

An IFS study of Low-Powered Radio Galaxies



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Chapter 1

Introduction

For millennia, Man has turned his gaze skyward in awe and wonder at the heavens and wondered at the processes that drive structure on the colossal scale of galaxies and galaxy clusters. As physicists we conduct experiments to better our understanding of the processes and laws that control the universe; an old joke goes that as astronomers, our experiment has already been conducted for us, it's called The Big Bang: our job is to use it to test the theories and laws of physics. Indeed, astronomy is a unique opportunity to test our knowledge at scales and limits that are not attainable here on Earth.


The most successful cosmological description of the universe is the dark energy (represented by Λ), cold dark matter (LCDM) coupled with inflation (Guth, 1981). This describes how the current state of the universe (69% Dark Energy (currently best described with a scalar field), 26% cold dark matter (a matter-like species which only interacts gravitationally), $\sim 5\%$ ordinary baryonic matter and $< 0.003\%$ radiation, complete with deep gravitational potential wells in which galaxies and galaxy clusters have formed) has evolved from an infinite and random spacetime 'foam'. Quantum fluctuations within this 'foam' cause regions to form which have the correct conditions for inflation to occur: this includes a quantum scalar field with a sufficient energy density to dominate the region (Linde, 1982; Albrecht & Steinhardt, 1982). This field causes the exponential expansion of space at superluminal speeds. Quantum fluctuations are rapidly expanded to beyond the scale of the horizon (distance that a massless particle can travel within the age of the universe). At this point they become frozen in as perturbations to the density profile of the universe (any perturbations prior to the start of inflation are stretched so thin that they can be considered completely negligible). Once the energy density of the inflationary field becomes too low, inflation ends, radiation dominates and the expansion slows to subluminal speeds. As the horizon expands at the speed of light, the density perturbations re-enter the horizon

(smallest first), providing seed for the gravitational collapse of dark matter. This process has hardwired hierarchical growth of structure (where the smallest structures, small galaxies, form first, which then coalesce into large galaxies, which form groups, clusters and eventually super-clusters) into the cosmos.

Having said that, we have said nothing the materials that we can directly observe: baryonic matter and radiation. Baryonic matter is far more complicated than dark matter and thus when we observe galaxies we see far more structure than simply smooth halos which it is believed that dark matter mostly exists as. The large plethora of shapes of galaxies are summed up in the Hubble Diagram (or tuning-fork), where galaxies are believed to begin life on one of the two prongs, working their way up towards the handle as the spiral arms get tighter, until a major merger removes all ordered motion to leave a very round, random motion (velocity dispersion) dominated elliptical galaxy 1.1 (note Hubble originally assumed evolution to occur in the opposite direction on the Hubble diagram). However, recent studies have shown that we cannot completely classify a galaxy by its morphology only (Cappellari et al., 2011; Sanchez et al., 2011; Arnold et al., 2013; Bryant et al., 2014) (though some broad trends with morphology do exist). A more physical representation might be by color: Baldry et al. (2004) showed that galaxies exist within two distinct regions (known as the blue cloud and red sequence) on a color–magnitude plot. The length of a star’s life is inversely proportional to it’s mass, which in turn is also followed by it’s temperature. Thus young stellar populations are observed as bright and blue, while old stellar populations are only apparent when there is a lack a younger population and have are characteristically redder. Since magnitude can be considered a proxy for stellar mass (the mass to light ratio of galaxies only varies by a factor of ~ 2 within morphological types and an order of magnitude across all galaxies (Faber & Gallagher, 1979)) the underlying physics is a star formation rate–mass bimodality. This is shown in figure 1.2 for the MPA-JHU sample¹ using Sloan Digital Sky Survey (SDSS) galaxies. These two classes roughly follow the Hubble diagram: spirals tend to exist in the cloud, while elliptical and lenticular (S0s) galaxies (which are collectively known as Early-type galaxies (ETGs)) mostly exist in the red sequence. This thesis has a particular focus on ETGs for reasons described below and more a detailed description of the our current understandings of ETGs is given in section 1.1.

Given their low star formation rates, ETGs are dominated by old stars, while spirals are dominated by young stars. The reasons for this dichotomy is a major part of modern astrophysics. ETGs are often the most massive galaxies, so they clearly

¹<http://wwwmpa.mpa-garching.mpg.de/SDSS/DR7/>



introduction/Hubble.png

Figure 1.1:

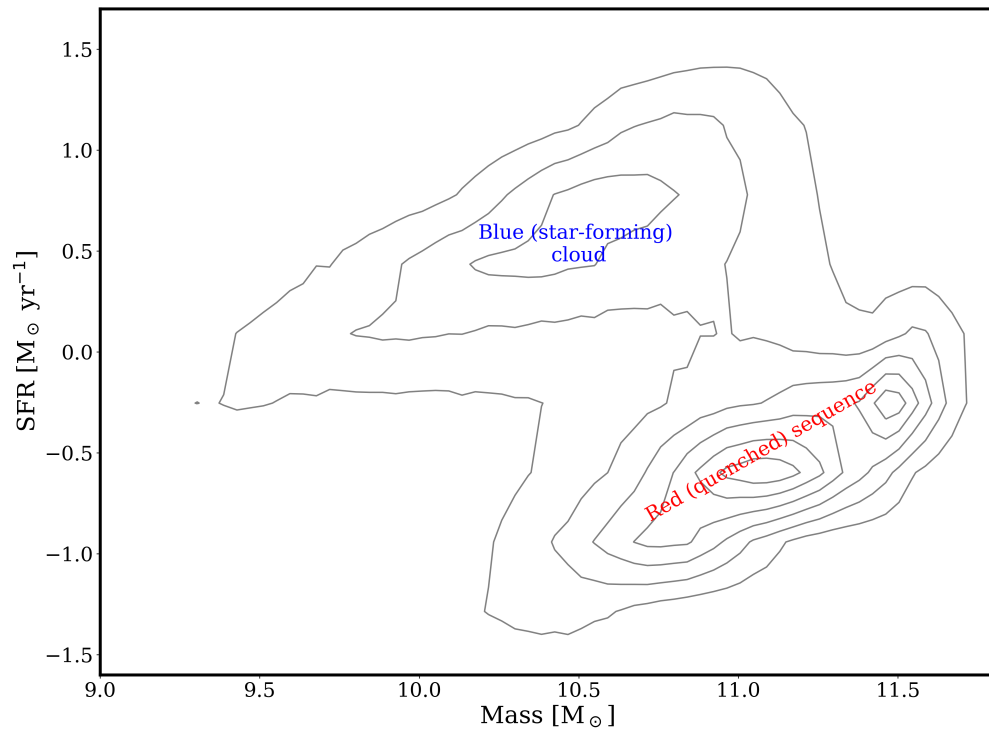


Figure 1.2: The dichotomy in the star-formation rate–mass plane for the MPA-JHU sample (i.e. SDSS DR7) galaxies (Kauffmann et al., 2003; Brinchmann et al., 2003; Salim et al., 2007).

must have undergone substantial star formation in their early history which has now stopped. They either must have run out of fuel for forming stars or have undergone some process which is preventing star formation. Many have been observed to contain large reservoirs of cold (molecular) gas, the material from which stars are formed: this seems to favor the latter.

Many studies now point to the fact that most (if not all) galaxies contain a super-massive black hole at the center. Many studies also show that there are correlations (known as scaling relations) between the properties of the central black hole and the host galaxy despite orders of magnitude differences between the sphere of gravitational influence of the black hole and the size of the host galaxy.

Many galaxies contain a bright point source at the center. These are known as Active Galactic Nuclei (AGN) and in extreme cases these can completely outshine the combined starlight of the host galaxy (these extreme AGN are known as quasars). AGN are understood to be the result of accretion of the inter-stellar medium (ISM) onto the central black hole and the energy emitted by the AGN is generally invoked to explain the scaling relationship as well as the quenching of star formation. This is known as AGN feedback.

AGNs are very varied, however much work has been done on a single unified model of AGNs, where the differences between individual observations can be explained by differing orientations of the AGN. The Unified Model by Antonucci (1993) describes AGN as made up of up to five components or regions surrounding the central black hole: (i) immediately around the black hole is an accretion disk. Surrounding this is (ii) a dust obscuring structure which is torus in shape. This blocks a side-on view of the black hole and accretion disk. Radiation from the accretion disk ionizes (iii) the region immediately above and below the disk. This emits in the optical regime with characteristic broad lines and gives rise to its name: the broad-line region (BLR). The BLR is also obscured by the torus. Above this is (iv) the narrow-line region (NLR), an area of more ordered and less ionized gas. Finally, a small proportion of AGN contain (v) jets of plasma traveling at relativistic speeds out from the poles of the accretion disk (the exact orientation of the jet may well also be dependent of the orientation of the spin of the black hole, though this remains an unknown). The nature of the jet is explored in more detail below. AGN viewed side-on, with the black hole, accretion disk and BLR obscured by the torus are known as Type 2 AGN, while AGN viewed face-on, sometime viewing directly down the jet, are known as Type 1. A schematic the unified model is shown in figure ??

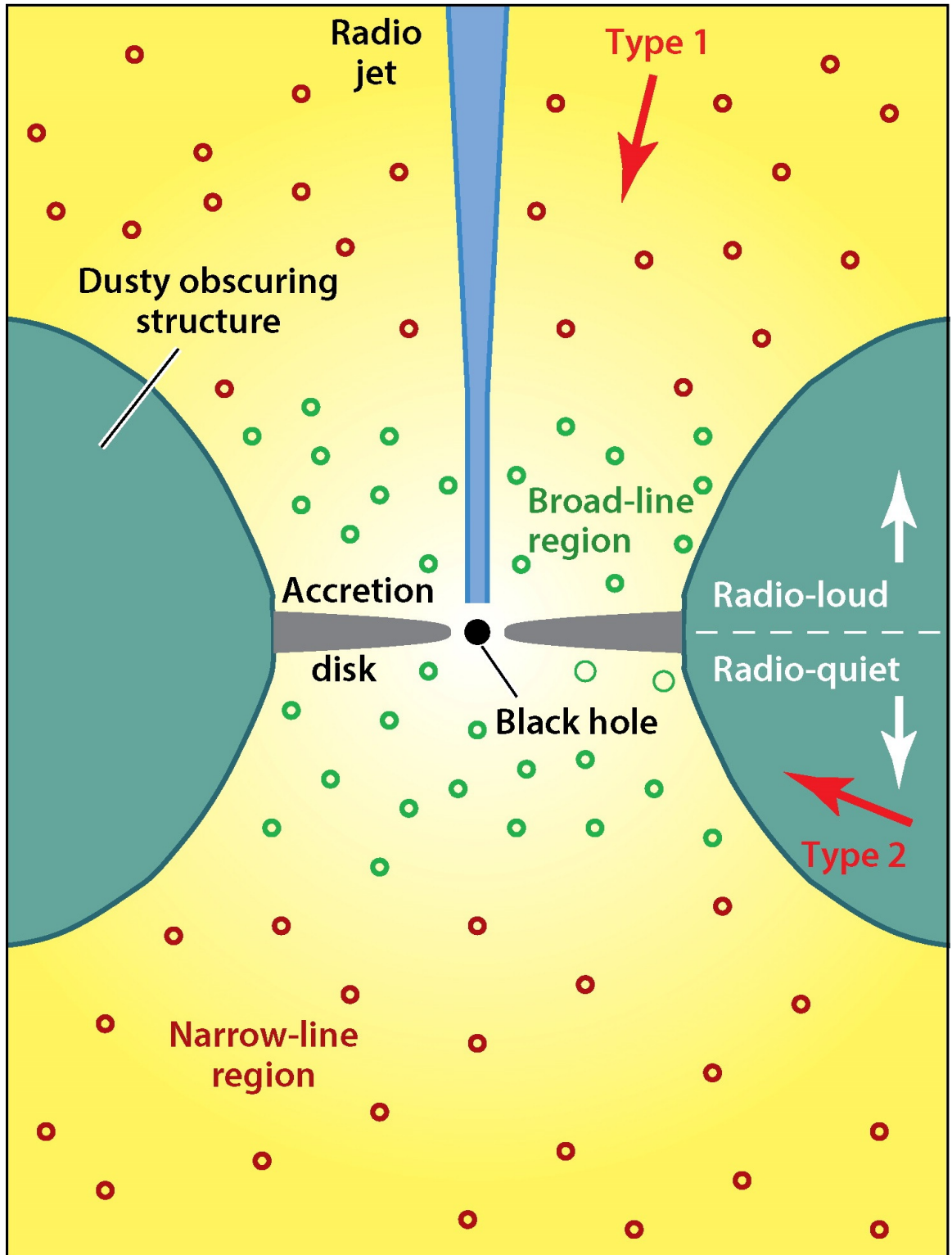


Figure 1.3: A schematic of the unified model of AGN. The primary source of radiation is the accretion disk. This in turn ionizes both the broad-line and the narrow-line regions. A dust torus obscures a side-on view (Type 2 AGN), while a jet sometimes occurs from the poles of the disk. Figure courtesy of Heckman & Best (2014)

Despite the unification of the many observed morphologies of AGN, it is becoming increasingly clear that AGNs do exist in two different modes, Radiative and Jet-mode, raising speculation about differing mechanisms and possibly fuel sources for each. More detail on both modes can be found in section 1.2.

Many galaxies also emit in the radio, although most are extremely faint. These galaxies are known as radio galaxies (RGs) and most radio emission is attributed to synchrotron emission (the broader term of free-free emission is sometimes used in the literature) from the jet within an AGN. However, Nyland et al. (2016) notes that in the fainter sources, the radio emission could be consistent with circumnuclear star formation. There is broad consensus within the literature that the most powerful RGs are the product of gas rich (wet), major mergers of galaxies (a major merger is defined as a ratio in mass closer than 3:1 between progenitor galaxies). These AGN tend to be extremely luminous: indeed despite clearly containing a very powerful jet, they are mostly radiative mode AGN. These RGs form a very small minority of the RG population and hence are not representative. Heckman & Best (2014) presents evidence that most RGs are fueled through secular processes. Exploring the nature and origin of the fuel, as well as the mechanism of accretion and effect of the resulting AGN feedback in RGs forms the basis of this thesis.

The rest of this chapter will cover the background required to study the RG population. Firstly, since RGs are mostly found in ETGs we will cover our current understanding of ETGs. This is mostly a summary of the recent Integral Field Spectroscopy (IFS) surveys including Atlas3D, MaNGA, MASSIVE, SAMI and Califa and can be considered a setting out of the control sample for our later investigations. After this we will give a more detailed summary of the current state of our understanding of AGN. This will include covering RGs.

1.1 Early Type Galaxies

In this section we mostly summarize Cappellari’s excellent review of IFS studies of ETGs (Cappellari, 2016).

The Atlas3D survey (Cappellari et al., 2011) found that morphologically categorizing ETGs was misleading and did not reflect any physical background. Instead they used spatially resolved kinematics to classify ETGs into fast or slow rotators, depending on their value of λ_R (Emsellem et al., 2011). λ_R is a new parameter that reflects

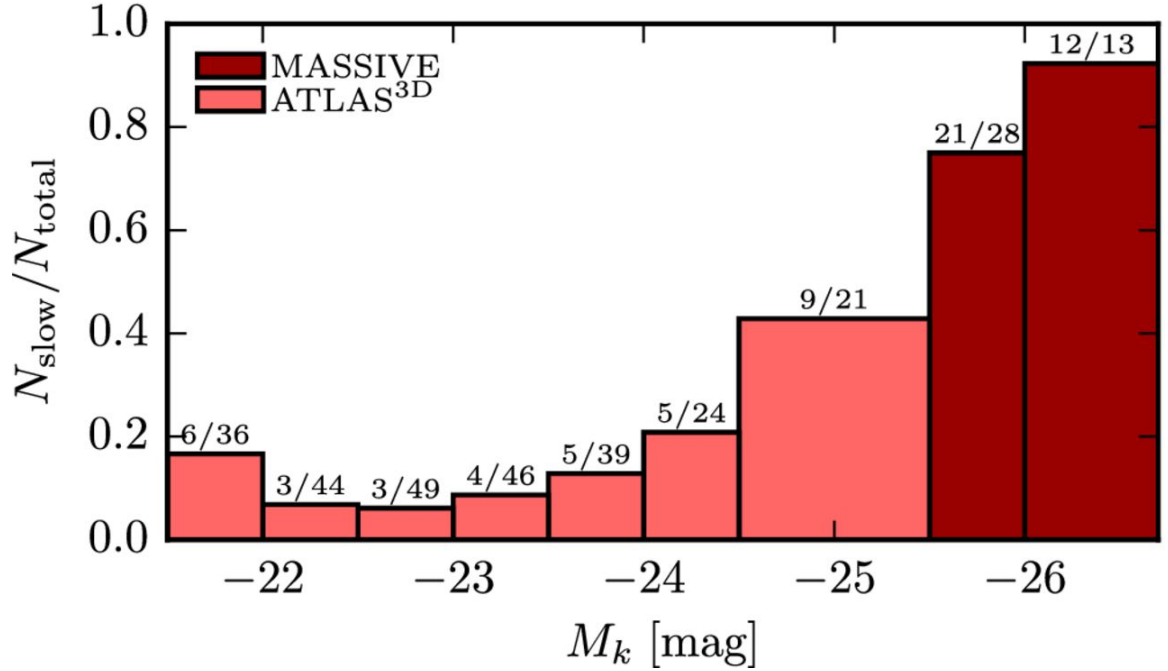


Figure 1.4: The fraction of slow rotators within the total number of galaxies within a given K-band magnitude (M_K). M_K is used here as a proxy for stellar mass. Figure courtesy of Veale et al. (2016)

the galaxies specific angular momentum and is defined as:

$$\lambda_R \equiv \frac{\langle R|V_*| \rangle}{\langle R\sqrt{V_*^2 + \sigma_*^2} \rangle} \quad (1.1)$$

where the angled brackets represent the flux weighted average with a radius, R , V_* is the stellar velocity and σ_* is the stellar velocity dispersion for a given spatial bin. The most recent definition for if a galaxy is considered a slow rotator is:

$$\lambda_{R_e} < 0.08 + \frac{\epsilon}{4} \text{ and } \epsilon < 0.4 \quad (1.2)$$

where ϵ is the ellipticity and λ_{R_e} the value of λ_R evaluated at the effective radius (R_e). It was found that most slow rotators are massive, with $M \gtrsim 10^{10.5} M_\odot$ and that they dominate the high mass end (Emsellem et al., 2011; Veale et al., 2016). Figure 1.4 shows the change in proportion of slow rotators as a function of stellar mass (for which K-band magnitude, M_K , is used as a proxy).

Atlas3D also found a raft of kinematic substructures.

The discovery of the fast/slow rotator scheme lead the Atlas3D group to re-evaluate the Hubble sequence.

1.2 Active Galactic Nuclei

Much of this section, makes use of the review by Heckman & Best (2014) and the references therein, and we would point readers in the direction of this review if more detail is required.

The three main questions that astronomers studying AGN focus on:

1. The origin of the ISM being accreted onto the black hole,
2. The process by which the ISM is accreted and,
3. Any processes by which the AGN might effect to the host galaxy (feedback).

The final point has been invoked (and indeed seems to be required) in numerous semi-analytic models and numerical simulations to successfully reproduce observations of massive galaxies (Di Matteo et al. (2005); Bower et al. (2005); Springel et al. (2005) etc.).

This section will first look at the different global parameters that can be measured for AGN. After this the two different modes (radiative and jet modes) will be explored.

1.2.1 AGN parameters

1.2.1.1 Bolometric Luminosity

Perhaps the most basic property of any AGN is it's total (bolometric) luminosity. Given the non-spherically symmetric nature of the unified model of AGN, estimating the bolometric luminosity can be a non-trivial task, particularly for obscured (Type 2) AGN.

1.2.1.2 Black Hole Mass

Thanks to scale relations, measuring the black hole mass, M_{BH} is relatively simple. In this work we use the mass – stellar velocity dispersion, σ_* relation (M-s) from McConnell & Ma (2013):

$$\log \left(\frac{M_{BH}}{M_{\odot}} \right) = (8.32 \pm 0.05) + (5.64 \pm 0.32) \log \left(\frac{\sigma_*}{200 \text{ km s}^{-1}} \right) \quad (1.3)$$

where σ_* is the luminosity weighted stellar velocity dispersion defined as:

$$\sigma_*^2 \equiv \frac{\int_{R_{\text{inf}}}^{R_e} [\sigma^2(r) + v^2(r)] I(r) dr}{\int_{R_{\text{inf}}}^{R_e} I(r) dr} \quad (1.4)$$

where $I(r)$, $v(r)$ and $\sigma(r)$ are the surface brightness profile, the stellar velocity and stellar velocity dispersion at a radius r . R_e is the effective radius: the radius which contains half of the total stellar light, while R_{inf} is the radius defining the sphere of influence of the black hole ($R_{\text{inf}} \equiv \frac{GM_{\text{BH}}}{\sigma^2}$). In other words, the direct gravitational effect of the black hole is removed from this measure of σ_* .

1.2.1.3 Eddington Ratio

The Eddington ratio, L/L_{Edd} , is the ratio of the Bolometric Luminosity to the maximum luminosity the AGN could possibly achieve while still in hydrostatic equilibrium (the Eddington limit), L_{Edd} . The Eddington limit for pure ionized hydrogen is:

$$\frac{L_{\text{Edd}}}{L_{\odot}} = \frac{4\pi G m_p c}{\sigma_T} \frac{M_{\text{BH}}}{M_{\odot}} = 3.3 \times 10^4 \frac{M_{\text{BH}}}{M_{\odot}} \quad (1.5)$$

where m_p is the mass of proton and σ_T is the Thomson scattering cross-section for an electron on a proton (the ratio σ_T/m_p is the opacity of a plasma).

1.2.2 Radiative Mode

Radiative mode AGN are thought to show actively growing black holes as accreted matter falls directly into the event horizon. In this case the energy expelled as radiation exceeds that expelled by any jet that may be present (most of the energy of a jet is kinetic): they are typically radiatively efficient with Eddington ratios of 1-100%. This means that the accretion disk must be geometrically thin. It is thought radiative mode AGN can also drive powerful interstellar winds, capable of terminating star formation and thus moving a galaxy from the blue cloud to the red sequence.

1.2.3 Jet Mode

In jet mode, the dominant form of emitted energy away from the black hole is the kinetic energy of the material that makes up the jet. Jet-mode AGN can be directly observed feeding-back to their host galaxies: they often terminate at shocks (hot spots), which make up the walls to cavities in the hot (X-ray) component of the host galaxy.

Jet mode AGN tend to be radiatively inefficient with Eddington ratios of $<1\%$. This implies that the inner most accretion disk must be geometrically thick, with an inflow time that is much less than the radiative cooling time scale. These are known as an advection-dominated accretion flow (ADAF). It is possible that there is a thin

disk beyond this (further from the black hole) but it must be truncated by the thick disk. Such a system is capable of launching a jet.

It is thought that the duty cycle of a jet is very short: it is turned on for very short periods of time, before turning off for an extended period. When the jet is turned off or only very weakly present, the galaxy may present as containing a low ionization nuclear emission-line region (LINER). It is worth noting that while the most powerful LINERs are likely to be low-luminosity AGN, many are not: they belong to a new category known as LIERs (the emission-line region is not nuclear - i.e. LINER but without the 'N') (Sarzi et al., 2005, 2010; Singh et al., 2013; Belfiore et al., 2016). Such emission is mostly attributed to ionization from radiation from post-asymptotic giant branch (AGB) stars.

As stated above, most RGs result from synchrotron emission from the relativistic plasma within the jet, moving in the magnetic field of the galaxy. This means that all such RGs can be considered a detection of a jet (we ignore the very low radio luminosity galaxies that result from high levels of circumnuclear star-formation). RGs are always jet mode AGN, except at the highest radio luminosities where the bright radio jet is dominated (in terms of energy output) by the brighter still accretion disk. These are radiative mode AGN.

There is still much speculation on the fuel (and origin) for jet mode AGN, though the consensus is swinging back in favor of an internal origin. Many people advocate hot gas as being the ultimate original fuel: some suggest that this could be directly accreted through Bondi accretion, a simple model where gas is accreted in a spherically symmetric manner directly from the hot phase onto the black hole; while some invoke a cooling mechanism to allow the x-ray gas to cool such that it condenses into molecular (cold) gas before accretion. Others point to a higher detection rate of dust and molecular (cold) gas in RGs compared to the general population suggesting that cold gas from either internal reservoirs or accreted from outside that gas may be the fuel.

Bibliography

- Albrecht A., Steinhardt P. J., 1982, *Physical Review Letters*, 48, 1220
- Antonucci R., 1993, *Annual Review of Astronomy and Astrophysics*, 31, 473
- Arnold J. A., et al., 2013, *The Astrophysical Journal*, Volume 791, Issue 2, article id. 80, 27 pp. (2014)., 791
- Baldry I. K., Glazebrook K., Brinkmann J., Ivezić Ž., Lupton R. H., Nichol R. C., Szalay A. S., 2004, *The Astrophysical Journal*, 600, 681
- Belfiore F., et al., 2016, *Monthly Notices of the Royal Astronomical Society*, 461, 3111
- Bower R. G., Benson A. J., Malbon R., Helly J. C., Frenk C. S., Baugh C. M., Cole S., Lacey C. G., 2005, *Monthly Notices of the Royal Astronomical Society*, Volume 370, Issue 2, pp. 645-655., 370, 645
- Brinchmann J., Charlot S., White S. D. M., Tremonti C., Kauffmann G., Heckman T., Brinkmann J., 2003, *Monthly Notices of the Royal Astronomical Society*, Volume 351, Issue 4, pp. 1151-1179., 351, 1151
- Bryant J. J., et al., 2014, *Monthly Notices of the Royal Astronomical Society*, Volume 447, Issue 3, p.2857-2879, 447, 2857
- Cappellari M., 2016, *Annu. Rev. Astron. Astrophys.*, 54, 1
- Cappellari M., et al., 2011, *Monthly Notices of the Royal Astronomical Society*, 413, 813
- Di Matteo T., Springel V., Hernquist L., 2005, *Nature*, 433, 604
- Emsellem E., et al., 2011, *Monthly Notices of the Royal Astronomical Society*, 414, 888

- Faber S. M., Gallagher J. S., 1979, Annual Review of Astronomy and Astrophysics, 17, 135
- Guth A. H., 1981, Physical Review D, 23, 347
- Heckman T. M., Best P. N., 2014, Annual Review of Astronomy and Astrophysics, 52, 589
- Kauffmann G., et al., 2003, Monthly Notices of the Royal Astronomical Society, 341, 33
- Linde A. D., 1982, Physics Letters B, 108, 389
- McConnell N. J., Ma C.-P., 2013, The Astrophysical Journal, 764, 184
- Nyland K., et al., 2016, Monthly Notices of the Royal Astronomical Society, 458, 2221
- Salim S., et al., 2007, The Astrophysical Journal Supplement Series, 173, 267
- Sanchez S. F., et al., 2011, Astronomy & Astrophysics, 538, A8
- Sarzi M., et al., 2005, Monthly Notices of the Royal Astronomical Society, 366, 1151
- Sarzi M., et al., 2010, Monthly Notices of the Royal Astronomical Society, 402, 2187
- Singh R., et al., 2013, Astronomy & Astrophysics, Volume 558, id.A43, 8 pp., 558
- Springel V., Di Matteo T., Hernquist L., 2005, The Astrophysical Journal, 620, L79
- Veale M., et al., 2016, Monthly Notices of the Royal Astronomical Society, Volume 464, Issue 1, p.356-384, 464, 356