Shaping the chameleon

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In conversation with C. Burrage

Contents

IIntroduction 1

I. INTRODUCTION

$$\phi = \begin{cases} \phi_{c} & 0 < r < S, \\ \frac{\rho r^{2}}{6m} + \frac{c}{r} + d & S < r < R, \\ \frac{\alpha}{r} + \phi_{\infty} & R < r \end{cases}$$
(1.1a)

Matching value and derivative at r = S:

$$\phi_{\rm c} = \frac{\rho S^2}{6m} + \frac{c}{S} + d, \qquad 0 = \frac{\rho S}{3m} - \frac{c}{S^2}$$
 (1.2)

Matching value and derivative at r = R:

$$\frac{\rho R^2}{6m} + \frac{c}{R} + d = \frac{\alpha}{R} + \phi_{\infty}, \qquad \frac{\rho R}{3m} - \frac{c}{R^2} = -\frac{\alpha}{R^2}. \tag{1.3}$$

We thus combine to obtain

$$c = \frac{\rho S^3}{3m}, \qquad \alpha = \frac{\rho}{3m} \left(S^3 - R^3 \right), \qquad d = \phi_c - \frac{\rho S^2}{2m}.$$
 (1.4)

Now, let all quantities aguire an angular perturbation,

$$X \to X_0 + \varepsilon X_1(\theta). \tag{1.5}$$

Collecting like powers of ε yields

$$\frac{\alpha_1}{\alpha_0} = \frac{\rho_1}{\rho_0} + \frac{3}{1 - \left(S_0/R_0\right)^3} \left[\frac{R_1}{R_0} - \left(\frac{S_0}{R_0}\right)^2 \frac{S_1}{R_0} \right]. \tag{1.6}$$

Note that α will be the magnitude of the radial component of the chameleon force outside the source. Note that if $S_0 \approx R_0$ then the denominator above will become very small, meaning that α_1/α_0 could become rather large.

The volume of the thin shell is

$$V = \frac{4}{3}\pi \left(R^3 - S^3\right). \tag{1.7}$$

Hence.

$$\mathcal{V} = \mathcal{V}_0 + 4\pi R_0^3 \left[\frac{R_1}{R_0} - \left(\frac{S_0}{R_0} \right)^2 \frac{S_1}{R_0} \right]$$
 (1.8)

Does this make sense? volume is an integrated quantity, and here we have angular-dependent terms on the RHS

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