 1$
- Which implies a critical frequency of:  $f = [(N e^2) / (4 \pi^2 m_e \epsilon_0)]^{0.5}$
- Typically  $N \sim 6 \times 10^{10} \text{ m}^{-3}$  which renders the atmosphere opaque at frequencies of  $f < \sim 2.2 \text{ MHz}$



## Optical Extinction

- The atmosphere is not 100% transparent to optical radiation
- Rayleigh scattering off small molecules:
  - Wavelength small compared to size of molecule,  $\lambda < \sim 0.1 \text{ r}$
  - Air molecules, e.g. O<sub>2</sub>, CO<sub>2</sub>, H<sub>2</sub>O, ...
  - Scattered intensity is strongly wavelength dependent:  $I \propto \lambda^{-4}$
- Mie scattering off larger molecules:
  - Wavelength comparable with size of molecule,  $\lambda > \sim 0.1 \text{ r}$
  - Aerosols, e.g. volcanic dust, car exhaust fumes, cigarette smoke
  - Scattered intensity is almost wavelength independent

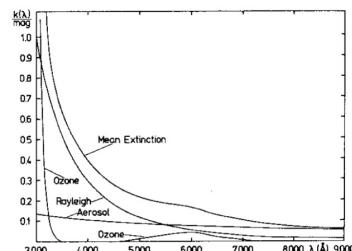
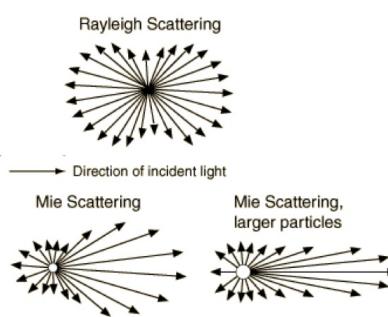
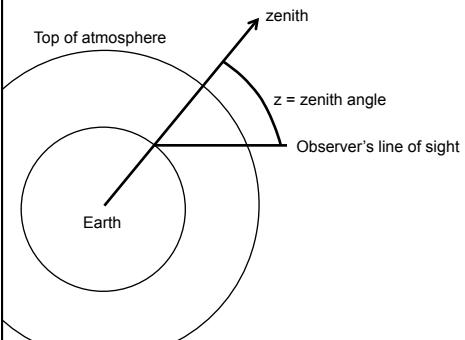


Fig. 1. Mean vertical extinction at Flagstaff, Arizona, in May-June 1976. The assumed ozone and Rayleigh contributions are shown separately



## Airmass

- Extinction depends on the position on the sky that is being observed



Define  $X$  to be the path length through the atmosphere

Define  $X=1$  at  $z=0$  (i.e. at zenith)

At  $z>0$ :  $\cos z = 1 / X$

i.e.:  $X = \sec z$

Assumes plane parallel atmosphere; a good approximation for  $X \leq 2$

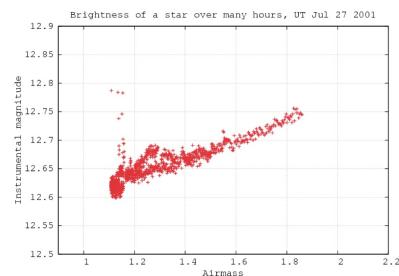
If optical depth depends only on path length, then:

$$\tau(X) = \tau_0 X$$

$$F_{\text{observed}} = F_{\text{top}} \exp(-\tau_0 X)$$

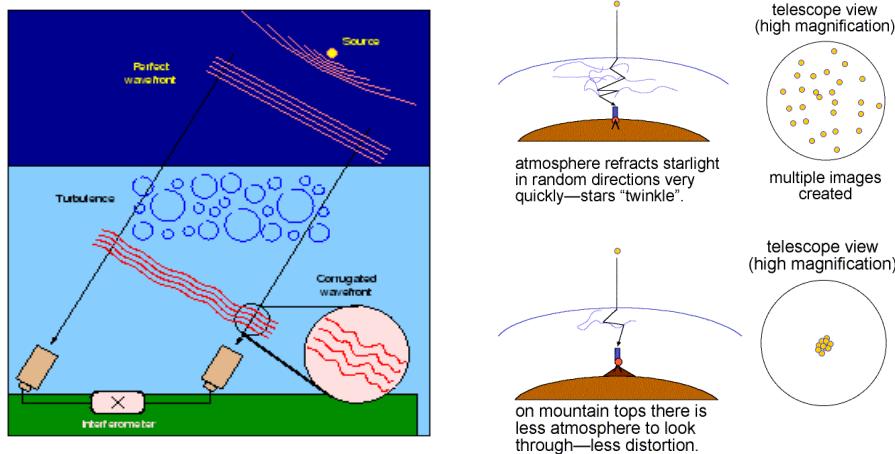
## Extinction Coefficients

- Re-writing the previous expression in terms of magnitudes gives observed magnitudes in terms of the exoatmospheric magnitude ( $m_0$ ), airmass, and the optical depth at zenith:
- $m = -2.5 \log_{10}(F_{\text{top}} \exp(-\tau_0 X)) = m_0 + 1.086 \tau_0 X$
- Extinction makes objects fainter, so their magnitude increases
- Extinction coefficient is then defined as:  $k = 1.086 \tau_0$
- $m_0$  and  $k$  can be measured by fitting a straight line to  $m$  as a function of  $X$



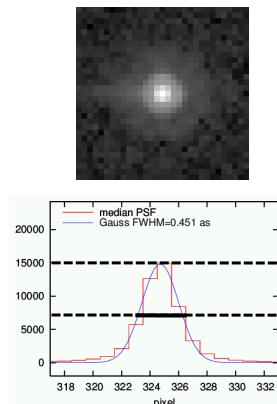
## Atmospheric Turbulence

- Turbulence causes the refractive index of the atmosphere to change randomly and rapidly – this blurs our view of stars and galaxies



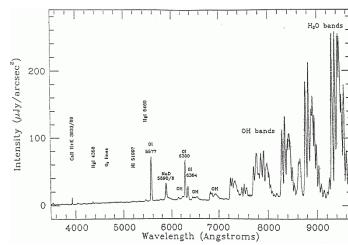
## “Seeing”

- In practice seeing blurs our view of stars and galaxies
- A star is blurred from a point to a smooth light distribution
- Astronomers measure seeing by fitting (e.g.) a Gaussian to the measured light distribution
- Seeing is defined as the full width at half maximum
- Typically FWHM = 1arcsec at 550nm
- It can be shown from Kolmogorov theory of turbulence that seeing has a gentle wavelength dependence: FWHM  $\propto \lambda^{-0.2}$



# Atmospheric Emission

- The atmosphere also emits radiation at optical wavelengths:
  - Rayleigh scattered moonlight and starlight
  - Emission from excited atoms and molecules in upper atmosphere (airglow)
- Airglow:
  - dominated by the hydroxyl radical OH<sup>-</sup>
  - Increases rapidly in strength beyond 700nm
- The sky brightness is a strong function of wavelength:



Filter	$\lambda(\mu\text{m})$	mag/arcsec <sup>2</sup>	photons/cm <sup>2</sup> /s/um/arcsec <sup>2</sup>
U	0.36	21.6	1.74x10e-2
B	0.44	22.3	1.76x10e-2
V	0.55	21.1	3.62x10e-2
R	0.64	20.3	5.50x10e-2
I	0.79	19.2	1.02x10e-1
J	1.23	14.8	2.49
H	1.66	13.4	4.20
K	2.22	12.6	3.98