

Electronics

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3rd February 2012

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Part I

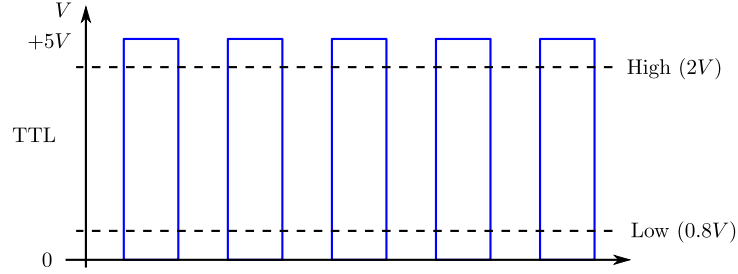
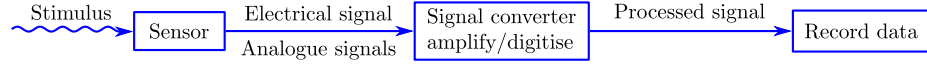
Overview

1 Electronics

Sensor: Device that measures a physical quantity and converts it into a signal, which can be detected and eventually read by a user, e.g. LDR, photomultiplier, thermistor.

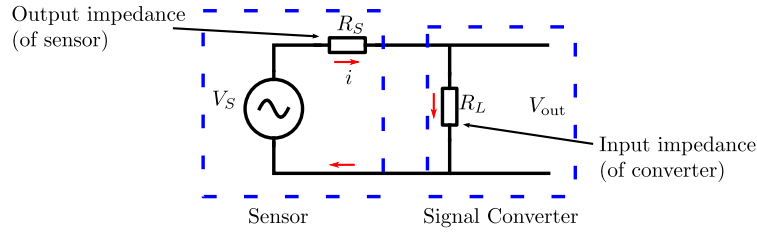
There are two types of transducer. First is a sensor, second, and actuator which is a device which converts a signal into some physical effect. For example, a microphone is a sensor, and a loudspeaker is an actuator.

1.1 Archetypal System



1.2 Sensor

A sensor can be represented by Thevenin's theorem. Thevenin's theorem says that any sensor can be represented by a voltage source and a source impedance.



1.3 Impedance Matching

The sensor signal which is measured is V_{out} (not V_S). Current flowing in the circuit is given by,

$$i = \frac{V_S}{R_S + R_L}$$

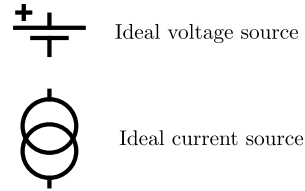
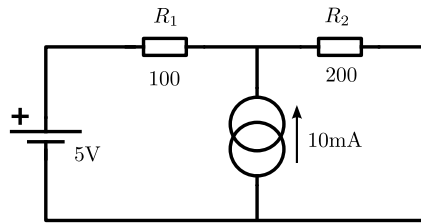
Therefore, the observed signal across the load R_L is

$$\begin{aligned} V_{out} &= iR_L \\ &= V_S \left(\frac{R_L}{R_S + R_L} \right) \end{aligned}$$

If $R_L \ll R_S$, then $V_{out} \ll V_S$. But if $R_L \approx R_S$, then $V_{out} \approx \frac{1}{2}V_S$

Ex

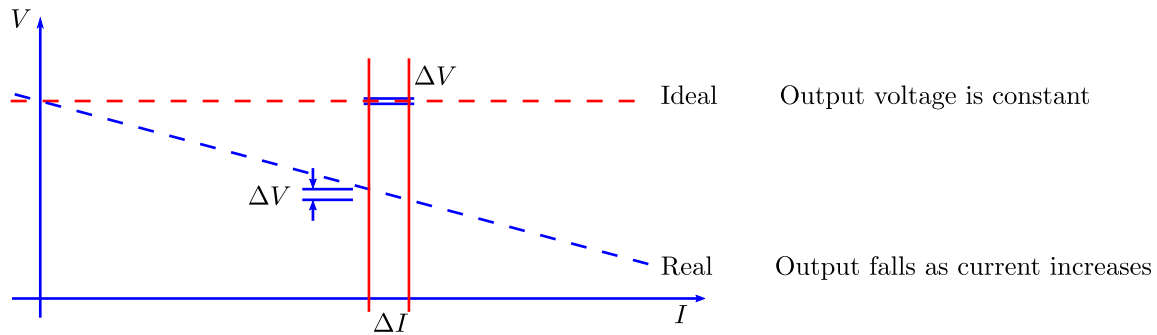
Consider the following circuit,



What is the voltage across R_2 ?

Superposition Consider the effect of each source separately then add the two effects.

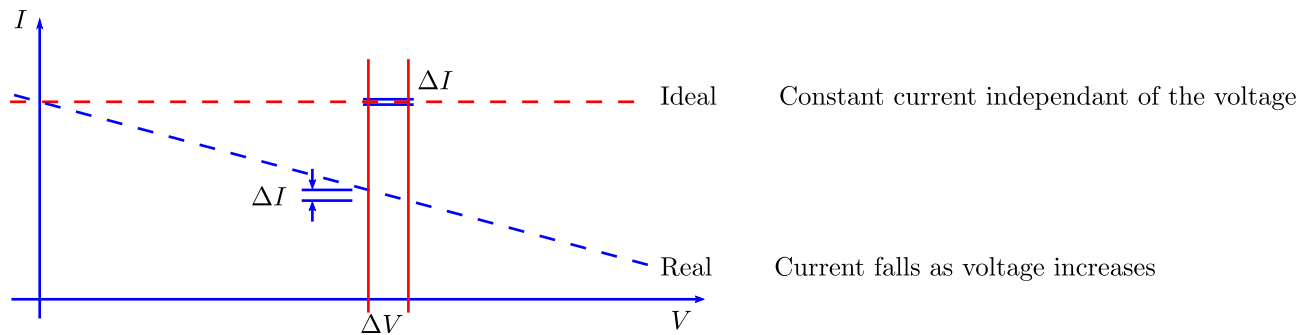
1.3.1 Voltage Source



Effective resistance, $R = \frac{\Delta V}{\Delta I}$. For the ideal source, $R_s = 0$ (i.e. short circuit)

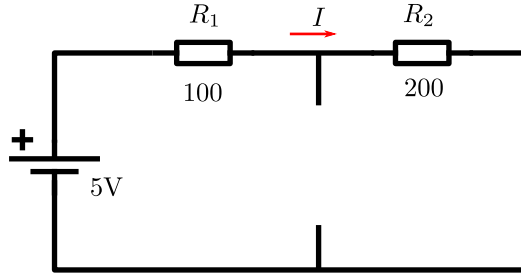
\therefore Consider the ideal voltage source to have zero internal resistance.

1.3.2 Current Source

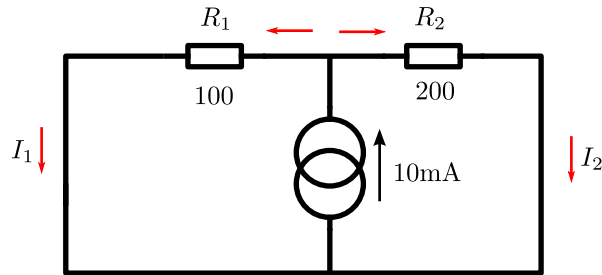


Effective resistance, $R = \frac{\Delta V}{\Delta I}$ For the ideal source, $R \rightarrow \infty$. \therefore Take the current source as having infinite internal resistance (i.e. open circuit).

1. Consider the effect of just the voltage source by switching the current source off, so $10\text{mA} \rightarrow 0$ and is replaced by an open circuit.



2. Consider the current source only by switching off the voltage source, so $5\text{V} \rightarrow 0$ and is replaced by a short circuit.



Using the superposition principle, we can say that the current through R_2 is the sum of these two,

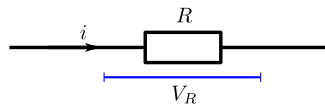
$$\begin{aligned} I &= \frac{5}{300} + \frac{10}{3} \times 10^{-3} \\ &= \frac{5}{300} + \frac{1}{50} = \frac{6}{300} = \frac{1}{50} \text{ Amps} \end{aligned}$$

\therefore Voltage across $R_2 = \frac{1}{50} \times 200 = 4\text{V}$.

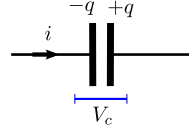
1.4 Review of Circuit Components

The most important three components that a general circuit can be composed of are the resistor, capacitor and the inductor.

1. **Resistance** By Ohm's law, $V_R = iR$



2. **Capacitance** $V_c = \frac{q}{C}$, where C is the capacitance given by $C = \frac{\epsilon_0 A}{d}$, where A is the area of the capacitor and d is the separation of the plates.



Now

$$q = \int i dt$$

$$\therefore V_c = \frac{\int i dt}{C} \quad \text{OR} \quad \frac{dV_c}{dt} = \frac{i}{C}$$

$$i = C \frac{dV_c}{dt}$$

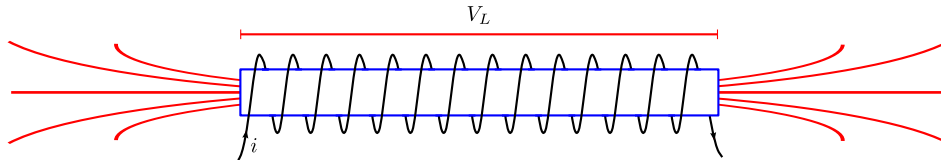
When V_c is constant, $i = 0$. Therefore, there is no current flows through a capacitor when a DC voltage is applied. Therefore, under DC conditions, the capacitor acts as an open circuit. But current does flow through a capacitor when the applied voltage across it changes,

$$\therefore i = C \frac{dV_c}{dt}$$

When $\frac{dV_c}{dt}$ is large, the current is large. For an infinitesimally small time interval, Δt , ΔV will be very small as well, but the ratio of $\frac{\Delta V}{\Delta t}$ may be significant. This implies then that the current, i may be appreciable. But as usual, the effective resistance, $R = \frac{\Delta V}{\Delta i}$.

For the infinitesimal changes, ΔV may be small and Δi may be large, such that $R \rightarrow 0$. Therefore a capacitor will act as a short circuit when instantaneous changes occur.

3. **Inductance "Choke"** An inductor is a coil of thick wire (so that the resistance, ideally, is zero) wrapped around a metal core.



If the current is constant, then $V_L = 0$. If the current changes, the magnetic field changes as well, so, by Faraday's Law, there is an induced electromotive force (back emf),

$$\epsilon = -\frac{d\phi}{dt}$$

where ϕ is the flux threading the coil. But $\phi \propto i$, which means

$$\phi = Li$$

$$\therefore \epsilon = -L \frac{di}{dt}$$

- (a) If $i = \text{constant}$, then $V = 0$. So the inductor acts as a short circuit under DC conditions.
 (b) For infinitesimally small changes, Δt will be small and Δi will be small also. But again, the ratio, $\frac{\Delta t}{\Delta i}$, may be significant, and so ΔL may be significant. The effective resistance is given by

$$R = \frac{\Delta V}{\Delta i}$$

$$\approx \frac{\text{finite}}{\text{zero}} \rightarrow \infty$$

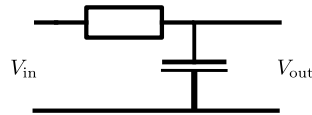
Therefore for instantaneous changes, the inductor behaves like an open circuit.

To summarise,

	R	C	L
Steady Conditions (DC)	R	open circuit	short circuit
Instantaneous Changes	R	short circuit	open circuit

2 Time Response

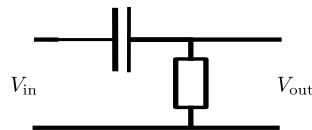
Consider the time response of an RC circuit. How does the circuit respond to DC conditions and instantaneous changes?



1. **DC** so the capacitor acts as an open circuit. $\frac{V_{out}}{V_{in}} = 1$
2. **Instantaneous changes** so the capacitor now acts as a short circuit. $\frac{V_{out}}{V_{in}} = 0$

Low frequency signals pass through unimpeded, whereas high frequency signals are blocked. This setup, then, acts as a "low pass" filter.

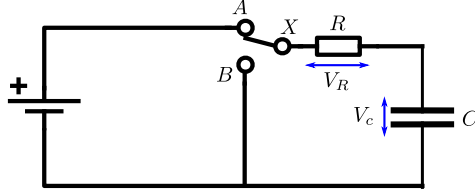
Consider the same circuit, but with the components swapped.



1. **DC**, the capacitor is still an open circuit, but now $\frac{V_{out}}{V_{in}} = 0$
2. **Instantaneous changes**, $\frac{V_{out}}{V_{in}} = 1$

Now, the low frequency signals are blocked and high frequency signals are unaffected leading to the "high pass" filter.

2.1 Quantitative Time Response of RC Circuit



Initially, the switch is at A, so $q = 0$, $V_R = 0$ and $V_c = 0$. At time $t = 0$, the switch is moved from A to B. At this instant, the capacitor C acts like a short circuit, and so a current flows,

$$i = \frac{V}{R}$$

And then, by Ohm's law, $V_R = V$ and $V_c = 0$. After a long time (under dc conditions) the capacitor acts like an open circuit, and so the current must fall to zero.

By Kirchhoff's laws,

$$\begin{aligned} V &= V_R + V_c \\ &= iR + q \frac{1}{C} \end{aligned}$$

Differentiating this gives

$$\begin{aligned} 0 &= R \frac{di}{dt} + \frac{1}{C} \frac{dq}{dt} \\ &= R \frac{di}{dt} + \frac{1}{C} i \\ \Rightarrow \int \frac{di}{i} &= \int \frac{1}{RC} dt \\ i &= \frac{V}{R} e^{-\frac{t}{RC}} \end{aligned}$$

This gives the exponential falloff as expected, and now,

$$\begin{aligned} \Rightarrow V_R &= iR = V e^{-\frac{t}{RC}} \\ \Rightarrow V_c &= V - V_R = V \left(1 - e^{-\frac{t}{RC}} \right) \end{aligned}$$

At time $t = 0$, the switch is moved from A to B and current flows through the components.

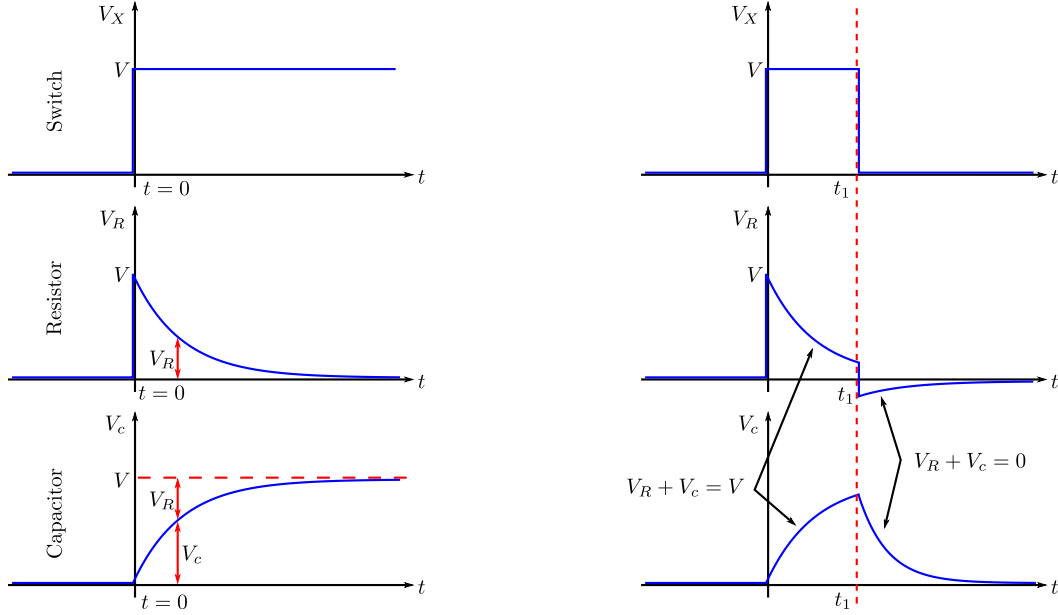
At time $t = t_1$, the switch is moved back from B to A, so the applied voltage falls back to zero. At this instant, no charge moves and therefore, V_c is unchanged. At $t \leq t_1$,

$$\begin{aligned} i &= \frac{V}{R} e^{-\frac{t_1}{RC}} \\ V_R &= V e^{-\frac{t_1}{RC}} \\ V_c &= V \left(1 - e^{-\frac{t_1}{RC}} \right) \end{aligned}$$

At $t \geq t_1$,

$$\begin{aligned} 0 &= V_R + V_c \\ V_R &= -V_c \\ V_R &= -V \left(1 - e^{-\frac{t_1}{RC}}\right) \end{aligned}$$

At $t = t_1$, V_R changes from $V_R = V \left(1 - e^{-\frac{t_1}{RC}}\right)$ to $V_R = -V \left(1 - e^{-\frac{t_1}{RC}}\right)$. So the change in $V_R = -V$.



When the switch is closed, the full change in voltage appears across the resistance.

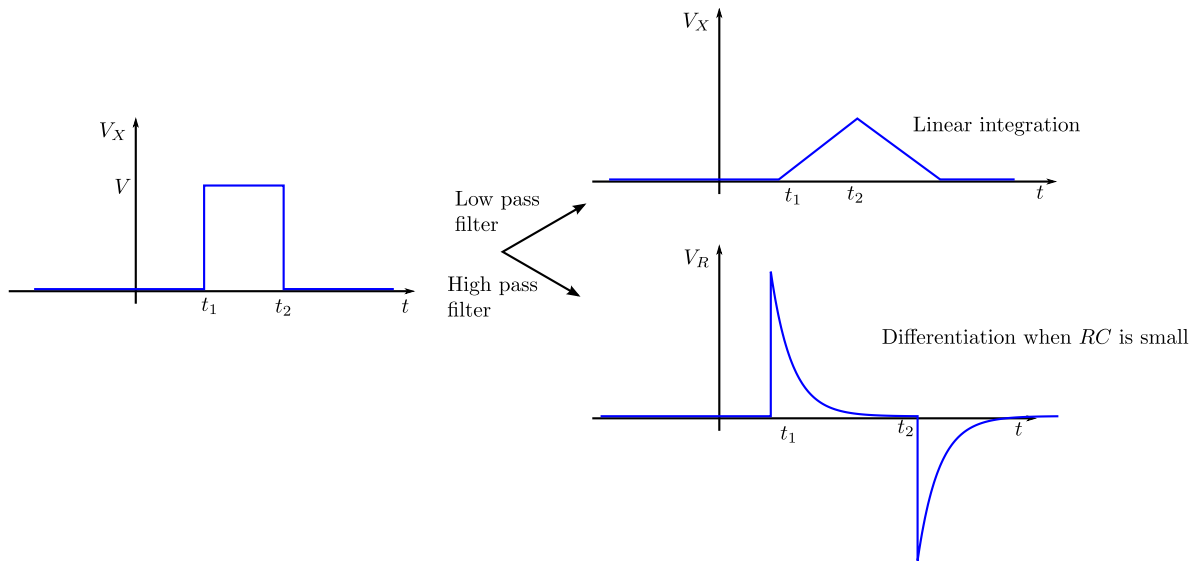
The low pass filter can be regarded then, as an integrator.

$$V_{\text{out}} = V_c = V \left(1 - e^{-\frac{t}{RC}}\right)$$

For small times, $t \ll RC$,

$$V_{\text{out}} = V \left(1 - \left(1 - \frac{t}{RC}\right)\right) = V \frac{t}{RC}$$

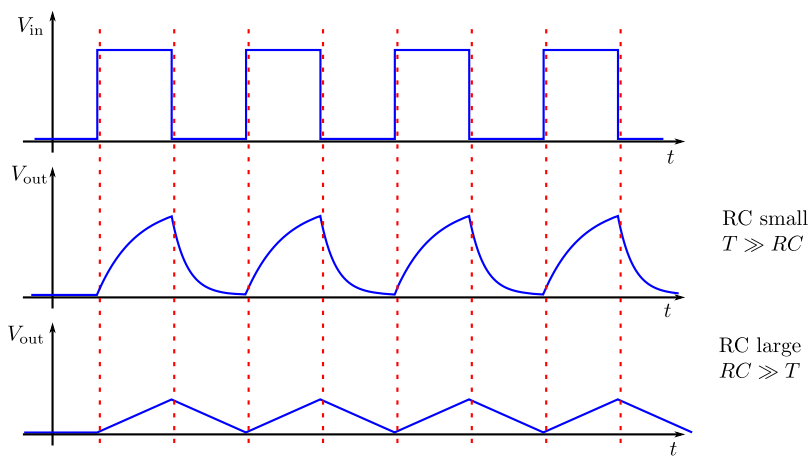
Therefore, V_{out} is approximately the integral of V_{in} with respect to t . Similarly, the high pass filter can be regarded as a differentiator.



2.2 Response to Pulses

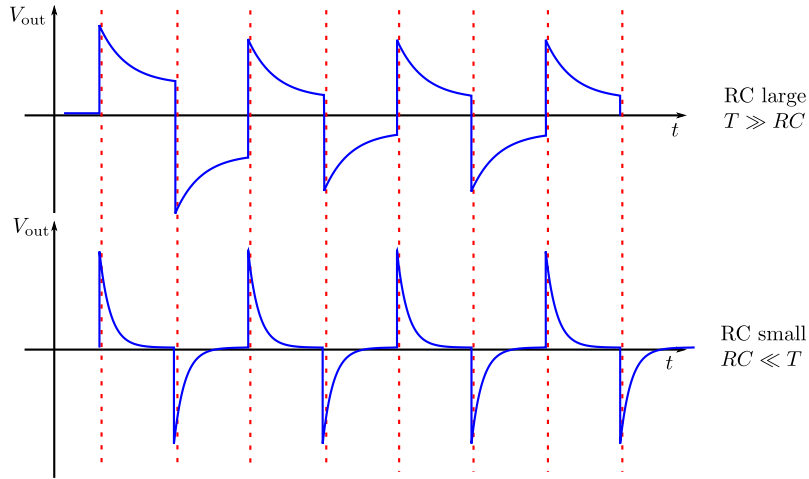
Apply now a train of pulses of time period T .

2.2.1 Low Pass Filter



2.2.2 High Pass Filter

When the same input pulses are input to V_{in} .



2.3 Transfer Functions

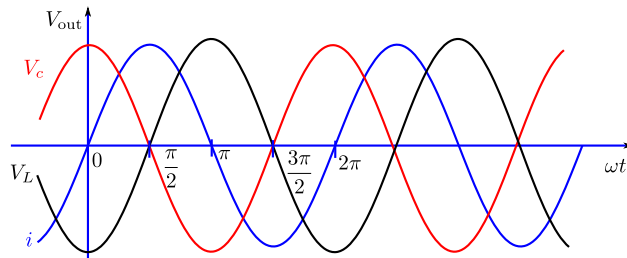
$$V_{\text{out}} = T \times V_{\text{in}}$$

Assume that the current is sinusoidal, such that,

$$i = I \sin(\omega t)$$

$$\begin{aligned} V_R &= I_0 R \sin(\omega t) & V_C &= -\frac{I}{\omega C} \cos(\omega t) & V_L &= I_0 \omega L \cos(\omega t) \\ & & &= \frac{I}{\omega C} \sin\left(\omega t - \frac{\pi}{2}\right) & &= I_0 \omega L \sin\left(\omega t + \frac{\pi}{2}\right) \end{aligned}$$

2.4 Phase Differences



V_C reaches its peak at a later time to i meaning V_C lags behind i by $\frac{\pi}{2}$. V_L reaches its peak at an earlier time, so V_L leads ahead of i by $\frac{\pi}{2}$. There is a dependence of the amplitude on frequency, ω , and there is a phase difference.

2.5 Move to Complex Notation

Now the current is given by

$$i = I_0 e^{j\omega t} \quad j = \sqrt{-1}$$

Also

$$\begin{aligned} V_R &= I_0 R e^{j\omega t} = iR \\ V_c &= \frac{I_0}{j\omega C} e^{j\omega t} = \frac{i}{j\omega C} = \frac{i}{\omega C} e^{-j\frac{\pi}{2}} \\ V_L &= LI_0 e^{j\omega t} = j\omega Li = i\omega L e^{j\frac{\pi}{2}} \end{aligned}$$

$e^{\pm j\frac{\pi}{2}}$ represents a phase lag and lead respectively of 90° .

Resistance is measured in ohms (Ω) and is real. Impedance is also measured in ohms (Ω) but is complex.

	R	C	L
Impedance, Z	$Z_R = R$	$Z_c = \frac{1}{j\omega C}$	$Z_L = j\omega L$

So in a circuit, just substitute

$$\begin{aligned} C &\rightarrow \frac{1}{j\omega C} \\ L &\rightarrow j\omega L \end{aligned}$$

Then apply Ohm's, Kirchhoff's laws etc.

2.6 RC Series Circuit as a Single Time Constant Filter

$$V_0 = TV_1$$

Potential divider,

$$\begin{aligned} T &= \frac{V_0}{V_1} = \frac{Z_c}{Z_c + Z_R} \\ &= \frac{\frac{1}{j\omega C}}{R + \frac{1}{j\omega C}} = \frac{1}{1 + j\omega CR} \end{aligned}$$

This is the voltage divider theorem. Any complex number can be expressed as an amplitude and a phase.

Complex Number: $T = a + jb \rightarrow re^{j\phi} \rightarrow r \cos \phi + jr \sin \phi$

$$\left. \begin{aligned} a &= r \cos \phi \\ b &= r \sin \phi \end{aligned} \right\} r = (a^2 + b^2)^{\frac{1}{2}}$$

$$\tan \phi = \frac{\sin \phi}{\cos \phi} = \frac{a}{b}$$

Transfer function (complex number) \rightarrow amplitude + phase

Ex

$$\begin{aligned} T &= \frac{1 - j\omega CR}{(1 - j\omega CR)(1 + j\omega CR)} \\ &= \frac{1 - j\omega CR}{1 + (\omega CR)^2} = Ae^{i\phi} \end{aligned}$$

where A is the amplitude and ϕ is the phase.

$$\begin{aligned} A &= \left[\left(\frac{1}{1 + (\omega CR)^2} \right)^2 + \left(\frac{\omega CR}{1 + (\omega CR)^2} \right)^2 \right]^{\frac{1}{2}} \\ &= \frac{1}{(1 + (\omega CR)^2)^{\frac{1}{2}}} \\ \tan \phi &= -\omega CR \end{aligned}$$

These data could be plotted linearly, but it is much more useful and easier to interpret if they are plotted on a double log plot, called a Bode plot.

2.7 Bode Plots

Bode plots exploit the complex impedance and involves simply plotting $\log_{10}(|T|)$ against $\log_{10}(\omega)$ or ϕ against $\log_{10}(\omega)$. The transfer function is defined as the output voltage divided by the input voltage, and is also complex.

When plotted in this way, the Bode plot is almost exactly two straight lines. There is a small amount of error at the connection of these lines since there should be a small amount of curvature here.

2.7.1 Decibels

Decibels are used to represent power ratios between two signals. For voltages,

$$\text{dB} = 20 \log_{10} \left(\frac{V_2}{V_1} \right)$$

and for powers,

$$\text{dB} = 10 \log_{10} \left(\frac{P_2}{P_1} \right)$$

since

$$P = \frac{V^2}{R}$$

$\frac{V_2}{V_1}$	dB
1000	60
100	40
10	20
$-\sqrt{2}$	3.01
1	0
$\frac{1}{\sqrt{2}}$	-3.01
0.1	-20
0.01	-40
0.001	-60

When $|T| = 20 \log_{10} \left(\frac{V_{out}}{V_{in}} \right) = -3\text{dB}$. This represents when the power has been reduced by a factor of 2. This occurs when $\omega = \omega_0$ and so ω_0 is called the half power frequency, or -3dB frequency.

Ex

High pass filter,

- Express the transfer function as a complex number,

$$\begin{aligned} T &= \frac{V_0}{V_1} = \frac{iR}{V_i} = \frac{R}{V_i} \frac{V_i}{R + \frac{1}{j\omega C}} \\ &= \frac{R}{R + \frac{1}{j\omega C}} = \frac{Rj\omega C}{1 + Rj\omega C} \end{aligned}$$

when $\omega \rightarrow \infty$, $T \rightarrow 1$ and when $\omega \rightarrow 0$, $T \rightarrow 0$, which is the usual high pass filter results, as expected.

- Express T as an amplitude, $|T|$, and a phase, ϕ .

$$\begin{aligned} T &= \frac{Rj\omega C}{1 + Rj\omega C} \times \frac{1 - Rj\omega C}{1 - Rj\omega C} \\ &= \frac{j\omega C R + \omega^2 C^2 R^2}{1 + \omega^2 C^2 R^2} \\ \Rightarrow |T| &= \frac{((\omega^2 C^2 R^2)^2 + \omega^2 C^2 R^2)^{\frac{1}{2}}}{1 + \omega^2 C^2 R^2} \\ &= \frac{\omega C R}{(1 + \omega^2 C^2 R^2)^{\frac{1}{2}}} \end{aligned}$$

So as $\omega \rightarrow \infty$, $|T| \rightarrow 1$, as $\omega \rightarrow 0$, $|T| \rightarrow 0$ and as $\omega \rightarrow \omega_0$, $|T| \rightarrow \frac{1}{\sqrt{2}}$ (-3dB) Phase difference between V_0 and V_1 is ϕ ,

$$\begin{aligned} \phi &= \tan^{-1} \left(\frac{\Im\{T\}}{\Re\{T\}} \right) \\ &= \tan^{-1} \left(\frac{\omega C R}{\omega^2 C^2 R^2} \right) = \tan^{-1} \left(\frac{1}{\omega C R} \right) \end{aligned}$$

when $\omega \rightarrow \infty$, $\phi \rightarrow 0$, $\omega \rightarrow \omega_0 = \frac{1}{CR}$, $\phi \rightarrow 45^\circ$ and $\omega \rightarrow 0$, $\phi \rightarrow 90^\circ$

Using these Bode plots, the more complex functions can be analysed simply by adding the phase shifts and amplitudes.

2.7.2 Cascading Transfer Functions

V_2 is both the output of T_1 and the input of T_2 . If $T = T_1 \times T_2$, then

$$\log_{10} T = \log_{10} T_1 + \log_{10} T_2$$

Therefore, to multiply the transfer functions, add the decibels of $|T_1|$ and $|T_2|$.

2.7.3 Phases

$$\begin{aligned}T_1 &= a + jb \\T_2 &= p + jq\end{aligned}$$

If

$$\begin{aligned}T &= T_1 T_2 \\&= (a + jb)(p + jq) \\&= (ap - bq) + j(bp + aq) \\ \Rightarrow \phi &= \tan^{-1} \left(\frac{bp + aq}{ap - bq} \right)\end{aligned}$$

But

$$\begin{aligned}\tan(\phi_1 + \phi_2) &= \frac{\tan \phi_1 + \tan \phi_2}{1 - \tan \phi_1 \tan \phi_2} \\&= \frac{b/a - q/p}{b/a \times q/p} \\ \phi &= \phi_1 + \phi_2\end{aligned}$$

When the transfer functions are multiplied, the phases are added together.

Ex Example of cascading transfer functions

The circuit is an effective high pass filter, but at low frequencies, the phase difference is -180° so the output is a negative multiple of the input.

2.7.4 Generalised Function

Any transfer function, $T(\omega)$ can be expressed as,

$$T = K \frac{\left(1 + \frac{j\omega}{Z_1}\right) + \left(1 + \frac{j\omega}{Z_n}\right) + \dots + \left(1 + \frac{j\omega}{Z_n}\right)}{\left(1 + \frac{j\omega}{P_1}\right) + \left(1 + \frac{j\omega}{P_2}\right) + \dots + \left(1 + \frac{j\omega}{P_n}\right)}$$

where K is a multiplicative constant, Z_n are transfer-function zeroes and P_n are transfer-function poles. To draw the Bode plot we draw the asymptotes for each zero and each pole and then add the various contributions.

- At every value of $\omega = Z_i$, the slope of the line increases by 20dB per decade and at every value of $\omega = P_i$ the slope of the line decreases by 20dB per decade.
- To correct a straight line Bode plot, at every zero, put a point 3dB above the line and at every pole put a point 3dB below the line.
- If K is positive, then the zero slope start line will be at 0 degrees, whereas if K is negative, then the start line will be at 180 degrees.

- At every value of $\omega = Z_i$ increase the slope of the line by 45 degrees per decade, beginning one decade below $\omega = Z_i$ and ending one decade above $\omega = Z_i$.
- At every value of $\omega = P_i$ decrease the slope of the line by 45 degrees per decade, beginning one decade before $\omega = P_i$ and ending one decade above $\omega = P_i$.

Part II

Operational Amplifiers

3 Op-amp Characteristics

The op-amp is a versatile building block in analogue electronic circuits. It was originally used to perform mathematical operations in analogue computers but are now included in a huge number of different applications because they are so cheap, easy to incorporate and perform very well under a vast range of conditions.

An IC op-amp (e.g. 741) consists of around 24 transistors, several resistors and a capacitor. We will not study transistors, so will treat the op-amp as a black box single unit.

3.0.5 Op-Amp Terminal

The op-amp has two input and one output terminals. It senses the difference between the voltages at its inputs, $V_+ - V_-$, and then multiplies this by a gain, A . The output, then becomes

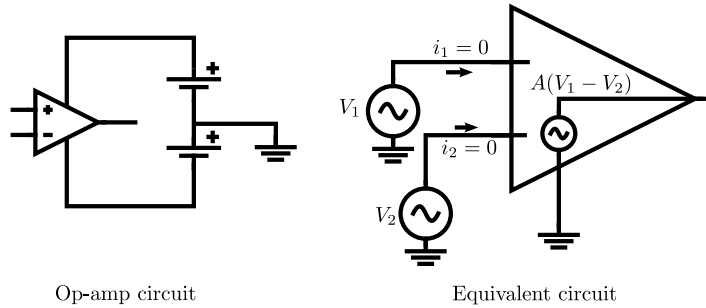
$$v_0 = a(v_+ - v_-)$$

All of the voltages are measured with respect to ground.

3.0.6 Ideal Properties

- The op-amp will be considered as an ideal component, meaning it draws no current through any of the input terminals, i.e. it has an infinite input impedance.
- The output is an ideal voltage source so has zero output impedance.
- It has infinite differential gain ($A \rightarrow \infty$) at all frequencies in the range $0 \leq \omega \leq \infty$.

Zero output impedance means that any amount of current can be taken without changing the voltage. This current cannot come from the inputs as they draw no current. Most IC op-amps require dc power to be supplied (not usually shown). Note that the $+$ and $-$ supplies have nothing to do with the $+$ and $-$ inputs.



The op-amp is a differential-input, single-ended-output and is only sensitive to voltage difference, any common signal is ignored. This property is called common-mode rejection. An ideal op-amp has infinite common-mode rejection. Op-amps are direct-coupled devices (amplifies signals whose frequency is as low as $\omega = 0$). The gain A is called the differential gain, or open-loop gain which is ideally infinite. They are never (rarely) used without feedback. Though this reduces the gain, it provides other benefits.

4 Op-Amp Golden Rules

- The open loop gain (A) is so high that a fraction of a milli-volt difference between the inputs will cause the output to reach maximum value (e.g. $\pm 15V$ of the power supply) So if the output signal is in the required range (i.e. $\pm 15V$), then it implies that the two inputs must be at the same value.

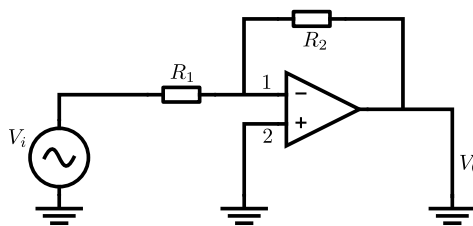
This then means that the voltage difference between the inputs is effectively zero.

- The op amp draws negligible current into the inputs, so no current flows into the device whatever voltage is drawn.

5 Applications

The op-amp is invariably used with feedback, in particular, negative feedback.

5.1 The Inverting Configuration



The resistor R_2 provides negative feedback. The output is fed back to the inverting terminal. The closed loop gain, G , of this configuration can be calculated as follows,

$$G = \frac{V_0}{V_i}$$

From the op-amp equation,

$$V_2 - V_1 = \frac{V_0}{A} \approx 0$$

The current in R_1 is,

$$i_1 = \frac{V_i - V_1}{R_1} \approx \frac{V_i}{R_1}$$

But the op-amp draws no current, so $i_2 = i_1$

$$\begin{aligned} V_0 &= 0 - i_1 R_2 \\ \Rightarrow G &= \frac{V_0}{V_i} = -\frac{R_2}{R_1} \end{aligned}$$

The inverting amplifier has an input impedance given by

$$Z_i = \frac{v_i}{i_1} = \frac{v_i}{v_i/R_1} = R_1$$

and an output impedance of

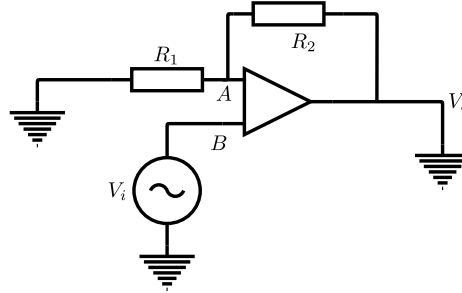
$$Z_0 = 0$$

The op-amp acts like an ideal voltage source.

5.1.1 Inverting Op-Amp Observations

- The feedback resistor R_2 closes the loop around the op-amp, hence G is known as the closed loop gain.
- The closed-loop gain is dependent on R_1 and R_2 only
- Gain can be chosen, and made as accurate as you like. The closed-loop gain is smaller than the open-loop gain, but is stable and predictable
- A consequence of the large (ideally infinite) open-loop gain is that the two input terminals track each other in potential.
- We speak of a virtual short-circuit between the input terminals. Don't make the mistake of actually shorting them though.
- This means that because the non-inverting terminal is grounded, the inverting terminal is a virtual ground.

5.2 The Non-Inverting Configuration



We can find the closed loop gain (or transfer function) G ,

From the op-amp equation

$$V_2 - V_1 = \frac{V_0}{A} \approx 0 \quad A \rightarrow \infty, V_1 \rightarrow V_2$$

Current in R_1

$$i_1 = \frac{0 - V_i}{R_1} = -\frac{V_i}{R_1} \quad V_2 = V_1 \rightarrow V_1 \rightarrow V_i$$

The op-amp draws no current, so

$$i_2 = i_1 \quad V_0 = V_i - i_1 R_2$$

$$\Rightarrow G = \frac{V_0}{V_i} = 1 + \frac{R_2}{R_1}$$

The input impedance is given by

$$Z_i = \frac{V_i}{i} = \frac{V_i}{0} = \infty$$

So the op-amp draws no current. The output impedance is given by

$$Z_0 = 0$$

So the op-amp acts like an ideal voltage source.

5.2.1 Non-Inverting Op-Amp Observations

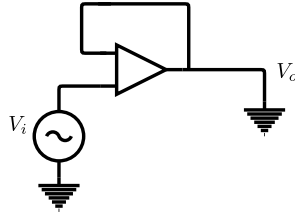
- Once again, the closed loop gain is dependant only on R_1 and R_2 , not on the properties of the op-amp itself.
- The closed loop gain is positive, hence non-inverting, so the output is in phase with the input.
- A consequence of the large (ideally infinite) open-loop gain is that the two terminals track each other in potential, as before.

This means that the non-inverting terminal is at the same potential as the inverting terminal, which is the applied voltage.

- Note that the inverting terminal cannot be connected directly to the ground, and this also assumes that the op-amp is working, i.e. the output is not saturated.

5.3 Unity Gain Buffer

The unity gain buffer, or voltage follower, has unity gain, so does not seem to be useful. Though it has a range of applications since it changes impedance.



Let $R_1 \rightarrow \infty$, so

$$G = \frac{V_0}{V_1} = 1 + \frac{R_2}{R_1} \rightarrow 1$$

This means that the input impedance of the ideal unity-buffer $Z_i = \infty$ so does not load the source. This is particularly useful for weak signals. Also $Z_o = 0$ so provides a lot of current without attenuating the signal.