

Measuring *Differences* in Momentum Scale as a Function of Track Parameters

Jim Pivarski

Texas A&M University

? March, 2010

Motivation

Jim Pivarski 2/14

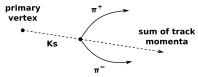




- ▶ Alignment optimizes tracks very well locally, but less so globally
 - any two regions that are rarely crossed by the same track can develop different momentum scales through weak modes
- ▶ Masses of J/ψ , Υ , and Z set the momentum scale, but with some complications
 - each daughter samples a different region of the tracker: needs to be untangled
 - shape is not symmetric due to final state radiation
 - backgrounds must be part of the fit
- ▶ In this talk, I'll present a procedure that can equalize momentum scales in all regions of the tracker, but not set the absolute scale
- ▶ How it fits into a complete track-correction program:
 - $\mathsf{alignment} \to \mathsf{this} \ \mathsf{correction} \to \mathsf{absolute} \ \mathsf{scale} \ \mathsf{correction}$
- Also provides uncertainties and correlations in the remaining bias, which are important for setting systematic errors on physics quantities



- Alignment makes module positions and track parameters mutually consistent (optimizes track shapes)
- Overall momentum scale correction adjusts observed resonance peaks to externally known values
- ▶ Another physics constraint: momentum direction of daughters of a decay-in-flight must be consistent with the flight direction



- sensitive to different momentum scales in different regions of the detector, but not absolute scale
- symmetric in the case of no misalignment (unlike mass peak)
- $K_S \to \pi^+\pi^-$ for low momenta, perhaps $B \to$ "all charged" for higher?
 - \triangleright boosted K_S are not sufficient because we need the daughters to sample different parts of the detector: need high mass

Untangling daughters (1/2)

Jim Pivarski 4/14





Same formalism as alignment (HIP for a single alignable):

(
$$N$$
 track residuals) = $(6 \times N \text{ matrix}) \cdot (6 - DOF \text{ parameters})$

 \triangleright N decays with angle difference $\Delta \vec{\phi}$, track-parameter corrections described by M parameters \vec{p} ($N \gg M$):

$$\left(\begin{array}{c} \Delta\phi_1\\ \vdots\\ \Delta\phi_N \end{array}\right) = A \cdot \left(\begin{array}{c} \rho_1\\ \vdots\\ \rho_M \end{array}\right)$$

- ightharpoonup A is the transformation from parameters \vec{p} to observables $\Delta \vec{\phi}$ and is therefore the derivative $\frac{\partial(\Delta\phi_i)}{\partial p_i}$, which can be computed numerically:
 - 1. for each decay i, compute nominal $\Delta \phi_i$
 - 2. for each parameter j, apply ϵp_i to all tracks for $\Delta \phi_i(\epsilon p_i)$

3.
$$A_{ij} = \frac{\Delta \phi_i - \Delta \phi_i (\epsilon p_j)}{\epsilon}$$



▶ The values of \vec{p} which minimize quadratic χ^2 (assumes Gaussian $\Delta \vec{\phi}$)

$$\chi^2 = \left(\Delta \vec{\phi} - A \cdot \vec{p}\right)^T \left(\sigma_{\Delta \phi}^{\ 2}\right)^{-1} \left(\Delta \vec{\phi} - A \cdot \vec{p}\right)$$

$$\text{are } \vec{p} = \left(A^T \left(\sigma_{\Delta \phi}^{\ 2}\right)^{-1} A\right)^{-1} \left(A^T \left(\sigma_{\Delta \phi}^{\ 2}\right)^{-1} \Delta \vec{\phi}\right) \qquad \sim \frac{\sum_{w_i x_i} w_i}{\sum_{w_i} w_i}$$

with uncertainty/covariance
$$\left(A^T \left(\sigma_{\Delta\phi}^2\right)^{-1} A\right)^{-1} \sim \frac{1}{\sum w_i}$$

- "convert $\Delta \vec{\phi}$ into \vec{p} space and compute weighted mean"
- the $\sigma_{\Delta\phi}^2$ weights are propagated uncertainties in $\Delta\vec{\phi}$
- \triangleright The final \vec{p} correct biases in tracking, but perhaps more importantly, the covariance provides a rigorous systematic error in tracking
- \triangleright Potential early application: Ξ^{\pm} mass measurement; tracking bias is the only systematic error

Parameterization of \vec{p}

Jim Pivarski 6/14



- ▶ Biases are small and slowly varying (largest alignment weak modes are global distortions)
- lacktriangle Expand general $ec{p}(\phi,\kappa,\cot\theta,d_{\mathsf{x}\mathsf{y}},d_{\mathsf{z}})$ function and keep low-order
 - ▶ Fourier-expand in ϕ : $\sin \phi$, $\cos \phi$, $\sin 2\phi$, $\cos 2\phi$
 - ▶ Taylor-expand in other parameters: κ , $\kappa \cot \theta$, $(\cot \theta)^2$
 - κ and d_{xy} are not sampled near zero due to physics signature: exclude κ^2 and d_{xy}^2 terms
- ▶ 33 terms for $\delta \kappa$

$$\delta\kappa = p_1 + p_2 \sin(\phi) + p_3 \cos(\phi) + p_4 \sin(2\phi) + p_5 \cos(2\phi) + p_6 \kappa + p_7 \cot \theta + p_8 d_{xy} + p_9 d_z + p_{10} \sin(\phi) \kappa + p_{11} \sin(\phi) \cot \theta + p_{12} \sin(\phi) d_{xy} + p_{13} \sin(\phi) d_z + p_{14} \cos(\phi) \kappa + p_{15} \cos(\phi) \cot \theta + p_{16} \cos(\phi) d_{xy} + p_{17} \cos(\phi) d_z + p_{18} \sin(2\phi) \kappa + p_{19} \sin(2\phi) \cot \theta + p_{20} \sin(2\phi) d_{xy} + p_{21} \sin(2\phi) d_z + p_{22} \cos(2\phi) \kappa + p_{23} \cos(2\phi) \cot \theta + p_{24} \cos(2\phi) d_{xy} + p_{25} \cos(2\phi) d_z + p_{26} \kappa \cot \theta + p_{27} \kappa d_{xy} + p_{28} \kappa d_z + p_{29} \cot \theta \cot \theta + p_{23} \cot \theta d_{xy} + p_{31} \cot \theta d_z + p_{32} d_{xy} d_z + p_{33} d_z d_z$$

- ▶ Same for $\delta\phi$ and δd_{xy} : 99 parameters, 99 × 99 matrix inversion
- lacksquare $\delta\cot heta$ and $\delta d_{
 m z}$ cannot be optimized; used to select primary vertex



- ▶ Identify $K_S \to \pi^+\pi^-$ by $\pi^+\pi^-$ invariant mass (it therefore cannot be used in the optimization)
- Use sideband for a background control sample; sideband events enter the minimization with negative weight, cancelling the effect of backgrounds under mass peak
- \triangleright Identify the correct primary vertex by propagating K_S to beamline in r-z plane ($\delta \cot \theta$ and δd_z cannot be optimized)



- \triangleright Require exactly one primary vertex in the road from K_S
- ▶ Require K_S vertex to be within $1.5 < |v_{xy}| < 3$ cm; away from beamspot and within first pixel layer
- From $K_S \to \pi^+\pi^-$, typical curvatures are $0.5 \le |\kappa| \le 5$ GeV⁻¹

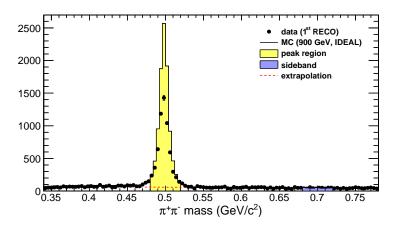
Feasibility study (1/6)

Jim Pivarski

8/14



- Quick check to make sure that there aren't any show-stoppers
- ▶ MC and 1st RECO data (non-optimal tracking: large APEs)



Upper sideband only: there can be signal events below peak due to final state radiation

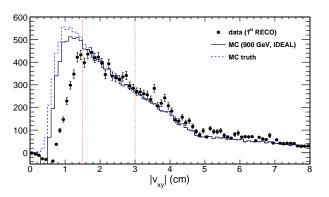
Feasibility study (2/6)

Jim Pivarski





- ► All plots from this point onward are background-subtracted and normalized within cut windows
- "MC truth" is matched to true K_S in MC (no sideband subtraction)



- ightharpoonup Could loosen track flight significance and lower $|v_{xy}|$ boundary
- ▶ Upper $|v_{xy}|$ boundary must be within first pixel layer so that resolution is smooth as a function of track d_{xy}

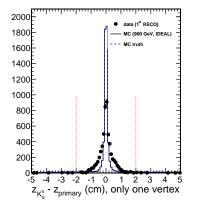
Feasibility study (3/6)

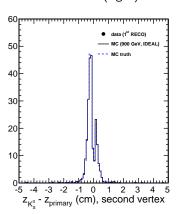
Jim Pivarski 10/14





 \triangleright Propagation of K_S to beamline in two cases: when there is only one primary vertex (left), and when there is also a second (right)





- ▶ In MC, 3.4% of events have a second primary vertex, in data, several do, but well beyond 5 cm (negligible pile-up plus beam-gas?)
- Efficiency of "only one vertex" requirement is dependent on luminosity

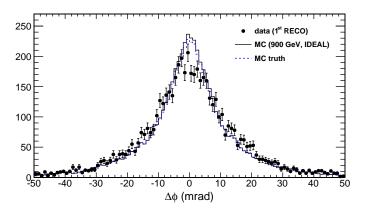
Feasibility study (4/6)

Jim Pivarski 11/14





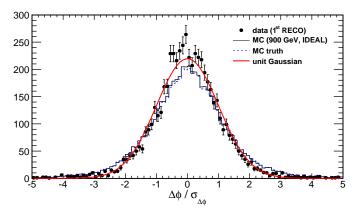
▶ Distribution of all $\Delta \phi$ is not Gaussian: is quadratic χ^2 the right estimator? Does the minimization need to be made non-linear (replace 99 × 99 matrix inversion with Minuit)?



 \blacktriangleright Some indication of $\Delta\phi$ variation in data? Remember, these are unrealistic APEs



▶ Calculation of $\sigma_{\Delta\phi}^2$ weights must include all correlations between two-track intersection and momentum sum, starting from 5×5 track parameter covariances

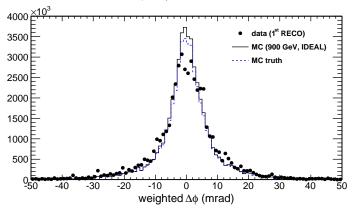


Uncertainty is slightly underestimated in MC, larger in data

Jim Pivarski



• $\Delta\phi$ histogram weighted by $1/\sigma_{\Delta\phi}{}^2$



- Much narrower, indicating that the largest outliers have the largest uncertainties
- ▶ Data are still broader than MC (and distribution still not Gaussian)

Conclusions

Jim Pivarski 14/14





- Proposal: multi-step track corrections
 - 1. alignment corrects track shapes, but can allow relative biases in regions of the tracker that are not connected by cosmic rays or collisions (assuming no resonance constraints)
 - 2. parameterize track parameter space and apply regional corrections to make momentum scale uniform, and also acquire matrix of uncertainties
 - 3. apply overall correction to set momentum scale with resonance masses, also with uncertainties
- ▶ Even if no track parameter biases are observed, uncertainty in bias is an important systematic for physics analyses
 - \blacktriangleright early analysis: CMS can improve Ξ^{\pm} mass measurement in an analysis with only one systematic uncertainty— tracking
- ▶ Demonstrated selection and $\Delta \phi$, $\sigma_{\Delta \phi}^2$ calculation with $K_S \to \pi^+ \pi^-$
 - next steps would be to calculate \vec{p} and its correlations; is it consistent with zero in MC? How much uncertainty per \sqrt{N} ?
- ▶ For most analyses, a higher-mass metastable particle would be more relevant: $B \rightarrow \text{all charged hadrons}$?