#### **Advanced Dynamics** & Automatic Control

## Nyquist Plots and Stability Margins

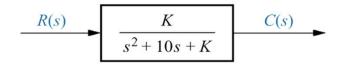
Dr. Mitch Pryor

THE UNIVERSITY OF TEXAS AT AUSTIN

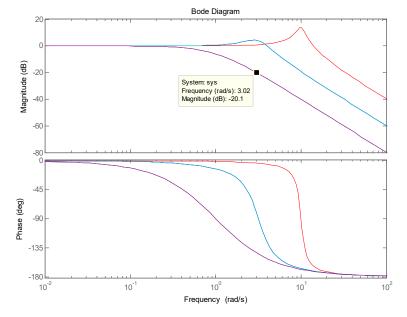
#### Lesson Objective

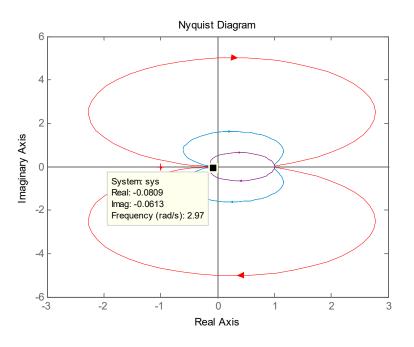
- Define Nyquist Plots
- Review how to create them
  - Learn how complex poles impact the Nyquist Plot
  - Learn how poles at the origin impact the Nyquist Plot
  - Learn how zeros at the origin impact the Nyquist Plot
- Summarize why they are useful
  - Define the Loop Transfer Function
  - Define the Stability Margins (gain and phase)

#### Nyquist plots vs. Bode Plots



- Bode Plots show the frequency vs amplitude and phase on different plots
- Nyquist (Polar) Plots display the frequency and amplitude on same plot
  - frequency (0 to infinity (or –infinity to infinity)) is the parameter that is plotted.
- Nyquist Stability Criterion defines a system's stability in the frequency domain.
  - The Nyquist plot changes with respect to a system variable in a way that controller designers can better see its impact on stability.
- Example plots K = 100 (red), 10 (green), 1 (blue)





#### Creating a Nyquist Plot

A simple example....

$$G(s) = G(j\omega) = \frac{1}{s+1} = \frac{1}{j\omega+1} = \frac{1}{j\omega+1} \left(\frac{1-j\omega}{1-j\omega}\right) = \frac{1-j\omega}{1+\omega^2}$$

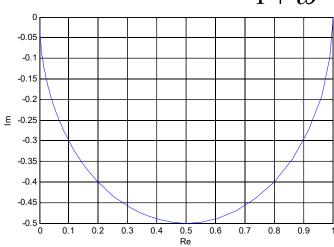
Separate this into the real and imaginary components...

$$\operatorname{Re}(G(j\omega)) = \frac{1}{1+\omega^2}$$

```
w = [0:.1:100];
for i=1:length(w)
    re(i) = 1/(1+w(i)*w(i));
    im(i) = -w(i)/(1+w(i)*w(i));
end

figure(1)
plot( re, im );
xlabel('Re'), ylabel('Im')
grid on;
```

$$\operatorname{Im}(G(j\omega)) = \frac{-\omega}{1+\omega^2}$$



#### Nyquist plots

$$G(s) = G(j\omega) = \frac{1}{s+1} = \frac{1}{j\omega+1} = \frac{1-j\omega}{1+\omega^2}$$

#### **Manually**

#### w = [0:.1:100];

```
for i=1:length(w)
    re(i) = 1/(1+w(i)*w(i));
    im(i) = -w(i)/(1+w(i)*w(i));
end
```

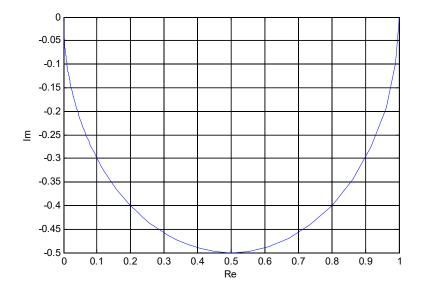
```
plot( re, im );
xlabel('Re'), ylabel('Im')
```

#### Or

```
for i=1:length(w)
    re(i) = real(1/(w(i)+1));
   im(i) = imag(1/(w(i) 1));
end
```

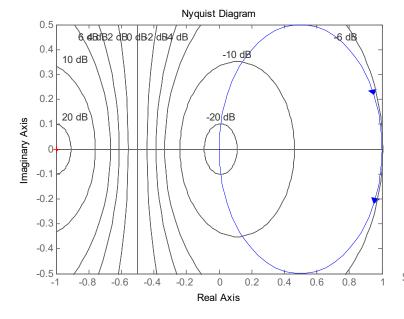
w = [0:.1:100];

```
plot( re, im );
xlabel('Re'), ylabel('Im')
```



#### **Using MATLAB**

```
num = [1]
den = [11]
sys = tf(num, den)
figure(2)
nyquist(sys)
```



### Nyquist for higher order systems

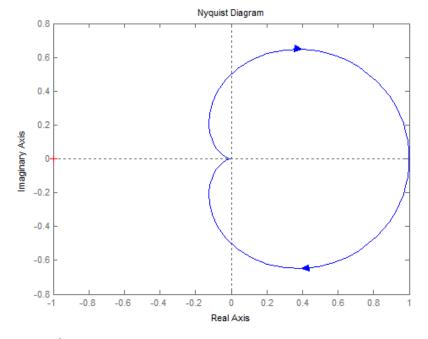
A 2<sup>nd</sup> order example....

$$G(j\omega) = \frac{1}{(s+1)^2} = \frac{1}{(j\omega+1)^2}$$
$$= \frac{1}{(j\omega+1)^2} \left(\frac{(1-j\omega)^2}{(1-j\omega)^2}\right)$$
$$= \frac{1-2j\omega-\omega^2}{\omega^4+2\omega^2+1}$$

Separate this into the real and imaginary components...

$$\operatorname{Re}(G(j\omega)) = \frac{1 - \omega^2}{\omega^4 + 2\omega^2 + 1}$$

$$\operatorname{Im}(G(j\omega)) = \frac{-2j\omega}{\omega^4 + 2\omega^2 + 1}$$



...without MATLAB it can start to get a little complicated...

But there are some simple rules/patterns.

### Nyquist for Higher Order Systems

Consider a general transfer function...

$$G(s) = \frac{(s - z_1) \cdots (s - z_m)}{(s - p_1) \cdots (s - p_n)}$$

with *m* zeroes and *n* poles. Assume (for the moment) that there are no poles at the origin..

$$G(j\omega) = \frac{(j\omega - z_1)\cdots(j\omega - z_m)}{(j\omega - p_1)\cdots(j\omega - p_n)}$$

As the frequency goes to zero, the value of G will be a finite real number (aka the *zero frequency gain*).

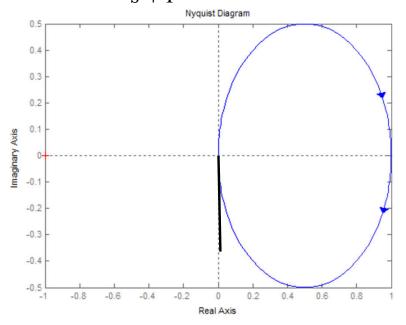
$$G(j\omega) = \frac{(-z_1)(-z_2)\cdots(-z_m)}{(-p_1)(-p_2)\cdots(-p_n)}$$

As the frequency becomes very large, the frequency terms will dominate.

$$G(j\omega) \approx \frac{(j\omega)^m}{(j\omega)^n} = \frac{1}{(j\omega)^{n-m}}$$

Systems of this form (where n>m) will approach 0 as frequency increases. The angle of approach is determined by n-m and is given by  $-90(n-m)^{\circ}$  For example:

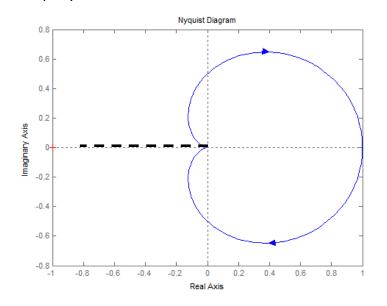
 $G(s) = \frac{1}{s+1} \Rightarrow n-m=1$ 



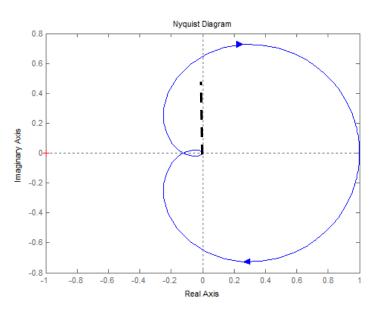
### Nyquist for Higher Order Systems

Some more examples showing the angle of approach is given by  $-90(n-m)^{\circ}$ .

$$G(s) = \frac{1}{(s+1)^2} \Rightarrow n-m=2$$

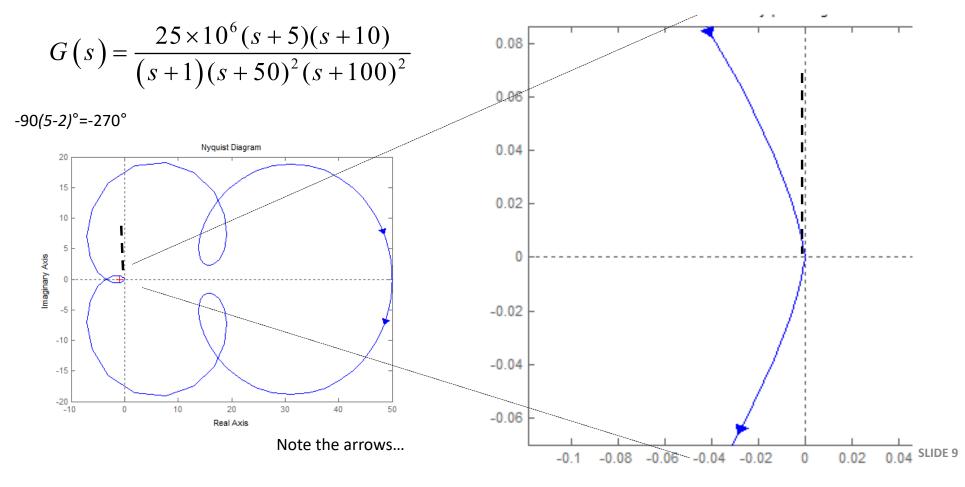


$$G(s) = \frac{1}{(s+1)^3} \Rightarrow n-m=3$$



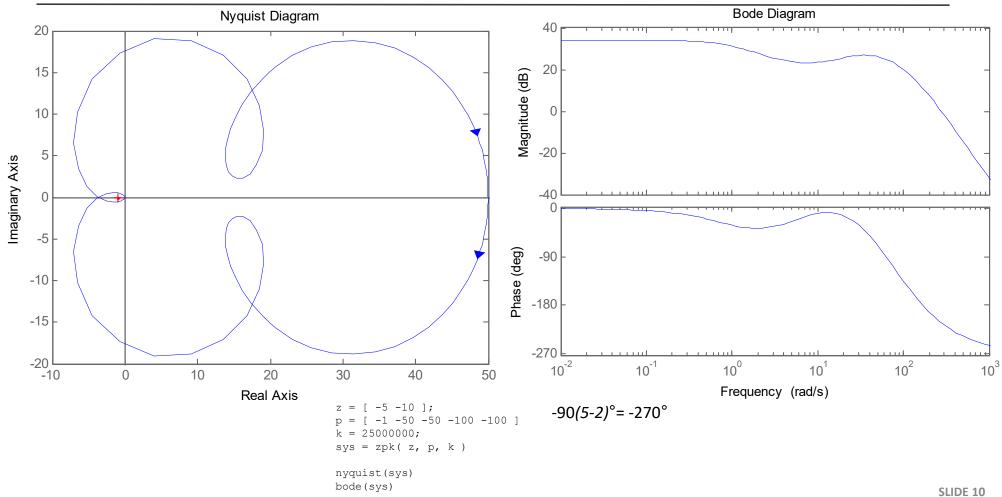
#### Nyquist for Higher Order Systems

One last example showing the angle of approach is given by  $-90(n-m)^{\circ}$ .



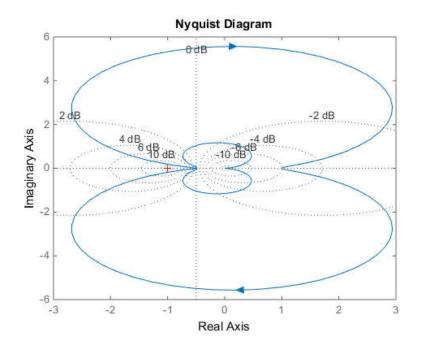
# Nyquist vs. Bode

$$G(s) = \frac{25 \times 10^6 (s+5)(s+10)}{(s+1)(s+50)^2 (s+100)^2}$$



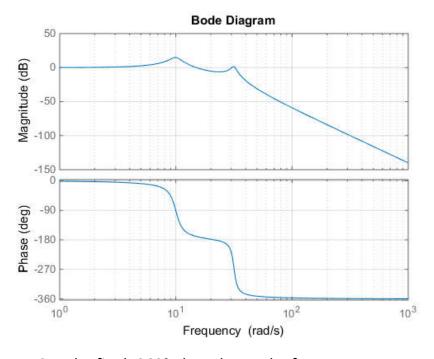
#### Special cases: Complex roots

$$G(s) = \frac{10^5}{(s^2 + 2s + 100)(s^2 + 3s + 1000)}$$



Still observe the increases in magnitude at times when the frequency increases.

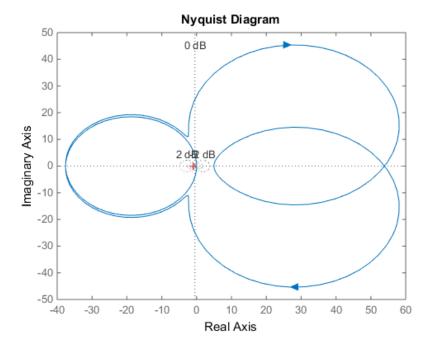
Looking at the Bode Plot for the system...



See the final -360° phase lag as the frequency goes to infinity as expected.

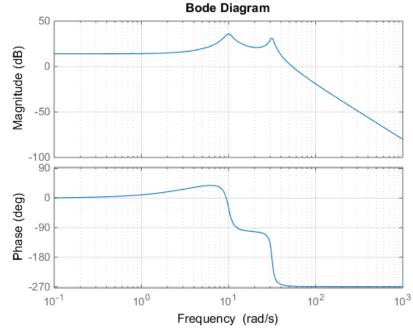
#### Special cases: Complex roots

$$G(s) = \frac{10^{5}(s+5)}{(s^{2}+2s+100)(s^{2}+3s+1000)}$$



Now considering more complex systems, it is less clear why we care about Nyquist plots. We are getting to that.

The impact of a zero we should recognize on the Bode plot, but the response becomes harder to follow on the Nyquist Plot



You see the bulge, the rapid phase change and the -270° phase lag as the frequency goes to infinity.

## Special cases: Zero(s) at Origin

#### Zero at the origin

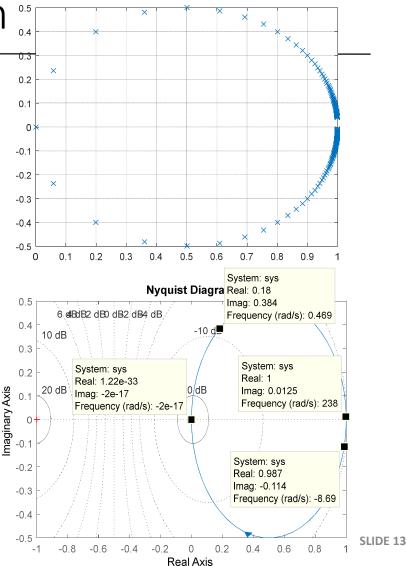
Similar procedure, but note the frequency of the real and imaginary components when the frequency is 0.

$$G(j\omega) = \frac{j\omega}{(j\omega+1)} \frac{(j\omega-1)}{(j\omega-1)}$$
$$= \frac{-\omega^2 - j\omega}{-\omega^2 - 1}$$
$$= \frac{\omega^2}{\omega^2 + 1} + j\frac{\omega}{\omega^2 + 1}$$

```
clear all;
w(1) = 0-100i;
for c=2:500
   w(c) = w(c-1) + (0+0.25*i);
end

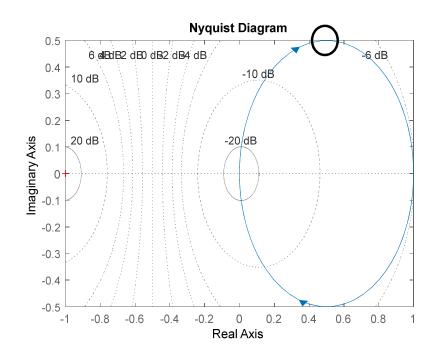
for i=1:length(w)
   re(i) = real( w(i)/(w(i)+1 ) );
   im(i) = imag( w(i)/(w(i)+1 ) );
end
figure(1); plot( re, im, 'x'); grid on;

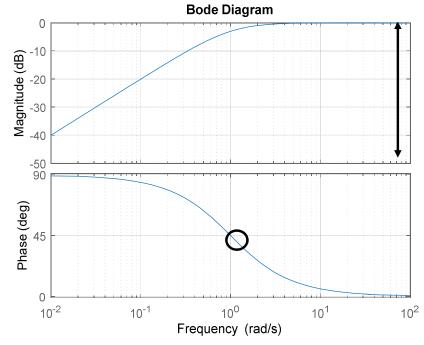
sys = zpk( [0], [-1], 1 );
figure(2); nyquist( sys ); grid on;
```



# Special Case – Zero(s) at the origin

$$G(s) = \frac{s}{(s+1)} \Rightarrow n-m=1$$





### Special cases – poles at origin

$$G(s) = \frac{1}{s(s+1)} \Rightarrow n-m = 2$$

Now our gain approaches infinity as the frequency approaches 0 instead of a finite value

$$G(j\omega) = \frac{1}{j\omega(j\omega+1)}$$

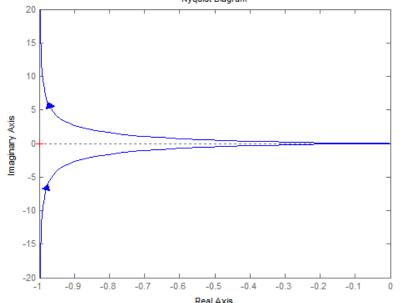
$$= \frac{1}{(j\omega-\omega^2)} \left(\frac{(j\omega+\omega^2)}{(j\omega+\omega^2)}\right)$$

$$= \frac{j\omega+\omega^2}{-\omega^4-\omega^2}$$

This one we can separate into real and imaginary components...

$$\operatorname{Re}(G(j\omega)) = \frac{\omega^2}{-\omega^2(\omega^2 + 1)} = \frac{1}{-(\omega^2 + 1)}$$

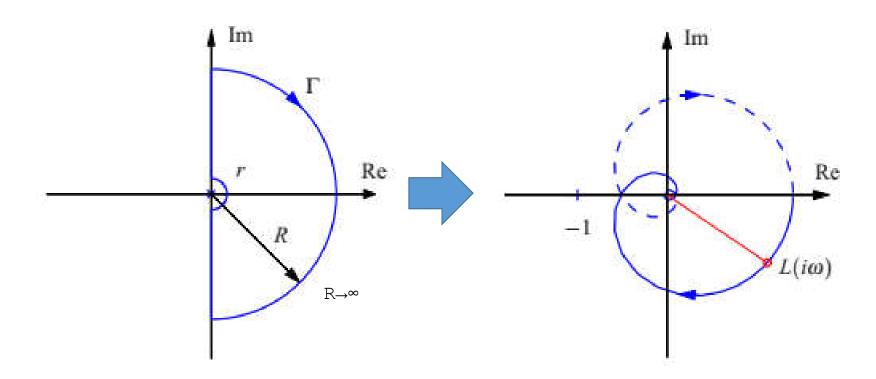
$$\operatorname{Im}(G(j\omega)) = \frac{\omega}{-\omega^2(\omega^2 + 1)} = \frac{1}{-\omega(\omega^2 + 1)}$$



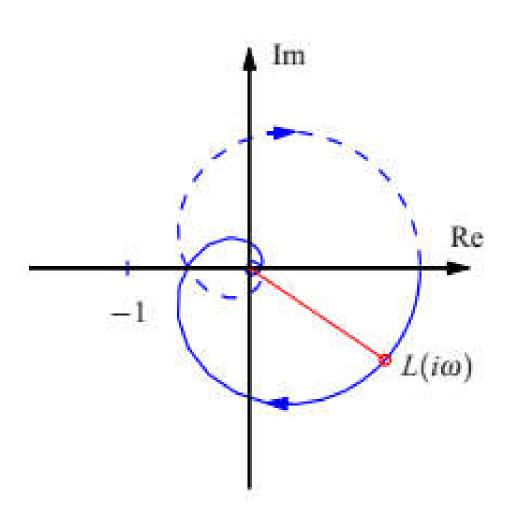
The high frequency asymptote is still -180  $^{\circ}$ 

# The Nyquist D Contour

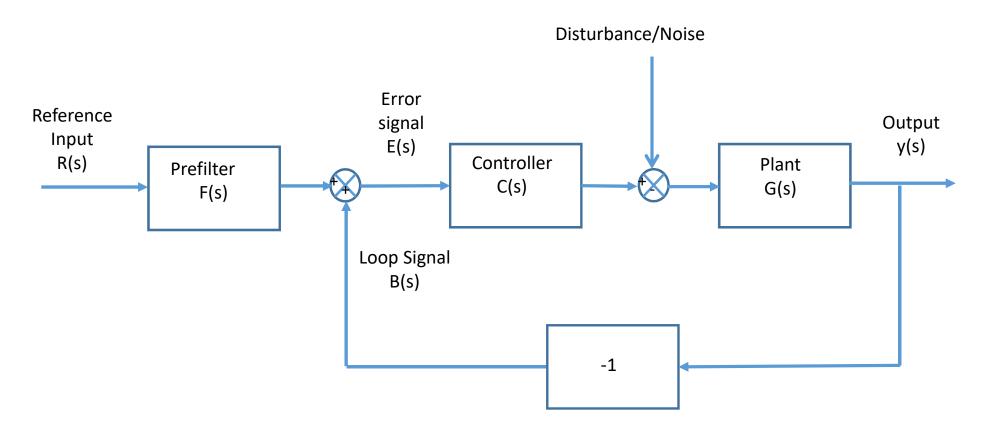
It has been helpful in when drawing the Nyquist to not just think of the frequencies from 0 to  $\infty$ , but from  $-\infty$  to  $\infty$ .



# So why are these useful?



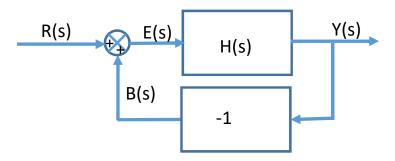
## Recall our closed system nomenclature...



<sup>\*</sup>not standardized. Nomenclature varies from textbook to textbook

### Recall our closed loop transfer function

Make a couple assumptions (no prefilter, no disturbance, etc.)



#### **Typical Control Blocks and Signals**

R(s) – Reference signal or desired output

E(s) - Error signal

H(s) - Often C(s)\*P(s)

C(s) – The control law that produces a u for the model or plant

P(s) – The model of the system to be controlled or plant

Y(s) – The output of the system

B(s) – Loop transfer signal (last signal value prior to feedback)

We get...

$$T_{cl}(s) = \frac{Y(s)}{R(s)} = \frac{C(s)P(s)}{1 + C(s)P(s)}$$

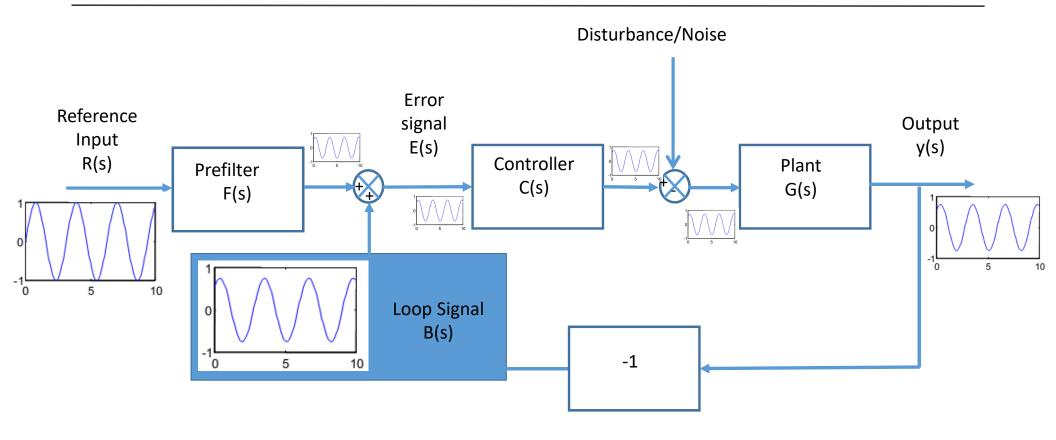
The system will be stable if all the roots of the CE are negative.

$$CE = \Delta(s) = 1 + C(s)P(s) = d_C(s)d_P(s) + n_C(s)n_P(s)$$

Given the calculation there is little insight into how to (re)design a system for stability.

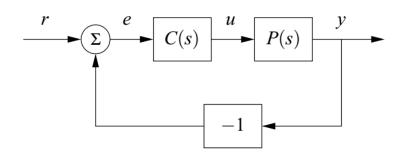
This is where the *Loop Transfer Function* **B(s)** can be helpful.

# Loop Transfer Function



<sup>\*</sup>not standardized. Nomenclature varies from textbook to textbook

### Previous method to determine stability...



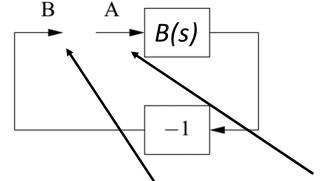
$$G_{yr}(s) = \frac{PC}{1 + PC} = \frac{n_p(s)n_c(s)}{d_p(s)d_c(s) + n_p(s)n_c(s)}$$

Find the Characteristic Equation...

$$\lambda(s) = d_p(s)d_c(s) + n_p(s)n_c(s)$$

Ensure all the roots have a negative real part.

If it is not stable, there is little information here that helps us pick a better C(s)



If we inject a signal with frequency  $\omega_o$  here...

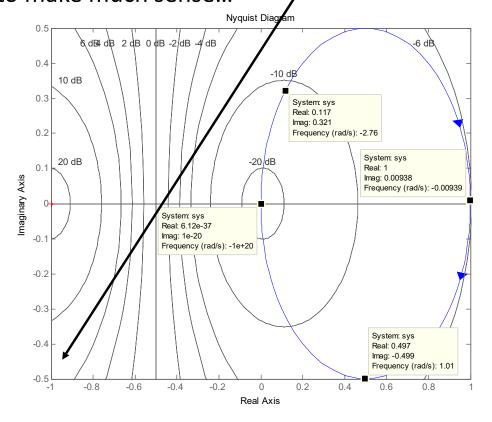
For a linear system we should get an output frequency  $\omega_o$  here...

It is reasonable that an oscillation can be maintained if the amplitude at B is the same as the amplitude at A. Thus...

$$B(i\omega) = -1$$

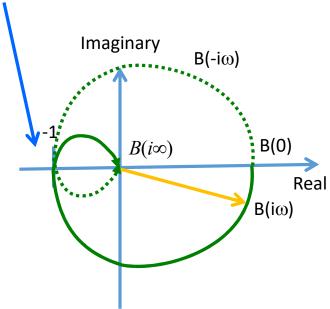
#### Back to our Nyquist Plot...

Notice that MATLAB ensures -1 is shown on the Nyquist Plot... even if that didn't seem to make much sense...



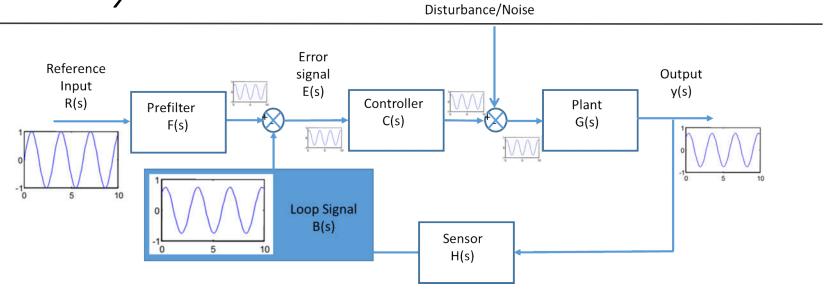
Plane of the Open Loop Transfer Function

$$B(i\omega_0) = \frac{n_c(s)}{d_c(s)} \frac{n_p(s)}{d_p(s)} = -1$$



-1 is called the critical point

### Loop Analysis



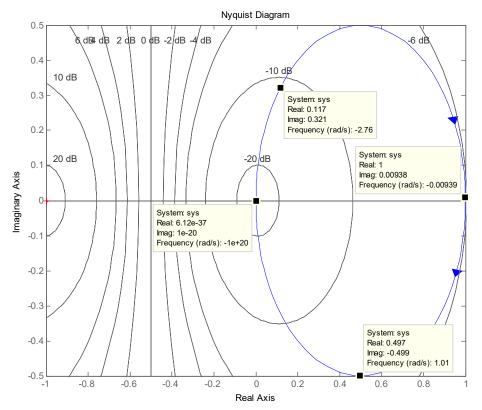
**Loop analysis**: trace how a sinusoidal signal propagates in the feedback loop and investigate if the propagating signal grows or decays.

This is the key concept supporting the **Nyquist Stability Theorem**.

Its key advantage over eigenanalysis and Lyapunov stability is it also provides insight into how stable we are since we can use it to define the **gain margin** and **phase margin** (aka the **stability margins**.)

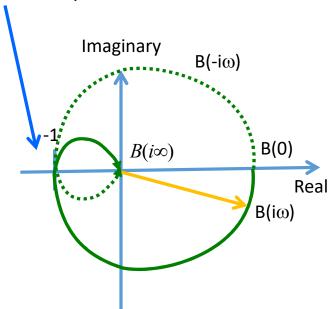
#### Back to our Nyquist Plot...

Notice that MATLAB ensures -1 is shown on the Nyquist Plot... even if that didn't seem to make much sense...



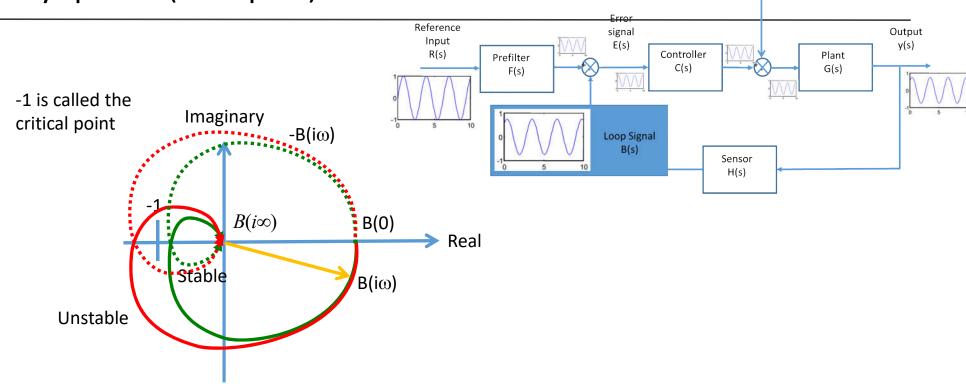
Plane of the Open Loop
Transfer Function

$$B(i\omega_0) = \frac{n_c(s)}{d_c(s)} \frac{n_p(s)}{d_p(s)} = -1$$



-1 is called the critical point

#### Nyquist (simple)Theorem



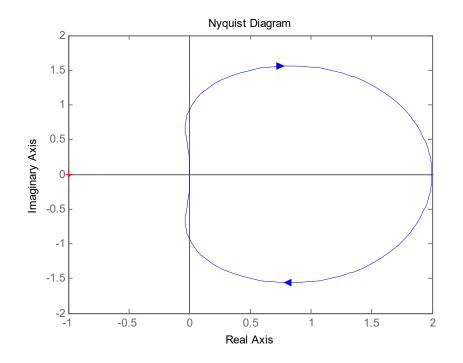
Nyquist (simple) Theorem: If the Loop Transfer Function  $B(i\omega)$  has no poles in the right hand side (except for simple poles on the imaginary axis) then the system is stable iff there are no encirclements of the critical point, -1.

Disturbance/Noise

#### Example, stable system

B(s) = 
$$\frac{K(s^2 + 2s + 2)}{s^3 + 2s^2 + 2s + 1}$$
  
B(j\omega) =  $\frac{1(j\omega^2 + j\omega s + 2)}{j\omega^3 + 2j\omega^2 + 2j\omega + 1}$ 

$$B(0) = 2$$
$$B(\infty) = B(-\infty) = 0$$

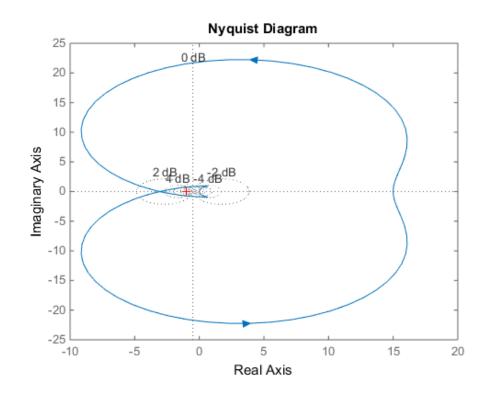


Will increasing K make this system unstable?

#### Example, unstable system

B(s) = 
$$\frac{K(s^2 + s + 3)}{s^3 + 6s^2 - 2s + 2}$$
  
B(j\omega) =  $\frac{10(j\omega^2 + j\omega + 3)}{j\omega^3 + 6j\omega^2 - 2j\omega + 2}$ 

$$B(0) = 15$$
$$B(\infty) = L(-\infty) = 0$$

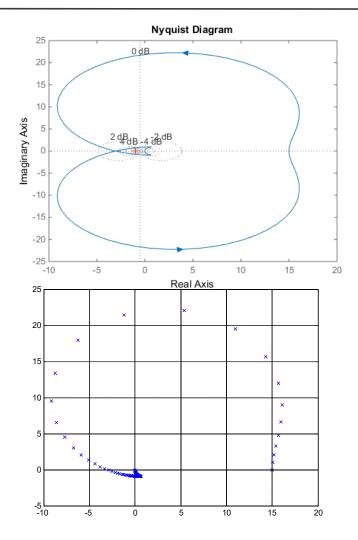


Will decreasing K make this system stable?

## Example, Unstable System (Old School)

B(s) = 
$$\frac{K(s^2 + s + 3)}{s^3 + 6s^2 - 2s + 2}$$
  
B(j\omega) =  $\frac{10(j\omega^2 + j\omega + 3)}{j\omega^3 + 6j\omega^2 - 2j\omega + 2}$ 

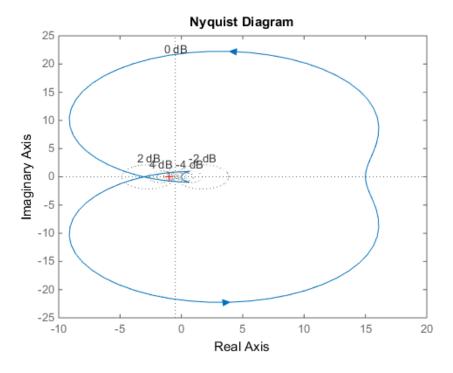
```
clear all;
f=[0:.05:500];
K = 10;
for j=1:length(f)
    w = 0 + 1i*f(j);
    Bnum = K*(w^2 + w + 3);
    Bden = w^3 + 6*w^2 - 2*w + 2;
    B = Bnum/Bden;
    re(j) = real(B);
    im(j) = imag(B);
end
plot( re, im, 'bx' );
```

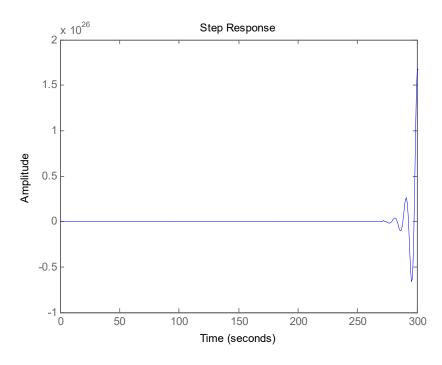


### Example, Unstable System

B(s) = 
$$\frac{K(s^2 + s + 3)}{s^3 + 6s^2 - 2s + 2}$$
  
B(j\omega) =  $\frac{10(j\omega^2 + j\omega + 3)}{j\omega^3 + 6j\omega^2 - 2j\omega + 2}$ 

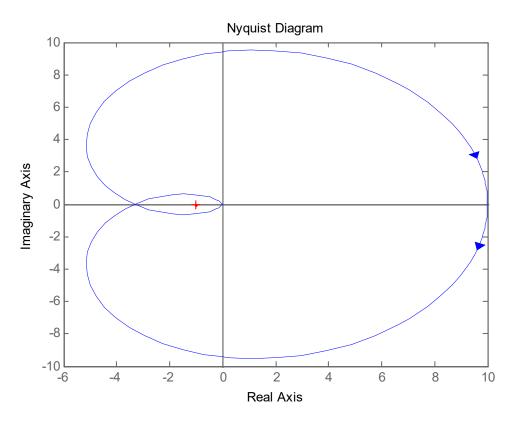




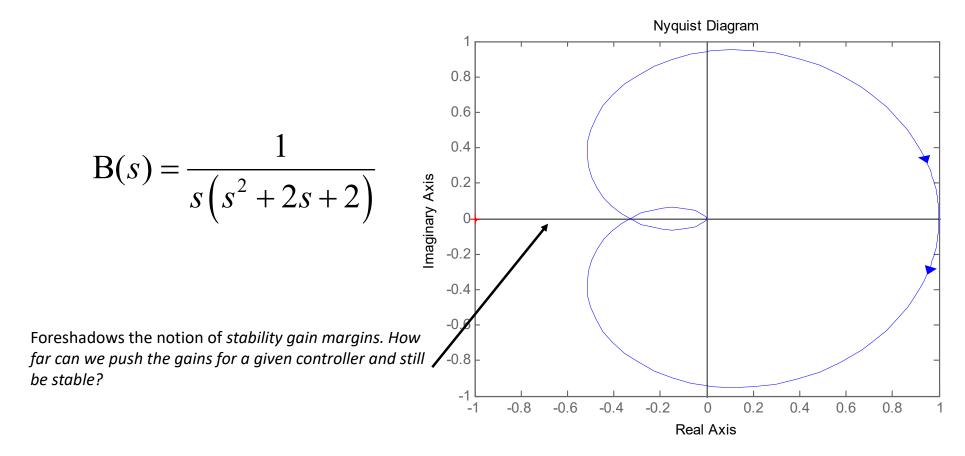


## Example, unstable system

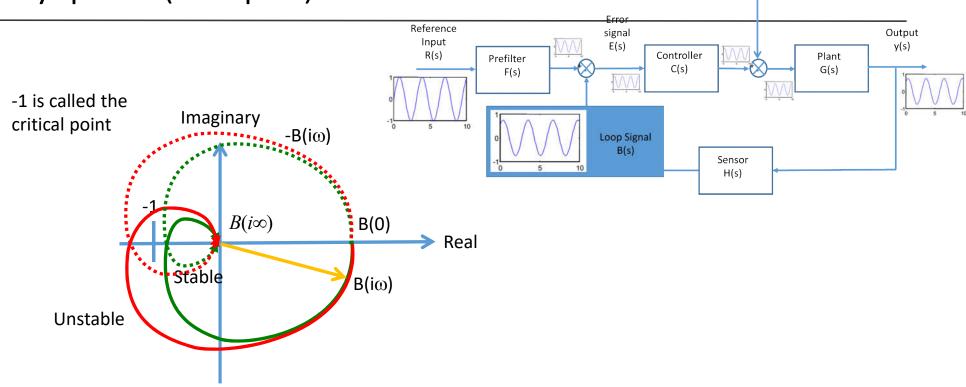
$$B(s) = \frac{10}{s\left(s^2 + 2s + 2\right)}$$



# But made stable by reducing the gain



#### Nyquist (simple)Theorem



Nyquist (simple) Theorem: If the Loop Transfer Function  $B(i\omega)$  has no poles in the right hand side (except for simple poles on the imaginary axis) then the system is stable iff there are no encirclements of the critical point, -1.

Disturbance/Noise

## Nyquist's (full) Stability Theorem

Nyquist (simple) stability Theorem required that B(s) have NO poles in the right half plane. In some cases this will not be true *even* when the closed loop transfer function T(s) is stable.

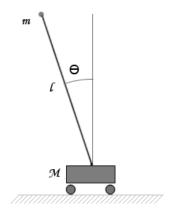
Thus, we need a more general (but slightly more complicated) theorem.

Nyquist Stability Theorem: Given a closed loop system with the loop transfer function B(s) with P poles in the right hand plane.

• Let N be the number of clockwise encirclements of -1 by  $B(i\omega)$  minus the counterclockwise encirclements of -1 by  $B(i\omega)$ . Then the closed loop system has Z=N+P zeros in the right half plane.

This requires that if a B(s) with P poles in the right hand plane, there must be P counter-clockwise encirclements of -1.

#### An example, Inverted Pendulum

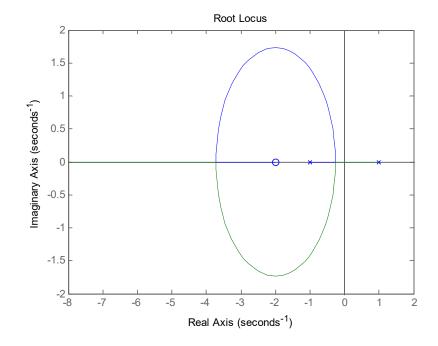


The old fashion way...

$$B(s) = \frac{k(s+2)}{s^2 - 1}$$

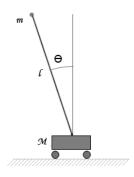
Note there is one pole in the RHP.

The Nyquist stability criterion dictates we should expect 1 counterclockwise encirclements of -1 if the system is stable...



...for what values of *k* is the system stable?

#### An example, Inverted Pendulum

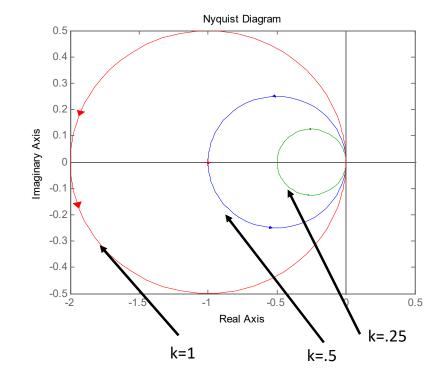


The old fashion way...

$$B(s) = \frac{k(s+2)}{s^2 - 1}$$

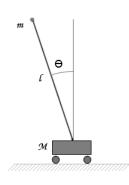
Note there is one pole in the RHP, we should expect 1 counterclockwise encirclement of - 1 if the system is stable....

...for what values of *k* is the system stable?



Does this result make physical sense for this system?

#### An example, Inverted Pendulum

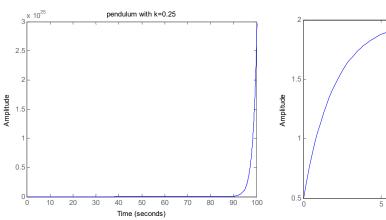


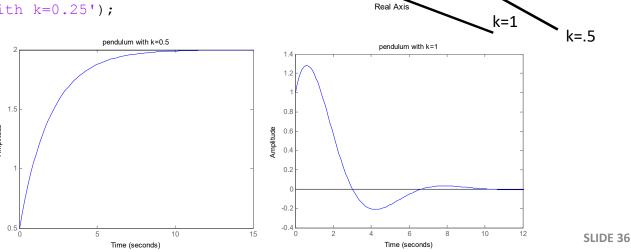
$$B(s) = \frac{k(s+2)}{s^2 - 1}$$

$$T_{cl}(s) = \frac{ks + 2k}{s^2 + ks + (2k - 1)}$$

The old fashion way...

```
sys = tf([k 2*k], [1 k 2*k-1]);
impulse(sys);
title('pendulum with k=0.25');
```





0.4

Imaginary Axis

-0.3

-0.4

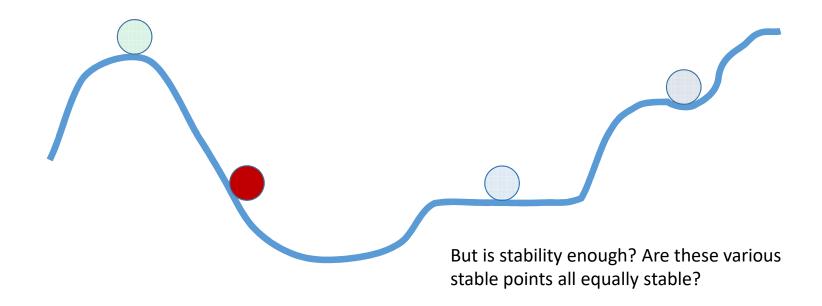
Nyquist Diagram

k=.25

### Stability Margins

Recall the definition of stability: A system's response to an input results in an output that stays arbitrarily near some value, a, for all of time greater than some value,  $t_{\rm f}$ 

$$||b-a|| < \delta \Rightarrow ||x(t;b)-x(t;a)|| < \varepsilon \text{ for all } t > 0$$



### Stability Margins Summary

 Margins are the range from the current system design to the edge of instability. We will determine...

#### Gain Margin

- How much can gain be increased?
- Formally: the smallest multiple amount the gain can be increased before the closed loop response is unstable.

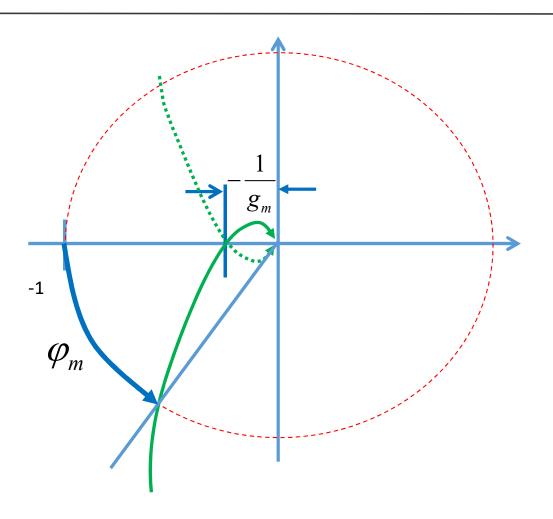
#### Phase Margin

- How much further can the phase be shifted?
- Formally: the smallest amount the phase can be increased before the closed loop response is unstable.

#### Stability Margin

How far is the system from the critical point?

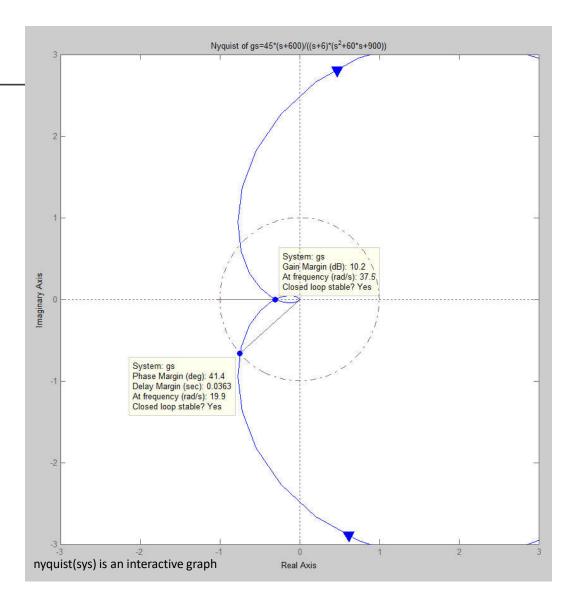
# Gain & Phase Margins



#### Using MATLAB

$$B(s) = \frac{45(s+600)}{(s+6)(s^2+60s+900)}$$

From the plot, the gain margin is 10.2 dB The gain multiple is  $G = 10^{0.05g_m} = 10^{0.05*10.2} = 3.2359$ 

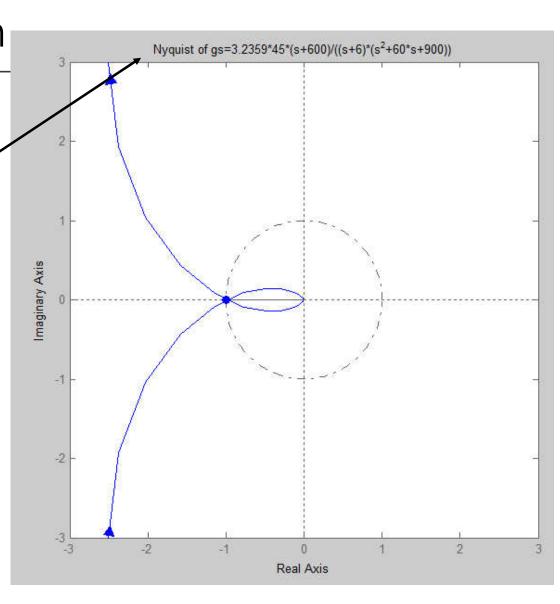


Using the gain margin

B(s) = 
$$\frac{(3.2359*45)(s+600)}{(s+6)(s^2+60s+900)}$$

Here the gain from the previous plot has been multiplied by 3.2359

With this gain, the system is neutrally stable.



## Using the gain margin

B(s) = 
$$\frac{(270)(s+600)}{(s+6)(s^2+60s+900)}$$

Here the gain has been changed to 270.

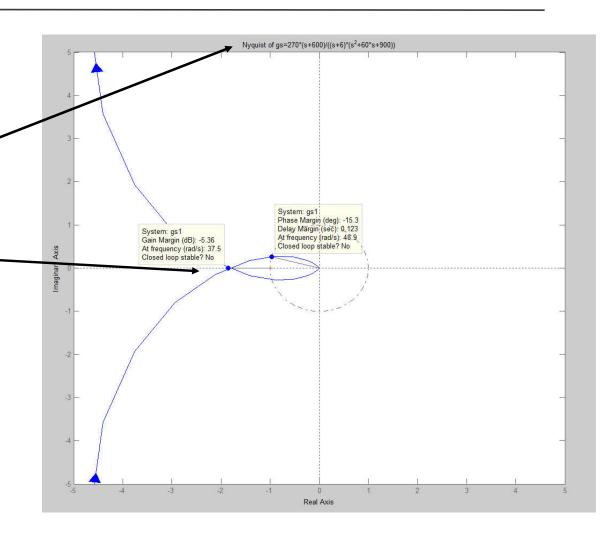
With this gain, the system is unstable.

With this gain margin is now -5.36 dB

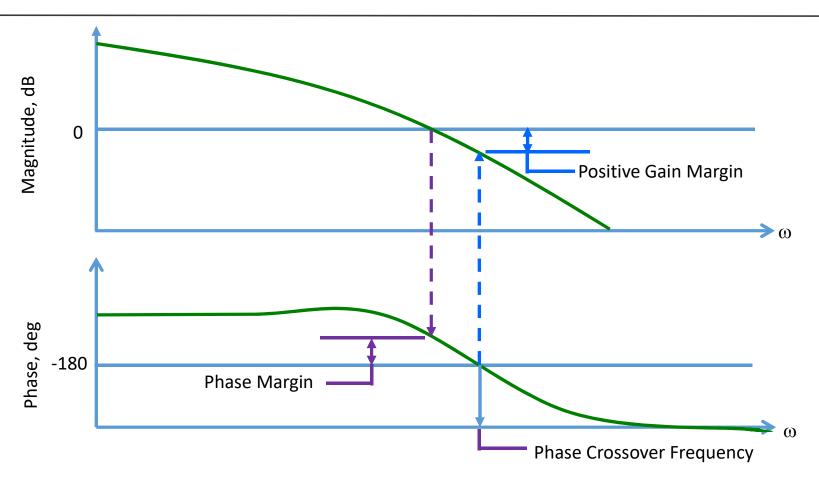
$$-5.36 = 20\log_{10} x$$
$$10^{\frac{-5.36}{20}} = x$$
$$0.54 = x$$

And this is the correct gain multiplier, since

$$0.5393 = \frac{45 * 3.2359}{270} = \frac{K_{neutral stable}}{K}$$

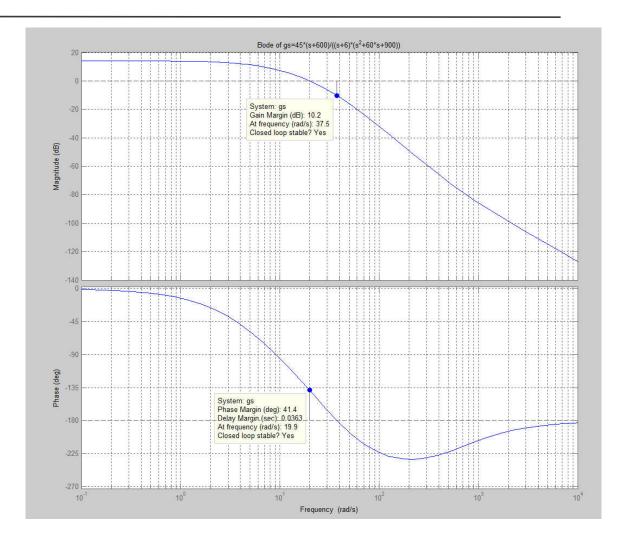


# Margins on a Bode Plot (for B(s)!)



# Back to our example...

$$B(s) = \frac{45(s+600)}{(s+6)(s^2+60s+900)}$$



#### Summary

- Examining the Loop Transfer Function B(s) is a common sense way to look at the response of a feedback system.
  - Can simplify stability analysis. Don't have to calculate T<sub>cl</sub>(S)
- By graphing real and imaginary components of the response as a function of frequency (Nyquist plot) we can determine the stability of the system.
- Nyquist plots can get complicated, but a combination of MATLAB, hand computations and symmetry make most problems solvable
- Using Nyquist plots we can roughly determine how stable we are by calculating the gain and phase margins.