

# Spectral Rigid Body Dynamics

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# Overview

Rigid Body Dynamics

Lagrangian Mechanics

Standard Collisions

Constraint Based Collisions

Fourier Methods

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An approximate model of low energy physics for stiff objects

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Pros:

- ▶ + Pretty accurate at human energy scales
- ▶ + Good for stiff materials (ie metals, plastics etc.)
- ▶ + Easy kinematic constraints (useful for mechanisms)
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Cons:

- ▶ - Inaccurate at extremely large energies
- ▶ - Bad for materials with low elastic modulus
- ▶ - Not always solvable! (See: Painleve's paradox)

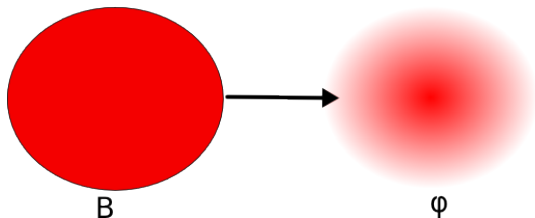
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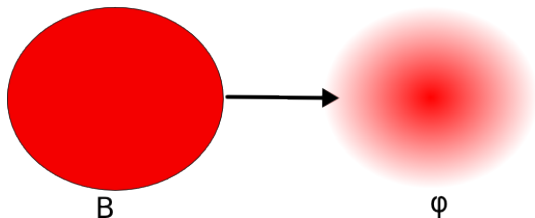
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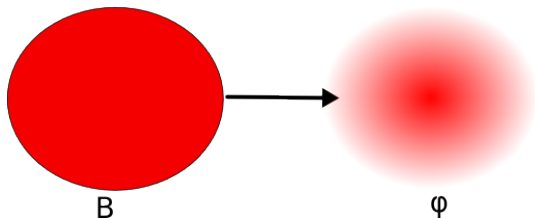
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$\varphi$  represents the mass distribution of  $B$

$\varphi(x) = 0$  indicates  $B$  does not occupy the space at  $x$

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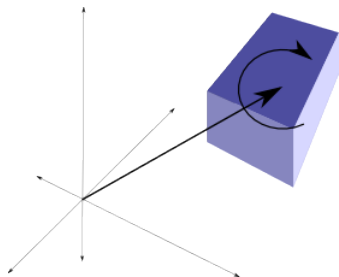
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Can be parameterized by a translation  $t$  and a rotation  $R$

$$\text{Matrix: } \begin{pmatrix} R & t \\ 0 & 1 \end{pmatrix}$$

$\binom{d+1}{2}$  degrees of freedom

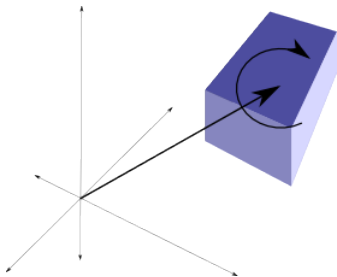
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**Motions of rigid objects  $\cong$  curves  $q(t) \subset SE(d)$**

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$$\mathcal{L}(q, \dot{q}, t) = T(\dot{q}) - U(q, t)$$

Where  $T(\dot{q}) = \frac{1}{2} \dot{q}^T M \dot{q}$  is the kinetic energy and  $U$  is a potential

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$$\operatorname{argmin}_{q: [t_0, t_1] \rightarrow SE(d)} \int_{t_0}^{t_1} L(q, \dot{q}, t) dt$$

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Newton's equations!

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A: Tensor sum

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Configuration space  $SE(d)^2 \cong SE(d) \oplus SE(d)$

Motion  $q(t) \cong q_i(t) \oplus q_j(t)$

Lagrangian  $L(q, \dot{q}, t) = L(q_i, \dot{q}_i, t) + L(q_j, \dot{q}_j, t)$

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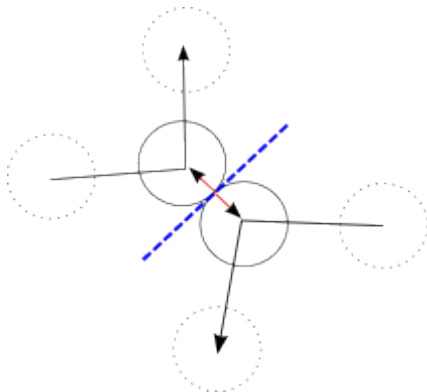
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Scales to  $n$  bodies, get Lagrangian in  $SE(d)^n$



# Collisions

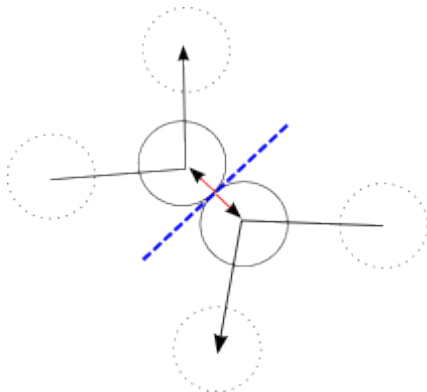
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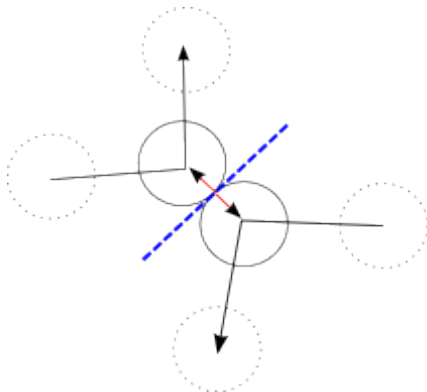
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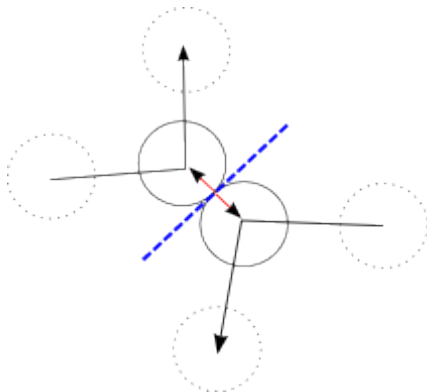
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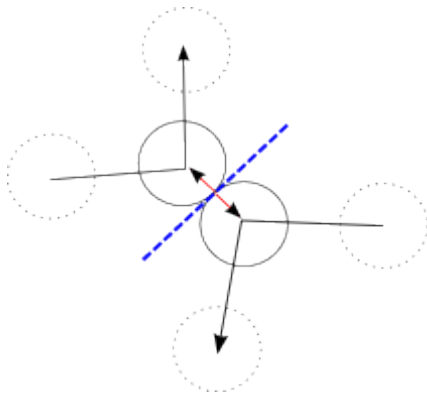
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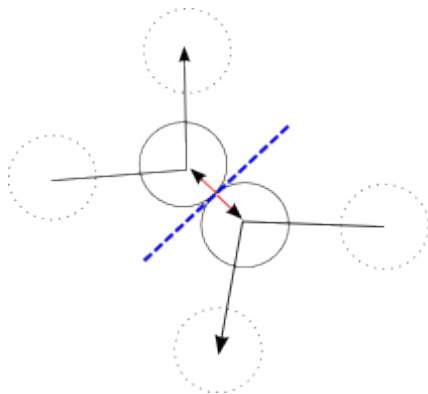
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+ Just like high school physics

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But can be made to work with enough hacking...

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*At all times no two solids intersect*

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Is this really all there is to it?

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Define

$$C_{i,j}(q_i, q_j) \stackrel{\text{def}}{=} \operatorname{vol} q_i A_i \cap q_j A_j$$

And so we replace the impact forces with a system of differentiable holonomic inequality constraints:

$$C_{i,j} \leq 0$$



# Equations of motion revisited

New problem:

$$\text{minimize } \int_{t_0}^{t_1} L(q, \dot{q}, t) dt$$

subject to  $C_{i,j}(q_i, q_j) \leq 0 \quad \forall t \in [t_0, t_1), i \neq j$

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Apply KKT conditions + Euler-Lagrange to get complementarity problem:

$$\frac{d}{dt} \left( \frac{\partial T(\dot{q}_i)}{\partial \dot{q}_i} \right) - \frac{\partial U(q, t)}{\partial q_i} + \sum_{j \neq i} \mu_{i,j} \frac{\partial C_{i,j}(q_i, q_j)}{\partial q_i} = 0$$

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Exactly elastic collision response!

Slack variables are impulse forces

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Need to compute:

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Observe:

$$\mathbf{1}_{A_i \cap A_j}(x) = \mathbf{1}_{A_i}(x) \mathbf{1}_{A_j}(x)$$

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So:

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Substitute  $q_j^{-1}q_i x \mapsto R(x - y)$  and let  $\widetilde{\mathbf{1}}_{A_j}(x) = \mathbf{1}_{A_j}(-x)$ :

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# Fourier Methods

Convolution?

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Convolution? Take a Fourier transform!

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Fix parameters  $q_i = (R_i, t_i)$ ,

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Solve for  $R, y$ ,

$$R = R_j R_i^{-1}$$

$$y = R_i R_j^{-1} t_j - t_i$$

Need to compute  $\frac{\partial C_{i,j}(R_i, t_i, R_j, t_j)}{\partial R_i}$ ,  $\frac{\partial C_{i,j}(R_i, t_i, R_j, t_j)}{\partial t_i}$

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Let  $t_i^k$  denote the  $k^{th}$  component of  $t_i$ , then

$$\begin{aligned}\frac{\partial C_{i,j}}{\partial t_i^k} &= \frac{\partial}{\partial t_i^k} \left( \int_{\mathbb{R}^d} \widehat{\mathbf{1}_{A_i}}(\omega) \overline{\widehat{\mathbf{1}_{A_j}}(R_j R_i^{-1} \omega)} e^{2\pi i \langle \omega, R_i R_j^{-1} t_j - t_i \rangle} d\omega \right) \\ &= \int_{\mathbb{R}^d} 2\pi i \langle \omega, v^k \rangle \widehat{\mathbf{1}_{A_i}}(\omega) \overline{\widehat{\mathbf{1}_{A_j}}(R_j R_i^{-1} \omega)} e^{2\pi i \langle \omega, R_i R_j^{-1} t_j - t_i \rangle} d\omega\end{aligned}$$

Where  $v^k$  denotes the  $k^{th}$  basis vector

# Translational Gradient

Start with the translational case first.

Let  $t_i^k$  denote the  $k^{th}$  component of  $t_i$ , then

$$\begin{aligned}\frac{\partial C_{i,j}}{\partial t_i^k} &= \frac{\partial}{\partial t_i^k} \left( \int_{\mathbb{R}^d} \widehat{\mathbf{1}_{A_i}}(\omega) \overline{\widehat{\mathbf{1}_{A_j}}(R_j R_i^{-1} \omega)} e^{2\pi i \langle \omega, R_i R_j^{-1} t_j - t_i \rangle} d\omega \right) \\ &= \int_{\mathbb{R}^d} 2\pi i \langle \omega, v^k \rangle \widehat{\mathbf{1}_{A_i}}(\omega) \overline{\widehat{\mathbf{1}_{A_j}}(R_j R_i^{-1} \omega)} e^{2\pi i \langle \omega, R_i R_j^{-1} t_j - t_i \rangle} d\omega\end{aligned}$$

Where  $v^k$  denotes the  $k^{th}$  basis vector

Conclusion: Translational gradient is just a multiplier

# Rotational Gradient

Parameterize  $R_i = \exp(\mathfrak{r})$ , where  $\mathfrak{r} \in \mathfrak{so}(d)$

In otherwords  $d \times d$  skew symmetric matrices,  $\mathfrak{r}_{k,l} = -\mathfrak{r}_{l,k}$

# Rotational Gradient

Parameterize  $R_i = \exp(\mathfrak{r})$ , where  $\mathfrak{r} \in \mathfrak{so}(d)$

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$$\frac{\partial C_{i,j}}{\partial \mathfrak{r}_{k,l}} = \frac{\partial}{\partial \mathfrak{r}_{k,l}} \left( \int_{\mathbb{R}^d} \widehat{\mathbf{1}_{A_i}}(\omega) \overline{\widehat{\mathbf{1}_{A_j}}(R_j \exp(-\mathfrak{r})\omega)} e^{2\pi i \langle \omega, \exp(\mathfrak{r})R_j^{-1}t_j - t_i \rangle} d\omega \right)$$

# Rotational Gradient

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# Rotational Gradient

Parameterize  $R_i = \exp(\mathfrak{r})$ , where  $\mathfrak{r} \in \mathfrak{so}(d)$

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Get two terms:

# Rotational Gradient

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Get two terms: **a multiplier (easy)**,

# Rotational Gradient

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Get two terms: a multiplier (easy), a gradient (can be precomputed).

# Rotational Gradient

Parameterize  $R_j = \exp(\mathbf{r})$ , where  $\mathbf{r} \in \mathfrak{so}(d)$

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Get two terms: a multiplier (easy), a gradient (can be precomputed).

Computationally not too bad, but still pretty messy in  $d$ -dimensional space.