

# Notes on the Fundamental Equations of PHASTA

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## 1 Shorthand

$$\phi_{,t} = \frac{\partial \phi}{\partial t} \tag{1}$$

$$\phi_{,i} = \frac{\partial \phi}{\partial x_i} \tag{2}$$

$$u_{i,t} = \frac{\partial u_i}{\partial t} \tag{3}$$

$$u_{i,j} = \frac{\partial u_i}{\partial x_j} \tag{4}$$

$$[\phi u_i]_{,j} = \frac{\partial \phi u_i}{\partial x_j} \tag{5}$$

## 2 Fundamental Fluid Equations

Unsteady Compressible Navier-Stokes (UCNS) equations:

### 2.1 Traditional Form

**Continuity**

$$\rho_{,t} + [\rho u_j]_{,j} = 0 \tag{6}$$

**Momentum**

$$[\rho u_i]_{,t} + [\rho u_i u_j]_{,j} + p_{,i} = \tau_{ij,j} + b_i \tag{7}$$

**Energy**

$$[\rho e_{tot}]_{,t} + [\rho e_{tot} u_j]_{,j} + [\rho u_j]_{,j} = [\tau_{ij} u_j]_{,j} + b_i u_j + r + q_{i,i} \tag{8}$$

## 2.2 Conservative Vectorized Form

$$\mathbf{U} \equiv \begin{Bmatrix} \rho \\ \rho u_1 \\ \rho u_2 \\ \rho u_3 \\ \rho e_{tot} \end{Bmatrix} = \begin{Bmatrix} \rho \\ \rho u_j \\ \rho e_{tot} \end{Bmatrix} \quad (9)$$

Flux Vector:

$$\begin{aligned} \mathbf{F}_i &= \underbrace{\begin{Bmatrix} \rho u_i \\ \rho u_i u_j \\ \rho u_i e_{tot} \end{Bmatrix}}_{\text{Advective Flux}} + \underbrace{\begin{Bmatrix} 0 \\ p \delta_{ij} \\ u_i p \end{Bmatrix}}_{\text{Diffusive Flux}} - \underbrace{\begin{Bmatrix} 0 \\ \tau_{ij} \\ \tau_{ij} u_j \end{Bmatrix}}_{\text{Diffusive Flux}} + \underbrace{\begin{Bmatrix} 0 \\ \mathbf{0} \\ q_{i,i} \end{Bmatrix}}_{\text{Diffusive Flux}} \\ &= \mathbf{F}_i^{\text{adv}} + \mathbf{F}_i^{\text{dif}} \end{aligned} \quad (10)$$

Also:

$$\mathbf{F}_i^{\text{adv}} = u_i \mathbf{U} + \begin{Bmatrix} 0 \\ p \delta_{ij} \\ u_i p \end{Bmatrix} \quad (11)$$

Source Vector:

$$\mathcal{F} = \begin{Bmatrix} 0 \\ b_j \\ b_j u_j + r \end{Bmatrix} \quad (12)$$

These terms combine together to form:

$$\mathbf{U}_{,t} + \mathbf{F}_{i,i} = \mathcal{F} \quad (13)$$

## 3 Finite Element Discretization

Before discretizing, we must first setup the UCNS equations. First, rearrange eq. (13) into a residual form:

$$\mathbf{U}_{,t} + \mathbf{F}_{i,i} - \mathcal{F} = \mathbf{0} \quad (14)$$

Next multiply by weight/test functions,  $\mathbf{W}$ . Note that  $\mathbf{0}$  is now just a scalar 0.

$$\mathbf{W} \cdot (\mathbf{U}_{,t} + \mathbf{F}_{i,i} - \mathcal{F}) = 0 \quad (15)$$

To create the **Weak Form** of the NS equations, integrate over the domain  $\Omega$ :

$$\int_{\Omega} \mathbf{W} \cdot (\mathbf{U}_{,t} + \mathbf{F}_{i,i} - \mathcal{F}) \, d\Omega = 0 \quad (16)$$

Next, perform integration by parts and Gauss's theorem on  $\mathbf{F}_{i,i}$ :

$$\int_{\Omega} \{ \mathbf{W} \cdot \mathbf{U}_{,t} - \mathbf{W}_{,i} \cdot \mathbf{F}_i - \mathbf{W} \cdot \mathcal{F} \} d\Omega + \int_{\Gamma} \mathbf{W} \cdot \mathbf{F}_i \cdot \hat{n}_i d\Gamma = 0 \quad (17)$$

where  $\Gamma$  is the boundary of the domain  $\Omega$  and  $\hat{n}_i$  is the normal unit vector of the boundary surface. Equation (17) represents the **Weak UCNS in IBP Form**.

### 3.1 Domain Discretization

Define nodes, points, and elements.

#### 3.1.1 Shape Function Decomposition

Define a set of functions  $N$  that are a basis for the weight functions. In the case of the **Galerkin Form**, the basis/shape functions  $N$  are the same for the weight function and the solution function. We can decompose some constant  $\phi$  into

$$\phi(\mathbf{x}) = \sum_{A=1}^{n_n} N_A(\mathbf{x}) \phi_A \quad (18)$$

where  $A$  is the index of each node and  $\phi_A$  is the value of  $\phi$  at each node  $A$ . This decomposition can be extrapolated to our unknowns:  $\mathbf{U}$ ,  $\mathbf{U}_{,t}$ , and  $\mathbf{W}$ :

$$\mathbf{U}(\mathbf{x}) = \sum_{A=1}^{n_n} N_A(\mathbf{x}) \mathbf{U}_A \quad (19a)$$

$$\mathbf{U}_{,t}(\mathbf{x}) = \sum_{A=1}^{n_n} N_A(\mathbf{x}) \mathbf{U}_{A,t} \quad (19b)$$

$$\mathbf{W}(\mathbf{x}) = \sum_{B=1}^{n_n} N_B(\mathbf{x}) \mathbf{W}_B \quad (19c)$$

Since we are working in Galerkin,  $N_A(\mathbf{x}) = N_B(\mathbf{x})$  for  $A = B$  (ie, at the same node).