

# Range Optimal Trajectories for an Aircraft Flying in the Vertical Plane

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Range optimal trajectories for an aircraft flying in the vertical plane are obtained from Pontryagin's minimum principle. Control variables are the load factor, which appears nonlinearly in the equations of motion, and the throttle setting, which appears only linearly. Both controls are subject to fixed bounds. Additionally, a dynamic pressure limit is imposed, which represents a first-order state-inequality constraint. For fixed flight time, initial coordinates, and final coordinates of the trajectory, the effect of the load factor limit on the resulting optimal switching structure is studied. All trajectories involve singular control along arcs with active dynamic pressure limit. For large flight times the optimal switching structures have not yet been found.

#### I. Introduction

REAT efforts are being undertaken to develop real-time, near-optimal feedback algorithms either for enhancement of aircraft performance by optimizing specified maneuvers or as autonomous guidance schemes for short- and mediumrange air-to-air missiles.<sup>1-5</sup> Open-loop control logics obtained by state-of-the-art optimization techniques are an important tool in testing the accuracy and finding the limits of such feedback laws. In a recent study<sup>6,7</sup> open-loop optimal control solutions in conjunction with perturbation techniques have been used directly to develop feedback algorithms. In this context minimum-time intercept trajectories, or, often equivalently, maximum-range trajectories for fixed flight time play an important role in modern air combat scenarios.

In the present paper Pontryagin's minimum principle  $(PMP)^{8-10}$  is applied to determine range optimal trajectories for an aircraft flying in the vertical plane. State variables are energy E, altitude h, and flight-path angle  $\gamma$ ; control variables are load factor n and throttle setting  $\eta$ . Control  $\eta$  appears only linearly in the equations of motion and is subject to fixed bounds  $0 \le \eta \le 1$ . Additionally, a dynamic pressure limit is imposed on the trajectory, which in the context of optimal control represents a first-order state inequality constraint.

For sufficiently large fixed final times a good approximation to the maximum-range trajectory can be obtained by breaking the flight into three phases: 1) climb to the dashpoint, 2) steady-state cruise at the dash-point, and 3) final transient to reach the terminal constraints. The dash-point is the point of maximum sustainable speed on the level-flight envelope. The associated state values can be found by solving the maximization problem

$$\max_{E, h, \gamma; n, \eta} \nu, \nu = \sqrt{2g(E - h)}$$

subject to the level-flight constraints

 $\gamma = 0$ 

 $\dot{\gamma} = 0$ 

 $\dot{E} = 0$ 

and the control constraints for n and  $\eta$ .

In the present paper we consider trajectories that start at the dash-point. The final conditions are chosen rather arbitrarily, but such that the aircraft is forced to dive to lower altitudes at final time. Furthermore, the final energy/altitude pair is located inside the level-flight envelope and off the dynamic pressure limit.

For fixed boundary conditions and fixed flight time we study the effect of the load factor limit  $|n| \le n_{\text{max}}$  on the structure of the optimal solution. Six different switching structures involving singular control on state constrained arcs are encountered if  $n_{\text{max}}$  is varied between  $n_{\text{max}} = \infty$  and  $n_{\text{max}} = 5$ . The switching structures identified in this paper appear to be of a general nature and are not restricted to the specific boundary conditions considered here. However, for the boundary conditions prescribed in this paper, the structure of the optimal solution has not yet been found for flight times greater than 62 s.

## II. Aircraft Model

The equations of motion of an aircraft flying in the vertical plane are

$$\dot{E} = (\eta T - D) \frac{v}{W} \tag{1}$$

$$\dot{h} = v \sin \gamma \tag{2}$$

$$\dot{\gamma} = \frac{g}{v} \left( \frac{L}{W} - \cos \gamma \right) \tag{3}$$

$$\dot{x} = v \cos \gamma \tag{4}$$

The specific energy E, replacing velocity  $\nu$ , the altitude h, the flight-path angle  $\gamma$ , and the range x are the state variables. Load factor n and the power setting  $\eta$  are the control vari-

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ables. The mass m of the vehicle and the gravitational acceleration g are assumed to be constant, namely

$$m = 16818 \text{ kg}$$

$$g = 9.80665 \text{ m/s}^2$$

The vehicle weight is given by W = mg. Velocity  $\nu$  is a short notation for  $\nu = \sqrt{2g(E - h)}$ . The air density  $\rho$  in kilograms per meter<sup>3</sup> is given by

$$\rho(h) = 1.225 e^{y}$$

$$y = -0.12122693 \ 10^{-3} \ h + r - 1.0228055$$

$$r = 1.0228055 e^{-z}$$

$$z = -3.48643241 \ 10^{-5} \ h + 3.50991865 \ 10^{-9} \ h^2$$

$$+ -8.33000535 \ 10^{-14} \ h^3 + 1.15219733 \ 10^{-18} \ h^4$$

The speed of sound in meters per second is given by

$$a(h) = 20.0468 \sqrt{\theta}$$

where the temperature  $\theta$  is given by

$$\theta = 292.1 - 8.87743 \ 10^{-3} \ h + 0.193315 \ 10^{-6} \ h^2 + 3.72 \ 10^{-12} \ h^3$$

In these expressions h is altitude in meters. The Mach number is given by M = v/a(h). The lift L, the drag D, and the maximum thrust T in Newtons, respectively, are given as functions of h, M, and n

$$q = \frac{1}{2}\rho(h)v^2S$$

$$L = Wn$$

$$D = q \left( C_{D0}(M) + K(M) \frac{W^2}{q^2} n^2 \right)$$

$$C_{D0} = \frac{a_4 M^4 + a_3 M^3 + a_2 M^2 + a_1 M + a_0}{b_4 M^4 + b_3 M^3 + b_2 M^2 + b_1 M + b_0}$$

$$K = \frac{c_4 M^4 + c_3 M^3 + c_2 M^2 + c_1 M + c_0}{d_5 M^5 + d_4 M^4 + d_3 M^3 + d_2 M^2 + d_1 M + d_0}$$

$$T(h, M) = e_5(M)h^5 + e_4(M)h^4 + e_3(M)h^3 + e_2(M)h^2$$

$$+ e_1(M)h + e_0(M)$$

where for  $i = 0, 1, \ldots, 5$ 

$$e_i(M) = (f_{i5}M^5 + f_{i4}M^4 + f_{i3}M^3 + f_{i2}M^2 + f_{i1}M + f_{i0})$$

The numerical values of the constants  $a_i$ ,  $b_i$ ,  $c_i$ ,  $d_i$ , and  $f_{ij}$  are given in Tables 1-3 and represent a high performance fighter-interceptor.

# III. Problem Formulation

The problem under consideration is that of finding control functions  $\eta(t)$  and n(t) that steer an aircraft from prescribed initial states: energy  $E_0$ , altitude  $h_0$ , and flight-path angle  $\gamma_0$  to

Table 1 Coefficients for  $C_{D0}$ -model

i	$a_i$		$b_i$			
0	-2.61059846050	10-2	+7.29821847445	10-1		
1	+8.57043966269	$10^{-2}$	-3.25219000620	$10^{0}$		
2	+1.07863115049	10-1	+ 5.72789877344	$10^{0}$		
3	-6.44772018636	$10^{-2}$	-4.57116286752	$10^{0}$		
4	+1.64933626507	$10^{-2}$	+1.37368651246	$10^{0}$		

prescribed final states: energy  $E_f$ , altitude  $h_f$ , and flight-path angle  $\gamma_f$  in prescribed flight time  $(t_f - t_0)$  (without loss of generality  $t_0 = 0$ ) such that the downrange x is maximized. Along the optimal trajectory a set of state and control constraints has to be satisfied. Explicitly the problem can be stated in Mayer form as follows:

$$\min - x(t_f) \tag{5}$$

subject to the state equations (1-4), the control constraints

$$-\eta \le 0 \tag{6}$$

$$\eta - 1 \le 0 \tag{7}$$

$$-n-n_{\max} \le 0 \tag{8}$$

$$+ n - n_{\text{max}} \le 0 \tag{9}$$

the state constraint

$$C_0(E, h, \gamma, x) := \nu - \nu_{\max}(h) \le 0$$
 (10)

and the boundary conditions

$$E(0) = 38029.207 \text{ m}$$
 (11a)

$$h(0) = 12119.324 \text{ m}$$
 (11b)

$$x(0) = 0 \text{ m}$$
 (11d)

(11c)

(11e)

$$h(t_f) = 942.292 \text{ m}$$
 (11f)

$$\gamma(t_f) = -0.2 \text{ rad} \tag{11g}$$

$$x(t_f)$$
 to be optimized (11h)

and the final time  $t_f$  prescribed, e.g.,

$$t_f = 60 \text{ s} \tag{12}$$

Table 2 Coefficients for K-model

 $\gamma(0) = 0$  rad

 $E(t_f) = 9000 \text{ m}$ 

i	ci		di			
0	+1.23001735612	10 <sup>0</sup>	+1.42392902737	10+1		
1	-2.97244144190	100	-3.24759126471	10+1		
2	+2.78009092756	10 <sup>0</sup>	+2.96838743792	10+1		
3	-1.16227834301	10 <sup>0</sup>	-1.33316812491	10+1		
4	+1.81868987624	10 <sup>-1</sup>	+2.87165882405	10+1		
5			-2.27239723756	10-1		

Table 3 Coefficients for thrust model

i	$f_{ij}, j=0$	$f_{ij}, j=1$	$f_{ij}, j=2$		
0	+0.11969995703 106	-0.35217318620 106	+0.60452159152 106		
1	-0.14644656421 10 <sup>5</sup>	+0.51808811078 105	- 0.95597112936 10 <sup>5</sup>		
2	$-0.45534597613 10^3$	+0.23143969006 104	- 0.38860323817 10 <sup>4</sup>		
3	+ 0.49544694509 10 <sup>3</sup>	$-0.2248231045510^4$	+ 0.39771922607 104		
4	- 0.46253181596 10 <sup>2</sup>	$+0.20894683419 10^3$	- 0.36835984294 10 <sup>3</sup>		
5	$+0.12000480258\ 10^{1}$	$-0.53807416658\ 10^{1}$	+0.94529288471 101		
i	$f_{ij}$ , $j=3$	$f_{ij}, j=4$	$f_{ij}, j=5$		
$\frac{i}{0}$	$f_{ij}$ , $j = 3$ $-0.43042985701 \ 10^6$	$f_{ij}$ , $j = 4$ + 0.13656937908 10 <sup>6</sup>	$f_{ij}, j = 5$ $-0.16647992124 10^5$		
$\frac{\overline{i}}{0}$					
_	-0.43042985701 10 <sup>6</sup>	+0.13656937908 106	-0.16647992124 10 <sup>5</sup>		
ĺ	-0.43042985701 10 <sup>6</sup> +0.83271826575 10 <sup>5</sup>	+ 0.13656937908 10 <sup>6</sup> - 0.32867923740 10 <sup>5</sup>	-0.16647992124 10 <sup>5</sup> +0.49102536402 10 <sup>4</sup>		
1 2	-0.43042985701 10 <sup>6</sup> +0.83271826575 10 <sup>5</sup> +0.12357128390 10 <sup>4</sup>	+0.13656937908 10 <sup>6</sup> -0.32867923740 10 <sup>5</sup> +0.55572727442 10 <sup>3</sup>	-0.16647992124 10 <sup>5</sup> +0.49102536402 10 <sup>4</sup> -0.23591380327 10 <sup>3</sup>		

Here  $n_{\text{max}}$  is a specified constant denoting the maximum allowed absolute value of the load factor n = L/W. In state constraint Eq. (10)  $v_{\text{max}}(h)$  is a specified function of altitude h. With  $v_{\text{max}}(h)$  chosen appropriately, this covers the important case of a dynamic pressure constraint. Boundary conditions (11a-11c) refer to the dash-point or high speed point. The boundary conditions (11e-11g) are picked more or less arbitrarily. The only important features are that  $h(t_f) < h(0)$  and  $v(t_f) = \sqrt{2g(E(t_f) - h(t_f))} < v_{\text{max}}(h(t_f))$ , i.e., in the altitude-velocity chart the prescribed final point of the trajectory is located to the left of the state constraint (10).

#### IV. Relaxed Problem Formulation

Existence theorems in optimal control theory require convexity of a certain velocity set or hodograph.9 Given a state equation  $\dot{x} = f(x, u), x \in \mathbb{R}^n, u \in \mathbb{R}^m$  with admissible controls  $u \in U \subset \mathbb{R}^m$ , the hodograph at some fixed state  $x_0 \in \mathbb{R}^n$  is defined as the set  $S = \{\dot{x} \in \mathbb{R}^n | \dot{x} = f(x_0, u), u \in U\}$  of possible state rates. For state equations (1-4) with controls  $[\eta, n] \in \mathbb{R}^2$ subject to the constraints equations (6-9), the hodograph is clearly nonconvex as indicated in Fig. 1. This nonconvexity results from the fact that with throttle  $\eta = 0$  the smallest energy rate ( = the largest negative energy rate) can be achieved only for  $n = \pm n_{\text{max}}$ . If, for example, an optimal control problem is formulated such that along some time interval the aircraft is required to dissipate energy at maximum rate while keeping the flight-path angle constant, then a near-optimal solution would be achieved by setting  $\eta = 0$  and by rapidly switching the load factor between  $n = \pm n_{\text{max}}$  such that, in the average,  $\dot{\gamma} = 0$  is achieved. In the limit the optimal solution would be achieved only by letting the load factor n chatter at an infinite rate. However, this limiting process is not well-defined in the sense that the limit load factor function of time is not Lebesgue integrable, and hence is not an admissible controller. Strictly speaking, a solution does not exist in this case for the problem at hand. It may be desirable, however, to determine the solution associated with the limiting case of chattering control. To do this it is necessary to convexize the hodograph. That means one has to add to the set of admissible state rates all those state rates that are located along straight lines in between any two admissible state rates of the original hodograph. For aircraft models with quadratic drag polar, such as the one used in this paper, this can be achieved by rewriting state equation (1) as

$$\dot{E} = \left[\delta(T - D + D_{\text{max}}) - D_{\text{max}}\right] \frac{v}{W} \tag{13}$$

with

$$D_{\max}(M, h) := D(M, h, n_{\max})$$

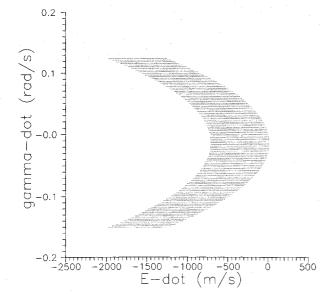


Fig. 1 Hodograph for unrelaxed problem.

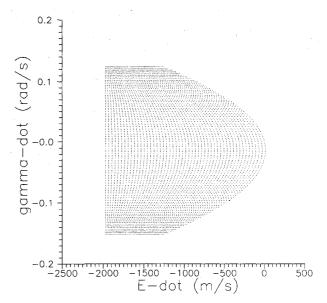


Fig. 2 Hodograph for relaxed problem.

The new control  $\delta$  replaces the old control  $\eta$  and is subject to the constraints

$$-\delta \le 0 \tag{14}$$

$$\delta - 1 \le 0 \tag{15}$$

Of course, the control  $\delta$  is rather artificial and lacks the physical meaning of a throttle parameter. It should be noted, however, that generally in optimal control the controls only play the role of bookkeeping parameters used to parameterize the set of admissible state rates. This arbitrariness of the controls becomes more clear by noting that under control transformations that leave the hodograph unchanged, the state and state rate functions of time associated with the optimal solution remain unchanged, whereas the optimal control function of time may become completely different.

To summarize, the relaxed control problem, which is used for the remainder of this paper, is given by

$$\min - x(t_f)$$

subject to the state equations (13) and (2-4), control constraints (14), (15), (8), and (9), state constraint (10), boundary conditions (11a-11g), and the final time  $t_f$  prescribed as in Eq. (12). That the hodograph of the relaxed dynamical system is equal to the convex hull of the hodograph of the original dynamical system can be seen by direct comparison (see Figs. 1 and 2). For the present problem, both hodographs can be constructed analytically by first determining the respective boundaries [by sequentially setting  $\eta = 0$ ,  $\eta = 1$ ,  $n = -n_{\text{max}}$ , and  $n = +n_{\text{max}}$  in Eqs. (13) and (2-4), and  $\delta = 0$ ,  $\delta = 1$ ,  $n = -n_{\text{max}}$ , and  $n = +n_{\text{max}}$  in Eqs. (14), (15), (8), and (9)] and then noting that all points in the interior of these boundaries also belong to the respective hodograph.

The state inequality constraint (10) being active on some time interval  $[\tau_1, \tau_2]$  {i.e.,  $C_0(E, h, \gamma, x) \equiv 0$  on  $[\tau_1, \tau_2]$ } is equivalent to

$$C_0(E, h, \gamma, x) \equiv 0 \quad \text{at} \quad t = \tau_1 \tag{16}$$

$$C_1(E, h, \gamma, x; n, \delta) \equiv 0$$
 on  $[\tau_1, \tau_2]$  (17)

where

$$C_1(E, h, \gamma, x; n, \delta) := \frac{\mathrm{d}}{\mathrm{d}t} C_0(E, h, \gamma, x)$$

= 
$$[\delta(T - D + D_{\text{max}}) - D_{\text{max}}]g/W - v \sin \gamma (v'_{\text{max}} + g/v)$$
 (18)

#### V. Minimum Principle

The relaxed optimization problem as stated earlier is solved by applying the PMP.<sup>8-10</sup> The PMP requires that at every point in time the controls are to be chosen such that the variational Hamiltonian

$$H(E, h, \gamma, x, \lambda_E, \lambda_h, \lambda_{\gamma}, \lambda_x, \delta, n) = \lambda_E \dot{E} + \lambda_h \dot{h} + \lambda_{\gamma} \dot{\gamma} + \lambda_x \dot{x}$$
(19)

is minimized subject to all control constraints. Let the vector valued function  $g: \mathbb{R}^6 \rightarrow \mathbb{R}^4$  be defined by

$$g_1(E, h, \gamma, x; \delta, n) = -\delta$$

$$g_2(E, h, \gamma, x; \delta, n) = \delta - 1$$

$$g_3(E, h, \gamma, x; \delta, n) = -n - n_{\text{max}}$$

$$g_4(E, h, \gamma, x; \delta, n) = +n - n_{\text{max}}$$
(20)

so that inequalities (14), (15), (8), and (9) can be written concisely as  $g \le 0$ . Then Lagrange multipliers  $\lambda_E$ ,  $\lambda_h$ ,  $\lambda_\gamma$ , and  $\lambda_x$  are solutions of the adjoint equations

$$\dot{\lambda}_{E} = -\frac{\partial H}{\partial E} - \sigma^{T} \frac{\partial g}{\partial E} - \mu \frac{\partial C_{1}}{\partial E}$$

$$\dot{\lambda}_{h} = -\frac{\partial H}{\partial h} - \sigma^{T} \frac{\partial g}{\partial h} - \mu \frac{\partial C_{1}}{\partial h}$$

$$\dot{\lambda}_{\gamma} = -\frac{\partial H}{\partial \gamma} - \sigma^{T} \frac{\partial g}{\partial \gamma} - \mu \frac{\partial C_{1}}{\partial \gamma}$$

$$\dot{\lambda}_{x} = -\frac{\partial H}{\partial x} - \sigma^{T} \frac{\partial g}{\partial x} - \mu \frac{\partial C_{1}}{\partial x}$$
(21)

where  $\sigma_i$ ,  $i=1,\ldots,4$  and  $\mu$  are multipliers associated with constraints  $g_i \leq 0$ ,  $i=1,\ldots,4$  and  $C_1 \leq 0$  [ $C_1$  as given in Eq. (18)], respectively. On time intervals where constraint Eq. (10) is not active [i.e.,  $C_0(E, h, \gamma, x) < 0$ ], multiplier  $\mu$  is identically

$$\mu = 0$$
 if  $C_0(t) < 0$  (22)

On these intervals multipliers  $\sigma_i$ ,  $i = 1, \ldots, 4$  are determined at each instant of time from the Kuhn-Tucker conditions<sup>8-11</sup> applied to the finite dimensional parameter optimization problem

$$(\delta, n) = \arg\min_{g \le 0} H \tag{23}$$

At times where state constraint (10) is active [i.e.,  $C_0$  (E, h,  $\gamma$ , x = 0, multipliers  $\sigma$  and  $\mu$  are determined from the Kuhn-Tucker conditions applied to the finite dimensional parameter optimization problem

$$(\delta, n) = \arg \min_{g \le 0, C_1 \le 0} H$$
 (24)

and components of multiplier vector  $\sigma$  are non-negative along intervals where the associated constraints are active, i.e.

$$\sigma_i \ge 0 \quad \text{if} \quad g_i = 0, \qquad i = 1, \dots, 4$$
 (26)

Note that the optimization problems Eqs. (23) and (24) simply represent the min-Hamiltonian operation required by the PMP.

# VI. Possible Control Logics

At each instant of time controls n and  $\delta$  are determined from the minimum principle given by Eq. (23) in case  $C_1 < 0$  and Eq. (24) in case  $C_1 = 0$ . Since the Hamiltonian H and the constraint functions g and  $C_1$  are smooth functions of their arguments  $\delta$  and n, the Kuhn-Tucker conditions imply that at the solution point  $[\delta^*, n^*]$  the following conditions have to be satisfied at each instant of time:

$$\frac{\partial}{\partial \delta} \left( H + \sigma^T g + \mu C_1 \right) = 0 \tag{27}$$

$$\frac{\partial}{\partial n} \left( H + \sigma^T g + \mu C_1 \right) = 0 \tag{28}$$

$$g \le 0 \tag{29}$$

$$C_1 \le 0 \tag{30}$$

$$\sigma \ge 0 \tag{31}$$

$$\mu \ge 0 \tag{32}$$

$$\begin{bmatrix} \Delta \delta \\ \Delta n \end{bmatrix} \begin{bmatrix} \frac{\partial^{2}(H + \sigma^{T}g + \mu C_{1})}{\partial \delta^{2}} & \frac{\partial^{2}(H + \sigma^{T}g + \mu C_{1})}{\partial \delta \partial n} \\ \frac{\partial^{2}(H + \sigma^{T}g + \mu C_{1})}{\partial n \partial \delta} & \frac{\partial^{2}(H + \sigma^{T}g + \mu C_{1})}{\partial n^{2}} \end{bmatrix} \begin{bmatrix} \Delta \delta \\ \Delta n \end{bmatrix} \geq 0 \quad (33)$$

for all  $(\Delta \delta, \Delta n) \in \mathbb{R}^2$  satisfying  $(\partial h/\partial \delta)\Delta \delta + (\partial h/\partial n)\Delta n = 0$ . Here vector function h contains exactly the active components of the inequality constraints  $g \le 0$  and  $C_1 \le 0$ . The set of active control constraints and the character of the solution of Eqs. (27-33) depends greatly on the direction of the multiplier vector  $[\lambda_E, \lambda_h, \lambda_{\gamma}, \lambda_x]^T$ , and through state constraint (10) also on the state itself. The explicit analysis for solving this finite dimensional constrained minimization problem as well as application of the generalized Legendre-Clebsch condition are given in Ref. 12. It is helpful to define

$$n_0 \equiv \frac{\lambda_{\gamma}}{\lambda_E} \frac{gq}{2v^2 WK} \tag{34}$$

$$n_1 = \frac{\lambda_E}{\lambda_x} \frac{v}{g} \cos^2 \gamma \frac{\partial (D_{\text{max}} v / W)}{\partial h} + 2 \cos \gamma$$
 (35)

$$n_1 = \frac{\lambda_E}{\lambda_x} \frac{v}{g} \cos^2 \gamma \frac{\partial (D_{\text{max}} v / W)}{\partial h} + 2 \cos \gamma$$

$$n_2 = \frac{q}{W} \sqrt{\frac{T - vW/g \sin \gamma (v'_{\text{max}} + q/v) - qC_{D0}}{qK}}$$
(35)

$$n_{3} = \cos \gamma - \frac{\left[ \left( \frac{\partial D_{\text{max}}}{\partial E} \left( \frac{v_{\text{max}}' v}{g} + 1 \right) + \frac{\partial D_{\text{max}}}{\partial h} \right) \frac{g}{W} + v_{\text{max}}'^{2} + v_{\text{max}}'' v} \right] v \sin \gamma}{\frac{g}{v} \left( v_{\text{max}}' v + g \right) \cos \gamma}$$
(37)

In both cases, as a consequence of the Kuhn-Tucker conditions, components of multiplier vector  $\sigma$  are zero along intervals where the associated constraint is not active

$$\sigma_i = 0 \quad \text{if} \quad g_i > 0, \qquad i = 1, \ldots, 4$$
 (25)

$$\delta_{1} \equiv \frac{D_{\text{max}} + \frac{vW}{g} \sin \gamma \left( v_{\text{max}}' + \frac{g}{v} \right)}{T - D \mid_{n = n_{\text{max}}} + D_{\text{max}}}$$
(38)

Here  $n_0$  is obtained from  $\partial H/\partial n = 0$ ,  $n_1$  solves  $\lambda_{\gamma} \equiv 0$  in the singular case 6b in the Appendix,  $n_2$  is implied by  $C_1 = 0$  with  $\delta = 1$ , and  $n_3$  is required for the singular control case 11b. The expression  $\delta_1$  stems from  $C_1 = 0$  with  $n = n_{\text{max}}$ . Then the different possible control logics are shown in the Appendix. A derivation of the results is given in Ref. 12.

#### VII. Transversality and Corner Conditions

All transversality and corner conditions are given such that the first variation of the cost function Eq. (5)  $J = -x(t_f)$  is zero.<sup>8-10</sup> With boundary conditions (11) this yields

$$\lambda_{\rm r} = -1 \tag{39}$$

In case of final time  $t_f$  to be minimized [i.e., cost function Eq. (5)  $J = -x(t_f)$  being replaced by  $J = t_f$ ], the associated boundary condition is

$$H(t_f) = 1 \tag{40}$$

The Hamiltonian H is continuous throughout the time interval  $[0, t_f]$ . At any corner point, say at time  $t_c$ , this yields an optimality condition on the switching time  $t_c$ , namely

$$H(t_c^+) - H(t_c^-) = 0 (41)$$

Here and later superscripts + and - denote evaluation just right and just left of the time under consideration, respectively. At the beginning, say  $t_1$ , of a state constrained are additional conditions are

$$C_0(t_1) = 0$$

$$\lambda_E^+ = \lambda_E^- - l_0 \frac{\partial C_0}{\partial E}$$

$$\lambda_h^+ = \lambda_h^- - l_0 \frac{\partial C_0}{\partial h}$$

$$\lambda_\gamma^+ = \lambda_\gamma^- - l_0 \frac{\partial C_0}{\partial \gamma}$$

$$\lambda_x^+ = \lambda_x^- - l_0 \frac{\partial C_0}{\partial \gamma}$$

$$\lambda_x^+ = \lambda_x^- - l_0 \frac{\partial C_0}{\partial \gamma}$$
(43)

where  $l_0$  is a constant multiplier. The end, say  $t_2$ , of a constrained arc is determined by the continuity of the Hamiltonian. The jump in multipliers in Eq. (43) is implied by the interior point condition Eq. (42).

#### VIII. Supplementary Optimality Conditions

Along constrained arcs we have the sign conditions<sup>14</sup>

$$\sigma_i \ge 0$$
 on arcs where  $g_i = 0$   $i = 1, \ldots, 4$  (44)

$$\mu \ge 0$$
 on arcs where  $C_0 = 0$  (45a)

$$\mu \le 0$$
 on arcs where  $C_0 = 0$  (45b)

Along singular arcs an additional optimality condition is the generalized Legendre-Clebsch condition. <sup>12,13</sup> This condition is already considered in the possible control logics stated in Sec. VI. The explicit analysis is given in Ref. 12.

#### IX. Switching Structures

Equation (5) is solved for fixed final time Eq. (12) subject to the equations of motion Eqs. (13) and (2-4) and the boundary conditions (11). As a first step only control constraints (14) and (15) and the state constraint (10) are enforced, whereas load factor limits (8) and (9) are neglected. The associated switching structure turns out to be

$$(S1) 1 - 7 - 11b - 7 - 1$$

where any number i in the sequence refers to case i of the possible control logics listed in the Appendix. The load factor n = L/W increases rapidly near the final time  $t_f$  and reaches a maximum value of approximately  $n_{\max} = 56.5886$ . Mathematically this value is perfectly reasonable, as will be explained heuristically in the next section. To make the solution meaningful from an engineering point of view lower values of  $n_{\max}$  have to be enforced. Starting with switching structure (S1) this procedure is done by reducing the load factor limit Eq. (9) in steps

$$|n| \le n_{\max}, \qquad n_{\max} = 56, 55, \ldots, 5$$

In the process we observe the following switching structures:

(S1) 
$$1-7-11b-7-1$$
 for  $n_{\text{max}} \in [56.6, \infty]$ 

(S2) 
$$1-7-11b-7-1-3$$
 for  $n_{\text{max}} \in [32.7, 56.5]$ 

(S3) 
$$1-7-11b-7-1-3-5$$
 for  $n_{\text{max}} \in [22.7, 32.6]$  (46)

(S4) 
$$1-7-11b-9-7-1-3-5$$
 for  $n_{\text{max}} \in [20.8, 22.6]$ 

(S5) 
$$1-7-11b-9-3-5$$
 for  $n_{\text{max}} \in [18.2, 20.7]$ 

(S6) 
$$1-7-11b-9-5$$
 for  $n_{\text{max}} \in [5.0, 18.1]$ 

# X. Large Load Factors near Final Time

Figure 3 shows the time history of the load factor along the solution without the bounds Eqs. (8) and (9) (i.e.,  $n_{\text{max}} = \infty$ ). It may be helpful to provide some explanation for why the peak-value occurs at the final time and how the peak-value depends on the boundary conditions. To this end, suppose the conditions (11e-11g) are replaced with

$$E(t_f) = E_f = \frac{v_{\text{max}}(h_f)^2}{2g} + h_f$$
 (47)

$$h(t_f) = h_f = 9000 (48)$$

$$\lambda_{\gamma}(t_f) = 0 \tag{49}$$

i.e., the prescribed final state lies on the dynamic pressure limit and the final flight path angle is free. Then numerical calculations show that the switching structure associated with the solution of this problem is given by switching structure (S1) of the previous section with the last two arcs deleted. Now, if boundary condition (47) is replaced by

$$E(t_f) = E_f - \Delta E \tag{50}$$

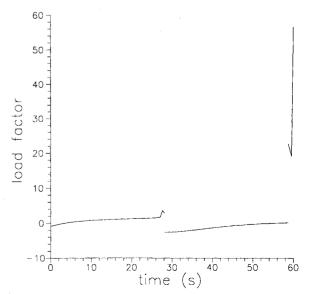


Fig. 3 Load factor n vs time t for  $n_{\text{max}} = \infty$  [switching structure (S1)].

for some  $\Delta E > 0$ , then, if the load factor is unbounded, the "optimal maneuver" for the aircraft would be to fly exactly as in the solution of the previous case (i.e., boundary conditions (47-49) until  $E(t_f) = E_f$  is reached and then to impulsively apply a high load factor  $n \to \infty$  on an infinitesimal time interval  $[t_f, t_f + \delta t_f]$ ,  $\delta t_f \to 0$ , such that the energy drops instantaneously to the prescribed value  $E_f - \Delta E$ . By noting that, in the dynamical equations, the load factor appears linearly in the  $\dot{\gamma}$ -equation and quadratically in the  $\dot{E}$ -equation, we expect that along this infinitesimal arc  $\delta E \sim n^2 \delta t$ , whereas  $\delta \gamma \sim n \delta t$ . Hence, with n and  $\delta t_f$  such that  $\delta E = -\Delta E$  we expect  $\delta \gamma \to 0$  for  $n \to \infty$ , and the flight-path angle does not change along this arc.

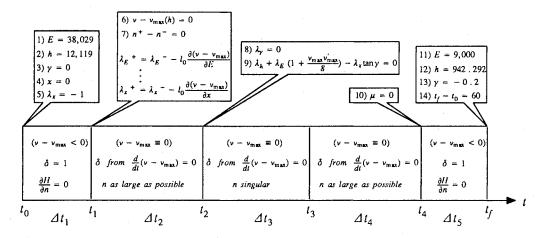
If the final flight-path angle is prescribed at a value different from the natural one, i.e., Eq. (49) is replaced by

$$\gamma(t_f) = \gamma_{\text{free}} + \Delta \gamma, \qquad \Delta \gamma \neq 0$$
 (51)

Then the dissipation of energy turns into a gradual process extending over a nonzero time interval, and the load factor remains finite. Paradoxically, nonzero  $\Delta \gamma$  results in a smaller peak value load-factor than does  $\Delta \gamma = 0$ .

# XI. Numerical Approach

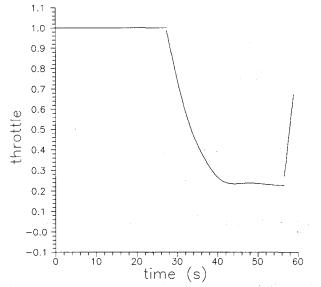
The switching structure, that is, the sequence of different control logics that actually solves a problem is not known in advance. For a given problem it has to be found by numerical experiments. Assuming a certain switching structure, the state and costate equations along with the boundary conditions, transversality conditions, and corner conditions implied by the assumed switching structure yield a multipoint boundary value problem (MPBVP). As an example case a schematic representation of the MPBVP associated with switching structure (S6) is given in Fig. 4. By inspection it is clear that the trajectory can be determined by simple forward integration if all parameters E(0), h(0),  $\gamma(0)$ , x(0),  $\lambda_E(0)$ ,  $\lambda_h(0)$ ,  $\lambda_{\gamma}(0)$ ,  $\lambda_x(0)$ ,  $l_0$ ,  $\Delta t_1$ ,  $\Delta t_2$ ,  $\Delta t_3$ ,  $\Delta t_4$ , and  $\Delta t_5$  are known. Basically, the numerical problem is to determine these 14 parameters such that all 14 conditions (numbered 1-14 in Fig. 4) are satisfied. This root finding problem is solved using routine ZSCNT of the IMSL subroutine library (version 9.2). In practice, forward integration causes the associated boundary value problem to be very badly conditioned. A remedy is to consider  $t_3$  as a new initial point and generate trajectories by successively integrating backward and forward, starting at switching time  $t_3$ , respectively. In an obvious way this generates a new set of parameters  $E(t_3)$ ,  $h(t_3)$ ,  $\gamma(t_3)$ ,  $\chi(t_3)$ ,  $\lambda_E(t_3)$ ,  $\lambda_h(t_3)$ ,  $\lambda_{\gamma}(t_3)$ ,  $\lambda_{\chi}(t_3)$ ,  $l_0$ ,  $\Delta t_1$ ,  $\Delta t_2$ ,  $\Delta t_3$ ,  $\Delta t_4$ , and  $\Delta t_5$  along with the conditions numbered 1-14 in Fig. 4. Noting that conditions 6, 8, and 10 in Fig. 4 can equivalently be enforced at time  $t_3$ , three unknowns, say  $E(t_3)$ ,  $\lambda_h(t_3)$ , and  $\lambda_{\gamma}(t_3)$  can be expressed in terms of the remaining 12



14 parameters:  $E(t_0)$ ,  $h(t_0)$ ,  $\gamma(t_0)$ ,  $x(t_0)$ ,  $\lambda_E(t_0)$ ,  $\lambda_h(t_0)$ ,  $\lambda_{\gamma}(t_0)$ ,  $\lambda_{x}(t_0)$ ,  $l_0$ ,  $\Delta t_1$ ,  $\Delta t_2$ ,  $\Delta t_3$ ,  $\Delta t_4$ ,  $\Delta t_5$ 

14 conditions: see above

Fig. 4 Schematic representation of the boundary value problem associated with switching structure (S6).





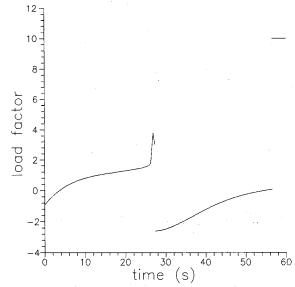


Fig. 6 Load factor n vs time t for  $n_{\text{max}} = 10$  [switching structure (S6)].

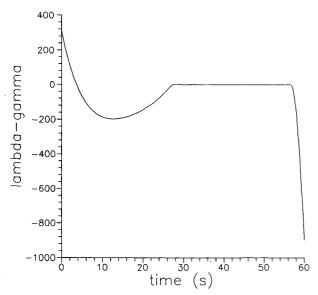


Fig. 7 Costate  $\lambda_{\gamma}$  vs time t for  $n_{\text{max}} = 10$  [switching structure (S6)].

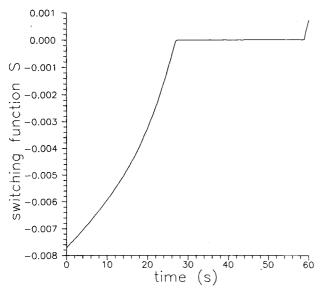


Fig. 8 Switching function S vs time t for  $n_{\text{max}} = 10$  [switching structure (S6)].

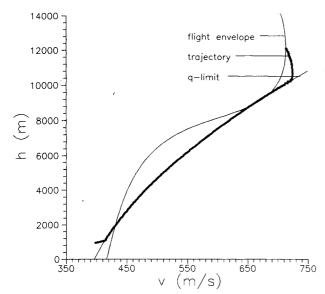


Fig. 9 Altitude-velocity chart for a typical extremal.

Table 4 Lengths of subarcs in seconds for selected trajectories associated with switching structures (S1)-(S6)

$n_{\max}$	$\Delta t_1$	$\Delta t_2$	$\Delta t_3$	$\Delta t_4$	$\Delta t_5$	$\Delta t_6$	$\Delta t_7$	$\Delta t_8$
free	27.622	0.525	30.614	0.796	0.443			
34	27.618	0.525	30.686	0.536	0.104	0.530		
23	27.491	0.521	30.539	0.451	0.015	0.515	0.567	
21	27.449	0.520	30.467	0.420	0.045	0.002	0.506	0.591
20	27,423	0.519	30.412	0.651	0.340	0.654		
10	26.856	0.504	29.231	2.31	1.098			

parameters  $h(t_3)$ ,  $\gamma(t_3)$ ,  $x(t_3)$ ,  $\lambda_E(t_3)$ ,  $\lambda_x(t_3)$ ,  $l_0$ ,  $\Delta t_1$ ,  $\Delta t_2$ ,  $\Delta t_3$ ,  $\Delta t_4$ , and  $\Delta t_5$ . Whereas it is only of minor importance that this reduces the number of parameters and conditions, it is very significant that this substitution ensures that 1) the characteristics of the singular arc, i.e., conditions 8 and 9 in Fig. 4, are satisfied along  $[t_2, t_3]$  and 2) the dynamic pressure limit  $\nu - \nu_{\text{max}}(h) = 0$  is satisfied along  $[t_1, t_4]$ . Note that both points hold true even before the root finding process converges.

#### XII. Numerical Results

As a general trend it is observed that all trajectories consist of mainly three phases. Phase 1 is the full thrust flight off the dynamic pressure limit (case 1) until dynamic pressure limit is reached. Phase 2 is the rapid descent with dynamic pressure limit active and singular control power (case 11b) until close to prescribed final altitude. Phase 3 is the rapid pitchup maneuver off the dynamic pressure limit with load factor on its upper limit and thrust first full, then zero.

In Table 4, the lengths of each arc in seconds are given for selected solutions with switching structures (S1)-(S6) [compare Eq. (46)]. For the case of  $n_{\text{max}} = 10$  (switching structure S6) time histories for throttle  $\eta$ , load factor n, Lagrange multiplier  $\lambda_{\gamma}$ , and switching function  $S = \lambda_{E} (v/W) + \mu (g/W)$  are given in Figs. 5-8, respectively. Figure 9 shows the altitude-velocity chart for this solution.

All switching structures found seem to be of some general nature in the sense that the same switching structures arise if initial or final coordinates of the trajectory are moderately changed. In this context trajectories starting at ground level with speed around takeoff velocity have been calculated for prescribed flight times over 200 s and final conditions as in Eq. (11). For long flight times [over 62 s for initial and final conditions as given in Eq. (11)] the obtained switching structures S1-S6 do not solve the problem. (Thrust oversaturates at the beginning of the singular thrust arc.) The correct switching structure for long flight times has not yet been found.

# XIII. Summary

Range optimal trajectories for an aircraft flying in the vertical plane are synthesized in the presence of a dynamic pressure limit (state inequality constraint) and a load factor limit (control inequality constraint). The trajectories start at the dash-point and end at prescribed final states inside the level-flight envelope, off the dynamic pressure limit. For a fixed flight time of 62 s the effect of the load factor limit on the resulting optimal switching structure is investigated. Six different switching structures are obtained with singular control along state constrained arcs always playing an important role. The switching structures identified in this paper appear to be of a general nature and are not restricted to the specific boundary conditions considered here. For long flight times greater than 62 s, however, and for the boundary conditions prescribed in this paper the optimal switching structure has not yet been found.

# **Appendix: Possible Control Logics**

Case 1: Constraint equation (10) is not active,  $\lambda_E < 0$ ,  $n_0 \in [-n_{\text{max}}, + n_{\text{max}}]$ :

$$n = n_0$$

$$\sigma_1 = 0$$

$$\sigma_2 = -\lambda_E (T - D + D_{\text{max}}) \frac{v}{W}$$

$$\sigma_3 = 0$$

$$\sigma_4 = 0$$

$$\mu = 0$$
(A1)

Case 2: Constraint equation (10) is not active,  $\lambda_E < 0$ ,  $n_0 < -n_{\text{max}}$ :

$$\delta = 1$$

$$n = -n_{\text{max}}$$

$$\sigma_1 = 0$$

$$\sigma_2 = -\lambda_E (T - D + D_{\text{max}}) \frac{v}{W}$$

$$\sigma_3 = \frac{\partial H}{\partial n}$$

$$\sigma_4 = 0$$

$$\mu = 0$$
(A2)

Case 3: Constraint equation (10) is not active,  $\lambda_E < 0$ ,  $n_0 > n_{\text{max}}$ :

$$\delta = 1$$

$$n = n_{\text{max}}$$

$$\sigma_1 = 0$$

$$\sigma_2 = -\lambda_E (T - D + D_{\text{max}}) \frac{v}{W}$$

$$\sigma_3 = 0$$

$$\sigma_4 = -\frac{\partial H}{\partial n}$$

$$\mu = 0$$
(A3)

Case 4: Constraint equation (10) is not active,  $\lambda_E > 0$ ,  $\lambda_{\gamma} > 0$ :

 $\delta = 0$ 

$$n = -n_{\text{max}}$$

$$\sigma_1 = \lambda_E (T - D + D_{\text{max}}) \frac{\nu}{W}$$

$$\sigma_2 = 0$$

$$\sigma_3 = \frac{\partial H}{\partial n}$$

$$\sigma_4 = 0$$

$$\mu = 0$$
(A5)

Case 5: Constraint equation (10) is not active,  $\lambda_E > 0$ ,  $\lambda_{\gamma} < 0$ :

$$\delta = 0$$

$$n = n_{\text{max}}$$

$$\sigma_1 = \lambda_E (T - D + D_{\text{max}}) \frac{v}{W}$$

$$\sigma_2 = 0$$

$$\sigma_3 = 0$$

$$\sigma_4 = -\frac{\partial H}{\partial n}$$

$$\mu = 0$$
(A6)

Case 6: Constraint equation (10) is not active,  $\lambda_E > 0$ ,  $\lambda_{\gamma} = 0$ .

In this case the controls are not determined uniquely by the minimum principle. Pointwise occurrence of this situation can be ignored. Assuming that  $\lambda_{\gamma} \equiv 0$  on some nonzero time interval yields control n after differentiating twice (singular control of first-order). Two cases have to be distinguished, namely  $\lambda_x = 0$  and  $\lambda_x \neq 0$ .

Case 6a: Constraint equation (10) is not active,  $\lambda_E > 0$ ,  $\lambda_{\gamma} = 0$ , and  $\lambda_{x} = 0$ .

Then necessarily  $\lambda_h \neq 0$  and

$$\delta = 0$$

$$\lambda_{\gamma} = 0$$

$$\cos \gamma = 0$$

$$n = 0$$

$$\sigma_{1} = \lambda_{E}(T - D + D_{\text{max}}) \frac{v}{W}$$

$$\sigma_{2} = 0$$

$$\sigma_{3} = 0$$

$$\sigma_{4} = 0$$

$$\mu = 0$$
(A9)

The generalized Legendre-Clebsch condition 12,13 implies

$$\sin \gamma > 0 \qquad \text{if} \qquad \lambda_h < 0$$
  
$$\sin \gamma < 0 \qquad \text{if} \qquad \lambda_h > 0$$

Case 6b: Constraint equation (10) is not active,  $\lambda_E > 0$ ,  $\lambda_{\gamma} = 0$ , and  $\lambda_{\chi} \neq 0$ :

$$\delta = 0$$

$$\lambda_{\gamma} = 0$$

$$\lambda_{h} - \lambda_{x} \tan \gamma = 0$$

$$n = n_{1}$$

$$\sigma_{1} = \lambda_{E} (T - D + D_{\text{max}}) \frac{v}{W}$$

$$\sigma_{2} = 0$$

$$\sigma_{3} = 0$$

$$\sigma_{4} = 0$$

$$\mu = 0$$
(A11)

The generalized Legendre-Clebsch condition implies  $\lambda_x < 0$ . The case  $\lambda_E \equiv 0$ ,  $\lambda_\gamma \equiv 0$  can be excluded. The case  $\lambda_E \equiv 0$ ,  $\lambda_\gamma \neq 0$  leads to first-order singular control in throttle  $\delta$ , which is rejected as nonoptimal by the generalized Legendre-Clebsch

condition. Case 7: Constraint equation (10) is active,  $\lambda_E + \mu(g/\nu) < 0$ ,  $\lambda_{\gamma} < 0$ :

$$\delta = 1$$

$$n = n_2$$

$$\sigma_1 = 0$$

$$\sigma_2 = -\left(\lambda_E + \mu \frac{g}{v}\right) (T - D + D_{\text{max}}) \frac{v}{W}$$

$$\sigma_3 = 0$$

$$\sigma_4 = 0$$

$$\mu = \frac{\lambda_\gamma}{n} \frac{q}{2KvW} - \lambda_E \frac{v}{g}$$
(A13)

Case 8: Constraint equation (10) is active,  $\lambda_E + \mu(g/\nu) < 0$ ,  $\lambda_{\gamma} > 0$ :

$$\delta = 1$$

$$n = -n_{2}$$

$$\sigma_{1} = 0$$

$$\sigma_{2} = -\left(\lambda_{E} + \mu \frac{g}{v}\right) (T - D + D_{\text{max}}) \frac{v}{W}$$

$$\sigma_{3} = 0$$

$$\sigma_{4} = 0$$

$$\mu = \frac{\lambda_{\gamma}}{n} \frac{q}{2KvW} - \lambda_{E} \frac{v}{g}$$
(A15)

The case where constraint equation (10) is active,  $\lambda_E + \mu(g/\nu) > 0$  and  $\lambda_{\gamma} \neq 0$  can be excluded. The case where constraint equation (10) is active,  $\lambda_E + \mu(g/\nu) = 0$  and  $\lambda_{\gamma} = 0$  is treated later.

Case 9: Constraint equation (10) is active,  $\lambda_E + \mu(g/\nu) = 0$ ,  $\lambda_\gamma < 0$ :

$$\delta = \delta_1 \text{ (from } C_1 = 0 \text{ with } n = n_{\text{max}})$$

$$n = n_{\text{max}}$$

$$\sigma_1 = 0$$

$$\sigma_2 = 0$$

$$\sigma_3 = 0$$
(A16)

$$\sigma_4 = -\lambda_\gamma \frac{g}{v}$$

$$\mu = -\lambda_E \frac{v}{g} \tag{A17}$$

Case 10: Constraint equation (10) is active,  $\lambda_E + \mu(g/\nu) = 0$ ,  $\lambda_{\gamma} > 0$ :

$$\delta = \delta_1 \text{ (from } C_1 = 0 \text{ with } n = n_{\text{max}})$$

$$n = -n_{\text{max}}$$

$$\sigma_1 = 0$$

$$\sigma_2 = 0$$

$$\sigma_3 = \lambda_\gamma \frac{g}{v}$$

$$\sigma_4 = 0$$

$$\mu = -\lambda_E \frac{v}{\sigma}$$
(A19)

Case 11: Constraint equation (10) is active,  $\lambda_E + \mu(g/\nu) = 0$ ,  $\lambda_{\gamma} = 0$ .

In this case the controls are not determined uniquely by the minimum principle. Pointwise occurrence of this situation can be ignored. Assuming that  $\lambda_{\gamma} \equiv 0$  on some nonzero time interval additional information has to be obtained from differentiating identity  $\lambda_{\gamma} \equiv 0$  (singular control). Two cases have to be distinguished, namely  $\lambda_{x} = 0$  and  $\lambda_{x} \neq 0$ .

Case 11a: Constraint equation (10) is active,  $\lambda_E + \mu(g/v) = 0$ ,  $\lambda_v = 0$ ,  $\lambda_v = 0$ .

In this case we have two possible control logics, namely

$$\delta = \frac{D_{\text{max}} + \frac{vW}{g} \sin \gamma \left( v_{\text{max}}' + \frac{g}{v} \right)}{T - D |_{n = \cos \gamma} + D_{\text{max}}}$$

$$\lambda_{\gamma} = 0$$

$$\sin \gamma = 0$$

$$n = \cos \gamma$$

$$\sigma_{1} = 0$$

$$\sigma_{2} = 0$$

$$\sigma_{3} = 0$$

$$\sigma_{4} = 0$$
(A20)

$$\mu = -\lambda_E \frac{v}{g} \tag{A21}$$

and

$$\delta = \frac{D_{\text{max}} + \frac{vW}{g} \sin \gamma \left(v_{\text{max}}' + \frac{g}{v}\right)}{T - D + D_{\text{max}}}$$

$$\lambda_{\gamma} = 0$$

$$\lambda_{h} + \lambda_{E} \left(1 + \frac{v_{\text{max}}'v}{g}\right) = 0$$

$$\sin \gamma = 0$$

$$n = \cos \gamma$$

$$\sigma_{1} = 0$$

$$\sigma_{2} = 0$$

$$\sigma_{3} = 0$$

$$\sigma_{4} = 0$$
(A22)

$$\mu = -\lambda_E \frac{v}{g} \tag{A23}$$

Here, Eqs. (A20) represent singular control of first-order (in control n). Equation (A22) represents a case of infinite-order singular control. In this case, control n is undetermined. Every control function of time n(t) is admissible, as long as it leads to state/costate time histories that satisfy all boundary, transversality, and switching conditions.

Case 11b: Constraint equation (10) is active,  $\lambda_E + \mu(g/\nu) = 0$ ,  $\lambda_x = 0$ , and  $\lambda_x \neq 0$ .

Then necessarily  $\cos \gamma \neq 0$  and

$$\delta = \frac{D_{\text{max}} + \frac{vW}{g} \sin \gamma \left( v'_{\text{max}} + \frac{g}{v} \right)}{T - D \mid_{n \text{ as below}} + D_{\text{max}}}$$

$$\lambda_{\gamma} = 0$$

$$\lambda_{h} - \lambda_{x} \tan \gamma + \lambda_{E} \left( 1 + \frac{v'_{\text{max}} v}{g} \right) = 0$$

$$n = \left( 1 - \frac{v'_{\text{max}} v}{g} \right) = \cos \gamma \qquad (A24)$$

$$\sigma_{1} = 0$$

$$\sigma_{2} = 0$$

$$\sigma_{3} = 0$$

$$\sigma_{4} = 0$$

$$\mu = -\lambda_{E} \frac{v}{g} \qquad (A25)$$

The generalized Legendre-Clebsch condition implies  $\lambda_x < 0$ . Case 12: Constraint equation (10) is active,  $\lambda_E + \mu(g/\nu) \neq 0$ ,  $\lambda_y = 0$ .

Then necessarily  $\cos \gamma \neq 0$  and

$$\delta = 0$$

$$\lambda_{\gamma} = 0$$

$$D_{\text{max}} \frac{g}{W} + \nu \sin \gamma \left(\nu_{\text{max}}' + \frac{g}{\nu}\right) = 0$$

$$n = n_3 \qquad (A26)$$

$$\sigma_1 = \left(\lambda_E + \mu \frac{g}{\nu}\right) (T - D + D_{\text{max}}) \frac{\nu}{W}$$

$$\sigma_{2} = 0$$

$$\sigma_{3} = 0$$

$$\sigma_{4} = 0$$

$$\mu = \frac{\lambda_{h} - \lambda_{x} \tan \gamma}{\nu'_{\text{max}} + \frac{g}{}}$$
(A27)

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