Analytical and Numerical Solution in The One-Dimensional Ising Model

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Concept

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1-1. Introduction to Ising Model

A simple but fundamental model that takes interaction only between nearest-neighbor spins into account

In Ising model, there are only up and down spins.

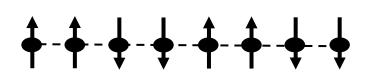


Fig. 1.1. The one-dimensional Ising model

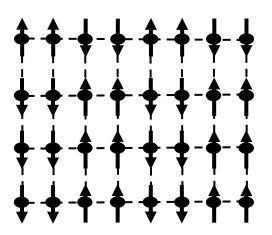


Fig. 1.2. The two-dimensional Ising model









Fig. 1.3. Parallel links

The energy of parallel link is -J

Fig. 1.4. Antiparallel links

The energy of antiparallel link is +J

$$N = N_+ + N_-$$

magnetization
$$M = \mu (N_{+} - N_{-})$$

interaction energy
$$E_i = -J(N_{++} + N_{--} - N_{+-} - N_{-+})$$

1-2. Partition Function

$$\begin{split} E(\{\sigma_i\}) &= -J \sum_{i=1}^N \sigma_i \sigma_{i+1} - \mu B \sum_{i=1}^N \sigma_i & \mu \text{: magnetic moment} \\ & J \text{: coupling constant} \\ & B \text{: external magnetic field} \\ & \beta = 1/k_B T \\ Z &= \sum_{\{\sigma_i\}} e^{-\beta E(\{\sigma_i\})} = \sum_{\sigma_1} \dots \sum_{\sigma_N} \exp[\beta \sum_{i=1}^N J \sigma_i \sigma_{i+1} + \frac{\mu}{2} (\sigma_i + \sigma_{i+1})] \end{split}$$

$$M(\sigma_{i}, \sigma_{i+1}) = \exp\left[\beta \sum_{i=1}^{N} J\sigma_{i}\sigma_{i+1} + \frac{\mu}{2}(\sigma_{i} + \sigma_{i+1})\right], \qquad A = \begin{pmatrix} e^{\beta+h} & e^{-\beta J} \\ e^{-\beta J} & e^{\beta(J+h)} \end{pmatrix}$$

$$Z = \sum_{\sigma} \dots \sum_{\sigma \in \mathcal{I}} A_{\sigma_1, \sigma_2, A_{\sigma_2, \sigma_3}} \dots A_{\sigma_{N-1}, \sigma_N} A_{\sigma_N, \sigma_1} = tr(A^N) = tr(A^N_{diag}) = \begin{pmatrix} \lambda_1 & 0 \\ 0 & \lambda_2 \end{pmatrix}^N = \lambda_1^N + \lambda_2^N$$

$$|A - \lambda I| = 0$$

$$\begin{vmatrix} e^{\beta(J+\mu B)} - \lambda & e^{-\beta J} \\ e^{-\beta J} & e^{\beta(J-\mu B)} - \lambda \end{vmatrix} = 0$$

$$\lambda = e^{\beta J} \cosh(\beta \mu B) \pm \sqrt{e^{2\beta J} \cosh^2(\beta \mu B) - 2\sinh(2\beta J)}$$
if $\lambda_1 > \lambda_2$ then $\lambda_1^N \gg \lambda_2^N$,

$$\mathbf{Z} \approx \lambda_1^{N} = \left[e^{\beta J} \cosh(\beta \mu B) + \sqrt{e^{2\beta J} \cosh^2(\beta \mu B) - 2 \sinh(2\beta J)} \right]^{N}$$

$$for N \to \infty$$

1-3. Energy, Heat Capacity, Magnetization and Susceptibility

Free energy

$$F = -k_B T \ln[Z] = -Nk_B T \ln[e^{\beta J} \cosh(\beta \mu B) + \sqrt{e^{2\beta J} \cosh^2(\beta \mu B) - 2 \sinh(2\beta J)}]$$

Entropy at B = 0

$$F = -k_B T ln[Z] = -Nk_B T ln[e^{\beta J} + e^{-\beta J}] = -Nk_B T ln[2cosh\beta J]$$

$$S = -\frac{\partial F}{\partial T} = Nk_B \left[\ln \left(2\cosh \frac{J}{k_B T} \right) - \frac{J}{k_B T} \left(\tanh \frac{J}{k_B T} \right) \right]$$

$$E = F + ST = -NJ \tanh \frac{J}{k_B T}$$

Heat Capacity at B = 0

$$C_V = \frac{dE}{dT} = \frac{NJ^2}{k_B T^2} \operatorname{sech}^2 \frac{J}{k_B T}$$

Magnetization

$$M = -\frac{\partial F}{\partial B} = \beta \mu N k_B T \frac{e^{\beta J} \sinh \beta \mu B + \frac{e^{2\beta J} \sinh \beta \mu B \cosh \beta \mu B}{\sqrt{e^{2\beta J} \cosh^2 \beta \mu B - 2 \sinh 2\beta J}}}{e^{\beta J} \cosh \beta \mu B + \sqrt{e^{2\beta J} \cosh^2 \beta \mu B - 2 \sinh 2\beta J}}$$

In a weak magnetic field $\sinh \beta \mu B \approx 1$, $\cosh \beta \mu B \approx 1$

$$M \approx N\mu^2\beta \ e^{2\beta J}B = \frac{N\mu^2}{k_BT} \ e^{\frac{2J}{k_BT}}B$$

Susceptibility

$$\chi = \mu_0 \frac{N\mu^2}{k_B T} e^{\frac{2J}{k_B T}} B$$

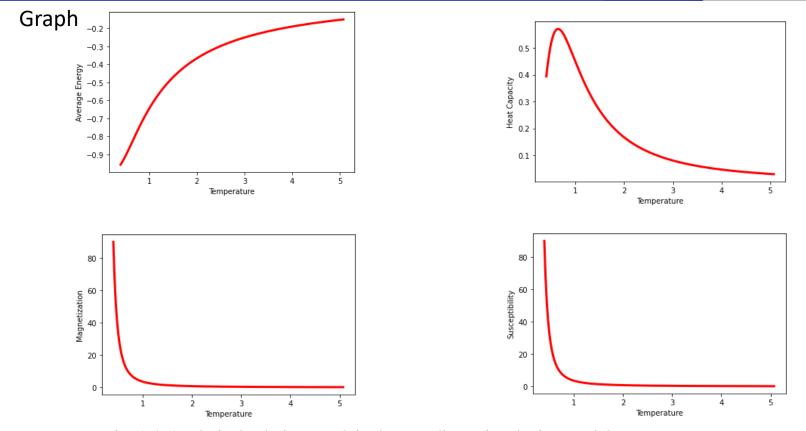


Fig. 1.4. Analytical solution graph in the one-dimensional Ising model

2-1. Monte Carlo Method

Monte Carlo Method is algorithm to get approximation using repeated random sampling.

For example, π approximation is caculated from Monte Carlo Method in picture.

$$\frac{\textit{the number of ball in Circle}}{\textit{the number of ball in Square}} = \frac{\textit{Area of Circle}}{\textit{Area of Square}}$$
$$= \frac{\pi \times 1^2}{2 \times 2}$$

Thus,
$$\pi = 4 \times \frac{\textit{the number of ball in Circle}}{\textit{the number of ball in Square}}$$

Fig. 2.1. pi approximation by Monte Carlo method

2-2. Markov Chain (MC)

Markov Chain Model is probabilistic model that present situation is only depended on last situation. After repeating, it reaches **Stationary Distribution**.

For example, π approximation can be calculated from Markov Chain.

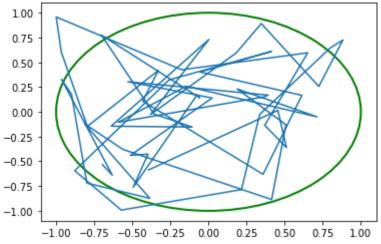


Fig. 2.2. pi approximation by Markov Chain

Start at (-1,-1) and Add random number between -1 and 1 each x and y coordinates.

If new coordinate is inside square (that x, y is from -1 to 1), new coordinate is saved.

If not, last coordinate is maintained.

And count the number of inside the circle.

$$\pi = 4 \times \frac{counted\ number}{total\ trial}$$

2-3. Markov Chain Monte Carlo (MCMC) and Metropolis Algorithm

Markov Chain + Monte Carlo

=

(Repeated random sampling)

+ (State only dependent on last state)

One simple example of Markov Chain Monte Carlo(MCMC) is Metropolis Algorithm.

Metropolis Algorithm

Select initial value:
$$x_0 \sim \pi(x)$$

Repeat:
$$i = 1, 2, ...$$

- Generate test sample $x^* \sim q(x_{i+1} | x_i)$

$$-\alpha(x^*, x_i) = \min\{1, \frac{\pi(x^*)}{\pi(x_i)}\}\$$

- Generate uniform random number: $u \sim U(0, 1)$

$$if \ u < \alpha \qquad x_{i+1} = x^*$$

$$else x_{i+1} = x_i$$

2-4. Application Metropolis Algorithm to The One-Dimensional Ising Model

Select initial value:
$$x_0 \sim \pi(a)$$

$$\pi(a) = \frac{e^{\beta E_a}}{Z}$$
, (Boltzmann distribution)

Flip and get $\pi(b)$

Fig. 2.3. Flip in

$$x^* \sim \pi(b) = \frac{e^{\beta E_b}}{z}$$

$$x^* \sim \pi(b) = \frac{e^{\beta E_b}}{Z}$$
$$-\alpha(x^*, x_i) = \min\{1, \frac{\pi(b)}{\pi(a)}\}$$

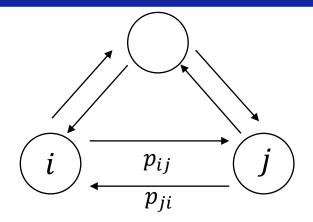
Generate uniform random

$$u \sim U(0,1)$$

$$if \ u < \alpha \qquad x_{i+1} = x^{3}$$

$$else \qquad x_{i+1} = x_{i}$$

2-5. Detailed balance



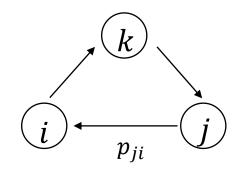


Fig. 2.5. Global balance diagram

Fig. 2.4. Detailed balance diagram
$$\pi_i p_{ij} = \pi_j p_{ji} \ orall i, j$$

The amount of probability flowing from i to j

= The amount of probability flowing from j to i

That represents no net flux of probability and **Stationary Distribution**

case 1)
$$\pi(a) > \pi(b)$$

$$p(a \rightarrow b) = \min \left[1, \frac{\pi(b)}{\pi(a)} \right] = \frac{\pi(b)}{\pi(a)}$$

$$\Rightarrow \pi(a)p(a \rightarrow b) = \pi(b)$$

$$p(b \to a) = \min \left[1, \frac{\pi(a)}{\pi(b)} \right] = 1$$

$$\Rightarrow \pi(b)p(b \to a) = \pi(b)$$

$$\therefore \pi(a)p(a \to b) = \pi(b)P(b \to a)$$

case 2)
$$\pi(a) < \pi(b)$$

$$p(a \rightarrow b) = \min \left[1, \frac{\pi(b)}{\pi(a)} \right] = 1$$

$$\Rightarrow \pi(a)p(a \rightarrow b) = \pi(a)$$

$$p(b \to a) = \min \left[1, \frac{\pi(a)}{\pi(b)} \right] = \frac{\pi(a)}{\pi(b)}$$

$$\Rightarrow \pi(b)p(b \rightarrow a) = \pi(a)$$

$$\therefore \pi(a)p(a \to b) = \pi(b)P(b \to a)$$

$$hamiltonian H = - J (N_{++} + N_{--} - N_{+-} - N_{-+})$$

$$\langle E \rangle = \frac{1}{N} \frac{1}{MC} \langle H \rangle$$

$$< M > = \frac{1}{N} \frac{1}{MC} | (N_{+} - N_{-}) |$$

$$< C_V> = \frac{\partial < E>}{\partial T} = \frac{\partial \beta}{\partial T} \frac{\partial < E>}{\partial \beta} = -\beta^2 \frac{\partial < E>}{\partial \beta} = \beta^2 \frac{\partial^2}{\partial \beta^2} log Z = \beta^2 (< E^2> - < E>^2)$$

$$\langle S \rangle = \frac{\partial \langle M \rangle}{\partial T} = \frac{\partial \beta}{\partial T} \frac{\partial \langle M \rangle}{\partial \beta} = -\beta^2 \frac{\partial \langle M \rangle}{\partial \beta} = \beta^2 (\langle M^2 \rangle - \langle M \rangle^2)$$

2-6. Python code

```
1 import random, math
2 import matplotlib.pyplot as plt
3 import numpy as np
4
5 global n, J, N, kb
6
7 n = 20
8 J = 1
9 kb = 1.3
```

Fig. 2.6. python code – part1(import and define)

```
1 def mk_matrix_1d():
2     data_init = np.zeros([n])
3     data_init = np.random.choice([-1,1], size = (n))
4     return data_init
```

Fig. 2.7. python code – part1(make random matrix)

```
1 def pi value 1d(data unit, beta):
       Energy = 0
      x = random.randint(0,(n-1))
       right = data unit[(x+1)%n]
       left = data unit[(x-1)%n]
       Energy += data unit[x] * (right + left)
10
       pi a = math.exp((+1)*beta*Energy)
11
       pi b = math.exp((-1)*beta*Energy)
       pi divided = pi b / pi a
13
14
       if pi divided <= 1:
15
           if random.random() < pi divided:</pre>
16
               data unit[x] = -data unit[x]
17
       elif pi divided > 1:
18
           data unit[x] = -data_unit[x]
19
20
       return data_unit
21
```

Fig. 2.8. python code – part3(flip)

```
1 def matrix repeat 1d(number, matrix matrix):
       for i in range(number):
           matrix matrix = pi value 1d(matrix matrix, 1/0.1)
       return matrix matrix
 Fig. 2.9. python code – part3(repeat)
 1 def cal temperature(matrix cal, beta):
      Sum Energy = 0
      Sum Energy square = 0
      Avg Energy = 0
      Avg Energy square = 0
      Sum M square = 0
      Sum M = 0
 9
      Avg M=0
10
      Avg M square = 0
11
      Sum Hc = 0
13
      Sum Sus = 0
14
15
      #Repeat N times
16
17
      for j in range(N):
18
          H = 0
19
          H square = 0
20
21
```

```
#Calculate Hamiltonian
23
           for x in range(n):
24
25
               right = matrix cal[(x+1)%n]
26
               left = matrix cal[(x-1)%n]
27
               h temp = -J * (matrix cal[x] * (right + left)) / 2
28
29
30
              H += h temp / (n)
              H square += h temp**2 / (n)
31
32
33
          Sum Hc += beta * beta * (H square - H*H)
34
35
          M = np.sum(matrix cal)
36
37
          Sum Energy += H
          Sum Energy square += H square
38
39
          Sum M += abs(M)
40
          Sum M square += M*M
          matrix cal = pi value 1d(matrix cal, beta)
41
```

Fig. 2.10. python code – part3(caculate)

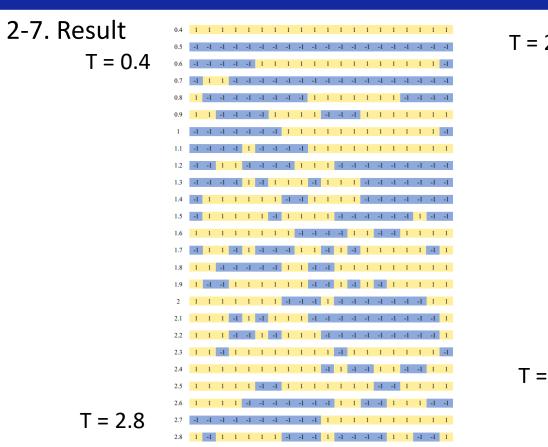


Fig. 2.11.(a) distribution diagram

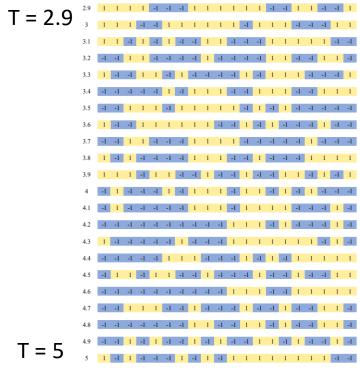


Fig. 2.11.(b) distribution diagram

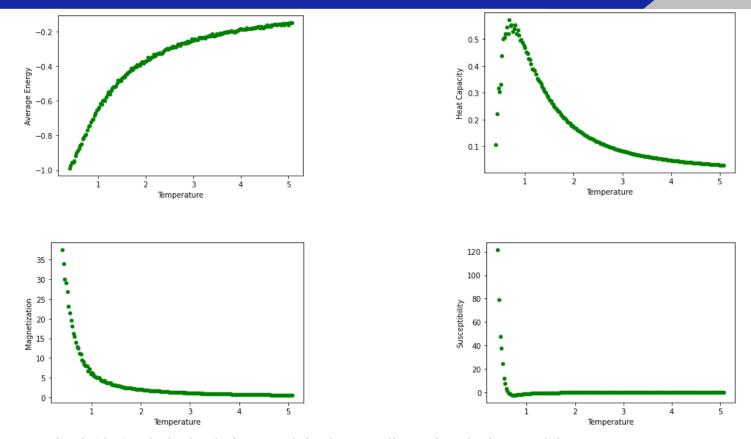


Fig. 2.12. Analytical solution graph in the one-dimensional Ising model

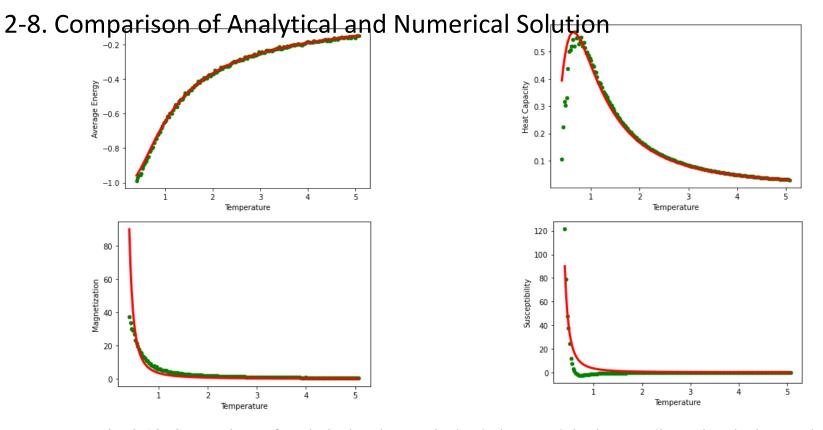


Fig. 2.13. Comparison of analytical and numerical solution graph in the one-dimensional Ising model

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