= topness , B ? = bottomness .

* Notation such as $173210 \pm 510 \pm 710$ denotes two types of measurement uncertainty . In the case of the top quark, the first uncertainty is statistical in nature, and the second is systematic.

= = Interacting quarks = =

As described by quantum chromodynamics, the strong interaction between quarks is mediated by gluons, massless vector gauge bosons. Each gluon carries one color charge and one anticolor charge. In the standard framework of particle interactions (part of a more general formulation known as perturbation theory), gluons are constantly exchanged between quarks through a virtual emission and absorption process. When a gluon is transferred between quarks, a color change occurs in both; for example, if a red quark emits a red? antigreen gluon, it becomes green, and if a green quark absorbs a red? antigreen gluon, it becomes red. Therefore, while each quark 's color constantly changes, their strong interaction is preserved.

Since gluons carry color charge , they themselves are able to emit and absorb other gluons . This causes asymptotic freedom : as quarks come closer to each other , the chromodynamic binding force between them weakens . Conversely , as the distance between quarks increases , the binding force strengthens . The color field becomes stressed , much as an elastic band is stressed when stretched , and more gluons of appropriate color are spontaneously created to strengthen the field . Above a certain energy threshold , pairs of quarks and antiquarks are created . These pairs bind with the quarks being separated , causing new hadrons to form . This phenomenon is known as color confinement : quarks never appear in isolation . This process of hadronization occurs before quarks , formed in a high energy collision , are able to interact in any other way . The only exception is the top quark , which may decay before it hadronizes .

= = = Sea quarks = = =

Hadrons contain, along with the valence quarks (q

- v) that contribute to their quantum numbers, virtual quark? antiquark (qq) pairs known as sea quarks (q
- s) . Sea quarks form when a gluon of the hadron 's color field splits; this process also works in reverse in that the annihilation of two sea quarks produces a gluon . The result is a constant flux of gluon splits and creations colloquially known as " the sea " . Sea quarks are much less stable than their valence counterparts , and they typically annihilate each other within the interior of the hadron . Despite this , sea quarks can hadronize into baryonic or mesonic particles under certain circumstances .

= = = Other phases of quark matter = = =

Under sufficiently extreme conditions , quarks may become deconfined and exist as free particles . In the course of asymptotic freedom , the strong interaction becomes weaker at higher temperatures . Eventually , color confinement would be lost and an extremely hot plasma of freely moving quarks and gluons would be formed . This theoretical phase of matter is called quark ? gluon plasma . The exact conditions needed to give rise to this state are unknown and have been the subject of a great deal of speculation and experimentation . A recent estimate puts the needed temperature at (1 @ .@ 90 \pm 0 @ .@ 02) \times 1012 kelvin . While a state of entirely free quarks and gluons has never been achieved (despite numerous attempts by CERN in the 1980s and 1990s) , recent experiments at the Relativistic Heavy Ion Collider have yielded evidence for liquid @-@ like quark matter exhibiting " nearly perfect " fluid motion .

The quark? gluon plasma would be characterized by a great increase in the number of heavier quark pairs in relation to the number of up and down quark pairs. It is believed that in the period prior to 10? 6 seconds after the Big Bang (the quark epoch), the universe was filled with quark?

gluon plasma, as the temperature was too high for hadrons to be stable.

Given sufficiently high baryon densities and relatively low temperatures? possibly comparable to those found in neutron stars? quark matter is expected to degenerate into a Fermi liquid of weakly interacting quarks. This liquid would be characterized by a condensation of colored quark Cooper pairs, thereby breaking the local SU (3) c symmetry. Because quark Cooper pairs harbor color charge, such a phase of quark matter would be color superconductive; that is, color charge would be able to pass through it with no resistance.