

$= 0$ and $A^2 =$

$\frac{m^2 k^2}{2} + 2 \frac{mEL^2}{2}$, giving five independent constants of motion. (Since the magnitude of A , hence the eccentricity e of the orbit, can be determined from the total angular momentum L and the energy E , only the direction of A is conserved independently; moreover, since A must be perpendicular to L , it contributes only one additional conserved quantity.)

This is consistent with the six initial conditions (the particle's initial position and velocity vectors, each with three components) that specify the orbit of the particle, since the initial time is not determined by a constant of motion. The resulting 1 @-@ dimensional orbit in 6 @-@ dimensional phase space is thus completely specified.

A mechanical system with d degrees of freedom can have at most $2d - 1$ constants of motion, since there are $2d$ initial conditions and the initial time cannot be determined by a constant of motion. A system with more than d constants of motion is called superintegrable and a system with $2d - 1$ constants is called maximally superintegrable. Since the solution of the Hamilton-Jacobi equation in one coordinate system can yield only d constants of motion, superintegrable systems must be separable in more than one coordinate system. The Kepler problem is maximally superintegrable, since it has three degrees of freedom ($d = 3$) and five independent constant of motion; its Hamilton-Jacobi equation is separable in both spherical coordinates and parabolic coordinates, as described below.

Maximally superintegrable systems follow closed, one @-@ dimensional orbits in phase space, since the orbit is the intersection of the phase @-@ space isosurfaces of their constants of motion. Consequently, the orbits are perpendicular to all gradients of all these independent isosurfaces, five in this specific problem, and hence are determined by the generalized cross products of all of these gradients. As a result, all superintegrable systems are automatically describable by Nambu mechanics, alternatively, and equivalently, to Hamiltonian mechanics.

Maximally superintegrable systems can be quantized using commutation relations, as illustrated below. Nevertheless, equivalently, they are also quantized in the Nambu framework, such as this classical Kepler problem into the quantum hydrogen atom.

$=$ Evolution under perturbed potentials $=$

The Laplace-Runge-Lenz vector A is conserved only for a perfect inverse @-@ square central force. In most practical problems such as planetary motion, however, the interaction potential energy between two bodies is not exactly an inverse square law, but may include an additional central force, a so @-@ called perturbation described by a potential energy $h(r)$. In such cases, the LRL vector rotates slowly in the plane of the orbit, corresponding to a slow apsidal precession of the orbit.