Spectral Rigid Body Dynamics

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May 5, 2010

Overview

Rigid Body Dynamics

Lagrangian Mechanics

Standard Collisions

Constraint Based Collisions

Fourier Methods

Rigid Body Dynamics

An approximate model of low energy physics for stiff objects

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- + Pretty accurate at human energy scales
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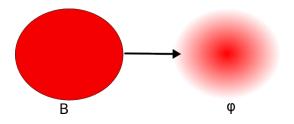
Cons:

- Inaccurate at extremely large energies
- Bad for materials with low elastic modulus
- Not always solvable! (See: Painleve's paradox)

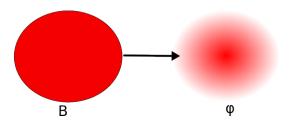


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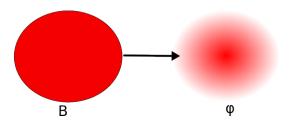


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 φ represents the mass distribution of B $\varphi(x)=0$ indicates B does not occupy the space at x

Transformations of rigid mass fields must preserve distance and handedness

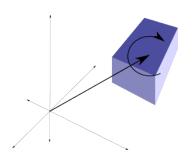
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Can be parameterized by a translation t and a rotation R

Matrix:
$$\begin{pmatrix} R & t \\ 0 & 1 \end{pmatrix}$$

 $\binom{d+1}{2}$ degrees of freedom

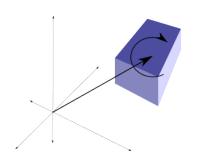
Tangent space: $\mathfrak{so}(d+1)$



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Motions of rigid objects \cong curves $q(t) \subset SE(d)$



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$$\mathcal{L}(q,\dot{q},t) = T(\dot{q}) - U(q,t)$$

Where $T(\dot{q}) = \frac{1}{2}\dot{q}^T M \dot{q}$ is the kinetic energy and U is a potential

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Physically plausible motions do minimal work



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$$M \ddot{q} = \nabla U$$

Newton's equations!



Multiple Bodies

Q: How to deal with multiple independent bodies?

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Let B_i, B_j be independent rigid bodies with motions q_i, q_j

Configuration space
$$SE(d)^2 \cong SE(d) \oplus SE(d)$$

Motion $q(t) \cong q_i(t) \oplus q_j(t)$
Lagrangian $L(q,\dot{q},t) = L(q_i,\dot{q}_i,t) + L(q_j,\dot{q}_j,t)$

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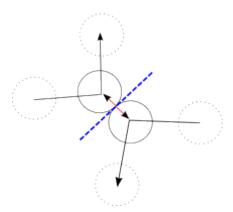
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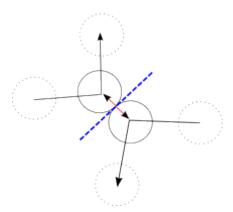
Scales to n bodies, get Lagrangian in $SE(d)^n$

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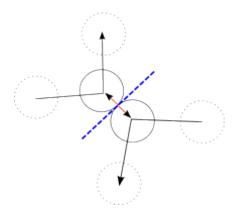


Standard method:

► Time step to point of impact

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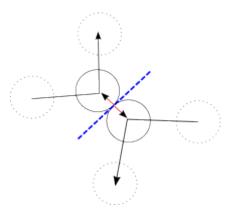
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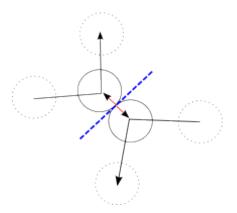
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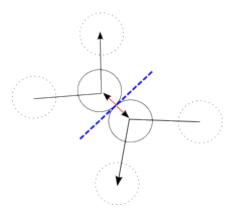
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- + Just like high school physics

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But can be made to work with enough hacking...

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At all times no two solids intersect

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Is this really all there is to it?

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Define

$$C_{i,j}(q_i,q_j)\stackrel{def}{\equiv} \operatorname{vol} \ q_iA_i\cap q_jA_j$$

And so we replace the impact forces with a system of differentiable holonomic inequality constraints:

$$C_{i,j} \leq 0$$

Equations of motion revisited

New problem:

minimize
$$\int\limits_{t_0}^{t_1}L(q,\dot{q},t)dt$$
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Apply KKT conditions + Euler-Lagrange to get complementarity problem:

$$\frac{d}{dt} \left(\frac{\partial T(\dot{q}_i)}{\partial \dot{q}_i} \right) - \frac{\partial U(q,t)}{\partial q_i} + \sum_{j \neq i} \mu_{i,j} \frac{\partial C_{i,j}(q_i, q_j)}{\partial q_i} = 0$$

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Exactly elastic collision response!

Slack variables are impulse forces



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Substitute $q_j^{-1}q_ix\mapsto Rx-y$ and let $\widetilde{\mathbf{1}_{A_j}}(x)=\mathbf{1}_{A_j}(-x)$:

$$\int_{\mathbb{R}^d} \mathbf{1}_{A_i}(x) \widetilde{\mathbf{1}}_{A_j}(y - Rx) dx = \int_{\mathbb{R}^d} \mathbf{1}_{A_i}(R^{-1}x) \widetilde{\mathbf{1}}_{A_j}(y - x) dx$$

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Convolution?

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Fix parameters $q_i = (R_i, t_i)$,

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Need to compute
$$\frac{\partial C_{i,j}(R,y)}{\partial y}$$
, $\frac{\partial C_{i,j}(R,y)}{\partial R}$
Or by symmetry: $C_{i,j}(q_i,q_j) = C_{j,i}(q_j,q_i)$

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Conclusion: Translational gradient is just a multiplier

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\dots = \int_{\mathbb{R}^d} -\left\langle \nabla \widehat{\mathbf{1}_{A_i}} (R^{-1}\omega), R^{-1} \mathfrak{r}_{k,l}\omega \right\rangle \overline{\widehat{\mathbf{1}_{A_j}}(\omega)} e^{2\pi i \langle \omega, y \rangle} d\omega$$

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\dots = \int_{\mathbb{R}^d} -\left\langle \nabla \widehat{\mathbf{1}_{A_i}} (R^{-1}\omega), R^{-1} \mathfrak{r}_{k,l}\omega \right\rangle \overline{\widehat{\mathbf{1}_{A_j}}(\omega)} e^{2\pi i \langle \omega, y \rangle} d\omega$$

Unfortunately not a Fourier multiplier.

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Can be precalculated with O(d) overhead.

In summary...

Unified method for dealing