Spectral Rigid Body Dynamics

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Overview

Rigid Body Dynamics

Lagrangian Mechanics

Standard Collisions

Constraint Based Collisions

Fourier Methods

Rigid Body Dynamics

An approximate model of low energy physics for stiff objects

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An approximate model of low energy physics for stiff objects

Pros:

- + Pretty accurate at human energy scales
- + Good for stiff materials (ie metals, plastics etc.)
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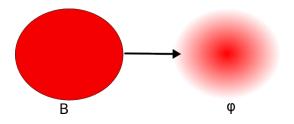
Cons:

- Inaccurate at extremely large energies
- Bad for materials with low elastic modulus
- Not always solvable! (See: Painleve's paradox)

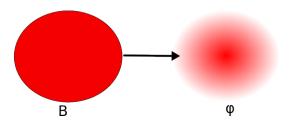


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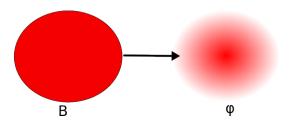


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 φ represents the mass distribution of B $\varphi(x)=0$ indicates B does not occupy the space at x

Transformations of rigid mass fields must preserve distance and handedness

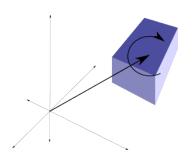
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Can be parameterized by a translation t and a rotation R

Matrix:
$$\begin{pmatrix} R & t \\ 0 & 1 \end{pmatrix}$$

 $\binom{d+1}{2}$ degrees of freedom

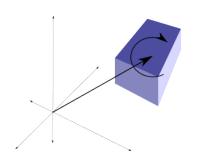
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Motions of rigid objects \cong curves $q(t) \subset SE(d)$



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Given a configuration curve q at time t, define the Lagrangian

$$\mathcal{L}(q,\dot{q},t) = T(\dot{q}) - U(q,t)$$

Where $T(\dot{q}) = \frac{1}{2}\dot{q}^T M \dot{q}$ is the kinetic energy and U is a potential

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Physically plausible motions do minimal work



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$$M \ddot{q} = \nabla U$$

Newton's equations!



Multiple Bodies

Q: How to deal with multiple independent bodies?

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A: Tensor sum

Let B_i, B_j be independent rigid bodies with motions q_i, q_j

Configuration space
$$SE(d)^2 \cong SE(d) \oplus SE(d)$$

Motion $q(t) \cong q_i(t) \oplus q_j(t)$
Lagrangian $L(q,\dot{q},t) = L(q_i,\dot{q}_i,t) + L(q_j,\dot{q}_j,t)$

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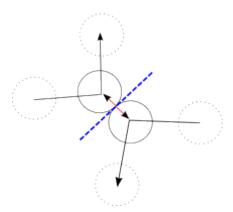
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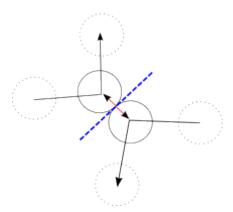
Scales to n bodies, get Lagrangian in $SE(d)^n$

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Need to keep them separated

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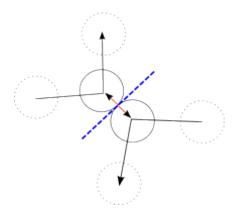


Standard method:

► Time step to point of impact

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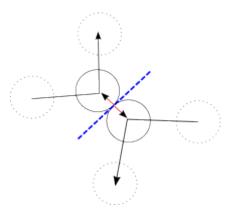
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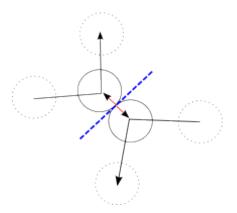
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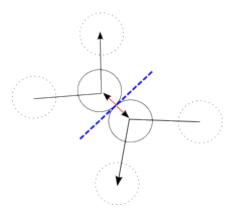
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- Iterate until all collisions are resolved
- + Just like high school physics

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Handling collisions this way is hard

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- Normal forces are ambiguous for curvy shapes
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But can be made to work with enough hacking...

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Minimal requirement for physical plausibility

At all times no two solids intersect

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Is this really all there is to it?

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Define

$$C_{i,j}(q_i,q_j)\stackrel{def}{\equiv} \operatorname{vol} \ q_iA_i\cap q_jA_j$$

And so we replace the impact forces with a system of differentiable holonomic inequality constraints:

$$C_{i,j} \leq 0$$

Equations of motion revisited

New problem:

minimize
$$\int\limits_{t_0}^{t_1}L(q,\dot{q},t)dt$$
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Apply KKT conditions + Euler-Lagrange to get complementarity problem:

$$\frac{d}{dt} \left(\frac{\partial T(\dot{q}_i)}{\partial \dot{q}_i} \right) - \frac{\partial U(q,t)}{\partial q_i} + \sum_{j \neq i} \mu_{i,j} \frac{\partial C_{i,j}(q_i, q_j)}{\partial q_i} = 0$$

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Exactly elastic collision response!

Slack variables are impulse forces



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Substitute $q_j^{-1}q_ix\mapsto R(x-y)$ and let $\widetilde{\mathbf{1}_{A_j}}(x)=\mathbf{1}_{A_j}(-x)$:

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Fourier Methods

Convolution?

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Fourier Methods

Convolution? Take a Fourier transform!

$$C_{i,j}(q_i,q_j) = (\mathbf{1}_{A_i} \star (\widetilde{\mathbf{1}_{A_j}} \circ R))(y) = \mathcal{F}^{-1}\left(\widehat{\mathbf{1}_{A_i}}\overline{\left(\widehat{\mathbf{1}_{A_j}} \circ R^T\right)}\right)$$