

HW 112 - Joseph Spectro

1) $U(x) = \frac{a}{x^2} + bx^2 \quad a, b \in \mathbb{R} > 0$

a) $\frac{dU}{dx} = -2ax^{-3} + 2bx$ want to find crit points, so $\frac{dU}{dx} = 0$

$0 = -2ax^{-3} + 2bx \Rightarrow 2bx = \frac{2a}{x^3} \Rightarrow bx^4 = a \Rightarrow x^4 = \frac{a}{b} \Rightarrow x = \pm \sqrt[4]{\frac{a}{b}}$

since we care about $x > 0$, disregard $-\sqrt[4]{\frac{a}{b}}$

$\frac{d^2U}{dx^2} = 6ax^{-4} + 2b \Rightarrow \frac{d^2U}{dx^2}(\sqrt[4]{\frac{a}{b}}) = 6a \frac{b}{a} + 2b = 6b + 2b = 8b > 0$

There is a min at $x = \sqrt[4]{\frac{a}{b}}$, second derivative test

b) Taylor expansion of $U(x) \approx U(x_0) + U'(x_0)(x-x_0) + \frac{U''(x_0)(x-x_0)^2}{2!} + \dots$

$U(x) \approx U(x_0) + U'(x_0)y + \frac{U''(x_0)}{2}y^2 + \dots$

$U(x) \approx U(\sqrt[4]{\frac{a}{b}}) + U'(\sqrt[4]{\frac{a}{b}})y + \frac{U''(\sqrt[4]{\frac{a}{b}})}{2}y^2 + \dots$

$\approx a\sqrt{\frac{b}{a}} + b\sqrt{\frac{a}{b}} + 0 + \frac{8b}{2}y^2 = a\sqrt{\frac{b}{a}} + b\sqrt{\frac{a}{b}} + 4by^2$

$k = 8b$

$C = a\sqrt{\frac{b}{a}} + b\sqrt{\frac{a}{b}}$

c) $\omega = \sqrt{k/m} = \sqrt{8b/m} \quad T = 2\pi / \sqrt{8b/m}$

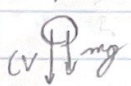
$T = 2\pi \cdot \sqrt{\frac{m}{8b}} = 2\pi \cdot \frac{1}{2} \sqrt{\frac{m}{2b}} = \pi \sqrt{\frac{m}{2b}}$

$T = \pi \sqrt{\frac{m}{2b}}$

k is the spring constant, so $\omega = \sqrt{k/m}$

HW #2 - Joseph Speltz

2)



$$F_{\text{drag}} = cV$$

$$mg > 0$$

$$K = \frac{c}{m}$$

$$F = cV - mg = m \frac{dv}{dt} \Rightarrow \frac{dv}{dt} = -Kv - g \Rightarrow \frac{dv}{-Kv - g} = dt$$

$$\Rightarrow \int_{v_0}^v \frac{dv}{-Kv - g} = \int_0^t dt \Rightarrow -\frac{1}{K} \ln \left(\frac{Kv + g}{Kv_0 + g} \right) = t + C \Rightarrow \ln \left(\frac{Kv + g}{Kv_0 + g} \right) = -Kt$$

$$\Rightarrow \frac{Kv + g}{Kv_0 + g} = e^{-Kt} \Rightarrow Kv + g = e^{-Kt} (Kv_0 + g) \Rightarrow \boxed{v = \frac{e^{-Kt} (Kv_0 + g) - g}{K}}$$

max height when $v = 0$, so

$$0 = \frac{e^{-Kt} (Kv_0 + g) - g}{K} \Rightarrow g = e^{-Kt} (Kv_0 + g) \Rightarrow \frac{g}{Kv_0 + g} = e^{-Kt}$$

$$\Rightarrow \ln \left(\frac{g}{Kv_0 + g} \right) = -Kt \Rightarrow t = -\frac{1}{K} \ln \left(\frac{g}{Kv_0 + g} \right)$$

now integrate velocity to find the position function

$$\frac{dy}{dt} = \frac{e^{-Kt} (Kv_0 + g) - g}{K} = \left(v_0 + \frac{g}{K} \right) e^{-Kt} - \frac{g}{K} \Rightarrow \int_0^y dy = \int_0^t \left(v_0 + \frac{g}{K} \right) e^{-Kt} dt - \int_0^t \frac{g}{K} dt$$

$$\Rightarrow y = \left(v_0 + \frac{g}{K} \right) \int_0^t e^{-Kt} dt - \frac{g}{K} t \Rightarrow y = \left(v_0 + \frac{g}{K} \right) \left(-\frac{1}{K} e^{-Kt} - 1 \right) - \frac{g}{K} t$$

$$\Rightarrow y(t=0) = \left(v_0 + \frac{g}{K} \right) \left(-\frac{1}{K} \left(\frac{g}{Kv_0 + g} \right) - 1 \right) + \frac{g}{K^2} \ln \left(\frac{g}{Kv_0 + g} \right)$$

$$= \left(v_0 + \frac{g}{K} \right) \left(\frac{-g}{K^2 v_0 + Kg} - 1 \right) + \frac{g}{K^2} \ln \left(\frac{g}{Kv_0 + g} \right) = \frac{-gv_0}{K^2 v_0 + Kg} - \frac{g^2}{K^3 v_0 + K^2 g} - \frac{v_0 + g}{K} + \frac{g}{K^2} \ln \left(\frac{g}{Kv_0 + g} \right)$$

$$\boxed{= \frac{g \ln \left(\frac{g}{Kv_0 + g} \right)}{K^2}}$$

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$$2b) \quad v(t=0) = \frac{-g v_0}{K^2 v_0 + K g} - \frac{g^2}{K^3 v_0 + K^2 g} - v_0 - \frac{g}{K} + g \ln\left(\frac{g}{(K v_0 + g)}\right) \frac{1}{K^2}$$

$$= \frac{-g K v_0}{K^3 v_0 + K^2 g} - \frac{g^2}{K^3 v_0 + K^2 g} - v_0 - \frac{g}{K} + \frac{g \ln(g/(K v_0 + g))}{K^2}$$

$$= \frac{-g^2 - g K v_0}{K^3 v_0 + K^2 g} - \frac{v_0 K^2}{K^2} - \frac{g K}{K^2} + \frac{g \ln(g/(K v_0 + g))}{K^2}$$

$$\boxed{v = \frac{-g^2 - g K v_0 + g \ln\left(\frac{g}{K v_0 + g}\right) - g K - v_0 K^2}{K^2}} \quad \text{where } K = c/m$$

$$c) \quad v = \frac{e^{-Kt}(K v_0 + g) - g}{K} \quad \text{exponentiated part } e^{-Kt}$$

$$e^{-x} = 1 - x + \frac{x^2}{2!} + \frac{x^3}{3!} + \dots \quad \text{around } x=0$$

$$e^{-Kt} = 1 - Kt + \frac{K^2 t^2}{2} + \frac{K^3 t^3}{6} + \dots = 1 - \frac{ct}{m} + \frac{c^2 t^2}{2m^2} + \frac{c^3 t^3}{6m^3} + \dots$$

if $c \ll 1$, the expression for e^{-Kt} simplifies to

$$1 - \frac{ct}{m} + \dots \approx 1 - \frac{ct}{m} \quad \text{plugging back into the expression for } v$$

$$v = \frac{(1 - Kt)(K v_0 + g) - g}{K} = \frac{K v_0 + g - K^2 v_0 t - K t g - g}{K}$$

$$\Rightarrow \frac{K v_0 - K^2 v_0 t - K t g}{K} = v_0 - \overset{\approx 0}{K v_0 t} - t g \approx v_0 - g t$$

$$\boxed{v \approx v_0 - g t}$$

This is what is expected because this is the velocity formula for a projectile thrown in the air

HW#2 - Joseph Speltz

B) $F(v, t) = -ktv^2$ $k > 0$ $\text{at } t=0, x=0 \text{ \& } v=v_0$

a) $F = -ktv^2 = m \frac{dv}{dt} \Rightarrow -\frac{k}{m} t dt = v^{-2} dv \Rightarrow \int_0^t -\frac{k}{m} t dt = \int_{v_0}^v v^{-2} dv$

$\Rightarrow -\frac{k}{2m} t^2 = -v^{-1} + v_0^{-1} \Rightarrow v^{-1} = v_0^{-1} + \frac{kt^2}{2m} \Rightarrow \frac{1}{v} = \frac{1}{v_0} + \frac{kt^2}{2m} \times \frac{v v_0 2m}{v v_0 2m}$

$\Rightarrow \frac{v v_0 2m}{v} = \frac{v v_0 2m}{v_0} + \frac{kt^2 v v_0 2m}{2m} \Rightarrow 2m v_0 = 2m v + kt^2 v v_0$

$\Rightarrow 2m v_0 = v(2m + kt^2 v_0) \Rightarrow \boxed{v = \frac{2m v_0}{2m + kt^2 v_0}}$

c) This shows the particle doesn't come to a stop because this is a continuous function with no discontinuities, so it goes on forever in the form of $\frac{1}{t^2}$, which asymptotes eventually

c) Integrate velocity to find x & then take limit

$\frac{dx}{dt} = \frac{2m v_0}{2m + kt^2 v_0} \Rightarrow dx = \frac{2m v_0}{2m + kt^2 v_0} dt \Rightarrow \int_0^{x_g} dx = 2m v_0 \int_0^t \frac{dt}{2m + kt^2 v_0}$

$\Rightarrow x_g = 2m v_0 \int_0^t \frac{1}{2m + \frac{kt^2 v_0}{2m}} = v_0 \int_0^t \frac{dt}{1 + \frac{kt^2 v_0}{2m}}$

let $u = \sqrt{\frac{kt v_0}{2m}} t$ $du = \sqrt{\frac{kt v_0}{2m}} dt$

$\Rightarrow dt = du \sqrt{\frac{2m}{kt v_0}}$ plugging back in $x_g = v_0 \int_0^{\sqrt{\frac{kt v_0}{2m}} t} \frac{1}{1 + u^2} du$

$\Rightarrow x_g = v_0 \sqrt{\frac{2m}{kt v_0}} \int_0^{\sqrt{\frac{kt v_0}{2m}} t} \frac{1}{1 + u^2} du$ derivative of $\tan(u)$ $\Rightarrow x_g = v_0 \sqrt{\frac{2m}{kt v_0}} \left(\tan(u) \right) \Big|_0^{\sqrt{\frac{kt v_0}{2m}} t}$

$x_g = v_0 \sqrt{\frac{2m}{kt v_0}} \left(\tan\left(\sqrt{\frac{kt v_0}{2m}} t\right) \right) \Big|_0^{\infty}$
 $\tan(\infty) = \frac{\pi}{2}$

$\boxed{x_g = v_0 \sqrt{\frac{2m}{kt v_0}} \cdot \frac{\pi}{2}}$

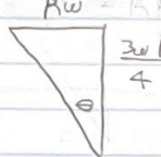
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a) $\omega = \text{ang vel @ } t=t_0$, $\phi = 90^\circ$, $r=R$, $\dot{r}(t=t_0) = \frac{3}{4}\omega R$, $\vec{F}_\theta = -k\vec{v}$

b) @ $t=t_0$ once the lead leaves, its position & \vec{F}_θ is only acting, so \vec{v} is constantly colinear w/ initial \vec{v}_F

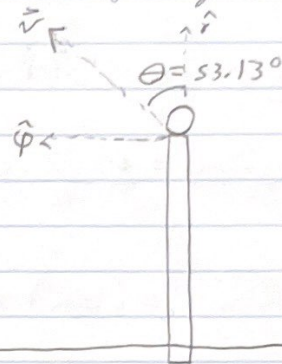
$$\vec{v}_F = \dot{r}\hat{r} + r\dot{\phi}\hat{\phi} = \frac{3}{4}\omega R\hat{r} + R\omega\hat{\phi}$$

$R\omega - k\vec{v}$



$$\tan\theta = \left(\frac{R\omega}{\frac{3}{4}\omega R}\right) = \left(\frac{4}{3}\right)$$

$$\theta = \arctan\left(\frac{4}{3}\right) = 53.13^\circ$$



The lead travels 53.13° counterclockwise off the \hat{r} vector @ $t=t_0$, trajectory stays on this line

$$\text{c) } \vec{F} = -k\vec{v} \Rightarrow m \frac{d\vec{v}}{dt} = -k\vec{v} \Rightarrow -\frac{m}{k} \frac{d\vec{v}}{\vec{v}} = dt \Rightarrow -\frac{m}{k} \int_{\vec{v}_0}^{\vec{v}} \frac{d\vec{v}}{\vec{v}} = \int_{t_0}^t dt$$

$$\Rightarrow -\frac{m}{k} (\ln(\vec{v}) - \ln(\vec{v}_0)) = t - t_0 \Rightarrow \ln\left(\frac{\vec{v}}{\vec{v}_0}\right) = -\frac{k}{m}(t - t_0)$$

$$\Rightarrow \frac{\vec{v}}{\vec{v}_0} = \exp\left(-\frac{k}{m}(t - t_0)\right) \Rightarrow \vec{v}(t) = \vec{v}_0 \left(\exp\left(-\frac{k}{m}(t - t_0)\right)\right)$$

$$\Rightarrow \vec{v}(t) = \exp\left(-\frac{k}{m}(t - t_0)\right) \left(\frac{3}{4}\omega R\hat{r} + R\omega\hat{\phi}\right)$$

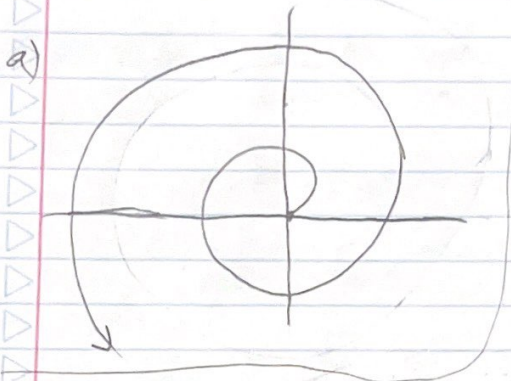
$$\vec{v}(t) = \left(\frac{3}{4}\omega R\hat{r} + R\omega\hat{\phi}\right) \exp\left(-\frac{k}{m}(t - t_0)\right)$$

* assuming \hat{r} & $\hat{\phi}$ are constant from point of lead leaving

This is never 0 because $\exp(A)$ can never be 0, it has an asymptote @ $t=\infty$ of 0, but it never reaches 0. A constant times an exponential can not be 0 ever.

HWK2 - cont - Joseph Specter

5) $r = b e^{kt}$



b) $r = b e^{kt}$

$\dot{r} = b k e^{kt}$

$\ddot{r} = b k^2 e^{kt}$

$\theta = ct$

$\dot{\theta} = c$

$\ddot{\theta} = 0$

$r(t) = r \hat{r} = b e^{kt} \hat{r}$

$\dot{r}(t) = \dot{r} \hat{r} + r \dot{\theta} \hat{\theta} = b k e^{kt} \hat{r} + b e^{kt} c \hat{\theta}$

$\ddot{r}(t) = (\ddot{r} - r \dot{\theta}^2) \hat{r} + (r \ddot{\theta} + 2 \dot{r} \dot{\theta}) \hat{\theta} = (b k^2 e^{kt} - b e^{kt} c^2) \hat{r} + (0 + 2 b k c e^{kt}) \hat{\theta}$

$r(t) = b e^{kt} \hat{r}$

$\dot{r}(t) = b k e^{kt} \hat{r} + b c e^{kt} \hat{\theta}$

$\ddot{r}(t) = (b k^2 e^{kt} - b c^2 e^{kt}) \hat{r} + 2 b k c e^{kt} \hat{\theta}$

c) $u \cdot v = |u| |v| \cos \theta$

$\dot{r} \cdot \ddot{r} = b^2 k^3 e^{3kt} - b^2 k c^2 e^{2kt} + 2 b^2 c^2 k e^{2kt}$
 $= b^2 k^3 e^{3kt} + b^2 c^2 k e^{2kt}$

$|\dot{r}| = \sqrt{b^2 k^2 e^{2kt} + b^2 c^2 e^{2kt}} = b e^{kt} \sqrt{c^2 + k^2}$

$|\ddot{r}| = \sqrt{b^2 k^4 e^{2kt} - 2 b^2 c^2 k^2 e^{2kt} + b^2 c^4 e^{2kt} + 4 b^2 c^2 k^2 e^{2kt}}$

$= \sqrt{b^2 k^4 e^{2kt} + b^2 c^4 e^{2kt} + 2 b^2 c^2 k^2 e^{2kt}}$

$= b e^{kt} \sqrt{k^4 + 2 k^2 c^2 + c^4} = b e^{kt} \sqrt{(c^2 + k^2)^2}$

$= b e^{kt} (c^2 + k^2)$

HW #2

5c) $\cos \theta = \frac{u \cdot v}{|u||v|}$

$$\cos \theta = \frac{b^2 K^3 e^{2\delta t} + b^2 c^2 K e^{2\delta t}}{b^2 e^{2\delta t} (b^2 + c^2) \sqrt{b^2 + c^2}}$$

$$= \frac{K^3 + c^2 K}{(b^2 + c^2)^{3/2}}$$

θ of w & \ddot{r} $\theta = \arccos \left(\frac{K^3 + c^2 K}{(b^2 + c^2)^{3/2}} \right)$

This expression is constant because it has no dependence on time & is only constants which never change

HW#2 - cont - Joseph Speduto

b) $\vec{R}(t) = b \sin \omega t \hat{i} + b \cos \omega t \hat{j} + ct^2 \hat{k}$ $b, c, \omega \in \mathbb{R}$

a) $r = \sqrt{x^2 + y^2} = \sqrt{b^2 \sin^2(\omega t) + b^2 \cos^2(\omega t)} = \sqrt{b^2 (\sin^2(\omega t) + \cos^2(\omega t))} = \sqrt{b^2} = b$

$\hat{r} = b \hat{i}$

$\Theta = \arctan\left(\frac{y}{x}\right) = \arctan\left(\frac{b \cos(\omega t)}{b \sin(\omega t)}\right) = \arctan(\cot(\omega t)) = \arctan(\tan(\frac{\pi}{2} - \omega t))$

$\Rightarrow \Theta = \frac{\pi}{2} - \omega t$

$R(t) = b \hat{r} + ct^2 \hat{k}$

$r(t) = b$

$\Theta(t) = \frac{\pi}{2} - \omega t$

$z(t) = ct^2$

b) $r = b$ $\Theta = \frac{\pi}{2} - \omega t$ $z = ct^2$

$\dot{r} = 0$

$\dot{\Theta} = -\omega$

$\dot{z} = 2ct$

$R = r \hat{r} + z \hat{k}$

$\dot{R} = \dot{r} \hat{r} + r \dot{\Theta} \hat{\Theta} + \dot{z} \hat{k} = r \dot{\Theta} \hat{\Theta} + \dot{z} \hat{k}$

$\ddot{r} = 0$

$\ddot{\Theta} = 0$

$\ddot{z} = 2c$

$\ddot{R} = (\ddot{r} - r \dot{\Theta}^2) \hat{r} + (r \ddot{\Theta} + 2 \dot{r} \dot{\Theta}) \hat{\Theta} + \ddot{z} \hat{k}$

$= -r \dot{\Theta}^2 \hat{r} + \ddot{z} \hat{k}$

$R(t) = b \hat{r} + ct^2 \hat{k}$

$\dot{R}(t) = -b\omega \hat{\Theta} + 2ct \hat{k}$

$\ddot{R}(t) = -b\omega^2 \hat{r} + 2c \hat{k}$

c) $|\ddot{R}| = \sqrt{b^2 \omega^4 + 4c^2}$ this is the same as $|\ddot{a}(t)| = \sqrt{b^2 \omega^4 + 4c^2}$ in cartesian words