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$$\frac{D}{Dt} \equiv \left( \frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right)$$

## Thermal Instability

$$\uparrow x_3 \quad \beta = -\frac{d\theta}{dz}$$

$$\partial_t \rho + \frac{\partial(\rho u_j)}{\partial x_j} = 0 \quad \text{continuity}$$

$$\rho \frac{D}{Dt} u_i = \rho \left( \frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) u_i = -g\rho \delta_{i3} + \frac{\partial \sigma_{ij}}{\partial x_j} \quad \text{Navier-Stokes EoM}$$

$$\text{stress tensor: } \sigma_{ij} = -p\delta_{ij} + \mu \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} - \frac{2}{3} \frac{\partial u_k}{\partial x_k} \delta_{ij} \right) + \lambda \frac{\partial u_k}{\partial x_k} \delta_{ij}$$

dynamic viscosity      bulk viscosity

write this before as

$$\rightarrow \partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} = -\frac{\nabla p}{\rho} + \nu \nabla^2 \mathbf{u}$$

Heat conduction:

$$\rho \frac{DE}{Dt} = \rho \left( \frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) E = \frac{\partial}{\partial x_j} \left( \overset{\text{temp}}{k} \frac{\partial \theta}{\partial x_j} \right) - \rho \frac{\partial u_j}{\partial x_j} + \Phi$$

internal E per unit mass      thermal conductivity

$$\text{rate of viscous dissipation per unit volume of fluid: } \Phi = \frac{1}{2} \mu \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)^2 + \left( \lambda - \frac{2}{3} \mu \right) \left( \frac{\partial u_k}{\partial x_k} \right)^2$$

$$\text{perfect liquid: } E = c\theta \quad (c \text{ is specific heat})$$

Boussinesq: density variation neglected everywhere except in buoyancy.

$$\therefore \rho = \rho_0 \{ 1 - \alpha(\theta - \theta_0) \}$$

$\hookrightarrow \mu, \kappa, c$  are constants

$\hookrightarrow$  neglect  $\lambda$

$$\hookrightarrow \frac{\partial u_j}{\partial x_j} = 0, \text{ as for an incompressible fluid } (\nabla \cdot \mathbf{u} = 0)$$

$$\therefore \sigma_{ij} = -p\delta_{ij} + \mu \left( \frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

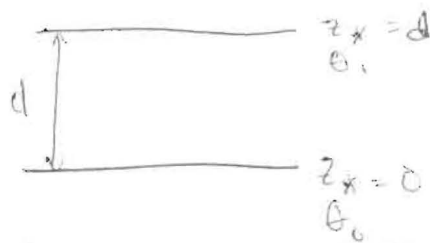
$$\frac{D}{Dt} u_i = -\frac{\partial}{\partial x_i} \left( \frac{p}{\rho_0} + g z \right) - \alpha g (\theta_0 - \theta) \delta_{i3} + \nu \frac{\partial^2}{\partial x_j^2} u_i$$

Heat eqn  $\rightarrow$ 

$$\frac{D}{Dt} \theta = \kappa \Delta \theta$$

$$\kappa = \frac{k}{\rho_0 c}$$

ratio of thermal conductivity  
to mean density \* specific heat



temp gradient  

$$\beta = \frac{(\theta_0 - \theta_1)}{d}$$

Boussinesq eqns:

instability  
when  $\theta_0 > \theta_1$

$$\begin{cases} \frac{\partial u_i}{\partial x_j} = 0 \\ \left( \frac{\partial}{\partial t} + \vec{u} \cdot \vec{\nabla} \right) u_i = - \frac{\partial}{\partial x_i} \left( \frac{P}{\rho_0} + g z \right) - \alpha g (\theta_0 - \theta) \delta_{i3} + \nu \frac{\partial^2 u_i}{\partial x_j^2} \\ \left( \frac{\partial}{\partial t} + \vec{u} \cdot \vec{\nabla} \right) \theta = \kappa \Delta \theta \end{cases}$$

Laplacean

$$u_* = 0$$

$$\theta_* = \theta_0 - \beta z_*$$

$$P_* = P_0 - g \rho_0 \left( z_* + \frac{1}{2} \alpha \beta z_*^2 \right) \quad \text{for } 0 \leq z_* \leq d$$

Small (primed) perturbations:

$$u_* = u'_*(\vec{x}_*, t_*)$$

$$\theta_* = \theta_*(z_*) + \theta'_*(\vec{x}_*, t_*)$$

$$p_* = P_*(z_*) + p'_*(\vec{x}_*, t_*)$$

$$\textcircled{1} \quad \nabla \cdot \vec{u} = 0 \implies \nabla_* \cdot \vec{u}'_* = 0$$

$$\frac{\partial}{\partial t} u'_* + \vec{u} \cdot \nabla u'_* = - \nabla \left( \frac{P}{\rho_0} \right) - \nabla (g z) - \alpha g (\theta_0 - \theta) \delta_{i3} + \nu \nabla^2 u'_*$$

$$\frac{\partial u'_*}{\partial t} = - \frac{1}{\rho_0} \nabla_* (P_*(z_*) + p'_*(\vec{x}_*, t_*)) - g \nabla_* z_* - \alpha g (\theta_0 - (\theta_*(z_*) + \theta'_*(\vec{x}_*, t_*))) \delta_{i3} + \nu \nabla_*^2 u'_*$$

$$\frac{\partial u'_*}{\partial t} = - \frac{1}{\rho_0} \nabla_* p'_* - \frac{1}{\rho_0} \nabla_* P_* - g \nabla_* z_* - \alpha g (\theta_0 - \theta_* - \theta'_*) \delta_{i3} + \nu \nabla_*^2 u'_*$$

$$\begin{aligned} \frac{\partial u'_*}{\partial t} &= - \frac{1}{\rho_0} \nabla_* p'_* - \frac{1}{\rho_0} \nabla_* (P_0 - g \rho_0 (z_* + \frac{1}{2} \alpha \beta z_*^2)) - \alpha g (\theta_0 - \theta_0 + \beta z_* - \theta'_*) \delta_{i3} \\ &= - \frac{1}{\rho_0} \nabla_* p'_* + g \nabla_* z_* + \frac{1}{2} g \alpha \beta \nabla_* z_*^2 - \alpha g \beta z_* + \alpha g \theta'_* \delta_{i3} + \nu \nabla_*^2 u'_* \end{aligned}$$

$$\textcircled{2} \quad \frac{\partial u'_*}{\partial t} = - \frac{1}{\rho_0} \nabla_* p'_* + \alpha g \theta'_* \hat{k} + \nu \nabla_*^2 u'_*$$

$$\frac{\partial}{\partial t} (\theta_0 - \beta z_*)' + u_* \cdot \nabla (\theta_0 - \beta z_*)' = \kappa \nabla_*^2 (\theta_0 - \beta z_*)' + \theta'_*$$

$$-\beta u_{z*} + \frac{\partial \theta'_*}{\partial t} + u_* \cdot \nabla_* \theta'_* = \kappa \nabla_*^2 \theta'_*$$

$$\textcircled{3} \quad \frac{\partial \theta'_*}{\partial t} - \beta u_{z*}' = \kappa \nabla_*^2 \theta'_*$$

So now we have the linearized, dimensional eqns.

$$\begin{cases} \nabla_{\mathbf{x}} \cdot \mathbf{u}_{\mathbf{x}}' = 0 \\ \frac{\partial \mathbf{u}_{\mathbf{x}}'}{\partial t_{\mathbf{x}}} = -\frac{1}{\rho_0} \nabla_{\mathbf{x}} P_{\mathbf{x}}' + \alpha g \Theta_{\mathbf{x}}' \hat{\mathbf{k}} + \nu' \nabla_{\mathbf{x}}^2 \mathbf{u}_{\mathbf{x}}' \\ \frac{\partial \Theta_{\mathbf{x}}'}{\partial t_{\mathbf{x}}} - \beta w_{\mathbf{x}}' = \kappa \nabla_{\mathbf{x}}^2 \Theta_{\mathbf{x}}' \end{cases}$$

scales: length:  $d$  time:  $\frac{d^2}{\kappa}$  temp diff:  $\Theta_0 - \Theta_1 = \beta d$

$$\therefore \vec{x} = \frac{\vec{x}_{\mathbf{x}}}{d} \quad t = \frac{\kappa t_{\mathbf{x}}}{d^2} \quad \vec{u} = \frac{d \vec{u}_{\mathbf{x}}'}{\kappa} \quad \Theta = \frac{\Theta_{\mathbf{x}}'}{\beta d} \quad P = \frac{d^2 P_{\mathbf{x}}'}{\rho_0 \kappa^2}$$

① so:  $\nabla \cdot \left( \frac{d \mathbf{u}_{\mathbf{x}}'}{\kappa} \right) = 0 \rightarrow \nabla \cdot \mathbf{u} = 0$   $\nabla_{\mathbf{x}} \rightarrow \frac{1}{d} \nabla$  ?

$\vec{x}_{\mathbf{x}} = d \vec{x} \quad \frac{d^2 t}{\kappa} = t_{\mathbf{x}} \quad \frac{\kappa \vec{u}}{d} = \vec{u}_{\mathbf{x}}' \quad \beta d \Theta = \Theta_{\mathbf{x}}' \quad \frac{\rho_0 \kappa^2 P}{d^2} = P_{\mathbf{x}}'$

$$\frac{\partial \left( \frac{\kappa \mathbf{u}}{d} \right)}{\partial \left( \frac{d^2 t}{\kappa} \right)} = \left( \frac{\kappa}{d} \right) \frac{\partial \mathbf{u}}{\partial t} = \frac{\kappa^2}{d^3} \frac{\partial \mathbf{u}}{\partial t} = -\frac{1}{\rho_0} \nabla \left( \frac{1}{d} \right) \left( \frac{\rho_0 \kappa^2}{d^2} \right) P + \alpha g \beta d \Theta \hat{\mathbf{k}} + \nu' \frac{1}{d^2} \nabla^2 \frac{\kappa \mathbf{u}}{d}$$

$$\frac{\kappa^2}{d^3} \frac{\partial \mathbf{u}}{\partial t} = -\frac{\kappa^2}{d^3} \nabla P + \alpha g \beta d \Theta \hat{\mathbf{k}} + \nu' \frac{\kappa}{d^3} \nabla^2 \mathbf{u}$$

$$\frac{\partial \mathbf{u}}{\partial t} = -\nabla P + \underbrace{\frac{\alpha g \beta d^4}{\kappa^2}}_{\text{Rayleigh number } R} \Theta \hat{\mathbf{k}} + \underbrace{\left( \frac{\nu'}{\kappa} \right)}_{\text{Prandtl number } Pr} \nabla^2 \mathbf{u}$$

②  $\therefore \frac{\partial \mathbf{u}}{\partial t} = -\nabla P + R Pr \Theta \hat{\mathbf{k}} + Pr \nabla^2 \mathbf{u}$

$$\frac{\partial (\beta d \Theta)}{\partial \left( \frac{d^2}{\kappa} t \right)} = \frac{\beta \kappa}{d} \frac{\partial \Theta}{\partial t} = -\beta \frac{\kappa}{d} w = \kappa \left( \frac{1}{d} \right)^2 \beta d \Theta = \frac{\kappa \beta^2}{d} \nabla^2 \Theta$$

$$\frac{\partial \Theta}{\partial t} - w = \frac{d}{\beta \kappa} \left( \frac{\kappa \beta^2}{d} \right) \nabla^2 \Theta$$

③  $\frac{\partial \Theta}{\partial t} - w = \nabla^2 \Theta$

Now we have perturbed, linearized, dimensionless eqns:

$$\begin{aligned} (1) & \quad \nabla \cdot \mathbf{u} = 0 \\ (2) & \quad \frac{\partial \mathbf{u}}{\partial t} = -\nabla p + RPr \theta \hat{\mathbf{k}} + Pr \nabla^2 \mathbf{u} \\ (3) & \quad \frac{\partial \theta}{\partial t} = -w = \nabla^2 \theta \end{aligned}$$

Take the curl of (2):

$$\nabla \times \left( \frac{\partial \mathbf{u}}{\partial t} \right) = \nabla \times (-\cancel{\nabla p} + RPr \theta \hat{\mathbf{k}} + Pr \nabla^2 \mathbf{u})$$

$$\frac{\partial}{\partial t} (\nabla \times \mathbf{u}) = RPr (\nabla \theta \times \hat{\mathbf{k}}) + Pr \nabla^2 (\nabla \times \mathbf{u})$$

$\vec{\omega} = \nabla \times \vec{u}$  : vorticity

$$\frac{\partial}{\partial t} (\omega) = RPr (\nabla \theta \times \hat{\mathbf{k}}) + Pr \nabla^2 \omega$$

→ take the curl again...

$$\nabla \times \left( \frac{\partial \omega}{\partial t} \right) = \nabla \times (RPr (\nabla \theta \times \hat{\mathbf{k}})) + \nabla \times (Pr \nabla^2 \omega)$$

$$\text{LHS: } \frac{\partial}{\partial t} (\nabla \times (\nabla \times \mathbf{u})) = \frac{\partial}{\partial t} ((\nabla \cdot \mathbf{u}) \nabla - \nabla \cdot \nabla \mathbf{u}) = -\frac{\partial}{\partial t} (\nabla^2 \mathbf{u})$$

$$\text{RHS: 1st term } RPr (\nabla \times (\nabla \theta \times \hat{\mathbf{k}})) = (\nabla \cdot \hat{\mathbf{k}} - \nabla \cdot \nabla \theta \hat{\mathbf{k}}) RPr$$

$$\text{2nd term } \nabla \times (Pr \nabla^2 \omega) = Pr (\nabla \times \nabla^2 (\nabla \times \mathbf{u})) = Pr (\nabla^2)^2 \mathbf{u} \quad \leftarrow \text{check}$$

$$\therefore -\frac{\partial}{\partial t} (\nabla^2 \mathbf{u}) = RPr (\nabla \cdot \hat{\mathbf{k}} - \nabla \cdot \nabla \theta \hat{\mathbf{k}}) + Pr (\nabla^2)^2 \mathbf{u}$$