Notes on the Fundamental Equations of PHASTA

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1 Shorthand

$$\phi_{,t} = \frac{\partial \phi}{\partial t} \tag{1}$$

$$\phi_{,i} = \frac{\partial \phi}{\partial x_i} \tag{2}$$

$$u_{i,t} = \frac{\partial u_i}{\partial t} \tag{3}$$

$$u_{i,j} = \frac{\partial u_i}{\partial x_j} \tag{4}$$

$$\left[\phi u_i\right]_{,j} = \frac{\partial \phi u_i}{\partial x_j} \tag{5}$$

2 Fundamental Fluid Equations

Unsteady Compressible Navier-Stokes (UCNS) equations:

2.1 Traditional Form

Continuity

$$\rho_{,t} + \left[\rho u_j\right]_{,j} = 0 \tag{6}$$

Momentum

$$[\rho u_i]_{,t} + [\rho u_i u_j]_{,j} + p_{,i} = \tau_{ij,j} + b_i$$
 (7)

Energy

$$[\rho e_{tot}]_{,t} + [\rho e_{tot} u_j]_{,j} + [\rho u_j]_{,j} = [\tau_{ij} u_j]_{,j} + b_i u_j + r + q_{i,i}$$
 (8)

2.2 Conservative Vectorized Form

$$\mathbf{U} \equiv \begin{cases} \rho \\ \rho u_1 \\ \rho u_2 \\ \rho u_3 \\ \rho e_{tot} \end{cases} = \begin{cases} \rho \\ \rho u_j \\ \rho e_{tot} \end{cases}$$
(9)

Flux Vector:

$$\mathbf{F}_{i} = \underbrace{\begin{cases} \rho u_{i} \\ \rho u_{i} u_{j} \\ \rho u_{i} e_{tot} \end{cases}}_{\text{Advective Flux}} + \underbrace{\begin{cases} 0 \\ p \delta_{ij} \\ u_{i} p \end{cases}}_{\text{Diffusive Flux}} + \underbrace{\begin{cases} 0 \\ \mathbf{0} \\ q_{i,i} \end{cases}}_{\text{Diffusive Flux}} \tag{10}$$

$$= \mathbf{F}_{i}^{\text{adv}} + \mathbf{F}_{i}^{\text{dif}}$$

Also:

$$\mathbf{F}_{i}^{adv} = u_{i}\mathbf{U} + \begin{cases} 0 \\ p\delta_{ij} \\ u_{i}p \end{cases}$$
 (11)

Source Vector:

$$\mathfrak{F} = \begin{cases} 0 \\ b_j \\ b_j u_j + r \end{cases} \tag{12}$$

These terms combine together to form:

$$\boldsymbol{U}_{,t} + \boldsymbol{F}_{i,i} = \boldsymbol{\mathcal{F}} \tag{13}$$

3 Finite Element Discretization

Before discritizing, we must first setup the UCNS equations. First, rearrange eq. (13) into a residual form:

$$\boldsymbol{U}_{,t} + \boldsymbol{F}_{i,i} - \boldsymbol{\mathcal{F}} = \boldsymbol{0} \tag{14}$$

Next multiply by weight/test functions, W. Note that 0 is now just a scalar 0.

$$\boldsymbol{W} \cdot (\boldsymbol{U}_{,t} + \boldsymbol{F}_{i,i} - \boldsymbol{\mathcal{F}}) = 0 \tag{15}$$

To create the **Weak Form** of the NS equations, integrate over the domain Ω :

$$\int_{\Omega} \boldsymbol{W} \cdot (\boldsymbol{U}_{,t} + \boldsymbol{F}_{i,i} - \boldsymbol{\mathcal{F}}) \, \mathrm{d}\Omega = 0$$
 (16)

Next, perform integration by parts and Gauss's theorem on $F_{i,i}$:

$$\int_{\Omega} \left\{ \boldsymbol{W} \cdot \boldsymbol{U}_{,t} - \boldsymbol{W}_{,i} \cdot \boldsymbol{F}_{i}(\boldsymbol{U}, \boldsymbol{U}_{,i}) - \boldsymbol{W} \cdot \boldsymbol{\mathcal{F}}(\boldsymbol{U}) \right\} d\Omega + \int_{\Gamma} \boldsymbol{W} \cdot \boldsymbol{F}_{i}(\boldsymbol{U}, \boldsymbol{U}_{,i}) \cdot \hat{n}_{i} d\Gamma = 0$$
(17)

where Γ is the boundary of the domain Ω and \hat{n}_i is the normal unit vector of the boundary surface. Equation (17) represents the **Weak UCNS in IBP Form**. The fact that \mathbf{F}_i and \mathbf{F} are functions of \mathbf{U} and $\mathbf{U}_{,i}$ is explicitly shown here.

3.1 Domain Discretization

Define nodes, points, and elements.

3.1.1 Shape Function Decomposition

Define a set of functions N that are a basis for the weight functions. In the case of the **Galerkin Form**, the basis/shape functions N are the same for the weight function and the solution function. We can decompose some constant ϕ into

$$\phi(\boldsymbol{x}) = \sum_{A=1}^{n_n} N_A(\boldsymbol{x})\phi_A \tag{18}$$

where A is the index of each node, ϕ_A is the value of ϕ at each node A, and n_n is the number of nodes used to discretize Ω . Note that ϕ on the LHS is continuos, but ϕ_A are discrete scalar values for every A. This decomposition can be extrapolated to our unknowns: U, $U_{,t}$, and W:

$$\boldsymbol{U}(\boldsymbol{x}) = \sum_{A=1}^{n_n} N_A(\boldsymbol{x}) \boldsymbol{U}_A$$
 (19a)

$$\boldsymbol{U}_{,t}(\boldsymbol{x}) = \sum_{A=1}^{n_n} N_A(\boldsymbol{x}) \boldsymbol{U}_{A,t}$$
 (19b)

$$\boldsymbol{W}(\boldsymbol{x}) = \sum_{B=1}^{n_n} N_B(\boldsymbol{x}) \boldsymbol{W}_B$$
 (19c)

Since we are working in Galerkin, $N_A(\mathbf{x}) = N_B(\mathbf{x})$ for A = B (ie, at the same node).

3.1.2 Discretize UCNS

Let's substitute the expressions in eq. (19) into the weak UCNS form in eq. (17).

$$\int_{\Omega} \left\{ \sum_{B=1}^{n_n} N_B(\boldsymbol{x}) \boldsymbol{W}_B \cdot \sum_{A=1}^{n_n} N_A(\boldsymbol{x}) \boldsymbol{U}_{A,t} - \sum_{B=1}^{n_n} N_{B,i}(\boldsymbol{x}) \boldsymbol{W}_B \boldsymbol{F}_i \left(\sum_{A=1}^{n_n} N_A(\boldsymbol{x}) \boldsymbol{U}_A, \sum_{A=1}^{n_n} N_{A,i}(\boldsymbol{x}) \boldsymbol{U}_A \right) - \sum_{B=1}^{n_n} \boldsymbol{W}_B \boldsymbol{\mathcal{F}} \left(\sum_{A=1}^{n_n} N_A(\boldsymbol{x}) \boldsymbol{U}_A \right) \right\} d\Omega
+ \int_{\Gamma} \sum_{B=1}^{n_n} \boldsymbol{W}_B \cdot \boldsymbol{F}_i \left(\sum_{A=1}^{n_n} N_A(\boldsymbol{x}) \boldsymbol{U}_A, \sum_{A=1}^{n_n} N_{A,i}(\boldsymbol{x}) \boldsymbol{U}_A \right) \hat{n}_i d\Gamma = 0 \quad (20)$$

Note that $U_{,i} = [\sum_{A=1}^{n_n} N_A(\boldsymbol{x}) \boldsymbol{U}_A]_{,i}$, but since \boldsymbol{U}_A is not a function of \boldsymbol{x} (it is only dependent on A), the sum and \boldsymbol{U}_A can be taken out of the brakets, leaving $N_A(\boldsymbol{x})$ to be derived: $\boldsymbol{U}_{,i} = \sum_{A=1}^{n_n} N_{A,i}(\boldsymbol{x}) \boldsymbol{U}_A$.

Since W_B is a vector of constants, we can take it out of the integral, along with it's respective summation:

$$\sum_{B=1}^{n_n} \mathbf{W}_B \left\{ \int_{\Omega} \left\{ N_B(\mathbf{x}) \cdot \sum_{A=1}^{n_n} N_A(\mathbf{x}) \mathbf{U}_{A,t} - N_{B,i}(\mathbf{x}) \mathbf{F}_i \left(\sum_{A=1}^{n_n} N_A(\mathbf{x}) \mathbf{U}_A, \sum_{A=1}^{n_n} N_{A,i}(\mathbf{x}) \mathbf{U}_A \right) - \mathbf{F} \left(\sum_{A=1}^{n_n} N_A(\mathbf{x}) \mathbf{U}_A \right) \right\} d\Omega + \int_{\Gamma} \mathbf{F}_i \left(\sum_{A=1}^{n_n} N_A(\mathbf{x}) \mathbf{U}_A, \sum_{A=1}^{n_n} N_{A,i}(\mathbf{x}) \mathbf{U}_A \right) \hat{n}_i d\Gamma \right\} = 0 \quad (21)$$

$$eq. (21) \Rightarrow \sum_{B=1}^{n_n} \mathbf{W}_B \cdot \mathbf{G}_B = 0 \quad (22)$$

If we assume that W_B is arbitrary, then $G_B = 0$.