Mètodes Numèrics 2ºn

20 de Gener de 2016

Nom i cognoms:

Problema 1.

(4.0 punts)

Considerem el problema $-\Delta u = f$ en $\Omega = (0,4) \times (0,2)$, amb les condicions de contorn:

$$u(0,y) = 0$$
, per a tot $y \in (0,2)$.

$$\frac{\partial u}{\partial y}(x,0) = 0$$
 i $\frac{\partial u}{\partial y}(x,2) = 2$, per a tot $x \in (0,4)$.

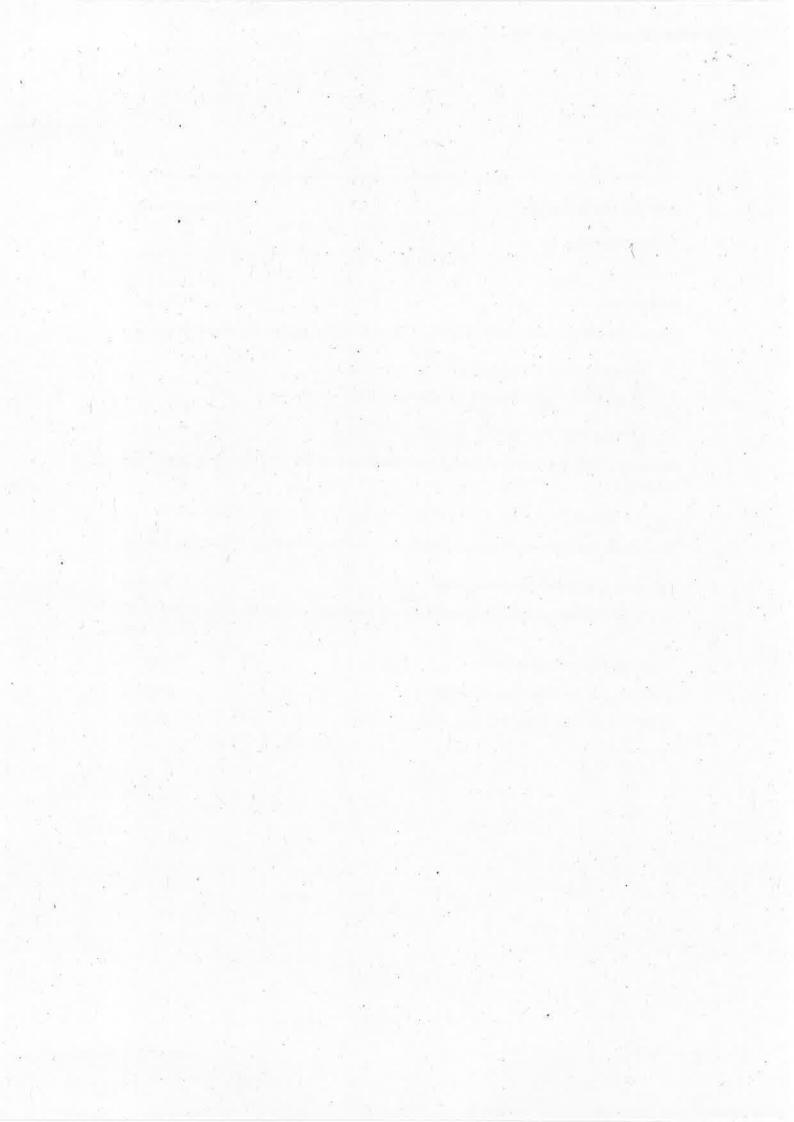
$$\frac{\partial u}{\partial x}(4,y) \text{ lineal en } y \in (0,2) \text{ , } \frac{\partial u}{\partial x}(4,0) = 0 \text{ i } \frac{\partial u}{\partial x}(4,2) = 2.$$

El resolem usant tres elements finits lineals triangulars Ω^1 , Ω^2 i Ω^3 que tenen els vèrtexs següents:

$$\Omega^1 = \{(0,0), (4,0), (4,1)\}, \quad \Omega^2 = \{(0,0), (4,1), (0,2)\}, \quad \Omega^3 = \{(4,1), (4,2), (0,2)\}.$$

Globalment, enumerem els nodes: $p_1 = (0,0)$, $p_2 = (0,2)$, $p_3 = (4,0)$, $p_4 = (4,1)$ i $p_5 = (4,2)$.

- (a) Escriviu la matriu de connectivitat.
- (b) Trobeu les matrius de rigidesa locals i, per a f constant, els vectors de càrregues locals. (1.0 punt)
- (c) Escriviu el sistema acoblat.
- (d) Expliciteu les condicions de contorn. (1.0 punt)
- (e) Si f = 1 i $U_3 = \frac{32084}{1683}$, trobeu U_4 i U_5 .



$$\Delta u = -\frac{\partial^2 u}{\partial x} - \frac{\partial^2 u}{\partial y} = f$$

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(a) Connectivity matrix
$$B = \begin{pmatrix} 134 \\ 142 \\ 452 \end{pmatrix}$$
 modes $= \begin{pmatrix} 0 & 0 \\ 0 & 2 \\ 1 & 0 \\ 4 & 2 \end{pmatrix}$

(6) T3-MN-FEMZD, page 48

521: Linear triangular rectangle element for Poisson equation

$$\frac{q_{11} = q_{22} = c, q_{12} = q_{21} = q_{00} = 0}{a^{2}} = \frac{q_{00} = 0}{a^{2}} = \frac{q_{00} = 0}{a^{2}} = \frac{q_{00} = q_{00}}{a^{2}} = \frac{q_{00}}{a^{2}} = \frac{q_{00}}{a^{2}}$$

SZ: T3-MN-FEMZD page 47.

$$\begin{vmatrix}
0 & 0 \\
4 & 1
\end{vmatrix} : \begin{cases}
\beta_1^2 = \lambda_2^2 + \lambda_3^2 = 1 - 2 = -1 \\
\beta_2^2 = \lambda_3^2 - \lambda_4^2 = 2 - 0 = 2
\end{vmatrix} \begin{cases}
\delta_2^2 = -(\lambda_3^2 - \lambda_4^2) = -(0 - 0) = 0 \\
\beta_3^2 = \lambda_1^2 - \lambda_2^2 = 0 - 1 = -1
\end{cases} \begin{cases}
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\end{cases}$$

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\delta_3 = \lambda_1^2 - \lambda_2^2 - \lambda_2^2 = 0 - 1 = -1
\end{cases} \begin{cases}
\delta_3 = -(\lambda_2^2 - \lambda_2^2) = -(0 - 0) = 0
\end{cases}$$

$$\begin{vmatrix}
\delta_3 = \lambda_1 - \lambda_2 - \lambda_$$

$$= \frac{1}{16} \left[\begin{pmatrix} 1 & -2 & 1 \\ -2 & 4 & 2 \\ 1 & -2 & 1 \end{pmatrix} + \begin{pmatrix} 16 & 0 & -16 \\ 0 & 0 & 0 \\ -18 & 0 & 16 \end{pmatrix} \right] = \frac{1}{16} \begin{pmatrix} 17 & -2 & -15 \\ -2 & 4 & -2 \\ 15 & -2 & 17 \end{pmatrix}$$

13: Same as R1, but now with a= 1, b=4. Therefore:

$$K^{3} = \frac{1}{16} \begin{pmatrix} 32 & -32 & 0 \\ -32 & 34 & -2 \\ 0 & -2 & 2 \end{pmatrix}$$

· Load vectors: T3-MN-FEM2D. pdf page 47, We assume fe f = ctit for all e=1,2,3.

Hence:
$$F^{2} = \frac{f_{2} Ae}{3} (1)$$
, and so $F^{1} = \frac{7}{3} f(1)$, $F^{2} = \frac{7}{3} f(2)$, $F^{3} = \frac{2}{3} f(1)$

(c) Global stiffmess matrix:
$$K = \begin{pmatrix} K_{11} + K_{11}^2 & K_{13}^2 & K_{12}^4 & K_{13}^4 + K_{12}^2 & 0 \\ K_{21} & K_{33}^2 + K_{33}^3 & 0 & K_{32}^2 + K_{31}^3 & K_{32}^2 \\ K_{31} & K_{32} & K_{22}^4 & K_{23}^4 & 0 \\ K_{41} & K_{42} & K_{43} & K_{33}^4 + K_{22}^2 + K_{41}^3 & K_{12}^3 \\ K_{51} & K_{52} & K_{53} & K_{54} \end{pmatrix} \times \begin{pmatrix} K_{54} \\ K_{51} \end{pmatrix} \times \begin{pmatrix} K_{52} \\ K_{53} \end{pmatrix} \times \begin{pmatrix} K_{54} \\ K_{54} \end{pmatrix} \times \begin{pmatrix} K_{52} \\ K_{54} \end{pmatrix} \times \begin{pmatrix} K_{54} \\ K_{54} \end{pmatrix} \times \begin{pmatrix} K_$$

with R=K. Moreover:

$$F = \begin{pmatrix} F_1 \\ F_2 \\ F_3 \\ F_3 \end{pmatrix} = \begin{pmatrix} F_1^4 + F_1^2 \\ F_3^2 + F_3^3 \\ F_2^4 \\ F_3 \end{pmatrix}, Q = \begin{pmatrix} Q_1^4 + Q_1^2 \\ Q_3^2 + Q_3^3 \\ Q_2^4 \\ Q_3^4 + Q_2^2 + Q_1^3 \\ Q_3^3 \end{pmatrix}$$

$$Q = \begin{pmatrix} Q_1^4 + Q_1^2 \\ Q_3^2 + Q_3^3 \\ Q_2^4 \\ Q_3^4 + Q_2^2 + Q_1^3 \\ Q_2^3 \end{pmatrix}$$

So, the coupled system easts

$$\frac{1}{16} \begin{pmatrix}
19 & -15 & -2 & -2 & 0 \\
-15 & 19 & 0 & -2 & -2 \\
-2 & 0 & 34 & -32 & 0 \\
-2 & -2 & -32 & 68 & -32 \\
0 & -2 & 0 & -32 & 34
\end{pmatrix} \begin{pmatrix}
U_1 \\
U_2 \\
U_3 \\
U_4 \\
U_5
\end{pmatrix} = \frac{1}{3} \begin{pmatrix} 6 \\ 6 \\ 2 \\ 8 \\ 2 \end{pmatrix} + \begin{pmatrix} Q_4 \\ Q_2 \\ Q_3 \\ Q_4 \\ Q_5 \end{pmatrix}$$

d) Essential B.C.: U1 = 0, U2=0

$$Q_{M,2}^{1}(s) = \left\langle \begin{pmatrix} a_{M} a_{12} \\ a_{N} a_{22} \end{pmatrix} \nabla u, M \right\rangle = \left\langle \nabla u, M \right\rangle = \left\langle \nabla u, M \right\rangle = \left\langle \frac{\partial u}{\partial x} | a_{12} a_{12} a_{13} a_{14} a_{15} \right\rangle = \left\langle \frac{\partial u}{\partial x} | a_{13} a_{14} a_{15} \right\rangle = \left\langle \frac{\partial u}{\partial x} | a_{15} a_{15}$$

$$\begin{array}{lll}
Q_{M,\Delta}^{3}(s) = \dots = & \langle \nabla u, M \rangle \Big|_{\Gamma_{\Delta}^{3}} = \left(\frac{\partial u}{\partial x} (4, 4+s), \frac{\partial u}{\partial y} (4, 4+s) \right) \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{\partial u}{\partial x} (4, 1+s) = 1+s, \quad 0 \leq s \leq 1 \\
Q_{M,\Delta}^{3}(s) = \dots = & \langle \nabla u, M \rangle \Big|_{\Gamma_{\Delta}^{3}} = \left(\frac{\partial u}{\partial x} (4+s, 2), \frac{\partial u}{\partial y} (4+s, 2) \right) \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \frac{\partial u}{\partial y} (4+s, 2) = 2, \quad 0 \leq s \leq 4
\end{array}$$

$$Q_{12}^{1} = \begin{cases} q_{M2}^{1}(s) & \forall_{21}^{1}(s) & ds = 1 \\ q_{M2}^{1}(s) & \forall_{21}^{1}(s) & ds = 1 \end{cases} \leq s(1-\frac{s}{h_{2}^{1}}) ds = \begin{bmatrix} \frac{1}{6} \end{bmatrix} : see T3-MN-FEM2D. pdf, page 60.$$

Aftermatively (*) $Q_{22}^1 = \frac{1}{3}.0 + \frac{1}{6} \times 1 = \frac{1}{6}$.

$$Q_{32}^{1} = \int_{0}^{h_{2}^{1}} q_{m,2}(s) \, Y_{32}^{1}(s) \, ds = \int_{0}^{h_{2}^{1} = 1} s \cdot \frac{h_{2}^{1} = 1}{s} \cdot Alternatively : Q_{32}^{1} = \frac{1}{3} \times 1 + \frac{1}{6}0 = \frac{1}{3}$$

$$Q_{M}^{3} = \begin{cases} h_{1}^{3} q_{M,1}^{3}(s) \cdot \psi_{11}^{3}(s) ds = \begin{cases} h_{1}^{3} = 1 \\ (1+s)(4-sh_{1}^{3}) ds = \begin{cases} 4-s^{2} ds = (s-s_{3}^{3}) \end{cases}^{1} = 3 \end{cases}$$

Alternatively: $Q_{11}^3 = \frac{1}{3} + \frac{2}{6} = \frac{2}{3}$

$$Q_{21}^{3} = \int_{0}^{h_{1}^{3}} q_{m,1}^{3}(s) \cdot \mathcal{V}_{21}^{3}(s) ds = \int_{0}^{h_{1}^{3}-1} (1+s) \cdot \frac{s}{h_{1}^{3}} ds = \int_{0}^{1} s(1+s) ds = \left(\frac{s}{2} + \frac{s^{3}}{3}\right) \Big|_{0}^{2} = \left[\frac{s}{2} + \frac{s^{3}}{3}\right] = \left[\frac{s}{2} + \frac{s}{3}\right] = \left[\frac{$$

Alternatively: $Q_{24}^3 = \frac{1}{6}1 + \frac{1}{3}Z = \frac{5}{6}$

$$Q_{12}^{3} = \int_{0}^{h_{2}} q_{M,2}(s) \psi_{1,2}(s) ds = \int_{0}^{h_{2}^{3}+4} Z \cdot \left(1 - \frac{s}{h_{2}^{3}}\right) ds = Z \left(\frac{4}{1 - \frac{s}{4}}\right) ds = Z \left(s - \frac{s^{2}}{8}\right) \Big|_{0}^{4}$$

$$= Z \left(4 - \frac{16}{8}\right) = 2(4 - z) = 4 \quad \text{Alternatively} : Q_{12}^{3} = \frac{1}{2} \times 2 \times 4 = 4.$$

(*) Remmark. We apply the "rule for linear flows". For example

$$Q_{33}^{e} = \alpha \left(1 - \frac{s_{e}}{h_{3}}\right) + \beta \frac{s_{e}}{h_{3}} = \alpha \frac{v_{33}^{e}(s)}{v_{33}^{e}(s)} + \beta \frac{v_{43}^{e}(s)}{v_{433}^{e}(s)} + \beta \frac{v_{433}^{e}(s)}{v_{433}^{e}(s)} + \beta \frac{v_{433}^{e}(s)}{v_{433}^{e}(s)} + \beta \frac{v_{433}^{e}(s)}{v_{433}^{e}(s)} + \beta \frac{v_{433}^{e}(s)}{v_{4333}^{e}(s)} + \beta \frac{v_{433$$

$$= \propto h_3 \frac{2!0!}{(2+0+1)!} + \beta h_3 \frac{1!1!}{3!} = h_3^e \left(\frac{\alpha}{3} + \frac{\beta}{6}\right).$$

(24) Let $p,q = D_1 1_1 2_1 3_{1111}$. Therefore $\int_{0}^{\infty} \left(\frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} \right)^{2} ds = \begin{cases} 0, & \text{if either } V_{ml}^{e}(s) = 0 \text{ or } V_{ml}^{e}(s) = 0 \end{cases}$ For, in the 2^{md} case: $\int_{0}^{\infty} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds = \int_{0}^{\infty} \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds + \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds + \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds + \frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \left(\frac{1}{1} \right)^{2} ds + \frac{1}{1} \left(\frac{$

$$Q_{13} = \int_{3}^{h_{3}} q_{13}^{e}(s) \Psi_{13}^{e}(s) ds = \int_{3}^{h_{3}} (x \Psi_{88}^{e}(s) + \beta \Psi_{13}^{e}(s)) \Psi_{18}^{e}(s) ds$$

$$= \alpha \int_{3}^{h_{3}} \Psi_{33}^{e}(s) \Psi_{13}^{e}(s) ds + \beta \int_{3}^{h_{3}} (\Psi_{18}^{e}(s))^{2} ds$$

$$= \alpha \int_{3}^{h_{3}} \Psi_{33}^{e}(s) \Psi_{13}^{e}(s) ds + \beta \int_{3}^{h_{3}} (\Psi_{18}^{e}(s))^{2} ds$$

$$= \alpha \int_{3}^{h_{3}} \Psi_{33}^{e}(s) \Psi_{13}^{e}(s) ds + \beta \int_{3}^{h_{3}} (\Psi_{18}^{e}(s))^{2} ds$$

$$= \alpha \int_{3}^{h_{3}} \Psi_{33}^{e}(s) \Psi_{13}^{e}(s) ds + \beta \int_{3}^{h_{3}} (\Psi_{18}^{e}(s))^{2} ds$$

____ Emd of Remark -

So, the essential B.C. turn out to be,

$$Q_{3} = Q_{22}^{1} = \frac{1}{6}$$

$$Q_{4} = Q_{32}^{1} + Q_{44}^{3} = \frac{1}{3} + \frac{3}{3} = 1$$

$$Q_{5} = Q_{24}^{3} + Q_{22}^{3} = \frac{1}{6} + 4 = \frac{29}{6}$$

(e) Take f=1. If U3 = 32084, we must solve the corresponding reduced system, i.e.:

$$\frac{1}{16} \begin{pmatrix} 68 & -32 \\ -32 & 34 \end{pmatrix} \begin{pmatrix} U_4 \\ U_5 \end{pmatrix} = \frac{1}{3} \begin{pmatrix} 2 \\ 2 \end{pmatrix} + \begin{pmatrix} 1 \\ 21 \\ 6 \end{pmatrix} - \frac{1}{16} \begin{pmatrix} -2 & -2 & -32 \\ 0 & -2 & 0 \end{pmatrix} \cdot \begin{pmatrix} U_4 = 0 \\ U_2 = 0 \\ U_3 = \frac{32084}{1683} \end{pmatrix}$$

$$\Leftrightarrow \begin{pmatrix} 17_4 & -2 \\ -2 & \frac{17}{8} \end{pmatrix} \begin{pmatrix} U_4 \\ U_5 \end{pmatrix} = \begin{pmatrix} \frac{70339}{1683} \\ 1083 \end{pmatrix}$$

$$A = \begin{vmatrix} 17_4 & -2 \\ -2 & \frac{17}{8} \end{vmatrix} = \frac{161}{32}, \quad \Delta U_4 = \begin{vmatrix} \frac{70339}{1683} & -2 \\ \frac{1683}{1683} & -\frac{79051}{1683}, \quad \Delta U_5 = \begin{vmatrix} 17 & \frac{10339}{1683} \\ \frac{17}{4} & \frac{10339}{1683} \end{vmatrix} = \frac{194014}{13464}$$

$$U_4 = \frac{\Delta U_4}{\Delta} = \frac{1964}{99}, \quad U_5 = \frac{\Delta V_5}{\Delta} = \frac{35780}{1683}$$

Finally the complete solution is given by:

$$U_4 = 0$$
, $U_2 = 0$, $U_3 = \frac{32084}{1683}$, $U_4 = \frac{1964}{99}$, $U_5 = \frac{35780}{1683}$.