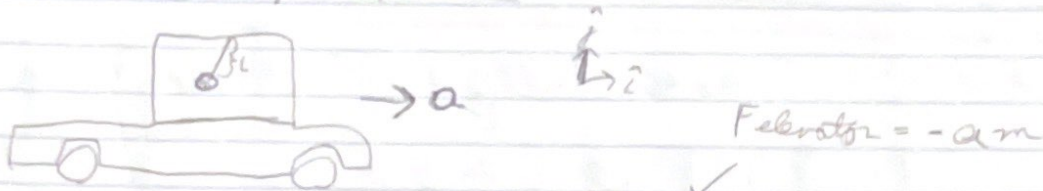


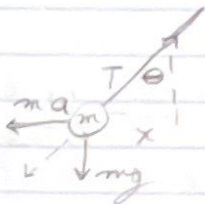
HW#9 - Joseph Spectro



a)



Real



Fictitious



$$\cos \theta = z/a$$

$$\sum F_{elast,x} = T \sin \theta - ma = 0 \Rightarrow T \sin \theta = ma \Rightarrow T = \frac{ma}{\sin \theta}$$

$$\sum F_{elast,y} = T \cos \theta - mg = 0 \Rightarrow T \cos \theta = mg \Rightarrow T = \frac{mg}{\cos \theta}$$

We can equate these 2 expressions for  $T$  & solve for  $\theta$

$$\frac{ma}{\sin \theta} = \frac{mg}{\cos \theta} \Rightarrow a = g \frac{\sin \theta}{\cos \theta} = g \tan \theta \Rightarrow \tan \theta = \frac{a}{g}$$

$$\therefore \theta = \arctan(a/g)$$

b) We know  $T = \frac{mg}{\cos \theta}$ , so if we know  $\theta$ , we can find  $T$

$$T = \frac{mg}{\cos \theta} \quad \text{w/ } \theta = \arctan(a/g)$$

## HW #9 - Joseph Spectro

c)



In this frame, we have  
a net  $\vec{a}$  to the bottom left,  
so we can call this the  
 $g_{eff}$  & w/ pendulums, we  
know the natural frequency  
is  $\sqrt{g/L}$

$$\therefore \omega_n = \sqrt{\frac{\sqrt{g^2 + a^2}}{L}} = \frac{(g^2 + a^2)^{1/4}}{L^{1/2}}$$

## HW#2 - Joseph Speed

2)



we have only Coriolis & Centrifugal  
 a/c  $\omega = 0$  & no relative translation  
 e/w frames.

$$a) \vec{F}_{cor} = -2m \vec{\omega} \times \vec{v} \quad \vec{F}_{cent} = -m \vec{\omega} \times (\vec{\omega} \times \vec{r})$$

$$\vec{\omega} = (0, 0, \omega) \quad \vec{v} = (v_r, 0, 0) \quad \vec{r} = (r, 0, 0)$$

$$\vec{F}_{cor} = -2m (0, \omega v_r, 0) = (0, -2m\omega v_r, 0)$$

$$\vec{F}_{cent} = -m \vec{\omega} \times (0, \omega r, 0) = -m (-\omega^2 r, 0, 0) = (m\omega^2 r, 0, 0)$$

$$\text{We know } |\vec{F}_a|^2 = |\vec{F}_{cent} + \vec{F}_{cor}|^2$$

$$|\mu m g|^2 = |(m\omega^2 r, -2m\omega v_r, 0)|^2$$

$$(\mu m g)^2 = (m\omega^2 r)^2 + (2m\omega v_r)^2 \Rightarrow \mu^2 m^2 g^2 = m^2 \omega^4 r^2 + 4m^2 \omega^2 v_r^2$$

$$\Rightarrow \mu^2 g^2 = \omega^4 r^2 + 4\omega^2 v_r^2 \Rightarrow \omega^4 r^2 = \mu^2 g^2 - 4\omega^2 v_r^2$$

$$r^2 = \frac{\mu^2 g^2}{\omega^4} - \frac{4v_r^2}{\omega^2} = \frac{1}{\omega^2} \left( \frac{\mu^2 g^2}{\omega^2} - 4v_r^2 \right)$$

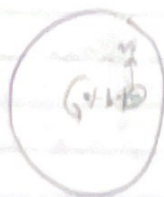
$$r = \sqrt{\frac{1}{\omega^2} \left( \left( \frac{\mu g}{\omega} \right)^2 - (2v_r)^2 \right)}$$

$$r_{max} = \frac{1}{\omega} \sqrt{\left( \frac{\mu g}{\omega} \right)^2 - (2v_r)^2}$$



# HW #9 - Joseph Speltz

a)



$$\vec{F}_{\text{cor}} = -2m\vec{\omega} \times \vec{v}$$

$$\vec{F}_{\text{cent}} = -m\vec{\omega} \times (\vec{\omega} \times \vec{r})$$

i) We know the speed is the speed in the reference frame plus the relative speed & this is equal to  $F_c$

$$F_c = \frac{mv(\vec{v}_r + \vec{\omega}b)^2}{b} = \mu mg$$

$$\Rightarrow \frac{(\vec{v}_r + \vec{\omega}b)^2}{b} = \mu g \Rightarrow b\mu g = (\vec{v}_r + \vec{\omega}b)^2 \Rightarrow \vec{v}_r + \vec{\omega}b = \sqrt{b\mu g}$$

$$\boxed{\vec{v}_r = \sqrt{b\mu g} - \vec{\omega}b}$$

ii) We know the following

$$\vec{r} = (b, 0, 0)$$

$$\vec{\omega} = (0, 0, \omega)$$

$$\vec{v} = (0, v_r, 0)$$

$$\vec{F}_{\text{cor}} = -2m(\omega v_r, 0, 0) = (-2m\omega v_r, 0, 0)$$

$$\vec{F}_{\text{cent}} = -m\omega \times (\omega \times \vec{r}) = -m(-\omega^2 b, 0, 0) = (m\omega^2 b, 0, 0)$$

We also know the drag has a centripetal force

$$F_c = F_{\text{cor}} + F_{\text{cent}} + F_{\text{centripetal}} \quad (\text{Equation net force})$$

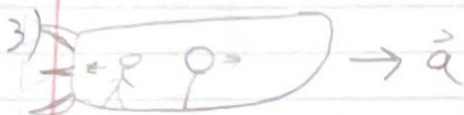
$$\mu mg = 2m\omega v_r + m\omega^2 b + m\frac{v_r^2}{b} \Rightarrow \mu g = 2\omega v_r + \omega^2 b + \frac{v_r^2}{b}$$

$$\Rightarrow b\mu g = 2\omega v_r b + \omega^2 b^2 + v_r^2 = (\vec{v}_r + \vec{\omega}b)^2$$

$$\Rightarrow \sqrt{b\mu g} = \vec{v}_r + \vec{\omega}b \Rightarrow \boxed{\vec{v}_r = \sqrt{b\mu g} - \vec{\omega}b}$$

## HW #2 - Joseph Spect

iii) They noted



a) If we consider the astronaut, air, & balloon are different densities w/  $\rho_{\text{air}} > \rho_{\text{air}} > \rho_{\text{balloon}}$ , we see that the bulk of the spaceship has a density of  $\rho_{\text{air}}$ . This volume of air will want to stay still as it has an inertia, but it is being pushed forward by the back wall. If we assume air is a relatively incompressible fluid, we can say the bulk of the volume accelerates the same. However, if something is more dense than air (humans), it will not be pushed as hard by the air behind it & thus experience less acceleration. However, the converse is true, if an object less dense than air (balloon) is pushed by air behind it, it will have less inertia than the air & accelerate forward more than the air.

b) Vacuums have 0 density & cannot provide support to anything w/ density, so

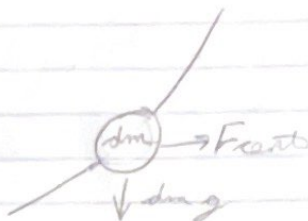
The balloon will move backward.

\*assuming an invincible balloon



# HW#9 - Joseph Specter

4)

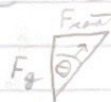


since we are in  
eq,  $v = \dot{u} = 0$ , so  
coriolis & azimuth  
are 0

$$\vec{F}_{cent} = -m \vec{\omega} \times (\vec{\omega} \times \vec{r}) = -m \vec{\omega} \times (0, \omega r, 0) = (m \omega^2 r, 0, 0)$$

$$\vec{F}_g = (0, 0, -mg)$$

if these are the only 2 forces, we can make  
a triangle



$$\text{We know } \tan \theta = \frac{(\omega^2 r \, dm)}{g \, dm}$$

& this is the slope, so we know @ this  
point these forces are equal as we  
are in eq, so

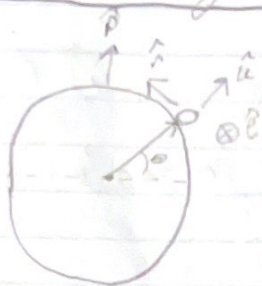
$$dm \omega^2 r \, dr = dm g \, dh \Rightarrow \omega^2 r \, dr = g \, dh$$

$$\omega^2 \int r \, dr = g \int dh \Rightarrow \frac{\omega^2 r^2}{2} = gh \Rightarrow h = \frac{\omega^2 r^2}{2g}$$

$$h = \frac{\omega^2 r^2}{2g}, \text{ \& this is equation of } V, \text{ so works}$$

# HW #9 - Joseph Specter

5)



$$\vec{r} = x\hat{e} + y\hat{n} + z\hat{u}$$

a)  $\hat{e} \perp \hat{p}$ , so no component in  $\hat{e}$  direction

if  $\theta = 0^\circ$ ,  $\hat{p} = \hat{n}$ , so most periodic function max @  $0^\circ$

if  $\theta = 90^\circ$ ,  $\hat{p} = \hat{u}$ , so most periodic function max @  $90^\circ$

$\hat{p} = \cos\theta\hat{n} + \sin\theta\hat{u}$ , so  $w\hat{p}$  is...

$$\vec{w} = w(\cos\theta\hat{n} + \sin\theta\hat{u})$$

b) Since we only have  $F_{cor}$  and  $F_g$ ,  $F_{net}$  is

$$F_{net} = F_{cor} + F_g$$

$$\hat{e} \quad \hat{n} \quad \hat{u}$$

$$F_g = (0, 0, -mg)$$

$$F_{cor} = -2m\vec{\omega} \times \vec{v}$$

$$F_{cor} = -2m\omega(\cos\theta\hat{n} + \sin\theta\hat{u}) \times (\dot{x}\hat{e} + \dot{y}\hat{n} + \dot{z}\hat{u})$$

$$\begin{vmatrix} \hat{e} & \hat{n} & \hat{u} \\ 0 & \cos\theta & \sin\theta \\ \dot{x} & \dot{y} & \dot{z} \end{vmatrix} = \hat{e}(\dot{y}\cos\theta - \dot{z}\sin\theta) - \hat{n}(-\dot{x}\sin\theta) + \hat{u}(-\dot{x}\cos\theta)$$

$$F_{cor} = -2m\omega(\dot{y}\cos\theta - \dot{z}\sin\theta, -\dot{x}\sin\theta, -\dot{x}\cos\theta)$$

# HW #9 - Joseph Specter

We know the zeroth order is just w/o rotation, so

$$\begin{array}{lll} \ddot{x}_0 = 0 & \ddot{y}_0 = 0 & \ddot{z}_0 = -g \\ \dot{x}_0 = 0 & \dot{y}_0 = 0 & \dot{z}_0 = -gt + A \\ x_0 = 0 & y_0 = 0 & z_0 = -\frac{1}{2}gt^2 + At + B \end{array} \quad \begin{array}{l} \text{This is the general} \\ \text{case w/o rotation} \end{array}$$

If we apply the conditions @  $t=0$  for  $\dot{y}$  &  $\dot{z}$ ...

$$\dot{y}(t=0) = A = v \quad \& \quad \dot{z}(t=0) = B = 0, \text{ so}$$

$$x_0 = 0 \quad y_0 = 0 \quad z_0 = vt - \frac{1}{2}gt^2$$

The solution for  $x_1, y_1, z_1$  is dependent on velocities, but the only non zero  $v$  or  $\omega$  is

$$\dot{z}_0 = v - gt \quad \dot{x}_0 = 0 = \dot{y}_0$$

$$\ddot{x}_1 = -2m\omega(v - gt)\cos\theta = -2m\omega(v - gt)\frac{1}{\sqrt{2}}, \quad \frac{1}{\sqrt{2}} = \frac{\sqrt{2}}{2}$$

$$= -\sqrt{2}m\omega(v - gt) = m\omega\sqrt{2}(gt - v)$$

$$\dot{x}_1 = \sqrt{2}m\omega\left(\frac{1}{2}gt^2 - vt + C\right) = 0 @ t=0, \text{ so } C=0$$

$$x_1 = \sqrt{2}m\omega\left(\frac{1}{6}gt^3 - \frac{1}{2}vt^2 + D\right) = 0 @ t=0, \text{ so } D=0$$

$$\text{we know } \ddot{y}_1(t=0) = \ddot{y}_1(t=0) \quad \& \quad \dot{y}_1(t=0) = \dot{y}_1(t=0) = 0$$

$$\therefore x(t) = \sqrt{2}\omega m\left(\frac{1}{6}gt^3 - \frac{1}{2}vt^2\right)$$

$$y(t) = 0$$

$$z(t) = vt - \frac{1}{2}gt^2$$