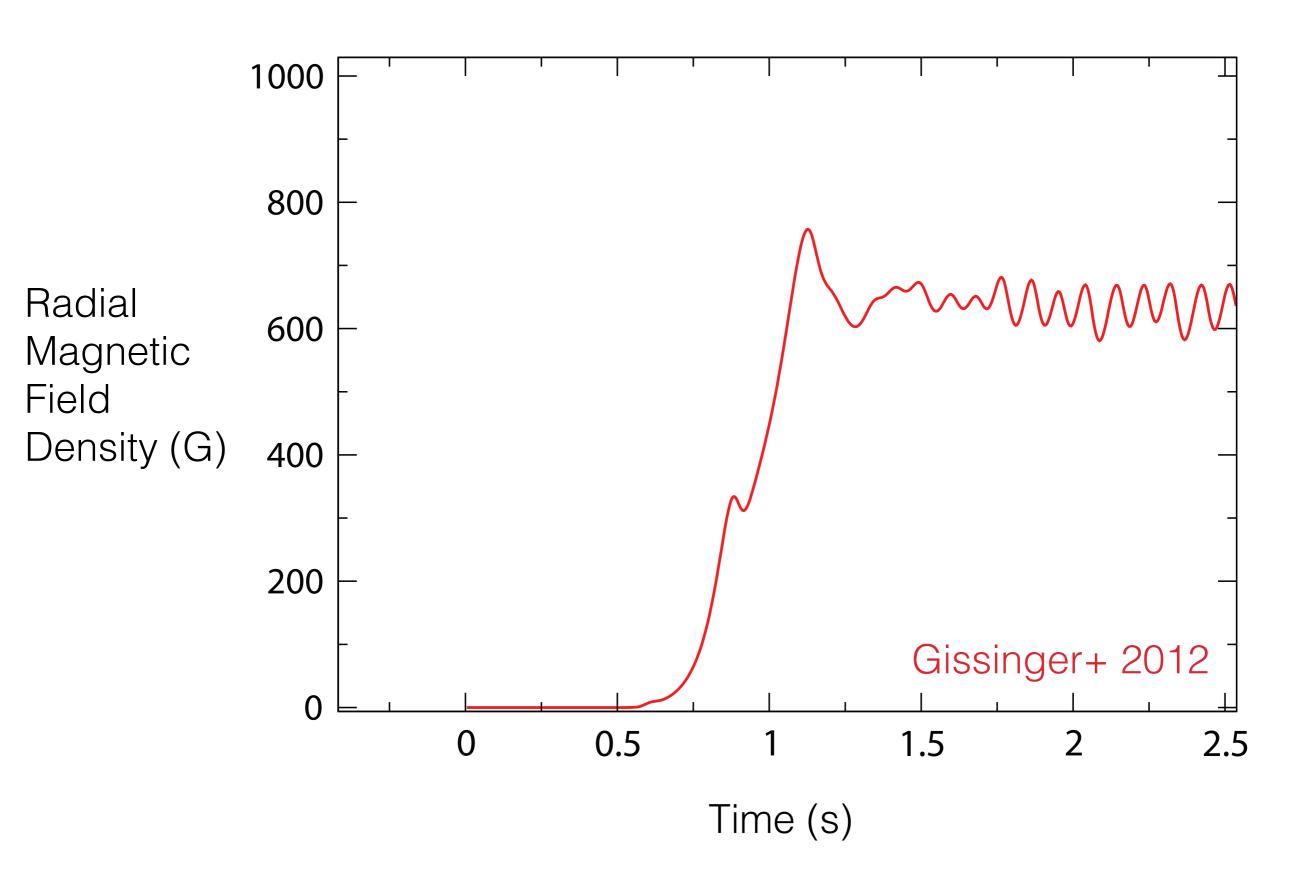
Exploring the saturation of the MRI via weakly nonlinear analysis

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The MRI evolves in stages.



set-up boundary conditions parameter range open questions, etc We solve the non-ideal MRI equations.

momentum

$$\partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} = -\frac{1}{\rho} \nabla P - \nabla \Phi + \frac{1}{\rho} \left(\mathbf{J} \times \mathbf{B} \right) - 2\Omega \times \mathbf{u} - \Omega \times (\Omega \times \mathbf{r}) + \nu \nabla^2 \mathbf{u}$$

induction

$$\partial_t \mathbf{B} = \nabla \times (\mathbf{u} \times \mathbf{B}) + \eta \nabla^2 \mathbf{B}$$

constraints

$$\nabla \cdot \mathbf{u} = 0$$

$$\nabla \cdot \mathbf{B} = 0$$

We solve the non-ideal MRI equations.

momentum

$$\partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} = -\frac{1}{\rho} \nabla P - \nabla \Phi + \frac{1}{\rho} \left(\mathbf{J} \times \mathbf{B} \right) - 2\Omega \times \mathbf{u} - \Omega \times (\Omega \times \mathbf{r}) + \nu \nabla^2 \mathbf{u}$$

magnetic

resistivity

$$\partial_t \mathbf{B} = \nabla \times (\mathbf{u} \times \mathbf{B}) + \eta \nabla^2 \mathbf{B}$$





$$\nabla \cdot \mathbf{u} = 0$$

$$\nabla \cdot \mathbf{B} = 0$$

We nondimensionalize and perturb the nonlinear MRI equations.

magnetic resistivity

microscopic viscosity

momentum

$$\partial_t \nabla^2 \Psi = \frac{2}{\beta} B_0 \partial_z \nabla^2 A + 2 \partial_z u_y + \frac{2}{\beta} J \left(A, \nabla^2 A \right) - J \left(\Psi, \nabla^2 \Psi \right) + \frac{1}{Re} \nabla^4 \Psi$$

$$\partial_t u = \frac{2}{\beta} B_0 \partial_z B_y - (2 - q) \Omega_0 \partial_z \Psi + \frac{2}{\beta} J(A, B_y) - J(\Psi, u_y) + \frac{1}{Re} \nabla^2 u_y$$

$$\partial_t A = B_0 \partial_z \Psi + J(A, \Psi) + \frac{1}{Rm} \nabla^2 A$$

$$\partial_t B_y = B_0 \partial_z u_y - q\Omega_0 \partial_z A + J(A, u_y) - J(\Psi, B_y) + \frac{1}{Rm} \nabla^2 B_y$$

momentum

viscous

$$\partial_t \nabla^2 \Psi = \frac{2}{\beta} B_0 \partial_z \nabla^2 A + 2 \partial_z u_y + \frac{2}{\beta} J \left(A, \nabla^2 A \right) - J \left(\Psi, \nabla^2 \Psi \right) + \frac{1}{Re} \nabla^4 \Psi$$

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$$\partial_t A = B_0 \partial_z \Psi + J\left(A, \Psi\right) + \frac{1}{Rm} \nabla^2 A$$
 resistive

$$\partial_t B_y = B_0 \partial_z u_y - q \Omega_0 \partial_z A + J(A, u_y) - J(\Psi, B_y) + \frac{1}{Rm} \nabla^2 B_y$$

momentum

viscous

$$\partial_t \nabla^2 \Psi = \frac{2}{\beta} B_0 \partial_z \nabla^2 A + 2 \partial_z u_y + \frac{2}{\beta} J \left(A, \nabla^2 A \right) - J \left(\Psi, \nabla^2 \Psi \right) + \frac{1}{Re} \nabla^4 \Psi$$

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shear

$$\partial_t A = B_0 \partial_z \Psi + J\left(A, \Psi\right) + \frac{1}{Rm} \nabla^2 A$$
 resistive

$$\partial_t B_y = B_0 \partial_z u_y - q \Omega_0 \partial_z A + J(A, u_y) - J(\Psi, B_y) + \frac{1}{Rm} \nabla^2 B_y$$

momentum

nonlinear

viscous

$$\partial_t \nabla^2 \Psi = \frac{2}{\beta} B_0 \partial_z \nabla^2 A + 2 \partial_z u_y + \frac{2}{\beta} J \left(A, \nabla^2 A \right) - J \left(\Psi, \nabla^2 \Psi \right) + \frac{1}{Re} \nabla^4 \Psi$$

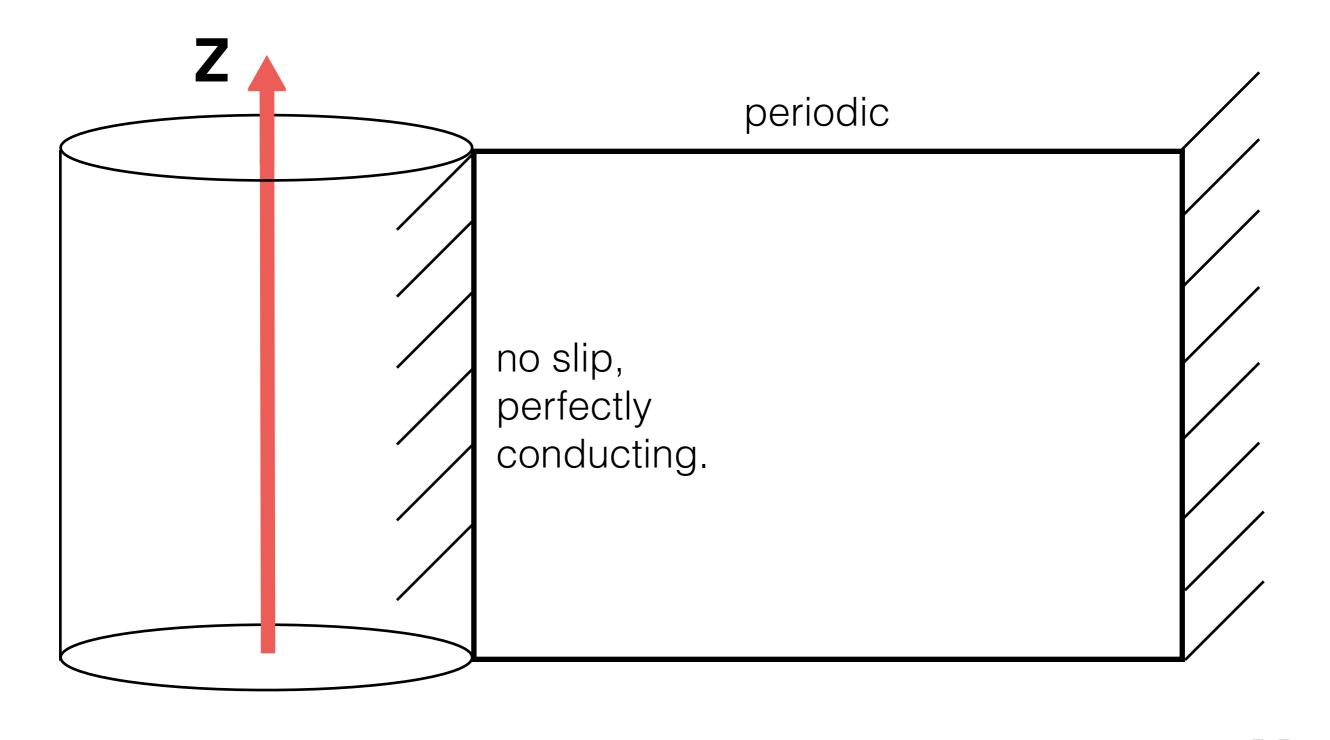
$$\partial_t u = \frac{2}{\beta} B_0 \partial_z B_y - \left(2 - q \right) \Omega_0 \partial_z \Psi + \frac{2}{\beta} J \left(A, B_y \right) - J \left(\Psi, u_y \right) + \frac{1}{Re} \nabla^2 u_y$$

shear

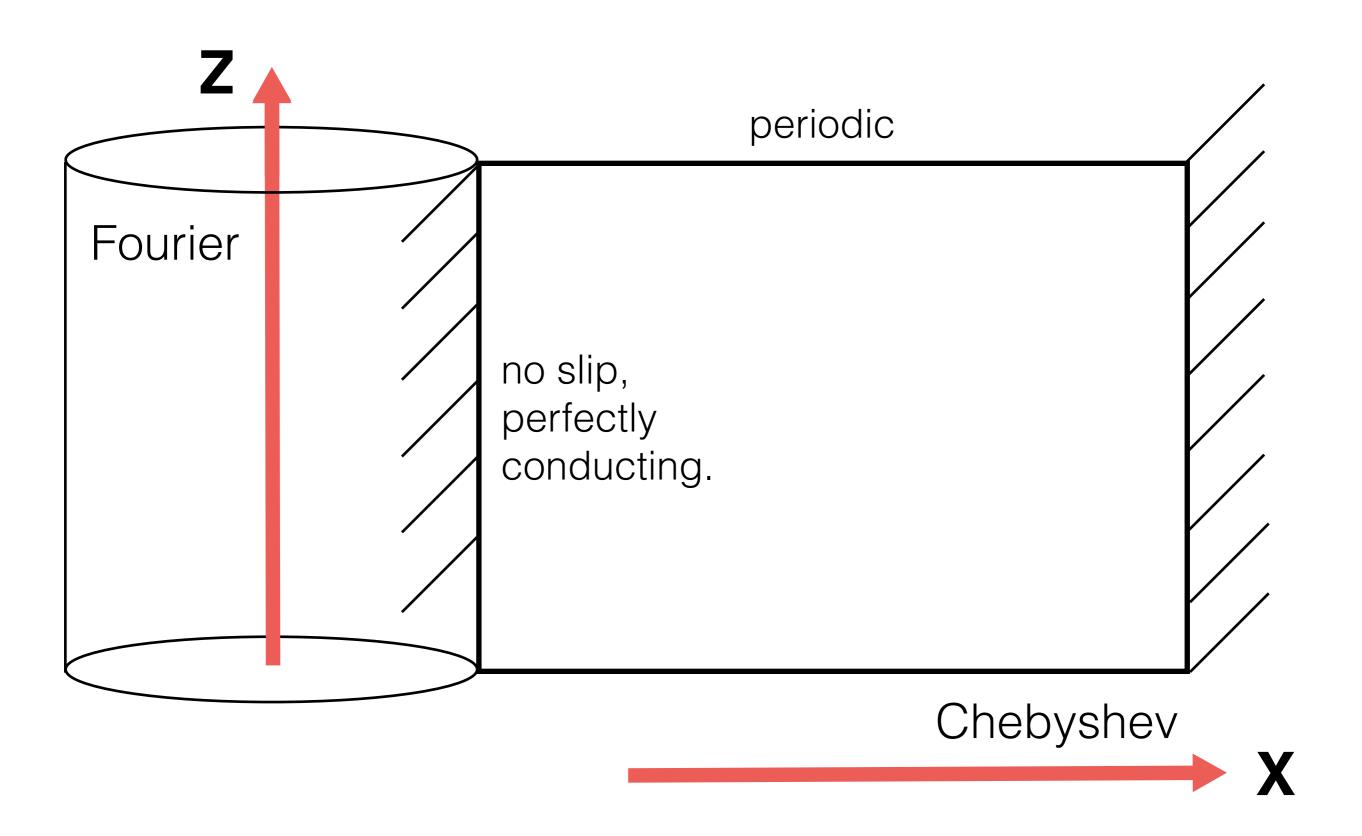
$$\partial_t A = B_0 \partial_z \Psi + J\left(A, \Psi\right) + \frac{1}{Rm} \nabla^2 A \quad \text{resistive}$$

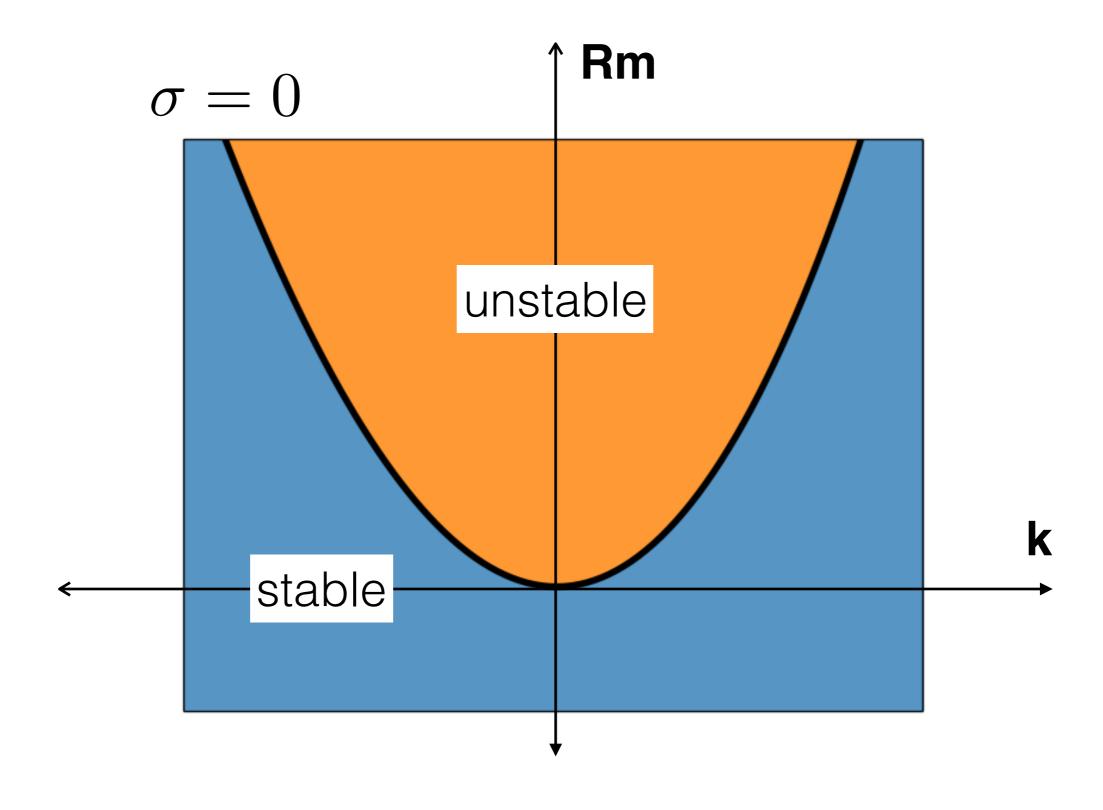
$$\partial_t B_y = B_0 \partial_z u_y - q \Omega_0 \partial_z A + J (A, u_y) - J (\Psi, B_y) + \frac{1}{Rm} \nabla^2 B_y$$

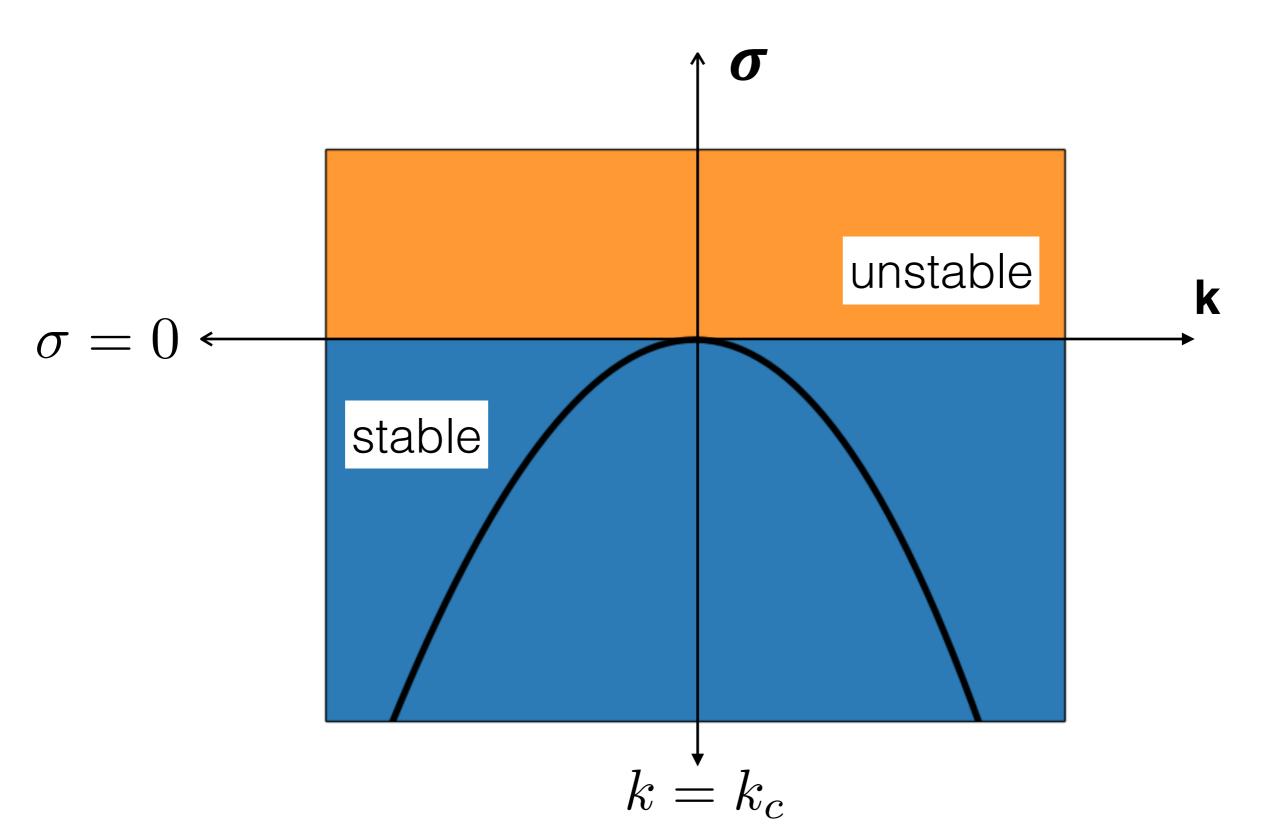
We use experimentally relevant boundary conditions.

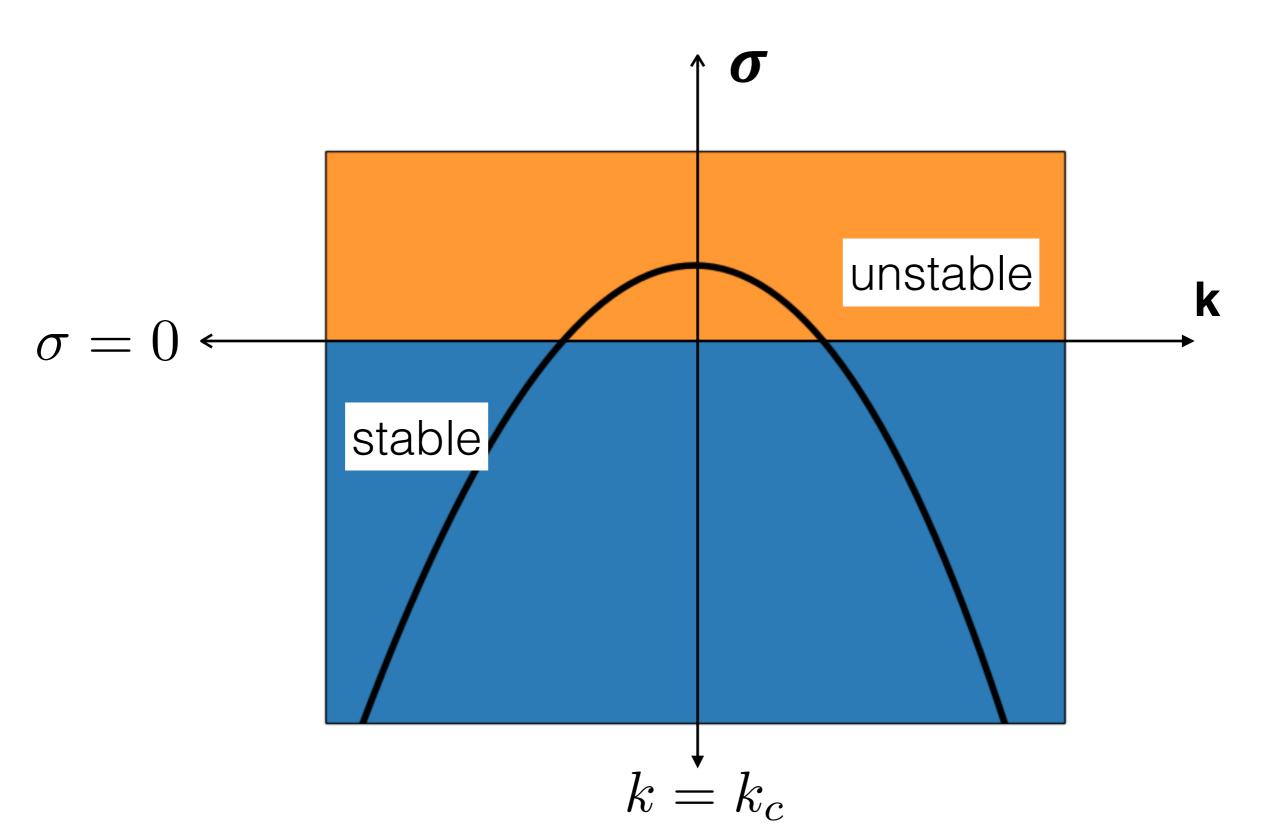


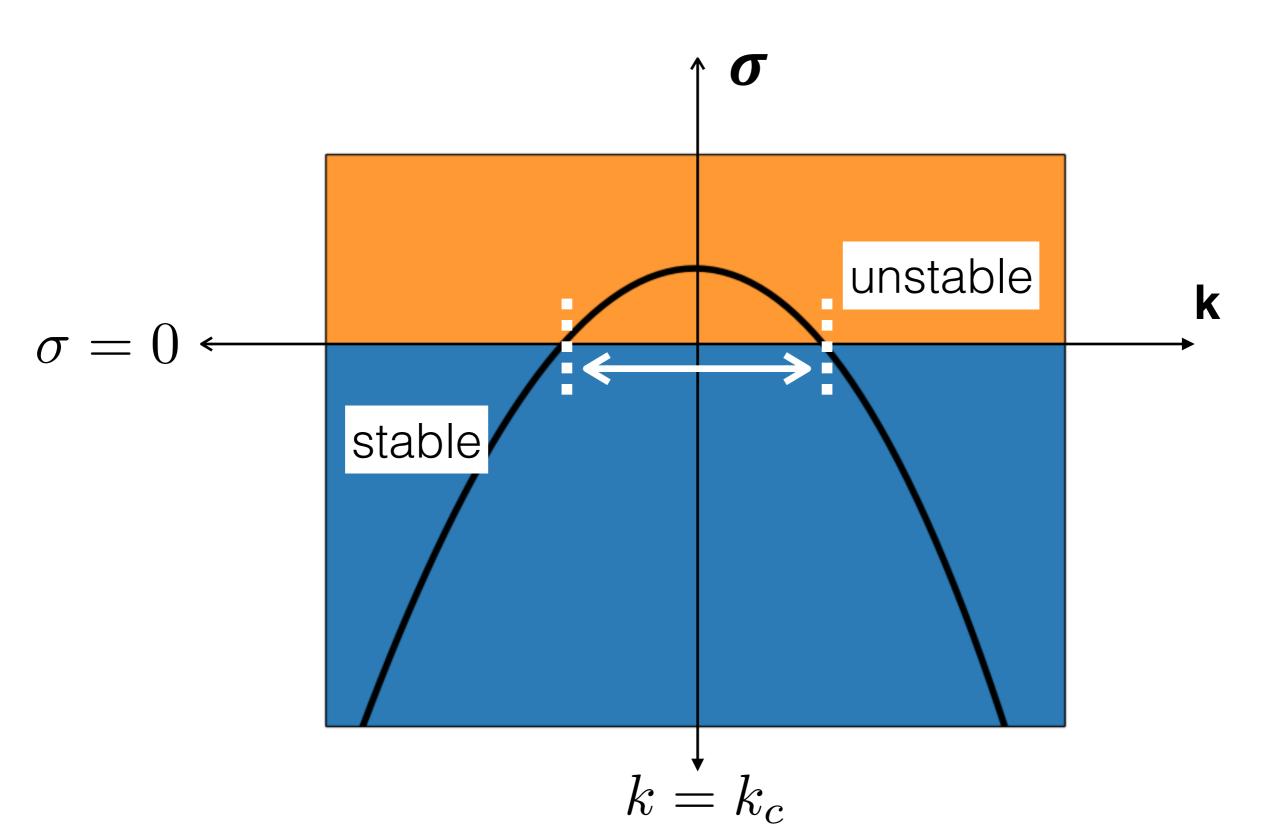
Dedalus is a general-purpose spectral code.



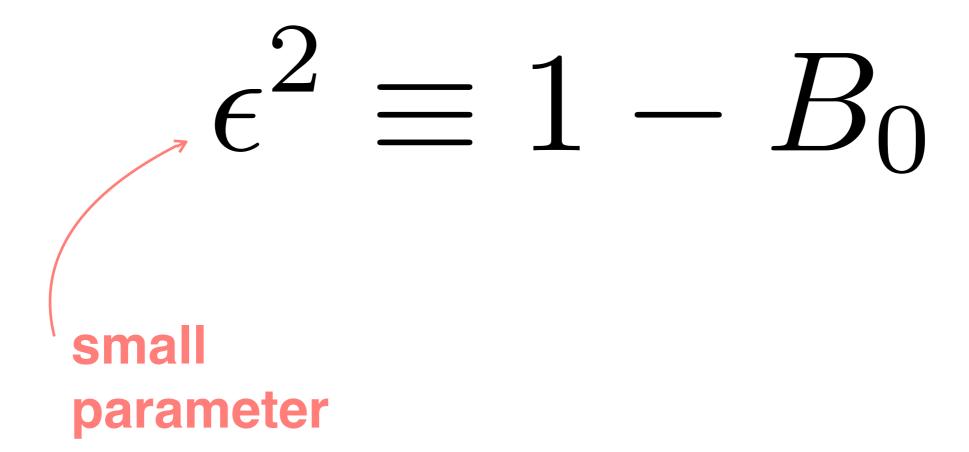




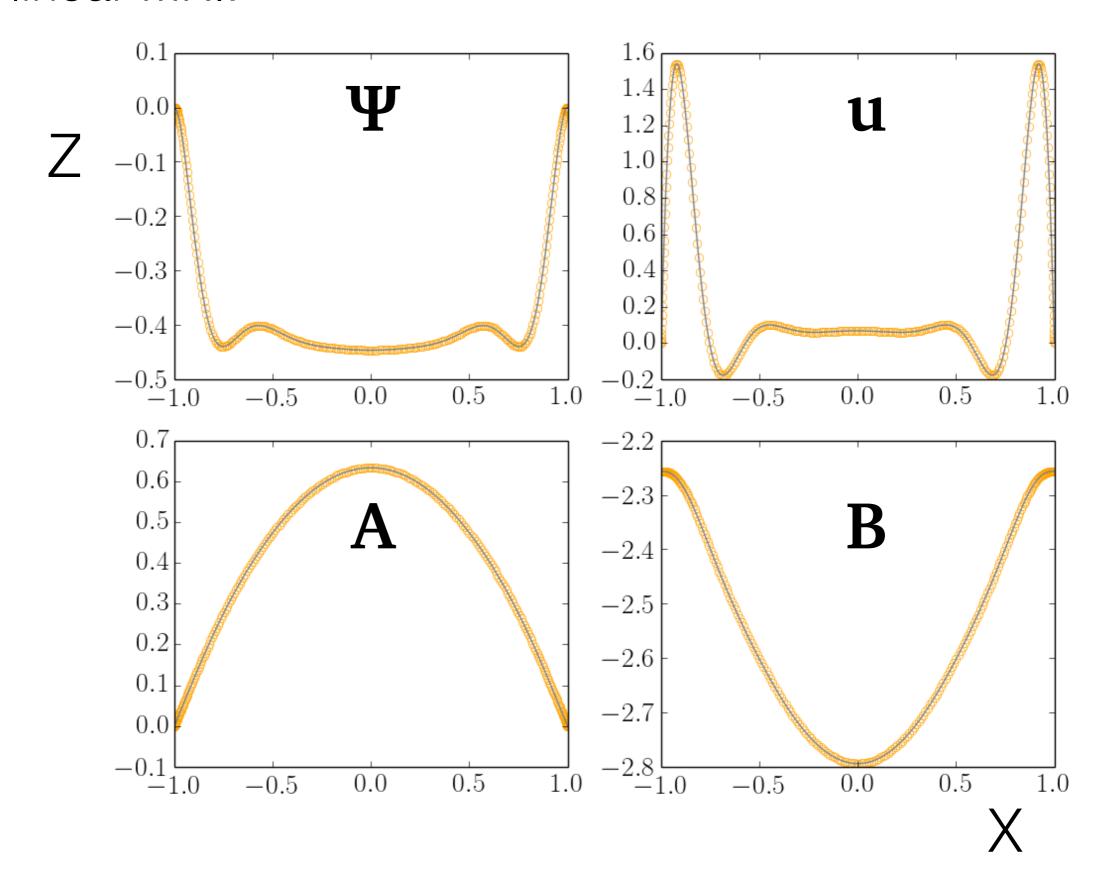




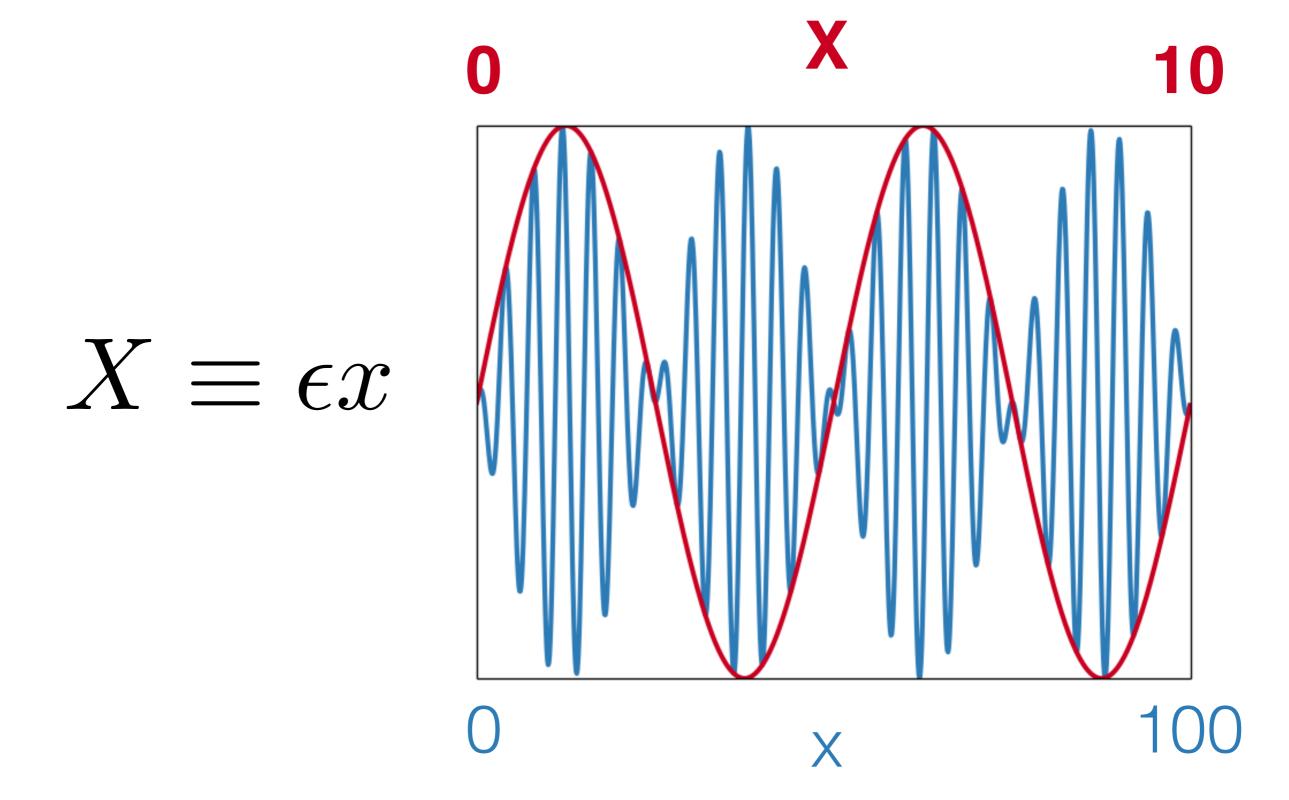
Tune the most unstable mode just over the threshold of instability.



Identify the most unstable mode of the linear MRI.



Multiscale analysis tracks the evolution of fast and slow variables.

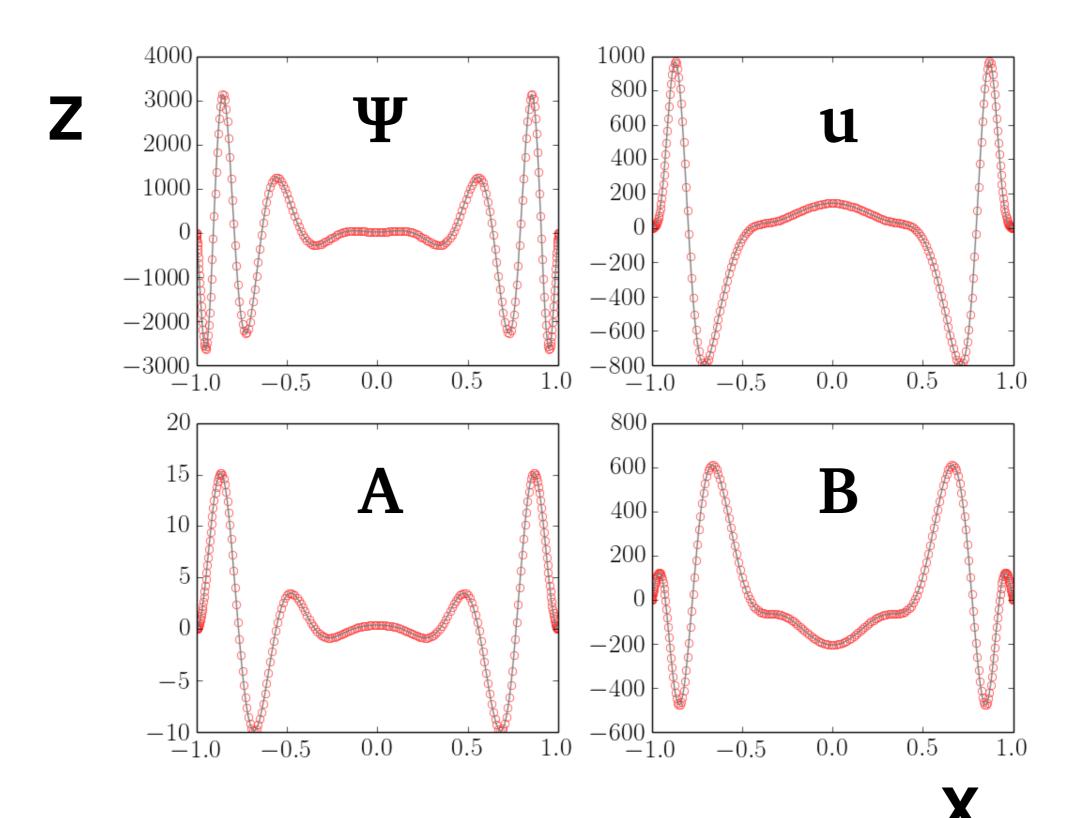


Equations are solved in a matrix formulation.

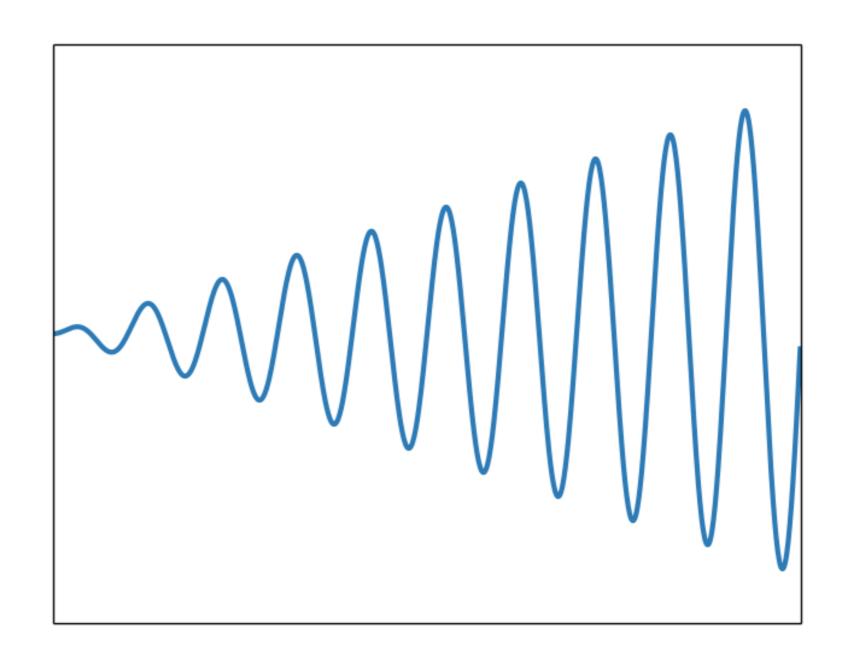
The fluid quantities are expanded in a perturbation series.

$$\mathbf{V} = \epsilon \mathbf{V_1} + \epsilon^2 \mathbf{V_2} + \epsilon^3 \mathbf{V_3} + \dots$$

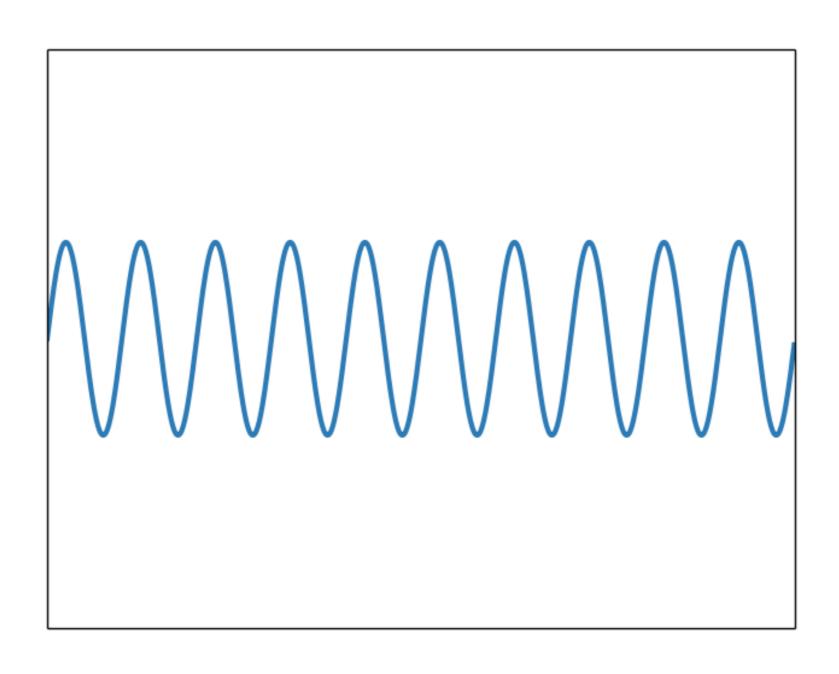
something about boundary layers?



The removal of secular terms yields solvability criteria.



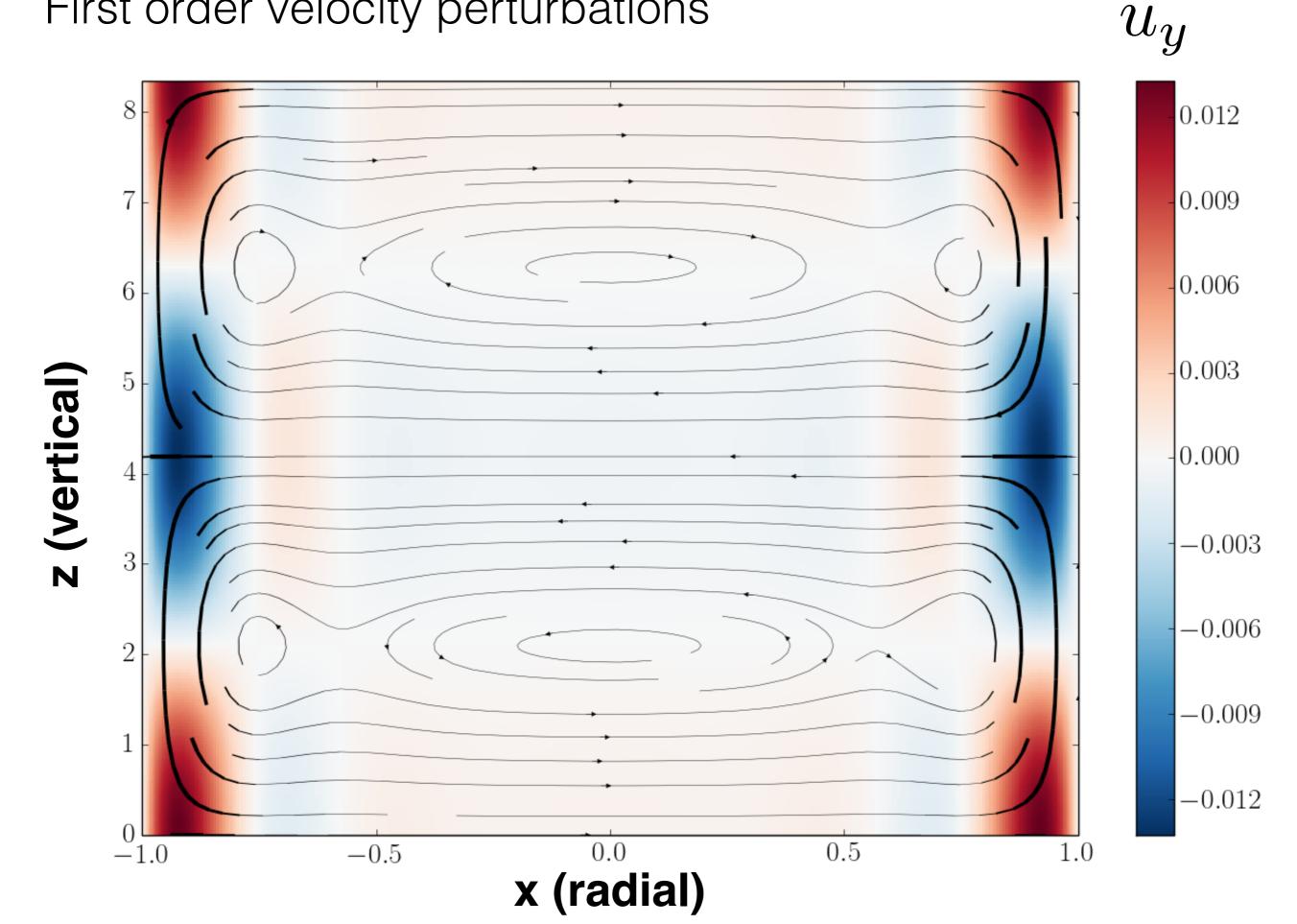
The removal of secular terms yields solvability criteria.



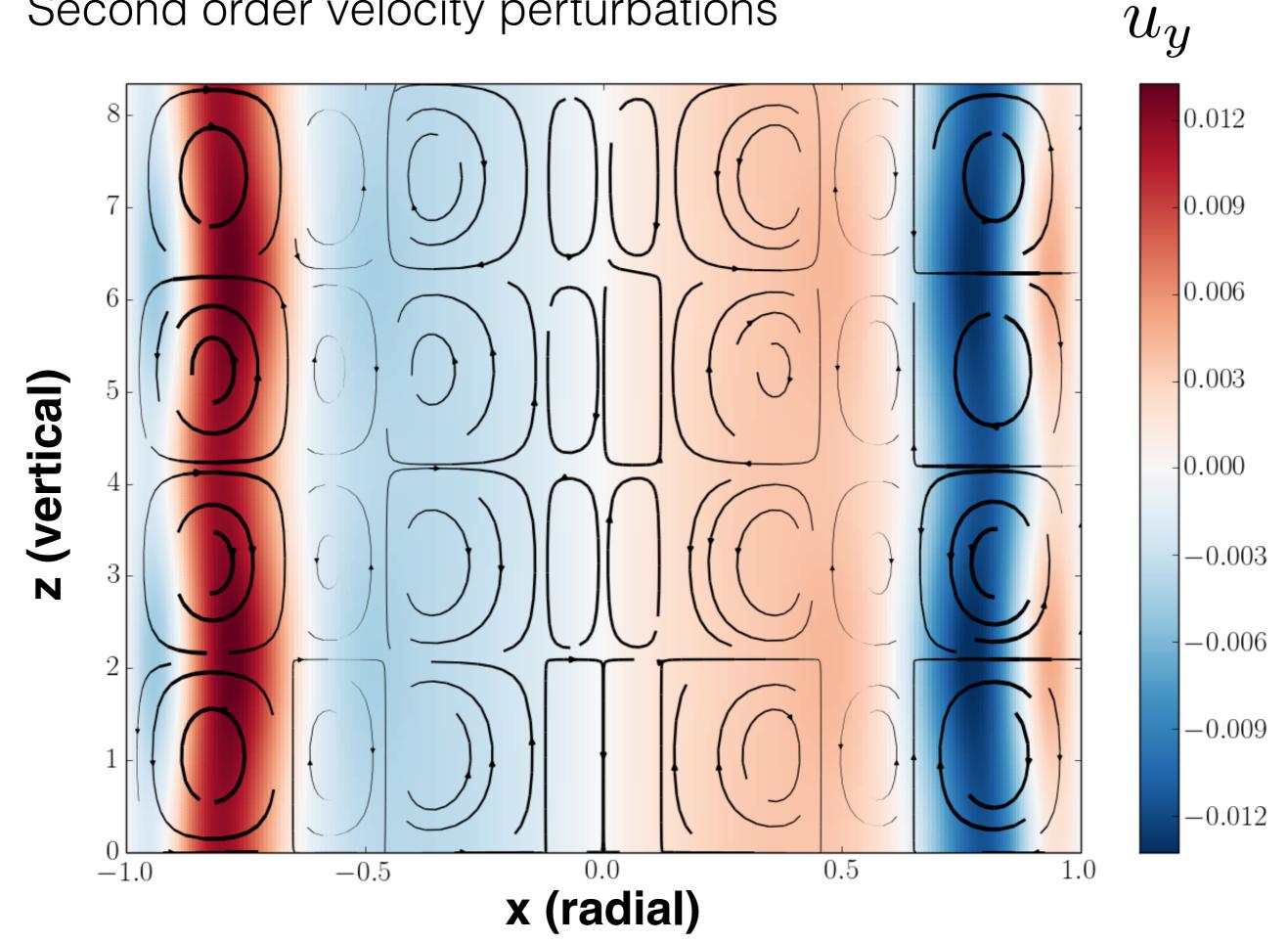
The result is an amplitude equation.

$$a\partial_T \alpha = -b\partial_Z \alpha - c\alpha \left| \alpha^2 \right| + h\partial_Z^2 \alpha + giQ^3 \alpha$$

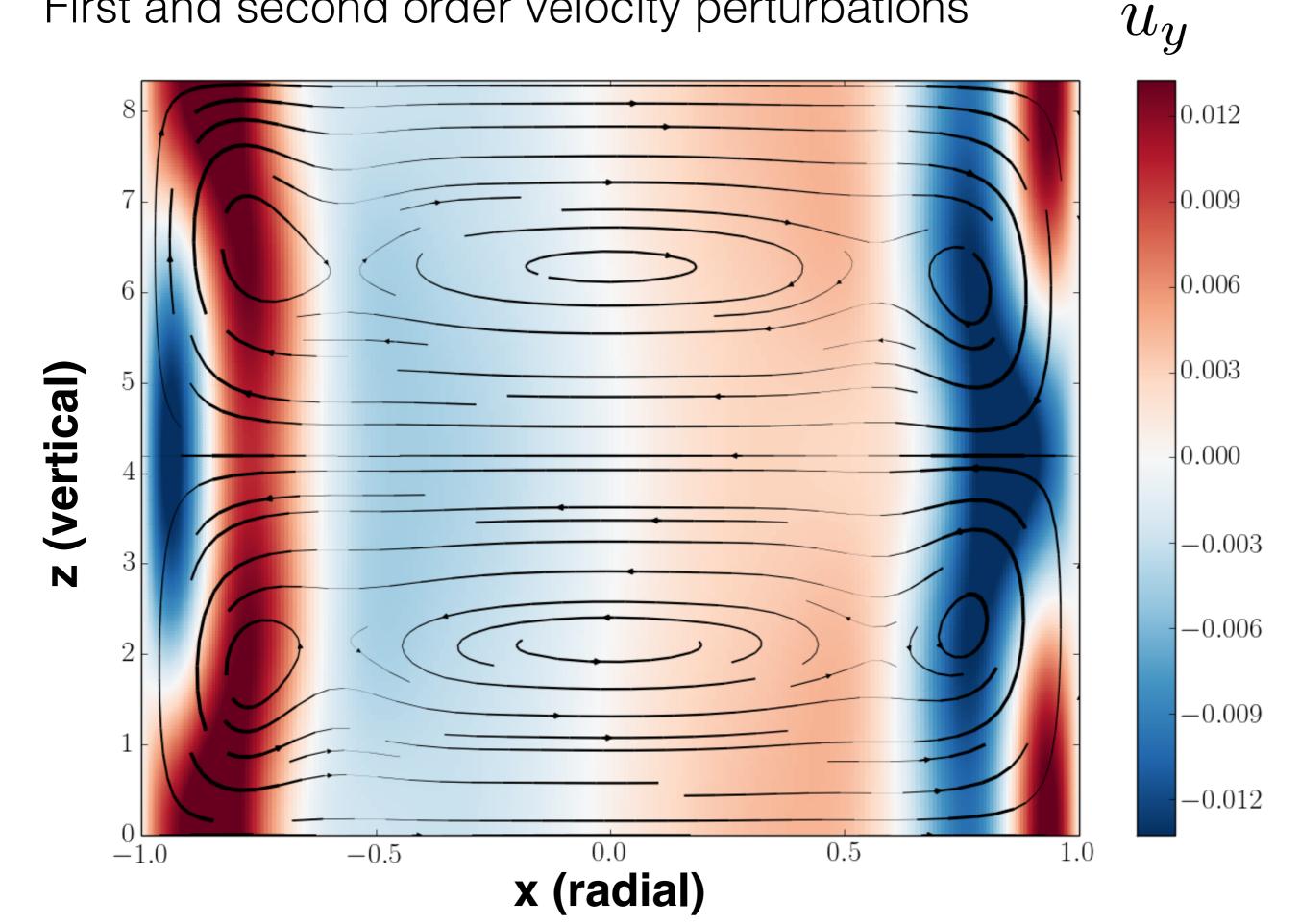
First order velocity perturbations



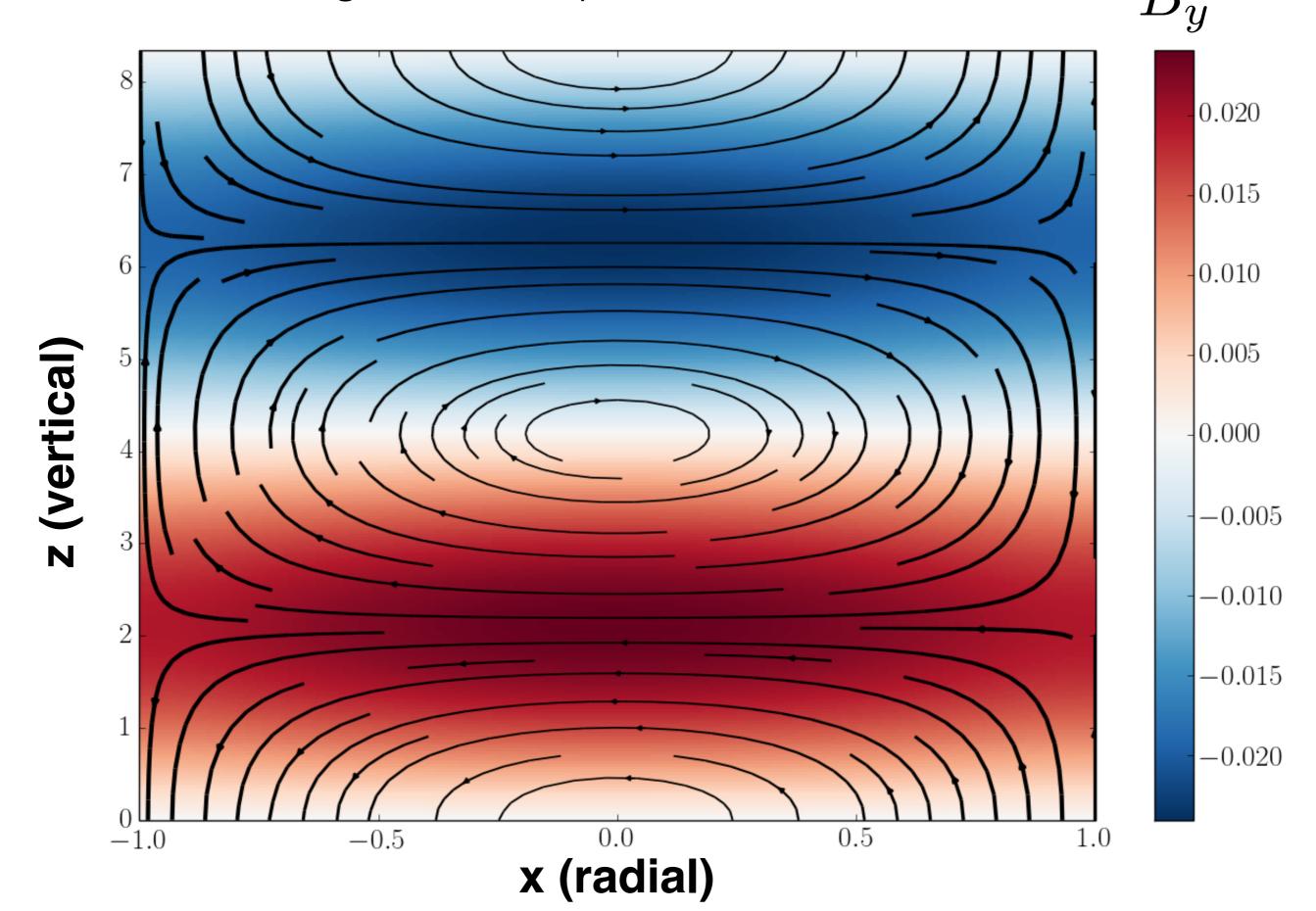
Second order velocity perturbations



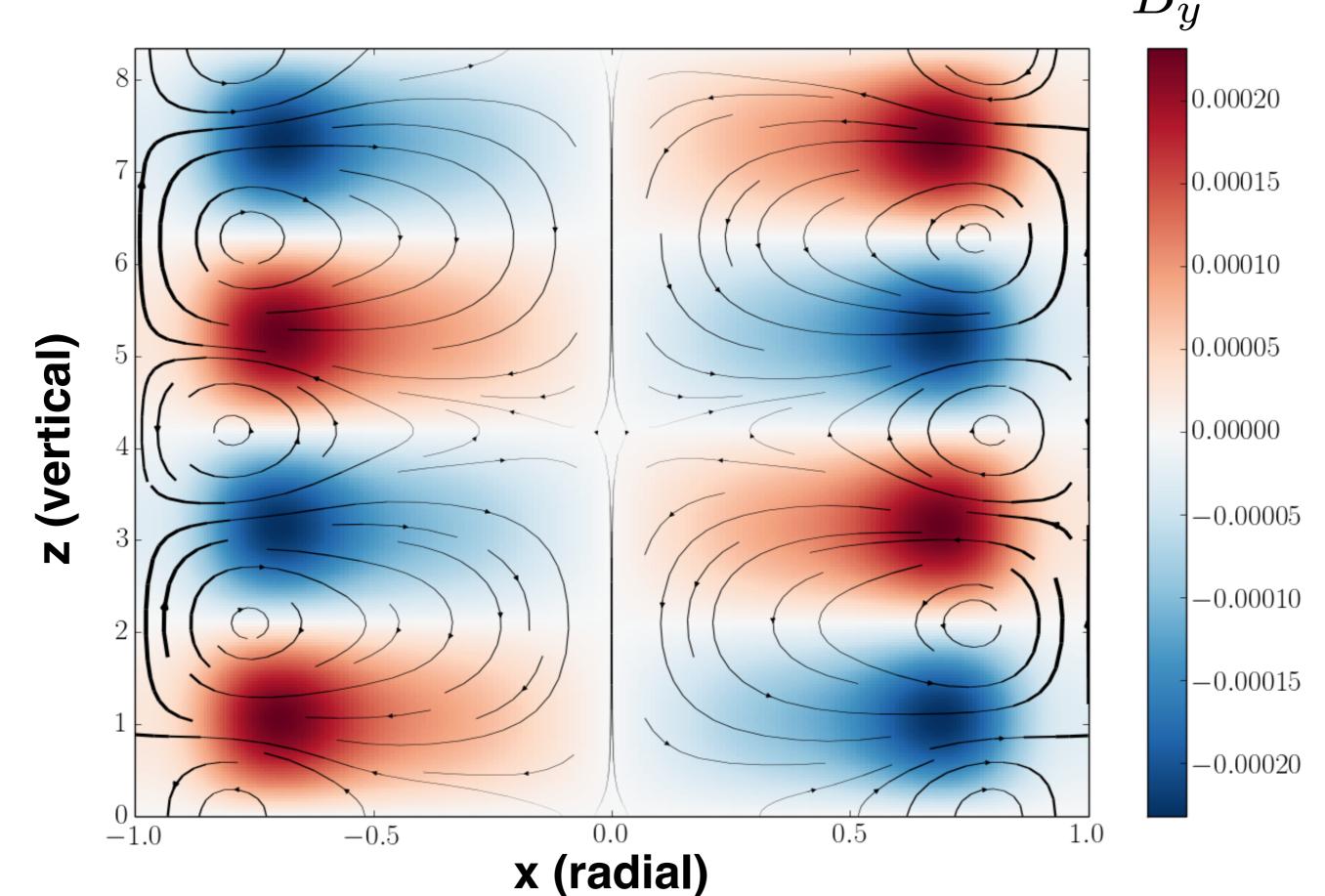
First and second order velocity perturbations



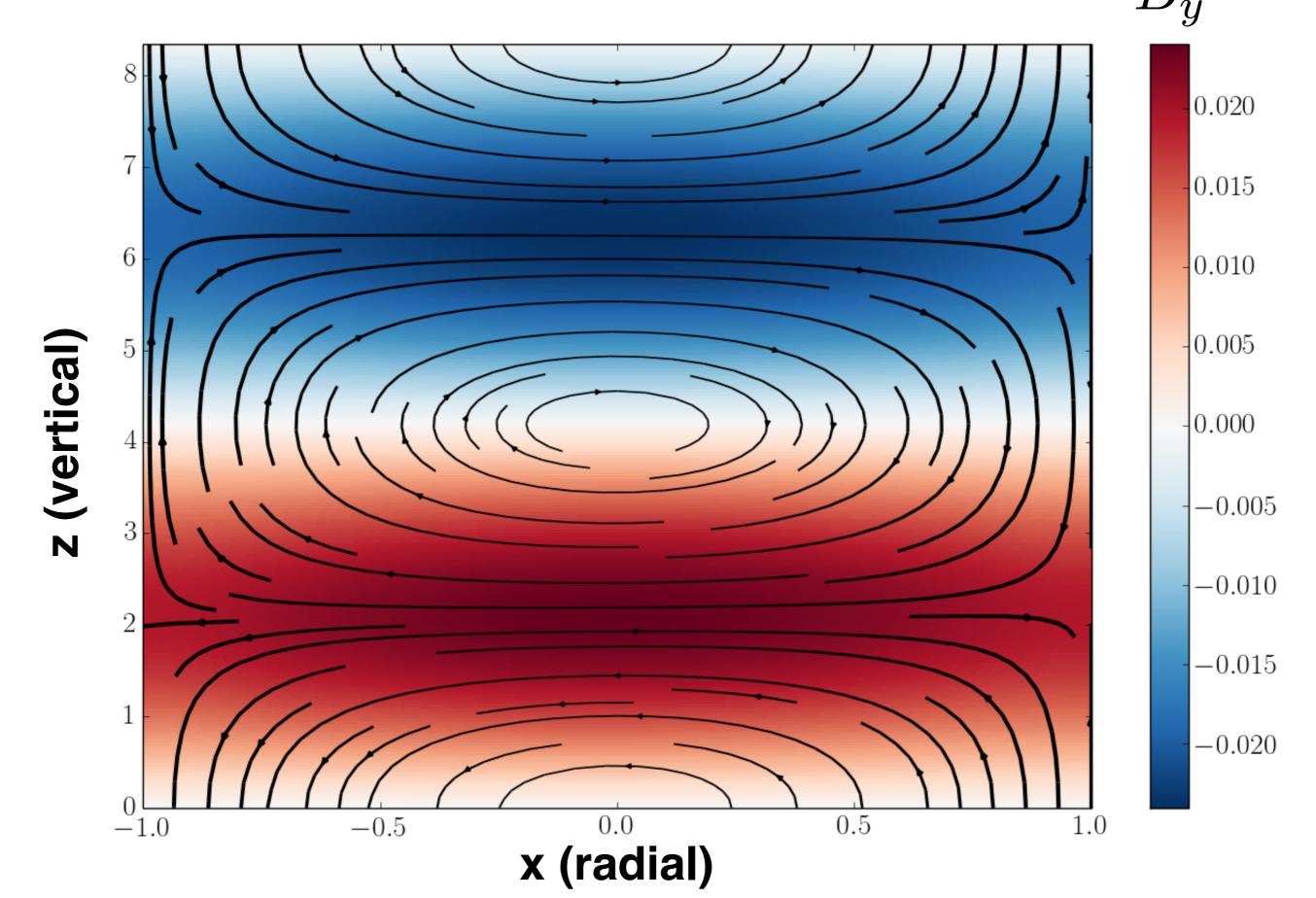
First order magnetic field perturbations



Second order magnetic field perturbations



First and second order magnetic field perturbations ${\cal B}_y$



Future work:

non-thin gap approximation helical MRI explore parameter space comparison to experiment