# Di-electron Widths of the $\Upsilon(1S)$ , $\Upsilon(2S)$ , and $\Upsilon(3S)$

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**CLEO Collaboration** 

#### Introduction

Even though Nuclear Strong force is "simpler" than Electroweak, it is an obstacle to understanding Electroweak

#### Nuclear Strong force (QCD)

- Highly symmetric
- One tunable parameter (+ quark masses)

Non-perturbative below 1 GeV

#### Electroweak interaction

- P, CP symmetries broken
- No obvious pattern in flavor-changing interactions:

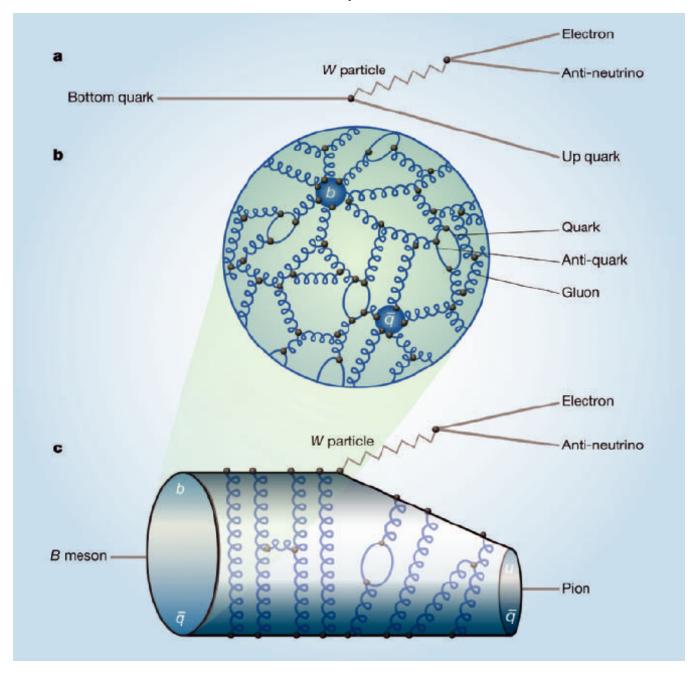
$$P(q_1 
ightarrow q_2) \propto \left| q_2 \cdot \left( egin{array}{ccc} V_{ud} & V_{us} & V_{ub} \ V_{cd} & V_{cs} & V_{cb} \ V_{td} & V_{ts} & V_{tb} \end{array} 
ight) \cdot q_1 
ight|^2$$

Perturbative

All measurements of quark properties must involve QCD

To learn more about electroweak interactions, we need to understand QCD better!

### A typical electroweak measurement: $b \rightarrow u$ process



#### Outline for this Talk

1. Show how  $V_{td}$  is obfuscated by QCD and how our knowledge of it is limited by our ability to compute QCD

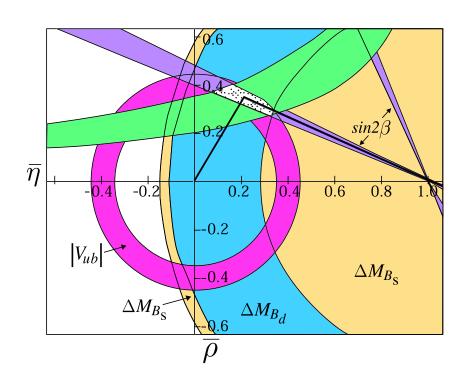
2. Introduce Lattice QCD as a tool which can help to compute the necessary parameter

3. Describe a CLEO experiment which tests this calculation: di-electron widths of  $\Upsilon(1S)$ ,  $\Upsilon(2S)$ ,  $\Upsilon(3S)$ 

# Measuring $V_{td}$

 $B\text{-}\bar{B}$  mixing:

$$B\left\{\begin{array}{c|c} b & t & d \\ \hline \hline w \lessgtr & \lessgtr w \\ \hline \overline d & \overline t & \overline b \end{array}\right\} \overline{B}$$



$$\Delta M_{B_d} = (known) \times (f_B^2 B_B) \times |V_{td}|^2 = 0.510 \pm 0.005 \text{ ps}^{-1}$$
, a 1% measurement!

But  $f_B$  is only known to 20% of itself

Hence the 20% uncertainty in  $V_{td}$  (blue band)

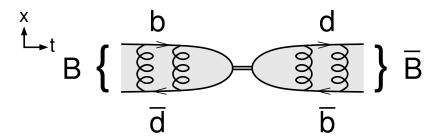
# What is $f_B$ ?

QCD corrections to b- $\bar{d}$ -electroweak "vertex"

On QCD length scales, B-mixing diagram

$$B\left\{\begin{array}{c|c} b & t & d \\ \hline \hline w & & w \\ \hline \overline{d} & \overline{t} & \overline{b} \end{array}\right\} \overline{B}$$

looks like this:

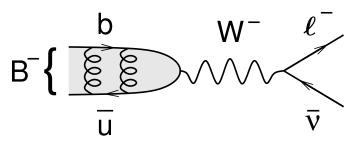


 $f_B$  expresses the probability that the  $ar{d}$  will fluctuate onto the b quark

That is, the value of the spatial wavefunction at the origin

# Determining $f_B$

Experimentally?  $B^- \to \ell^- \bar{\nu}$ 



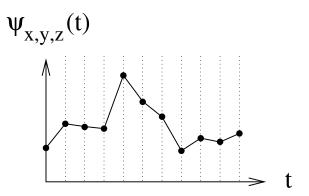
$$\bullet \Gamma(B^- \to \ell^- \bar{\nu}) = \frac{G_F^2}{8\pi} \quad |V_{ub}|^2 \quad m_l^2 M_B \left(1 - \frac{m_l^2}{M_B^2}\right)^2 \quad f_B^2$$
small small

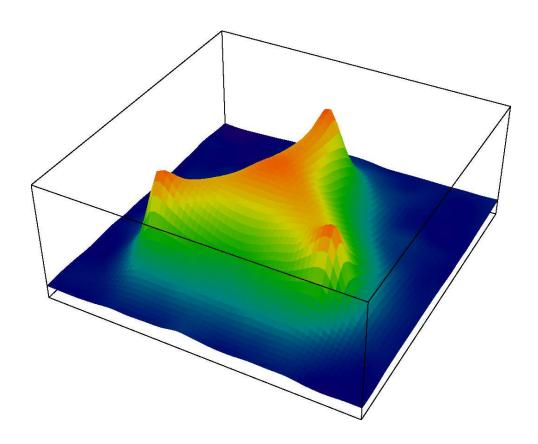
• 
$$\mathcal{B}(B^- \to \tau^- \bar{\nu}) < 1.8 \times 10^{-4}$$
 at 90% C.L. (253 fb $^{-1}$  at Belle 2005)

Theoretically? Need non-perturbative techniques. . .

# Lattice QCD

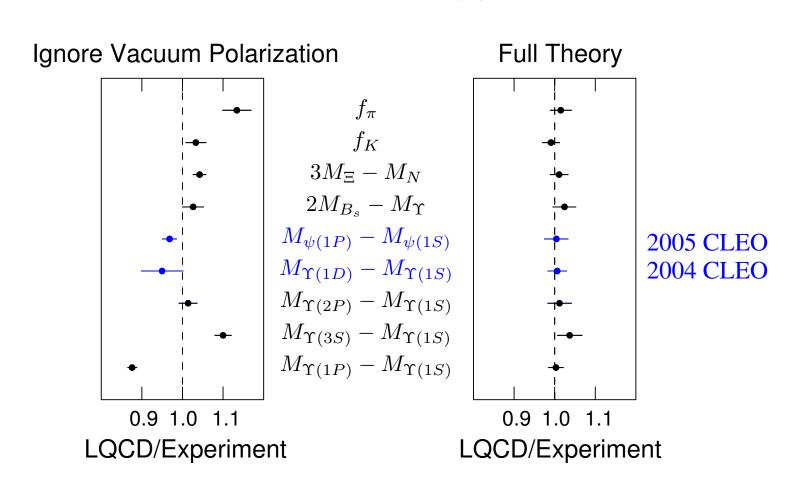
- Evaluate path integral with Monte Carlo integration
- Very computationally intensive





# Recent Breakthrough (c. 1999)

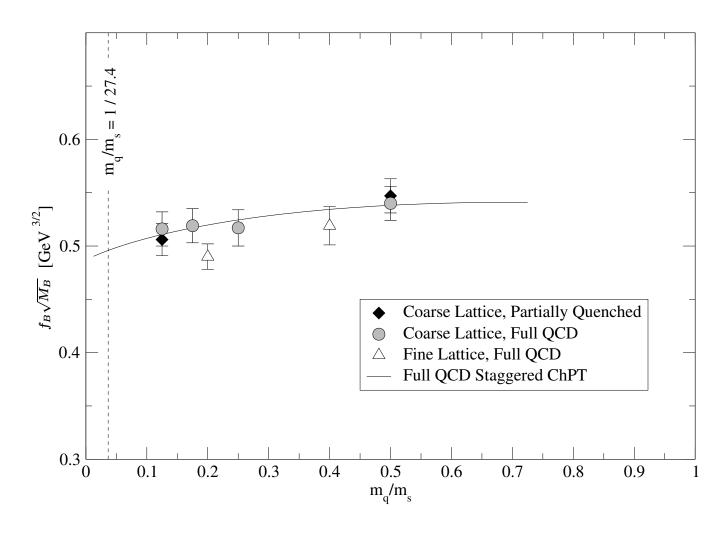
Allows for calculation of vacuum polarization  $\overline{u}, \overline{d}, \overline{s}$ 



# $f_B$ from Lattice QCD

$$f_B = 216 \pm 9 \pm 19 \pm 4 \pm 6 \text{ MeV}$$

Phys. Rev. Lett. 95, 212001 (2005)

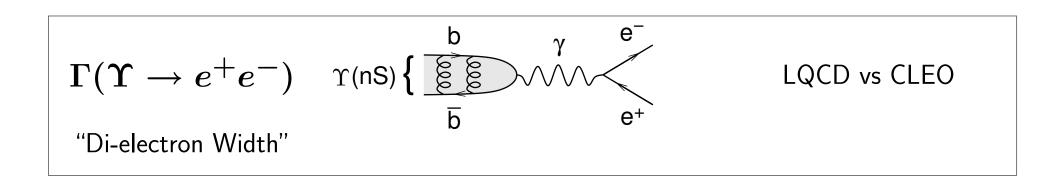


But how can we verify this?

### Test with Processes that Differ by One Quark Flavor

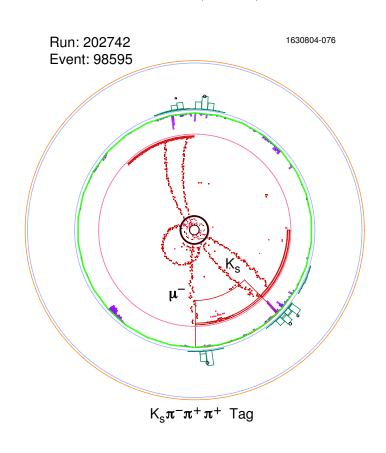
$$f_B$$
 B- $\{$   $\mathbb{R}^-$  LQCD only

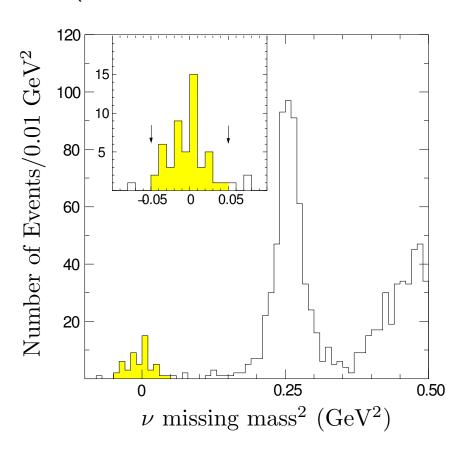
$$f_D$$
  $\mathsf{D}^+\{ egin{array}{c} \mathsf{C} & \mathsf{W}^+ & \ell^- \ \hline \mathsf{d} & \mathsf{V} \ \end{array} \}$  LQCD vs CLEO



## A Brief Look at $f_D$

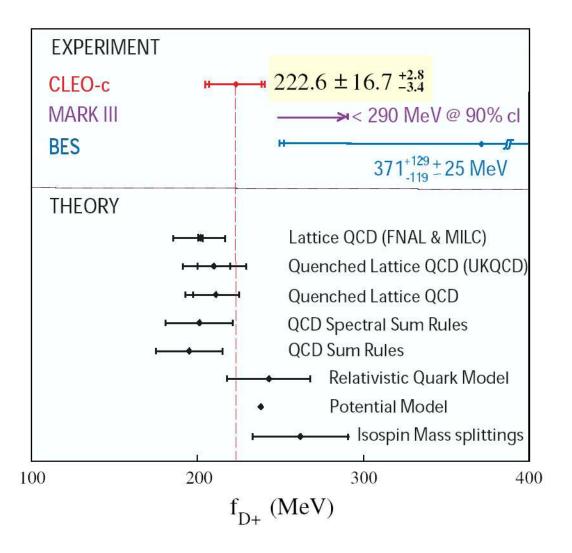
CLEO-c: 281 pb $^{-1}$  at  $\psi(3770)=3$  million  $D^+D^-$  (30 times MARK-III, 8.5 times BES-II)





50 events 
$$-$$
 2.8 background  $=$  47.2  $\pm$  7.1  $^{+0.3}_{-0.8}$   $\mathcal{B}(D^+ \to \mu^+ \nu) = (4.40 \pm 0.66 \, ^{+0.09}_{-0.12}) \times 10^{-4}$   $f_{D^+} = (222.6 \pm 16.7 \, ^{+2.8}_{-3.4}) \; \text{MeV}$ 

### A Brief Look at $f_D$



Projected final precision: 4.5% on  $f_D$  and 4.5% on  $f_{D_s}$ 

# Di-electron widths of $\Upsilon(1S)$ , $\Upsilon(2S)$ , $\Upsilon(3S)$

- Three high-precision measurements (1.5%, 1.8%, and 1.8%)
- Largely share systematics
- See top of screen for an outline

Di-electron width  $\Gamma_{ee}$  = rate of  $\Upsilon \to e^+e^- = \Gamma \times \mathcal{B}_{ee}$ 

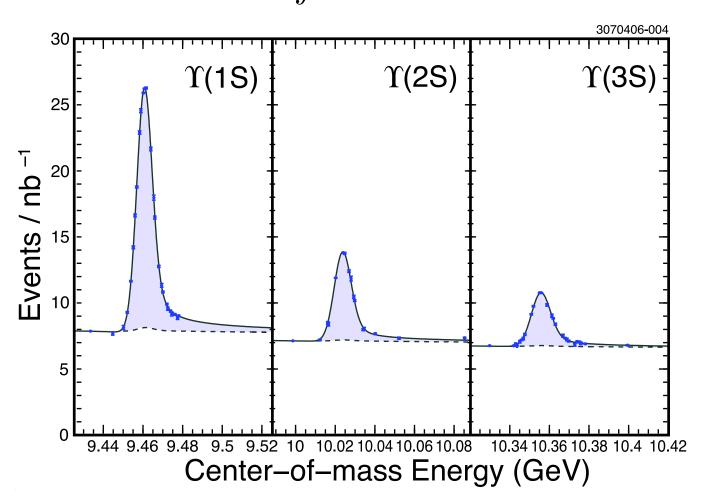
Cannot obtain  $\Gamma_{ee}$  from  $\mathcal{B}_{ee}$  because  $\Gamma$  is unresolvable

Instead, determine 
$$\Upsilon(nS)$$
 {  $b$   $\gamma$   $e^-$  from  $e^+$  from  $e^+$   $b$   $\Upsilon(nS)$ 

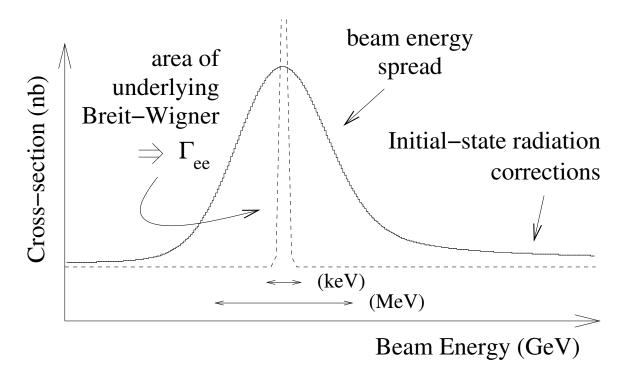
$$\Gamma_{ee} = \frac{M_{\Upsilon}^2}{6\pi^2} \int \sigma(e^+e^- \to \Upsilon) dE$$

Scan  $e^+e^-$  collision energies across  $M_\Upsilon$ , measure cross-section  $\sigma(E)$ , and integrate

$$\Gamma_{ee} = \frac{M_{\Upsilon}^2}{6\pi^2} \int \sigma(e^+e^- \to \Upsilon) dE$$

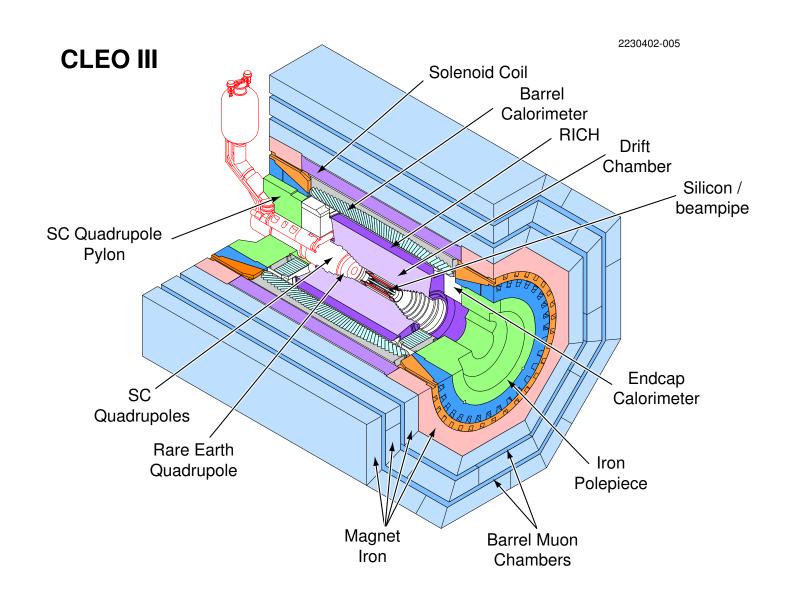


### Anatomy of an \( \gamma \) Lineshape

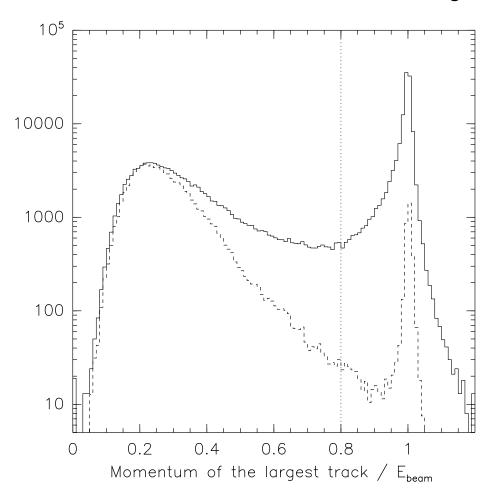


- Breit-Wigner resonance is convoluted by beam energy spread
- Further spread by initial-state radiation  $(e^+e^- \to \gamma \Upsilon)$
- Flat backgrounds

Simulate all effects with a fit function, report Breit-Wigner area only

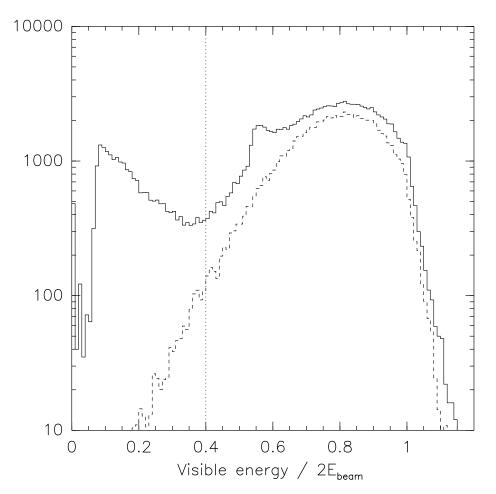


Solid = data, dashed = scaled Monte Carlo, log scale



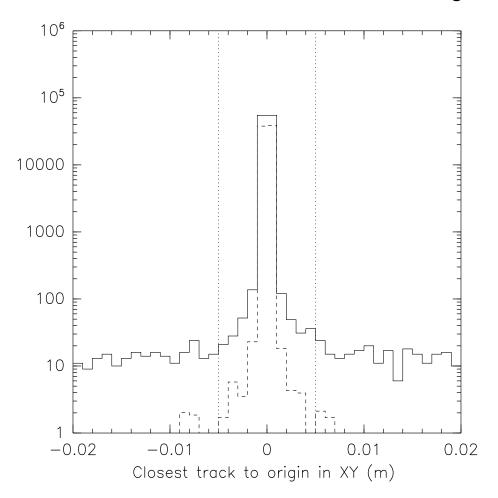
- 1. Bhabhas
- 2.
- 3.
- 4.

Solid = data, dashed = scaled Monte Carlo, log scale



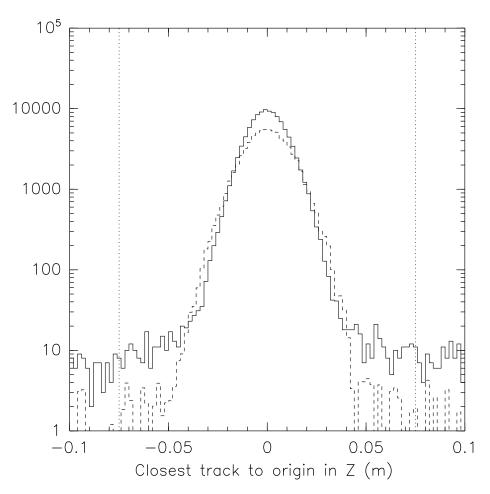
- 1. Bhabhas
- 2. Two-photon fusion
- 3.
- 4.

Solid = data, dashed = scaled Monte Carlo, log scale



- 1. Bhabhas
- 2. Two-photon fusion
- 3. Cosmic rays
- 4.

Solid = data, dashed = scaled Monte Carlo, log scale



- 1. Bhabhas
- 2. Two-photon fusion
- 3. Cosmic rays
- 4. Beam-gas

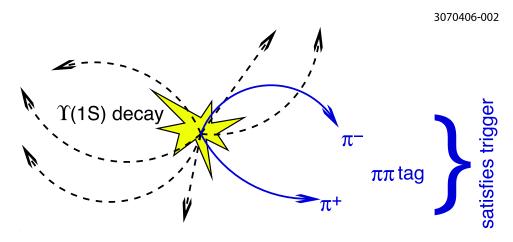
Event selection rejects  $e^+e^-$ ,  $\mu^+\mu^-$ , and 43% of  $\tau^+\tau^-$ 

Correct apparent cross-section with 
$$\frac{1}{1-2.43\mathcal{B}_{\mu\mu}}$$

Define hadronic efficiency = probability that non-leptonic decays pass cuts and trigger

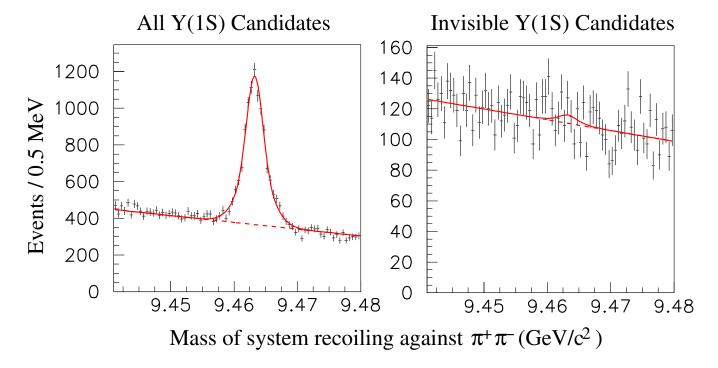
### Data-Derived Efficiency Study

- Select  $\Upsilon(2S) \to \pi^+\pi^-\Upsilon(1S)$  events by the  $\pi^+\pi^-$  only
- ullet Count how many  $\Upsilon(1S)$  decays pass cuts and trigger

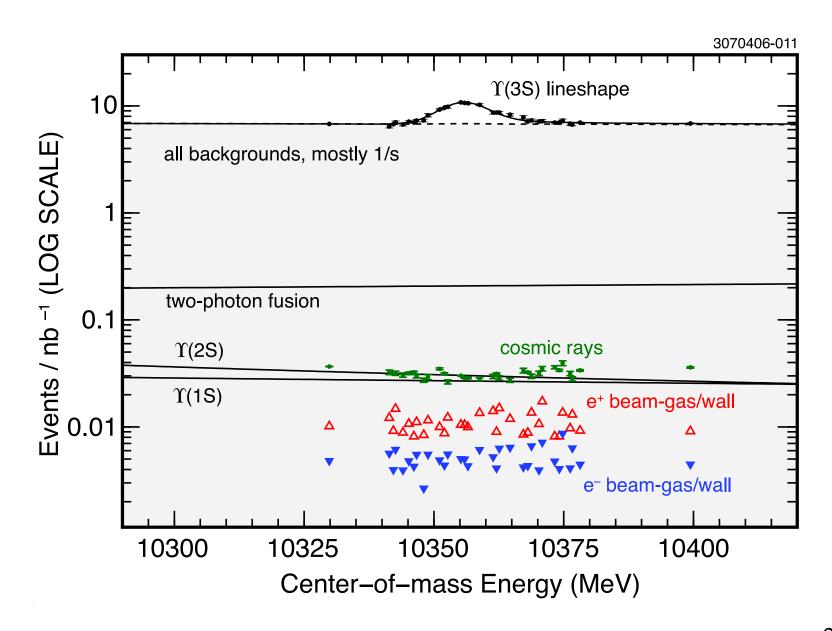


### Data-Derived Efficiency Study

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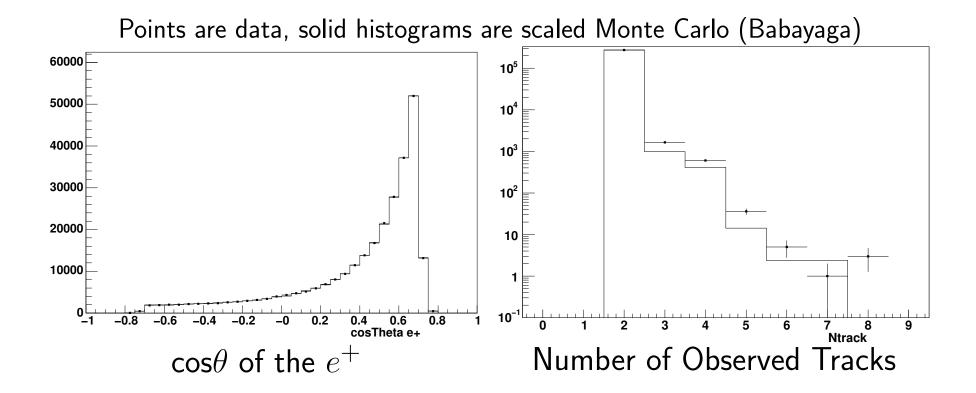
- $\bullet$   $\Upsilon(1S)$  hadronic efficiency is 97.8%  $\pm$  0.5%
- 90% upper limit on invisible  $\Upsilon(1S)$  decays is  $\mathcal{B}_{\mathsf{inv}} < 1.0\%$



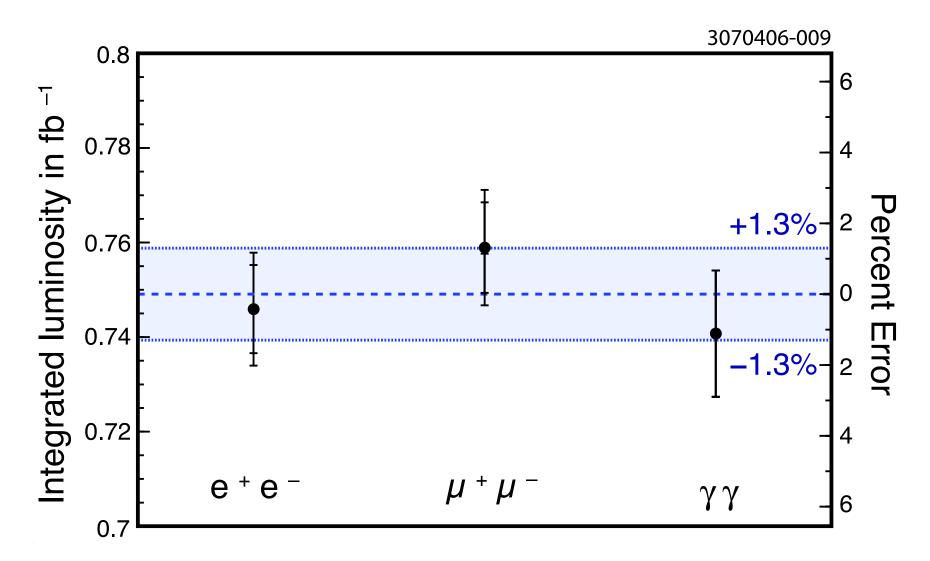
results

theory

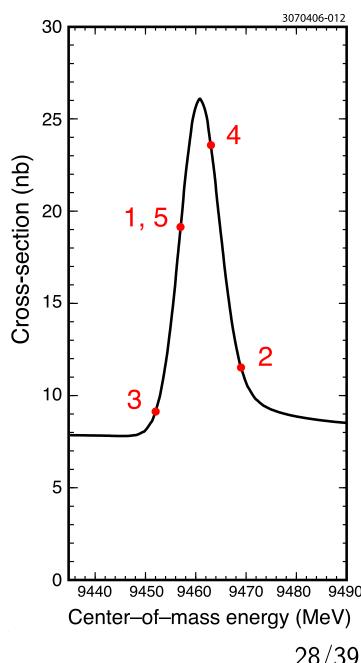
Integrated luminosity = observed Bhabhas / efficiency-weighted Bhabha cross-section

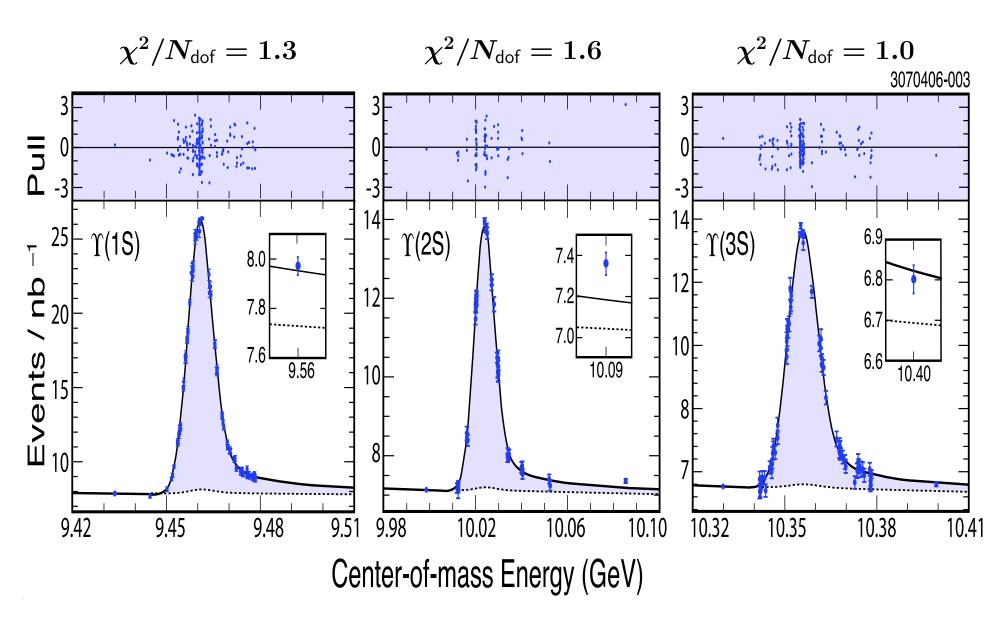


Overall scale of  $e^+e^-$  calculation is checked by  $\mu^+\mu^-$  and  $\gamma\gamma$ 



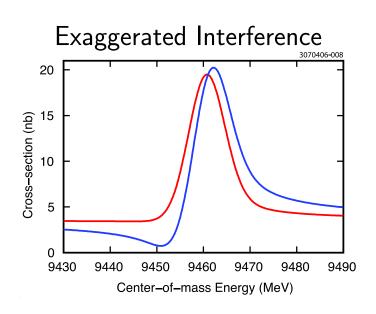
- Beam energy determined by dipole magnet measurement
- Calibration drifts with time (0.5 MeV/month)
- Each resonance completely scanned in 48 hours (repeated scans for statistical precision)
- Measurements alternated above and below resonance peak
- Repeated point of high slope (1 & 5): convert cross-section reproducibility into beam energy reproducibility
- $\bullet \Rightarrow 0.07$  MeV uncertainty in center-of-mass differences, 0.2% in  $\Gamma_{ee}$

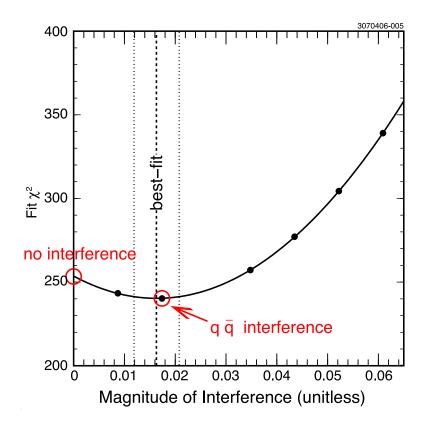




### Lineshape Distortions

- $\bullet$  Non-Gaussian beam energy spread? No, not observed with 0.3% statistical precision
- ullet Variable beam energy spread? Yes, we observed 1% variation in a month
- Interference between  $e^+e^- \to \Upsilon \to \text{hadrons}$  and  $e^+e^- \to \text{hadrons}$ ? Yes!





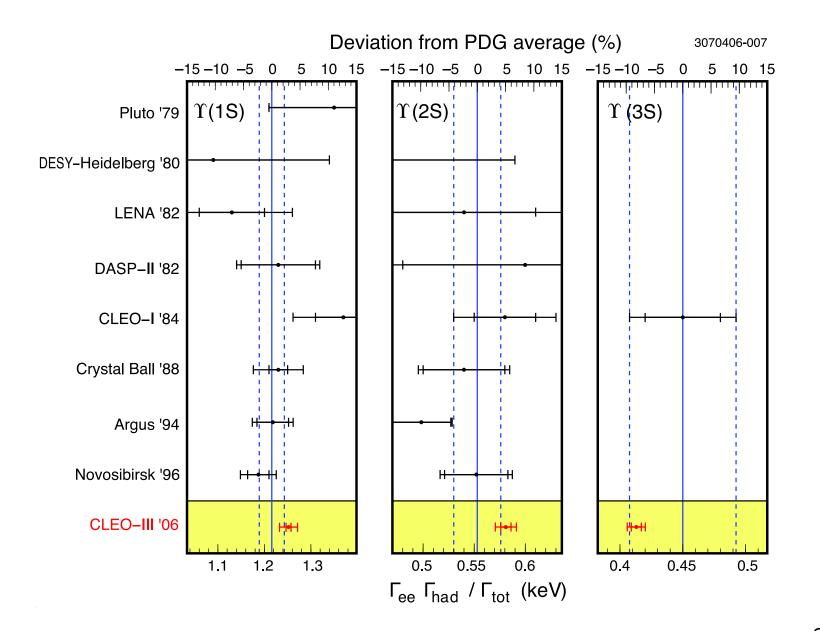
#### \*Common to all resonances

Contribution to $\Gamma_{ee}$	$\Upsilon(1S)$	$\Upsilon(2S)$	$\Upsilon(3S)$
Correction for leptonic modes	0.2%	0.2%	0.3%
Hadronic efficiency*	0.5%	0.5%	0.5%
$Xe^+e^-$ , $X\mu^+\mu^-$ correction	0	0.15%	0.13%
Overall luminosity scale*	1.3%	1.3%	1.3%
Bhabha/ $\gamma\gamma$ inconsistency	0.4%	0.4%	0.4%
Beam energy measurement drift	0.2%	0.2%	0.2%
Fit function shape	0.1%	0.1%	0.1%
$\chi^2$ inconsistency	0.2%	0.6%	0
Total systematic uncertainty	1.5%	1.6%	1.5%
Statistical uncertainty	0.3%	0.7%	1.0%
Total	1.5%	1.8%	1.8%

$\Gamma_{ee}(1S)$	=	$1.354\pm0.004\pm0.020$ keV	1.5%
$\Gamma_{ee}(2S)$	=	$0.619\pm0.004\pm0.010$ keV	1.8%
$\Gamma_{ee}(3S)$	=	$0.446\pm0.004\pm0.007$ keV	1.8%
$\Gamma_{ee}(2S)/\Gamma_{ee}(1S)$	_	$0.457\pm0.004\pm0.004$ keV	1.2%
$\Gamma_{ee}(3S)/\Gamma_{ee}(1S)$	=	$0.329\pm0.003\pm0.003\;{ m keV}$	1.3%
$\Gamma_{ee}(3S)/\Gamma_{ee}(2S)$	=	$0.720\pm0.009\pm0.007\;{ m keV}$	1.6%
$\Gamma(1S)$	=	54.4 $\pm$ 0.2 $\pm$ 0.8 $\pm$ 1.6 keV	3.3%
$\Gamma(2S)$	=	$30.5\pm0.2\pm0.5\pm1.3$ keV	4.6%
$\Gamma(3S)$	=	$18.6\pm0.2\pm0.3\pm$ $0.9$ keV $0.9$ keV	5.2%

 $\Gamma_{ee}$ : J.L. Rosner et~al. (CLEO Collaboration) Phys. Rev. Lett. **96**, 092003 (2006)

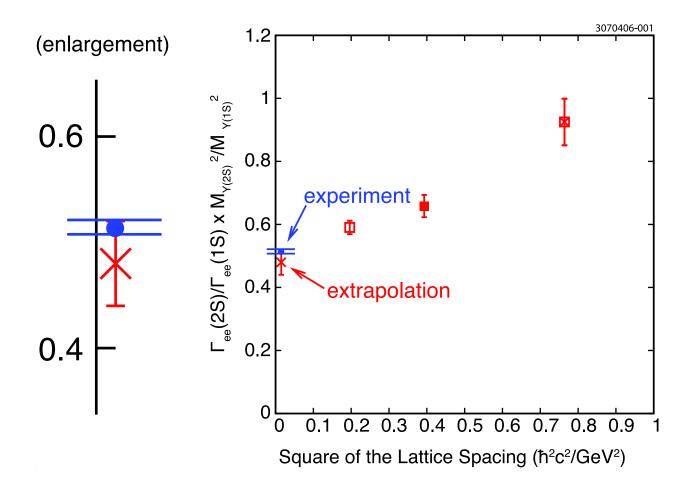
 $\mathcal{B}_{\mu\mu}$ : G.S. Adams et~al. (CLEO Collaboration), Phys. Rev. Lett. 94, 012001 (2005)



### Lattice QCD Calculations of $\Gamma_{ee}$

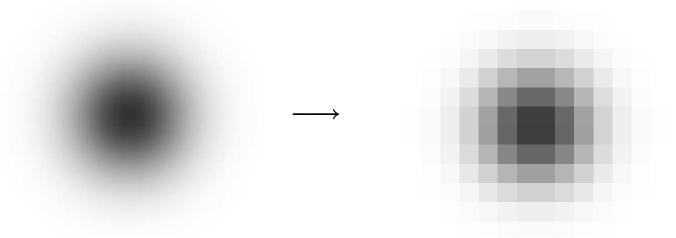
Summer 2005: 10%-level prediction of  $\Gamma_{ee}(2S)/\Gamma_{ee}(1S)$  ratio

Later this summer: few-percent ratios, 10% absolute  $\Gamma_{ee}$ 



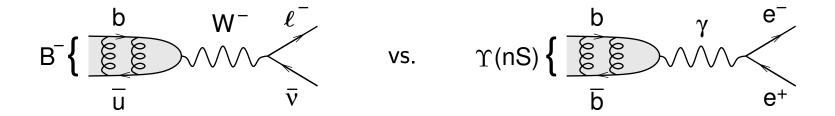
# Why the Steep Dependence on Lattice Spacing?

Decay constants sample wavefunction at the origin, which is discretized



 $\Upsilon$  is a small meson, making discretization more severe

### How Relevant is this Test to $f_B$ ?

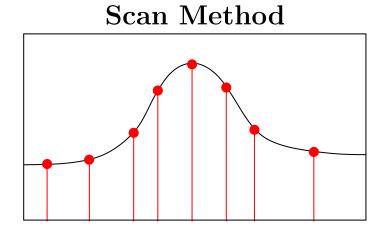


- ullet  $\Gamma_{ee}$  and  $f_B$  calculations share the same NRQCD action and staggered-quark formalism
- ullet Discretization errors in  $f_B$  are smaller than discretization errors in  $\Gamma_{ee}$
- ullet Though  $\Upsilon$  couples to vector current, the factor this introduces cancels in ratios of  $\Gamma_{ee}$

Compliments  $f_D$  and  $\Gamma(\psi \to e^+e^-)$  tests

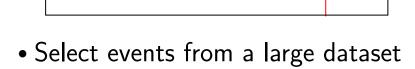
# CLEO has also measured $\Gamma_{ee}$ for $J/\psi$ and $\psi(2S)$

Measured with initial-state radiation, rather than scans



- Requires dedicated scans
- Measure inclusive cross-section
- Requires detailed understanding of the lineshape

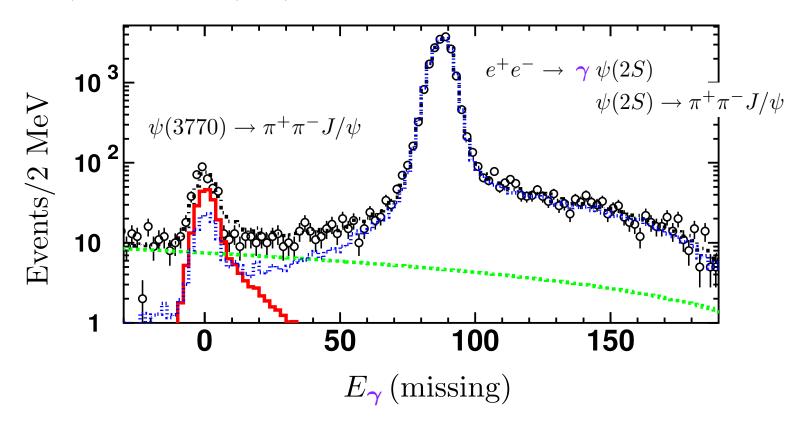




radiate photon

- Distinctive final states  $(\pi^+\pi^-J/\psi)$
- Requires precise knowledge of  $\mathcal{B}(\pi^+\pi^-J/\psi)$

# $\Gamma_{ee}$ for $J/\psi$ and $\psi(2S)$



$$\Gamma_{ee}(J/\psi) = 5.68 \pm 0.11 \pm 0.13 \; ext{keV} \; 3.0\%$$
 $\Gamma_{ee}(\psi(2S)) = 2.54 \pm 0.03 \pm 0.11 \; ext{keV} \; 4.5\%$ 
 $\Gamma_{ee}(\psi(2S))/\Gamma_{ee}(J/\psi) = 0.45 \pm 0.01 \pm 0.02 \; ext{keV} \; 5.0\%$ 

N.E. Adam (CLEO Collaboration) Phys. Rev. Lett. 96 082004 (2006) and

G.S. Adams (CLEO Collaboration) Phys. Rev. D73 051103 (2006)

### Summary

CLEO provides key tests of Lattice QCD relevant for  $f_B$ 

	heavy-heavy	heavy-light
bottom	$\Gamma_{ee}$ for $\Upsilon$	$f_B$
charm	$\Gamma_{ee}$ for $\psi$	$f_D$

 $\Gamma_{ee}$  for  $\Upsilon(1S)$ ,  $\Upsilon(2S)$ ,  $\Upsilon(3S)$ , and ratios have better than 2% precision

Few-percent calculations of  $\Gamma_{ee}$  from Lattice QCD are expected this summer

If they compare favorably, we can have greater confidence in  $V_{td}$  determinations and our understanding of QCD in general