

Feynman diagrams in equiv. enumerative geometry

Enumerative geometry is the study of counting geometric objects.

For me, this is counting certain curves into a special kind of variety

In particular K-theory, which is an invariant much like homology, allows us to use algebraic topology to study these counts in interesting ways.

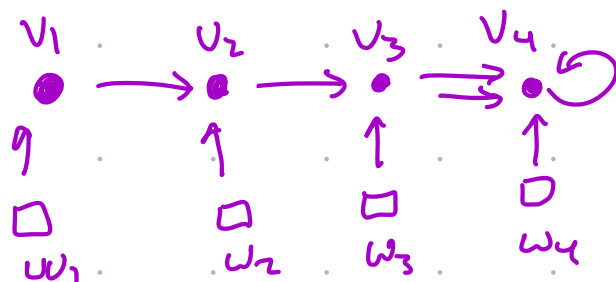
- Plan:
- What kinds of varieties are we going to study?
 - What is the K-theory of these objects & why should you care? (Bethe equations, interesting algebraic actions)
 - What are we counting?
 - How to diagrammatically model these objects (i.e. Feynman diagrams)

Part 1: Nakajima Varieties

In the 90's Hiroku Nakajima discovered an amazing connection between QFT & certain geometric objects arising from graph theory.

Def: A **quiver** is a directed graph

Ex:



Taking this quiver we may assign

$$\begin{aligned} \{ \text{vertices } v_i \} &\longleftrightarrow \{ \text{vector spaces } V_i \mid \dim V_i = v_i \} \\ \{ w_i \} &\longleftrightarrow \{ W_i \mid \dim W_i = w_i \} \\ \{ \text{edges } v_i \rightarrow v_j \} &\longleftrightarrow \{ \text{Hom}(V_i, V_j) \} \end{aligned}$$

$$\text{Rep}_Q(v) = \bigoplus_{i \rightarrow j} \text{Hom}(V_i, V_j) \oplus \bigoplus_i \text{Hom}(W_i, V_i)$$

Quiver Rep Theory studies group actions of these spaces

We want a geometric object. It turns out the way to do this is

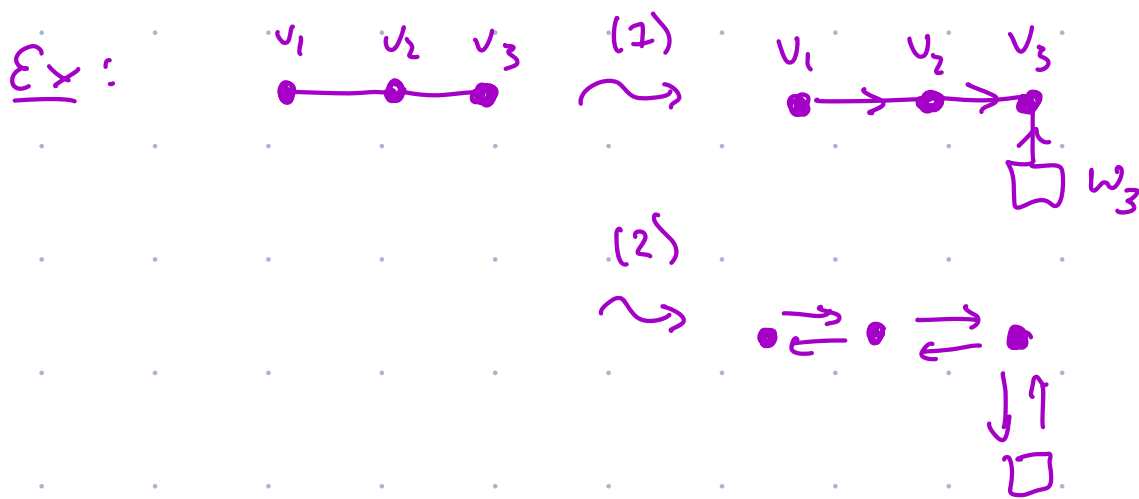
1) Add "framing vertices" and arrows

↳ Vert. space W_i w/ $\dim W_i = n_i$

2) Double all arrows

See Ginzburg's notes for precise reason

Answer: Nonempty quivers & Cotangent bundles



Now the rep space

$$T^* \text{Rep}_q(v, w) = \bigoplus_{i \rightarrow j} \text{Hom}(V_i, V_j) \oplus \bigoplus_i \text{Hom}(W_i, V_i)$$



\oplus duals

Actually cotangent bundle

Group action $G_v = \prod_i GL(v_i)$

$$G_w = \prod GL(w_i)$$

\exists Moment Map

$$\mu: T^* \text{Rep}_k(v, w) \longrightarrow \mathfrak{g}^* = (\text{Lie } G_v)^*$$

Traverse the arrows to get the formula

The last important piece is something called a **Stability parameter**

Again it's a super important part but why would take too long. It's a Map

$$\Theta: G_v \longrightarrow \mathbb{C}^*$$
$$g_i \mapsto \prod \det(g_i)$$

Changing $\Theta \longleftrightarrow$ Different Nakajima Variety

So what is a Nakajima Variety?

"twisted GIT quotient"

$$X = \mu^{-1}(0) //_{\Theta} G_v$$

If you've never seen these before here are some examples:

Ex: $T^*Gr(k, n)$, T^* Partial Flag, $H^1b^1(\mathbb{C}^2)$

$\hookrightarrow w = (0, \dots, 0, n)$

$\theta = \pm (1, \dots, 1)$

There is a torus action $T = \underbrace{A}_{\text{Max torus of Aut}(X)} \times \underbrace{\mathbb{C}_\hbar^\times}_{\text{Scaling symplectic form}} \curvearrowright X$

$(\uparrow, \downarrow) = (\uparrow, \hbar \downarrow)$

Scale opposite arrows

What are we counting?

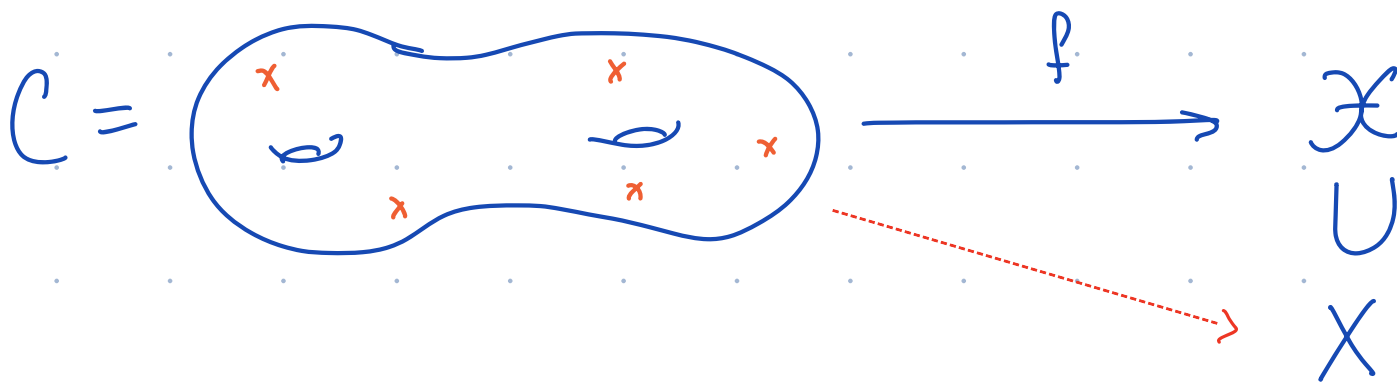
A: Quasimaps

Nakajima Varieties

Stable locus
Smooth, Symplectic

$\subseteq \mathcal{X} = [\text{Rep } G_V / G_V]$
dense open subset

A quasimap



"Def": Maps $f: C \rightarrow X$ are called
 quasihomomorphisms if $f(p) \in X$ for all but finitely many
 pts $p \in C$. Pts $p \in C$ where $f(p) \notin X$
 are called **Singularities of f**

$$QM^d(X) = \left\{ C = \mathbb{P}^1 \xrightarrow{f} X \mid \deg(f) = d \right\}$$

$d \in H_2(X, \mathbb{Z})$

\exists well-defined $\hat{\mathcal{G}}_{\text{vir}}^d \in K_G(QM^d(X))$

We are really in each of
 formulas counting the "number of degree
 d quasihomomorphisms"

$$\chi(\hat{\mathcal{G}}_{\text{vir}}^d \otimes \dots)$$

↑
natural
classes
on $QM^d(X)$

[= analog of $\int (\text{volume form})$
 QM^d in Cohomology]

There's an evaluation map

$$\begin{array}{ccc} \mathcal{QM}^d(X) & \xrightarrow{\text{ev}_p} & \mathcal{X} \\ \downarrow & & \downarrow \\ f & \mapsto & f(p) \end{array}$$

If $\tau \in \mathcal{X}_T(\mathcal{X})$ then $\text{ev}_p^*(\tau) \in K_T(\mathcal{QM}^d(X))$

There are useful shorthand notation for describing
 presimples "Feynman Diagrams"



- Component of source
 curve $C = \mathbb{P}^1$



- $\text{ev}_p^*(\tau)$ descendant
 insertion

$$G = \text{GL}(n)$$

$$\begin{aligned} K([X] = T^*V/G) &\simeq K(\text{pt}/G) \\ &\simeq K_G(\text{pt}) \\ &= \text{Rep}(G) \\ &= \text{Chow}(G) \end{aligned}$$

Another structure

$$QM^d(X)$$

\cup

$$QM_{ns_p}^d(X) = \{ f \in QM^d : f(p) \in X \}$$



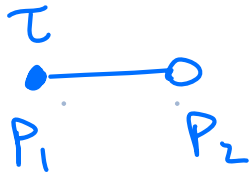
X



$$= ev_{P,*} : K_G(QM_{ns_p}) \rightarrow K_G(X)$$

+ Sum over all degrees

E_X :



$$V^{(\tau)}(z) = \sum_d ev_{P_2,*} (QM_{ns_p}^d, ev_{P_1}^*(\tau) \otimes \hat{Q}_{v,r}^d \otimes \dots) z^d$$

$\in K_{T \times \mathbb{C}_q^*}(X) [z]_{loc}$

Vertex function w/ descendent τ

Important pts:

$$1) \quad G = T \times \mathbb{C}_q^*$$

Relation
of \mathbb{P}^1
Symmetry

2) $ev_{P*} : K(GM_{ns,p}) \rightarrow K(x)$ not defined

it's nonproper. To fix this we need to
use equivariant localization, this is where \mathbb{C}_1^*
comes in.

This is why we land in localized equivariant
K-theory.

Claim: $\begin{array}{c} \circ \text{---} \circ \\ P_1 \quad P_2 \end{array}$ would only get constant germs



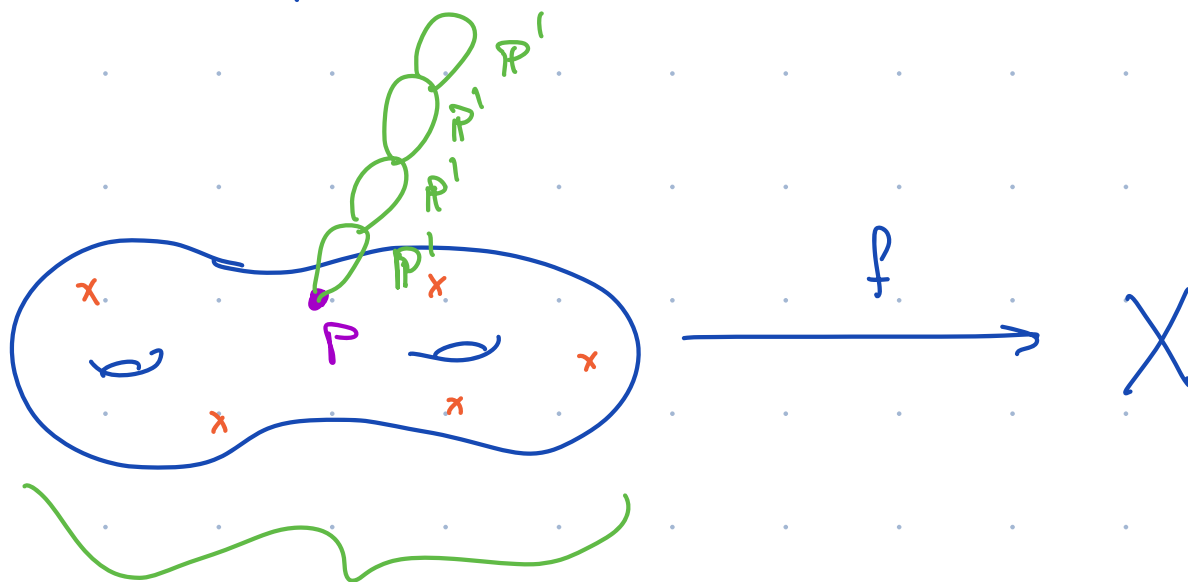
No singularities on $\mathbb{P}^1 \setminus \{P_1, P_2\}$

So we need localization, but the upshot is that this
means $V^{(r)}$ is easy to compute w/ localization

In fact, these are solutions of certain q -difference
equations (q -differential in H_T) & generalize
not ϕ_q hypergeometric functions.

How do we avoid localization?

Relative Quasimaps



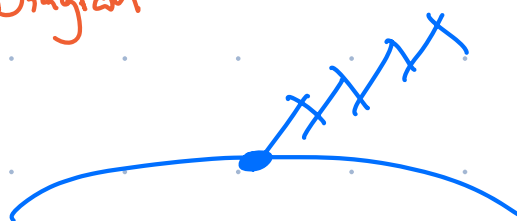
$$\tilde{C} \xrightarrow{\pi} C$$

collapsing the bubbles

$$QM^d_{rel\,p} = \{ \tilde{C} \dashrightarrow X \}$$

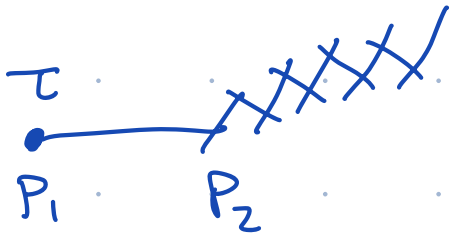
$$\begin{array}{c} \text{proper} \downarrow \text{ev}_p \\ X \end{array}$$

Feynman Diagram



$$:= \text{ev}_{p_x} + \text{sum of all degrees}$$

\mathcal{E}_X :



$$\hat{V}^{(\tau)}_{(z)} := \sum_d \text{ev}_{P_2, *}(Q\mathcal{M}^d(X)_{\text{rel } P_2}, \text{ev}_{P_1}^*(\tau) \otimes \hat{Q}_{\text{vir}}^{\dagger} \otimes \dots) z^d$$

$$\in K_T(X)[[z]]$$

nonlocalized K -theory

This is very difficult to compute since we don't have localization, and is the main focus of my research
Many deep properties:

Conjecture: This is a rational function

Thm [Pandharipande-Pixton, '13]: This is rational for the
Cohomological version in the case of certain stable pairs

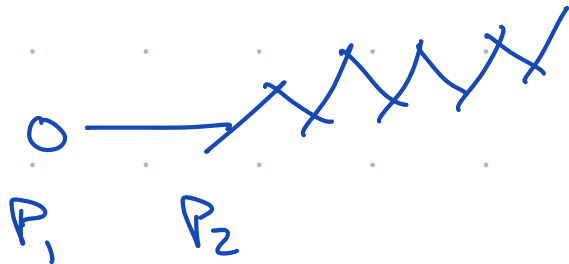
Thm [Smirnov, '16]: This is true for $X = \mathcal{M}(n, r)$

Thm [A-Smirnov, '24] This is true for $X = \text{Hilb}^n(\mathbb{C}^2)$
and compute an explicit combinatorial formula for this $\hat{V}^{(\tau)}$

Then [A - Dinkins, To appear soon] This is true for

$X = \text{Hilb}^n(\mathbb{C}^2/\mathbb{Z}_n\mathbb{Z})$ and compute an explicit combinatorial formula for $\widehat{V}^{(n)}$

Ex:



$$\sum_d \text{ev}_{P_1, *}\otimes \text{ev}_{P_2, *} \left(\underbrace{\mathcal{QM}^d(X)}_{\substack{\text{ns } P_1 \\ \text{rel } P_2}}, \widehat{\mathcal{O}}_{\text{vir}}^d \otimes \dots \right) z^d \in K_T^{\otimes 2}(X)_{\text{loc}}[[z]]$$

$\Psi(a, z)$

"Capping operator"

$\Psi(a, z)$ is the fund. solution matrix of q -diff equations

in both a & z

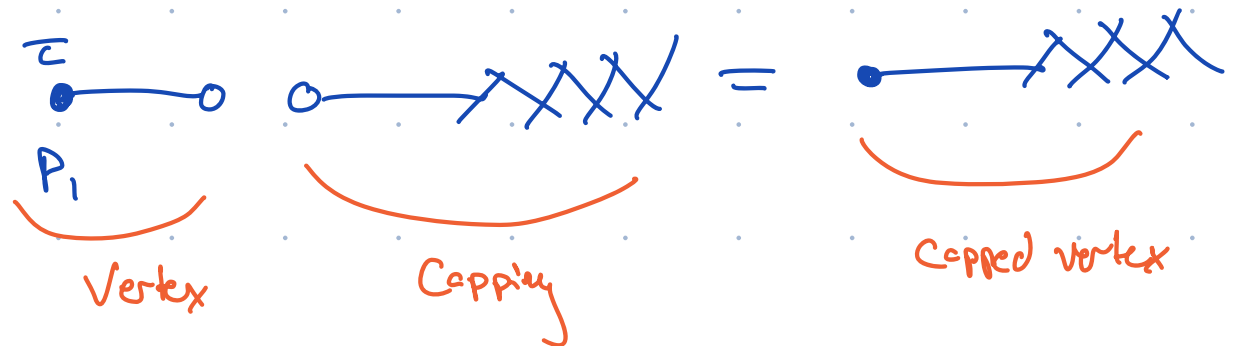
$q(z) \leftarrow \rightarrow$ dynamical equations

" a is a spectral parameter arising from RT"

There is a nice symmetry between a & z called

3D-Mirror Symmetry

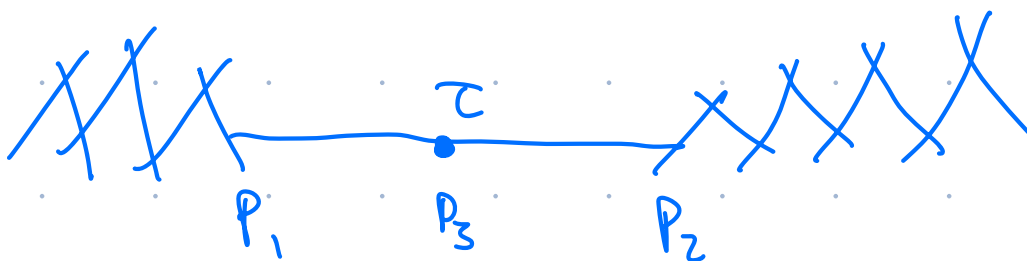
Why are diagrams useful? Bare vertex is easy to compute, capped isn't b/c no localization but



\Rightarrow

Very Very important thm in K-theoretic Enum. Geo proven by Okounkov. It connects quiver geometry, geometric rep theory & Mathematical physics into one equation.

Ex:



$$G = \sum_d \text{ev}_{P_1, x} \otimes \text{ev}_{P_2} (QM_{\text{rel } P_1, P_2}^d(x), \text{ev}_{P_3}^*(\tau) \otimes \hat{\omega}_{\text{irr}}^d \otimes \dots) z^d$$

"Glue Matrix"

Important Connection to RT: $G = \tau * -$

quantum multiplication by τ , these operators

commute & form Bethe algebra

In particular they're $H_\tau =$ Hamiltonian whose eigenvalues are controlled by τ for an XXZ-spin chain

Integrable Spin chains & Connections to geometry

$$U_q(\hat{\mathfrak{g}})$$

quantum
loop group

$$V(a)$$

fund rep of $U_q(\hat{\mathfrak{g}})$

$$\text{Fock} = V(a_1) \otimes \dots \otimes V(a_n)$$

\exists system of commuting Hamiltonians $H_1, H_2, \dots \curvearrowright \text{Fock}$

$$\text{Fock} = \bigoplus_k \text{Fock}_k$$

\uparrow weight k subspace

$$\text{Fock}_k \simeq K_T(X_k)$$

stable
envelope

X_k is a Nakajima variety

Under stable envelopes operators of quantum mult
 \longleftrightarrow Hamiltonians