

# Hard Probes used in Heavy Ion Collisions to Study QCD at Extreme Energy Densities

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on behalf of the CMS Collaboration

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- Heavy ion collisions produce a new kind of fluid called a quark-gluon plasma
- It survives such a short time that we cannot use traditional methods to study its properties
- Instead, we look for well-known particle physics processes ("probes") to see how they are modified by the medium
- "Hard" → high-energy, perturbative QCD production of dijets, bound states, and electroweak bosons



one of the two partons loses energy in the medium in the medium



bound particles dissociate

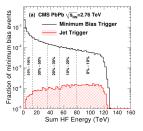


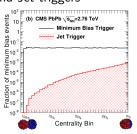
electroweak not strongly affected

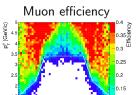


0.05

#### Minimum Bias and Jet triggers





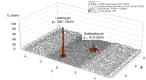


- $\eta$  coverage:  $\pm 5.2$  for jets,  $\pm 2.4$  for muons
- High-resolution tracking up to 100's of GeV
- Particle flow jets (get charged particle momenta from tracks)
- $\blacktriangleright$  150  $\mu b^{-1}$  Pb-Pb (Nov 2011) and 231  $nb^{-1}$  of 2.76 TeV pp (Mar 2012)





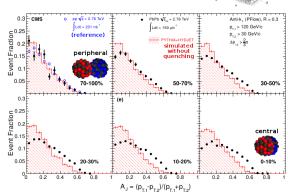
- ► Hard process produces two partons, one loses energy in the medium ("quenched")
- Dramatic enough at the LHC to see it in event displays



Quantify with

$$A_{J} = rac{p_{\mathrm{T,1}} - p_{\mathrm{T,2}}}{p_{\mathrm{T,1}} + p_{\mathrm{T,2}}}$$
 and  $p_{\mathrm{T,2}}/p_{\mathrm{T,1}}$ 

 Calorimeter-only and particle-flow jet algorithms yield similar results



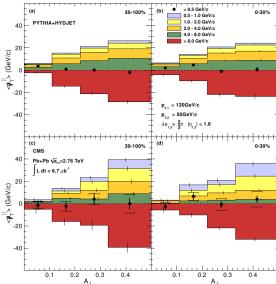
PRC 84 (2011) 024906 and PLB 712 (2012) 176



$$p_{\mathrm{T}}^{\parallel} = \sum_{i} -p_{\mathrm{T}}^{\mathrm{i}} \cos{(\phi_{\mathrm{i}} - \phi_{\mathrm{Leading Jet}})}$$

versus

$$A_J = \frac{p_{\text{T,1}} - p_{\text{T,2}}}{p_{\text{T,1}} + p_{\text{T,2}}}$$

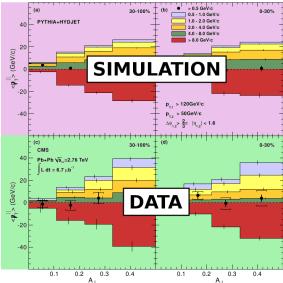


PRC 84 (2011) 024906



$$p_{\mathrm{T}}^{\parallel} = \sum_{i} -p_{\mathrm{T}}^{\mathrm{i}} \cos{\left(\phi_{\mathrm{i}} - \phi_{\mathrm{Leading Jet}}\right)}$$

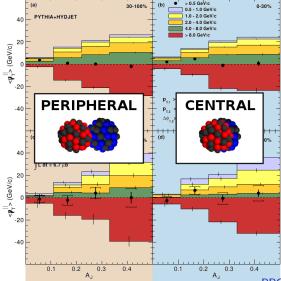
$$A_J = \frac{p_{\rm T,1} - p_{\rm T,2}}{p_{\rm T,1} + p_{\rm T,2}}$$





$$p_{\mathrm{T}}^{\parallel} = \sum_{i} -p_{\mathrm{T}}^{\mathrm{i}} \cos{(\phi_{\mathrm{i}} - \phi_{\mathrm{Leading Jet}})}$$

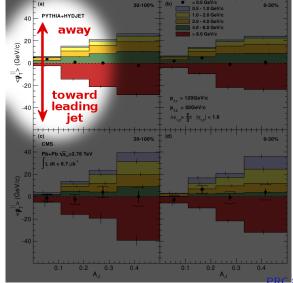
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$$p_{\mathrm{T}}^{\parallel} = \sum_{\mathrm{r}} -p_{\mathrm{T}}^{\mathrm{i}} \cos{\left(\phi_{\mathrm{i}} - \phi_{\mathrm{Leading Jet}}\right)}$$
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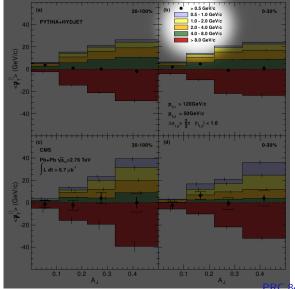
$$A_J = \frac{p_{\text{T,1}} - p_{\text{T,2}}}{p_{\text{T,1}} + p_{\text{T,2}}}$$





$$p_{\mathrm{T}}^{\parallel} = \sum_{\mathrm{i}} -p_{\mathrm{T}}^{\mathrm{i}}\cos\left(\phi_{\mathrm{i}} - \phi_{\mathrm{Leading\ Jet}}
ight)$$
 versus

 $A_J = \frac{p_{\rm T,1} - p_{\rm T,2}}{p_{\rm T,1} + p_{\rm T,2}}$ 

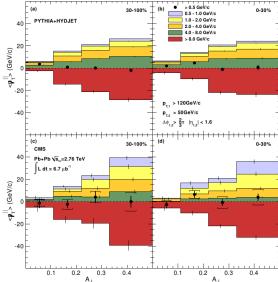


PRC 84 (2011) 024906



$$p_{\mathrm{T}}^{\parallel} = \sum_{i} -p_{\mathrm{T}}^{\mathrm{i}} \cos{(\phi_{\mathrm{i}} - \phi_{\mathrm{Leading Jet}})}$$

$$A_J = \frac{p_{\text{T,1}} - p_{\text{T,2}}}{p_{\text{T,1}} + p_{\text{T,2}}}$$

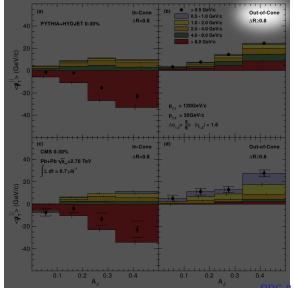




$$p_{ ext{T}}^{\parallel} = \sum_{ ext{i}} -p_{ ext{T}}^{ ext{i}} \cos{(\phi_{ ext{i}} - \phi_{ ext{Leading Jet}})}$$
 versus

S

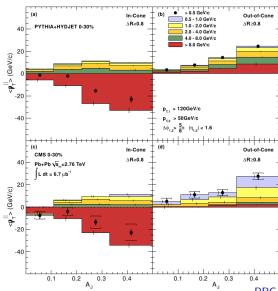
$$A_J = \frac{p_{\text{T,1}} - p_{\text{T,2}}}{p_{\text{T,1}} + p_{\text{T,2}}}$$



## Missing $p_T$ went into large-angle particles



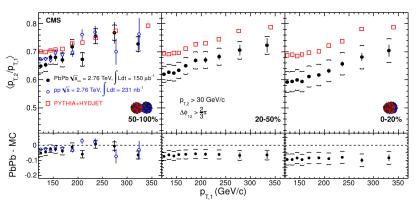
$$p_{
m T}^{||} = \sum_{
m i} -p_{
m T}^{
m i} \cos{(\phi_{
m i} - \phi_{
m Leading\ Jet})}$$
 versus



 $A_J = \frac{p_{\text{T,1}} - p_{\text{T,2}}}{p_{\text{T,1}} + p_{\text{T,2}}}$ 



ightharpoonup vertical axis:  $\frac{\text{subleading } p_T}{\text{leading } p_T}$  fraction, horizontal axis: leading  $p_T$ 

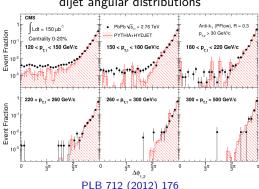


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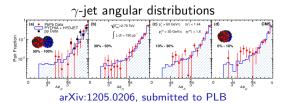


- Photon-jet pairs produced in  $qg \rightarrow \gamma q$  and  $q\bar{q} \rightarrow \gamma g$
- The photon does not interact with the medium, making it a good probe of the initial state of the parton that produced the jet
- Dijets and photon-jet correlations both show large energy loss, rather than angle decorrelation

#### dijet angular distributions

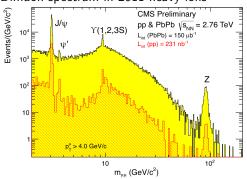


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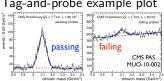




#### Upsilon candidate



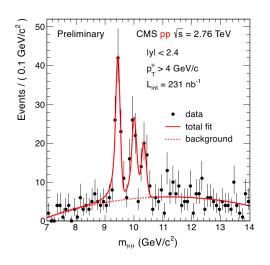
#### Tag-and-probe example plot



### Upsilon suppression, state by state

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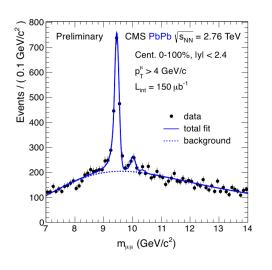




► CMS can resolve the three  $\Upsilon$  states

PRL 107 (2011) 052302 https://twiki.cern.ch/twiki/bin/view/CMSPublic/PhysicsResultsHIN11011



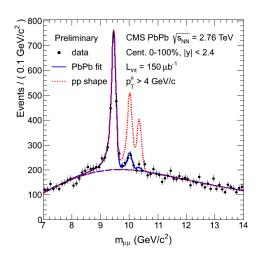


- ► CMS can resolve the three ↑ states
- In heavy ion collisions, all are suppressed, but  $\Upsilon(2S)$  and  $\Upsilon(3S)$  more than  $\Upsilon(1S)$

PRL 107 (2011) 052302

https://twiki.cern.ch/twiki/bin/view/CMSPublic/PhysicsResultsHIN11011



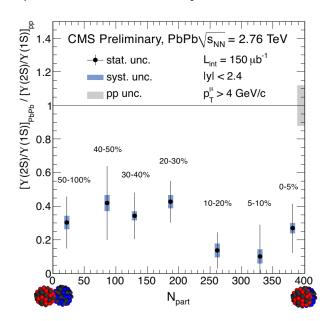


- ► CMS can resolve the three ↑ states
- In heavy ion collisions, all are suppressed, but  $\Upsilon(2S)$  and  $\Upsilon(3S)$  more than  $\Upsilon(1S)$
- Significance of suppression: 5.4σ
- ▶ Double ratio  $\frac{nS/1S|_{\text{Pb Pb}}}{nS/1S|_{\text{pp}}}$ 2S/1S = 0.21 ± 0.07 ± 0.02 3S/1S < 0.1 (95% C.L.)

PRL 107 (2011) 052302

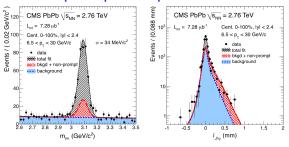
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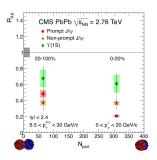




# CMS

### and non-prompt samples

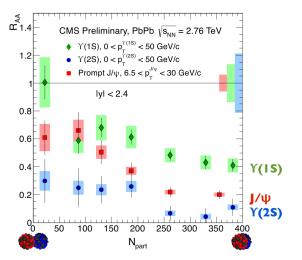




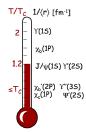
- ▶ In  $B \to J/\psi\,X$ , the secondary vertex from the long-lived B, far from the quark-gluon plasma, can be distinguished from prompt  $J/\psi$  production
- ▶  $I_{\mathrm{J/\psi}}$  is the proper flight distance along the dimuon momentum axis  $(I_{\mathrm{J/\psi}} < 0$  part of the distribution is purely resolution)
- $P_{AA} = \frac{\mathcal{L}_{pp}}{T_{AA}N_{MB}} \frac{N_{PbPb}(Q\bar{Q})}{N_{pp}(Q\bar{Q})} \cdot \frac{\varepsilon_{pp}}{\varepsilon_{PbPb}}, \text{ the nuclear modification factor}$  JHEP 1205 (2012) 063



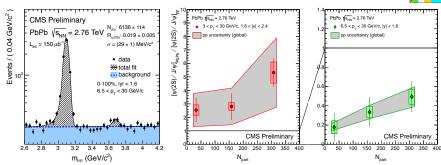
- ▶ Suppression of  $\Upsilon(1S)$ , prompt  $J/\psi$ ,  $\Upsilon(2S)$ , and  $\Upsilon(3S)$  follow the pattern of their binding energies
- ▶ Supports the idea that the medium is screening their potentials



state	J/ψ	χ <sub>c</sub>	ψ'(2S)
Mass(GeV)	3.10	3.53	3.69
ΔE (GeV)	0.64	0.20	0.05
T <sub>d</sub> /T <sub>c</sub>	2.1	1.16	1.12
state	Y(1S)	Y(2S)	Y(3S)
Mass(GeV)	9.46	10.0	10.36
ΔE (GeV)	1.10	0.54	0.20



PLB 178, 416 (1986)



- ▶  $J/\psi$  appears more suppressed than  $\psi(2S)$  in forward, low  $p_T$  region
- lacksquare  $\psi(2S)$  is more suppressed than  $J/\psi$  at midrapidity, higher  $ho_T$
- ▶ Note, however, that uncertainties (mostly from pp reference) are large: the highest point is

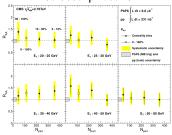
 $5.32\,\pm\,1.03$  (stat.)  $\pm\,0.79$  (syst.)  $\pm\,2.58$  (pp), or 1.5 sigma from unity

Colorless probes are unaffected ~

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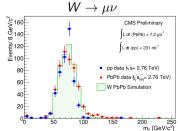




 $Z \to \mu\mu$   $\begin{array}{c} Z \to \mu\mu\\ \text{O 15} \\ \text{O 15} \\ \text{D 20} \\ \text{D 10} \\ \text{$ 

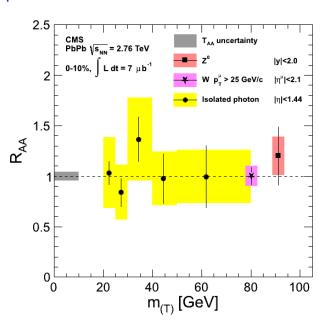
PRL 106 (2011) 212301

PLB 710 (2012) 256



arXiv:1205.6334, submitted to PLB







- ▶ Jet quenching:
  - Asymmetry in dijets is primarily due to momentum loss into low-p<sub>T</sub>, out-of-cone tracks, rather than angular decorrelation or lost particles
  - ▶ The effect has little momentum dependence
  - Jet-photon correlations lead to the same conclusion as jet-jet
- Suppression of quarkonia:
  - $\blacktriangleright$  Measured Upsilon states separately, prompt and non-prompt  $J/\psi$
  - ▶ The  $\Upsilon(1S)$ ,  $J/\psi$ ,  $\Upsilon(2S)$ , and  $\Upsilon(3S)$  follow the pattern of Matsui-Satz screening
  - ▶ The  $\psi(2S)$  may actually be less suppressed than the  $J/\psi$  in the forward, low- $p_T$  region, though the uncertainties are large
- Colorless probes:
  - Production and development of γ, W, and Z are unaffected by the strongly interacting medium
  - ► Supports the interpretation of jet quenching and suppression of quarkonia as effects that happen *after* they are produced