

# Notes on the Fundamental Equations of PHASTA

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## 1 Shorthand

$$\phi_{,t} = \frac{\partial \phi}{\partial t}$$

$$\phi_{,i} = \frac{\partial \phi}{\partial x_i}$$

$$u_{i,t} = \frac{\partial u_i}{\partial t}$$

$$u_{i,j} = \frac{\partial u_i}{\partial x_j}$$

$$[\phi u_i]_{,j} = \frac{\partial \phi u_i}{\partial x_j}$$

## 2 Theory

### 2.1 Fundamental Fluid Equations

Unsteady Compressible Navier-Stokes (UCNS) equations:

#### 2.1.1 Traditional (Strong) Form

**Continuity**

$$\rho_{,t} + [\rho u_j]_{,j} = 0 \tag{1}$$

**Momentum**

$$[\rho u_i]_{,t} + [\rho u_i u_j]_{,j} + p_{,i} = \tau_{ij,j} + b_i \tag{2}$$

**Energy**

$$[\rho e_{tot}]_{,t} + [\rho e_{tot} u_j]_{,j} + [\rho u_j]_{,j} = [\tau_{ij} u_j]_{,j} + b_i u_j + r + q_{i,i} \tag{3}$$

**Constitutive Equations**

$$q_i = -\kappa T_{,i} \tag{4}$$

### 2.1.2 Conservative Vectorized Form

$$\mathbf{U} \equiv \begin{Bmatrix} \rho \\ \rho u_1 \\ \rho u_2 \\ \rho u_3 \\ \rho e_{tot} \end{Bmatrix} = \begin{Bmatrix} \rho \\ \rho u_j \\ \rho e_{tot} \end{Bmatrix} \quad (5)$$

Flux Vector:

$$\begin{aligned} \mathbf{F}_i &= \underbrace{\begin{Bmatrix} \rho u_i \\ \rho u_i u_j \\ \rho u_i e_{tot} \end{Bmatrix}}_{\text{Advective Flux}} + \underbrace{\begin{Bmatrix} 0 \\ p \delta_{ij} \\ u_i p \end{Bmatrix}}_{\text{Diffusive Flux}} - \underbrace{\begin{Bmatrix} 0 \\ \tau_{ij} \\ \tau_{ij} u_j \end{Bmatrix}}_{\text{Diffusive Flux}} + \underbrace{\begin{Bmatrix} 0 \\ \mathbf{0} \\ q_{i,i} \end{Bmatrix}}_{\text{Diffusive Flux}} \\ &= \mathbf{F}_i^{\text{adv}} + \mathbf{F}_i^{\text{dif}} \end{aligned} \quad (6)$$

Also:

$$\mathbf{F}_i^{\text{adv}} = u_i \mathbf{U} + \begin{Bmatrix} 0 \\ p \delta_{ij} \\ u_i p \end{Bmatrix} \quad (7)$$

Source Vector:

$$\mathcal{F} = \begin{Bmatrix} 0 \\ b_j \\ b_j u_j + r \end{Bmatrix} \quad (8)$$

These terms combine together to form:

$$\mathbf{U}_{,t} + \mathbf{F}_{i,i} = \mathcal{F} \quad (9)$$

## 2.2 Finite Element Discretization

Before discretizing, we must first setup the UCNS equations. First, rearrange eq. (9) into a residual form:

$$\mathbf{U}_{,t} + \mathbf{F}_{i,i} - \mathcal{F} = \mathbf{0} \quad (10)$$

Next multiply by weight/test functions,  $\mathbf{W}$ . Note that  $\mathbf{0}$  is now just a scalar 0.

$$\mathbf{W} \cdot (\mathbf{U}_{,t} + \mathbf{F}_{i,i} - \mathcal{F}) = 0 \quad (11)$$

To create the **Weak Form** of the NS equations, integrate over the domain  $\Omega$ :

$$\int_{\Omega} \mathbf{W} \cdot (\mathbf{U}_{,t} + \mathbf{F}_{i,i} - \mathcal{F}) d\Omega = 0 \quad (12)$$

Next, perform integration by parts (IBP) and Gauss's theorem on  $\mathbf{F}_{i,i}$ :

$$\int_{\Omega} \{ \mathbf{W} \cdot \mathbf{U}_{,t} - \mathbf{W}_{,i} \cdot \mathbf{F}_i(\mathbf{U}, \mathbf{U}_{,i}) - \mathbf{W} \cdot \mathcal{F}(\mathbf{U}) \} d\Omega + \int_{\Gamma} \mathbf{W} \cdot \mathbf{F}_i(\mathbf{U}, \mathbf{U}_{,i}) \cdot \hat{n}_i d\Gamma = 0 \quad (13)$$

where  $\Gamma$  is the boundary of the domain  $\Omega$  and  $\hat{n}_i$  is the normal unit vector of the boundary surface. Equation (13) represents the **Weak UCNS in IBP Form**. The fact that  $\mathbf{F}_i$  and  $\mathcal{F}$  are functions of  $\mathbf{U}$  and  $\mathbf{U}_{,i}$  is explicitly shown here.

### 2.2.1 Domain Discretization

Define nodes, points, and elements.

**Shape Function Decomposition** Define a set of functions  $N$  that are a basis for the weight functions. In the case of the **Galerkin Form**, the basis/shape functions  $N$  are the same for the weight function and the solution function. We can decompose some constant  $\phi$  into

$$\phi(\mathbf{x}) = \sum_{A=1}^{n_n} N_A(\mathbf{x})\phi_A \quad (14)$$

where  $A$  is the index of each node,  $\phi_A$  is the value of  $\phi$  at each node  $A$ , and  $n_n$  is the number of nodes used to discretize  $\Omega$ . Note that  $\phi$  on the LHS is continuous, but  $\phi_A$  are discrete scalar values for every  $A$ . This decomposition can be extrapolated to our unknowns:  $\mathbf{U}$ ,  $\mathbf{U}_t$ , and  $\mathbf{W}$ :

$$\mathbf{U}(\mathbf{x}) = \sum_{A=1}^{n_n} N_A(\mathbf{x})\mathbf{U}_A \quad (15a)$$

$$\mathbf{U}_t(\mathbf{x}) = \sum_{A=1}^{n_n} N_A(\mathbf{x})\mathbf{U}_{A,t} \quad (15b)$$

$$\mathbf{W}(\mathbf{x}) = \sum_{B=1}^{n_n} N_B(\mathbf{x})\mathbf{W}_B \quad (15c)$$

Since we are working in Galerkin,  $N_A(\mathbf{x}) = N_B(\mathbf{x})$  for  $A = B$  (ie, at the same node).

**Discretize UCNS** Let's substitute the expressions in eq. (15) into the weak UCNS form in eq. (13).

$$\begin{aligned} \int_{\Omega} \left\{ \sum_{B=1}^{n_n} N_B(\mathbf{x})\mathbf{W}_B \cdot \sum_{A=1}^{n_n} N_A(\mathbf{x})\mathbf{U}_{A,t} - \sum_{B=1}^{n_n} N_{B,i}(\mathbf{x})\mathbf{W}_B \mathbf{F}_i \left( \sum_{A=1}^{n_n} N_A(\mathbf{x})\mathbf{U}_A, \sum_{A=1}^{n_n} N_{A,i}(\mathbf{x})\mathbf{U}_A \right) \right. \\ \left. - \sum_{B=1}^{n_n} \mathbf{W}_B \mathcal{F} \left( \sum_{A=1}^{n_n} N_A(\mathbf{x})\mathbf{U}_A \right) \right\} d\Omega \\ + \int_{\Gamma} \sum_{B=1}^{n_n} \mathbf{W}_B \cdot \mathbf{F}_i \left( \sum_{A=1}^{n_n} N_A(\mathbf{x})\mathbf{U}_A, \sum_{A=1}^{n_n} N_{A,i}(\mathbf{x})\mathbf{U}_A \right) \hat{n}_i d\Gamma = 0 \quad (16) \end{aligned}$$

Note that  $\mathbf{U}_{,i} = [\sum_{A=1}^{n_n} N_A(\mathbf{x})\mathbf{U}_A]_{,i}$ , but since  $\mathbf{U}_A$  is not a function of  $\mathbf{x}$  (it is only dependent on  $A$ ), the sum and  $\mathbf{U}_A$  can be taken out of the brackets, leaving  $N_{A,i}(\mathbf{x})$  to be derived:  $\mathbf{U}_{,i} = \sum_{A=1}^{n_n} N_{A,i}(\mathbf{x})\mathbf{U}_A$ .

Since  $\mathbf{W}_B$  is not a function of space ( $\mathbf{x}$ ), we can take it out of the integral, along with its

respective summation:

$$\sum_{B=1}^{n_n} \mathbf{W}_B \left\{ \int_{\Omega} \left\{ N_B(\mathbf{x}) \cdot \sum_{A=1}^{n_n} N_A(\mathbf{x}) \mathbf{U}_{A,t} - N_{B,i}(\mathbf{x}) \mathbf{F}_i \left( \sum_{A=1}^{n_n} N_A(\mathbf{x}) \mathbf{U}_A, \sum_{A=1}^{n_n} N_{A,i}(\mathbf{x}) \mathbf{U}_A \right) - \mathfrak{F} \left( \sum_{A=1}^{n_n} N_A(\mathbf{x}) \mathbf{U}_A \right) \right\} d\Omega + \int_{\Gamma} \mathbf{F}_i \left( \sum_{A=1}^{n_n} N_A(\mathbf{x}) \mathbf{U}_A, \sum_{A=1}^{n_n} N_{A,i}(\mathbf{x}) \mathbf{U}_A \right) \hat{n}_i d\Gamma \right\} = 0 \quad (17)$$

$$eq. (17) \Rightarrow \sum_{B=1}^{n_n} \mathbf{W}_B \cdot \mathbf{G}_B = 0 \quad (18)$$

If we assume that  $\mathbf{W}_B$  is arbitrary, then eq. (18) is true only if  $\mathbf{G}_B = \mathbf{0}$ . Thus,  $\mathbf{G}_B$  can be treated as  $5 \times n_n$  non-linear ODEs.

$$\mathbf{G}_B = \mathbf{G}_B(\tilde{\mathbf{U}}, \tilde{\mathbf{U}}_{,t}), \quad \tilde{\mathbf{U}} = \{\mathbf{U}_A\}$$

### 2.3 Local/Element level Formulation

For brevity, we will abbreviate the inputs of  $\mathbf{F}_i$  and  $\mathfrak{F}$  to their continuous inputs:  $\mathbf{F}_i(\mathbf{U}, \mathbf{U}_i)$  and  $\mathfrak{F}(\mathbf{U})$ .

$$\mathbf{G}_b^e = \int_{\Omega^e} \left\{ N_b^e(\mathbf{x}) \left[ \sum_{a=1}^{n_{en}} N_a^e(\mathbf{x}) \mathbf{U}_{a,t}^e - \mathfrak{F}(\mathbf{U}) \right] - N_{b,i}^e(\mathbf{x}) \mathbf{F}_i(\mathbf{U}, \mathbf{U}_i) \right\} d\Omega^e + \int_{\Gamma^e} N_b^e(\mathbf{x}) \mathbf{F}_i(\mathbf{U}, \mathbf{U}_i) \hat{n}_i d\Gamma^e \quad (19)$$

Notes:

1.  $n_{en}$  is the number of nodes per element
2.  $^e$  is the element ID number ( $e = \{1 \dots n_n\}$ ),
3.  $_a$  is the node ID number relative to the current element,  $a = \{1 \dots n_{en}\}$
4.  $\mathbf{U} = \sum_{a=1}^{n_{en}} N_a^e(\mathbf{x}) \mathbf{U}_a^e$ ,  $\mathbf{U}_i = \sum_{a=1}^{n_{en}} N_{a,i}^e(\mathbf{x}) \mathbf{U}_a^e$
5.  $\Gamma^e \subset \Gamma$ . In words,  $\Gamma^e$  is a subset (small part) of the global  $\Gamma$  (ie. is not the boundary of  $\Omega^e$ )

**Assembly (Local to Global)** The values at the local level must be brought up to the global level, in a process called assembly. This will be denoted by an  $\mathbf{A}$ . For a 1-D problem with 3 elements (and 4 nodes), this takes the form of:

$$\mathbf{G}_B = \mathbf{A}_{e=1}^{n_{el}} \mathbf{G}_b^e \Rightarrow \begin{cases} \mathbf{G}_1 = \mathbf{G}_1^1 \\ \mathbf{G}_2 = \mathbf{G}_2^1 + \mathbf{G}_1^2 \\ \mathbf{G}_3 = \mathbf{G}_2^2 + \mathbf{G}_1^3 \\ \mathbf{G}_4 = \mathbf{G}_2^3 \end{cases} \quad (20)$$

Each  $\mathbf{G}_a^e$  in the array represents a single  $N_a$  shape function.

### 3 Random Notes

RHS and LHS in the code refer to the Newton's method variation (see page 12b and 13 in notes).

RHS is the stabilized residual term:

$$\hat{\mathbf{G}}_B \quad (21)$$

and LHS is the mass matrix

$$\sum \frac{\partial \hat{\mathbf{G}}_B}{\partial \mathbf{Y}} \quad (22)$$

#### 3.1 Term Definition

Term	Definition
rlyi	$\mathbf{A}_i \mathbf{Y}_i$
ri	building of residual RHS (I think)
EGmass	mass matrix, $\frac{\partial \mathbf{G}}{\partial \mathbf{Y}}$

#### 3.2 Meaning of n...

Term	Definition	Source/Relevant Reference
nsd	number of spacial dimensions	common/common.h#343
nflow	number of flow variables (ie. size of $\mathbf{Y}$ )	?
nshape	number of interior element shape functions	common/common.h#444
ngauss	number of interior element integration points	common/common.h#447
npro	number of elements processed in a single call of e3.f	Jansen lecture
npro	number of virtual processors for the current block	common/common/h#586
nen	maximum number of element nodes	common/common.h#341
nQpt	number of quadrature points per element	common/shp4t.f#14
nshl	number of shape functions per element	common/genblkPosix.f#70
nshg	global number of shape functions	common/common.h#354
nenl	number of element nodes for current block	common/common.h#382
nedof	total number of degrees of freedom	common/e3.f#35,344