

# HW #4

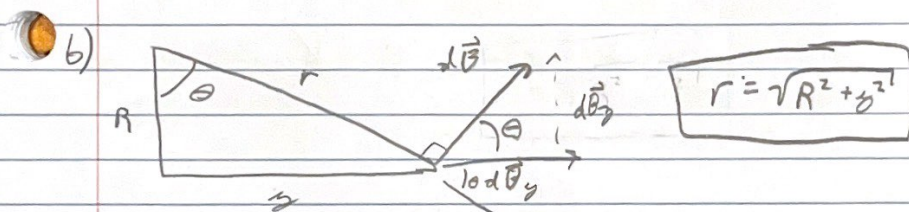
$$1) \vec{d\vec{B}} = \frac{\mu_0 d \vec{\ell} \times \vec{r}}{4\pi r^3}$$

a) find  $\vec{B}$  @ center of loop, Dist. same @  $r=R$ , find  $\vec{B}$

$$\vec{d\vec{B}} = \frac{\mu_0 d \vec{\ell} \times \vec{r}}{4\pi r^3} \quad \vec{r} \rightarrow -\vec{r} \quad \& \quad d\vec{\ell} = R d\phi \hat{\phi}$$

$$\vec{d\vec{B}} = \frac{\mu_0 d}{4\pi R^2} \begin{vmatrix} \hat{r} & \hat{\phi} & \hat{z} \\ 0 & R d\phi & 0 \\ -1 & 0 & 0 \end{vmatrix} = \frac{\mu_0 d}{4\pi R^2} (R d\phi) \hat{z} = \frac{\mu_0 d}{4\pi R} d\phi \hat{z}$$

$$\vec{B} = \int \vec{d\vec{B}} = \int \frac{\mu_0 d}{4\pi R} d\phi \hat{z} = \frac{\mu_0 d}{4\pi R} \hat{z} \int_0^{2\pi} d\phi = \frac{\mu_0 d}{4\pi R} (2\pi) \hat{z} = \boxed{\frac{\mu_0 d}{2R} \hat{z}} \quad \text{on axis}$$



$$c) \sin \theta = \frac{z}{r} \quad \sin \theta = \frac{d\vec{B}_z}{d\vec{B}} \Rightarrow d\vec{B} \sin \theta = d\vec{B}_z$$

$$\cos \theta = \frac{R}{r} \quad \cos \theta = \frac{d\vec{B}_y}{d\vec{B}} \Rightarrow d\vec{B} \cos \theta = d\vec{B}_y$$

$$d) \sin \theta = -\frac{d\vec{B}_z}{d\vec{B}'} \Rightarrow -d\vec{B}' \sin \theta = d\vec{B}_z = -d\vec{B}_z'$$

$$\cos \theta = \frac{d\vec{B}_y}{d\vec{B}'} \Rightarrow d\vec{B}' \cos \theta = d\vec{B}_y = d\vec{B}_y'$$

The  $\hat{z}$  components of  $d\vec{B}$  &  $d\vec{B}'$  will cancel out

$\therefore$  particle will only feel force in  $\hat{y}$  (on-axis)



HW #4 - cont

e)  $\cos \theta = \frac{A}{r} = \frac{R}{\sqrt{R^2 + z^2}}$

f)  $d\vec{B} = \frac{\mu_0 dI}{4\pi r^2} d\vec{\ell} \times \hat{r} = \frac{\mu_0 dI}{4\pi r^2} \begin{vmatrix} \hat{r} & \hat{\phi} & \hat{z} \\ 0 & R d\phi & 0 \\ -\cos\theta & 0 & 0 \end{vmatrix} = \frac{\mu_0 dI}{4\pi r^2} d\phi R \cos\theta \hat{z}$

$\Rightarrow \frac{\mu_0 dI R}{4\pi r^2} \frac{R \hat{z}}{\sqrt{R^2 + z^2}} \int_0^{2\pi} d\phi = \frac{\mu_0 dI R^2}{4\pi} \frac{1}{(R^2 + z^2)^{3/2}} \cdot 2\pi \hat{z}$

$\Rightarrow \vec{B} = \frac{\mu_0 dI R^2}{2(R^2 + z^2)^{3/2}} \hat{z}$

# HW #4

2) uniform  $\vec{B} = B_0 \hat{z}$

mirror  $\vec{B} = B_0 \left[ 1 + \left( \frac{z}{a_0} \right)^2 \right] \hat{z}$

a) mass  $m$ , charge  $q$ , find  $v_{||}(z)$  for  $I \ll I_c$ , assume particle mirrors @  $z = \pm z_t$

II) mirror ratio

$$B_{max} = B(z_t) = B_0 \left[ 1 + \left( \frac{z_t}{a_0} \right)^2 \right]$$

$$B_{min} = B(0) = B_0 [1 + 0] = B_0$$

$$\Rightarrow R_m = \frac{B_0 [1 + (z_t/a_0)^2]}{B_0} = 1 + (z_t/a_0)^2$$

$$I) \quad m \dot{\vec{r}} = q \vec{v} \times \vec{B} = q \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ v_x & v_y & v_z \\ 0 & 0 & B_0 \end{vmatrix} = q B_0 (v_y \hat{x} - v_x \hat{y})$$

$$m \dot{v}_z = 0 \Rightarrow \boxed{v_z = v_{z0} = v_{||,0}}$$

$$II) \quad \frac{1}{2} m v_{\perp}^2 + \frac{1}{2} m v_{||}^2 = C_0 \Rightarrow v_{\perp}^2 + v_{||}^2 = C$$

$$\Rightarrow v_{||} = \sqrt{C - v_{\perp}^2}, \quad (v_{||}(\pm z_t) = 0 = \sqrt{C - v_{\perp}^2} \Rightarrow C = v_{\perp}^2(\pm z_t))$$

reflects  $\uparrow$

$$v_{||} = \sqrt{v_{\perp}^2(\pm z_t) - v_{\perp}^2(z)}$$

$$\mu = \frac{1}{2} m \frac{v_{\perp}^2}{B} \Rightarrow v_{\perp}^2 = \frac{2\mu B}{m} \quad v_{\perp}^2(z) = \frac{2\mu}{m} B(z)$$

$$\Rightarrow v_{\perp}^2(z) = \frac{2\mu}{m} \left[ B_0 + B_0 \left( \frac{z}{a_0} \right)^2 \right], \quad v_{\perp}^2(\pm z_t) = \frac{2\mu}{m} \left[ B_0 + B_0 \left( \frac{\pm z_t}{a_0} \right)^2 \right]$$

$$\Rightarrow v_{||} = \sqrt{\frac{2\mu B_0}{m} \left[ 1 + \left( \frac{\pm z_t}{a_0} \right)^2 \right] - \frac{2\mu B_0}{m} \left[ 1 + \left( \frac{z}{a_0} \right)^2 \right]}$$



HW#4

$$\Rightarrow v_{||} = \sqrt{\frac{2\mu B_0}{ma_0^2} (z_t - z)}$$

b) show  $F_{||} \propto \frac{\partial B}{\partial z} \hat{z}$

$$F_{||} = m \frac{dv_{||}}{dt}, \quad v_{||} = \frac{dz}{dt} \Rightarrow F_{||} = m \frac{dv_{||}}{dt} = m \frac{dv_{||}}{dz} \left[ \frac{dz}{dt} \right]$$

$$F_{||} = m v_{||} \frac{dv_{||}}{dz} = \frac{m}{2} \frac{d(v_{||}^2)}{dz} = \frac{m}{2} \frac{d}{dz} \left[ \frac{2\mu B_0}{ma_0^2} (z_t^2 - z^2) \right]$$

$$\Rightarrow F_{||} = m \left( \frac{2\mu B_0}{ma_0^2} (-2z) \right) = \frac{-2\mu B_0}{a_0^2} z = F_{||}$$

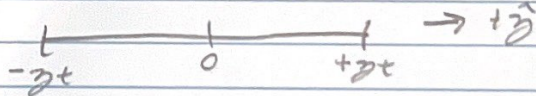
$$\frac{\partial v}{\partial z} = \frac{B_0}{a_0^2} 2z$$

$$F_{||} \neq \mu \quad \therefore F_{||}(z) \propto \frac{\partial B}{\partial z}$$

c) I)  $F_{||}(z) = 0$  as  $v_z \parallel B_z$

The force does nothing to the particle

$$\text{II) } F_{||}(z) = \frac{-2\mu B_0}{a_0^2} z$$



as  $z \rightarrow +z_t$ , the force gets stronger in  $-\hat{z}$  forcing particle back to center

as  $z \rightarrow -z_t$ , the force gets stronger in  $+\hat{z}$  forcing particle back to center



# HW# 4

$$2d) F_{||} = m \frac{dv_{||}}{dt} = \frac{-2\mu B_0 z}{a_0^3} = m \ddot{z}$$

$$\Rightarrow \boxed{\frac{d^2 z}{dt^2} + \underbrace{\left( \frac{2\mu B_0}{m a_0^3} \right)}_{\omega^2} z = 0} \quad \omega / \quad \omega^2 = \frac{2\mu B_0}{m a_0^3}$$

Period

$$\frac{2\pi}{T} = \omega = \sqrt{\frac{2\mu B_0}{m a_0^3}}$$

$\Rightarrow$

$$\boxed{\therefore T = 2\pi \sqrt{\frac{m a_0^3}{2\mu B_0}}}$$

#### HW#4

3) find trajectory w/  $B(90,1)$  &  $E=(E_0 \sin(\omega t), 0, 0)$

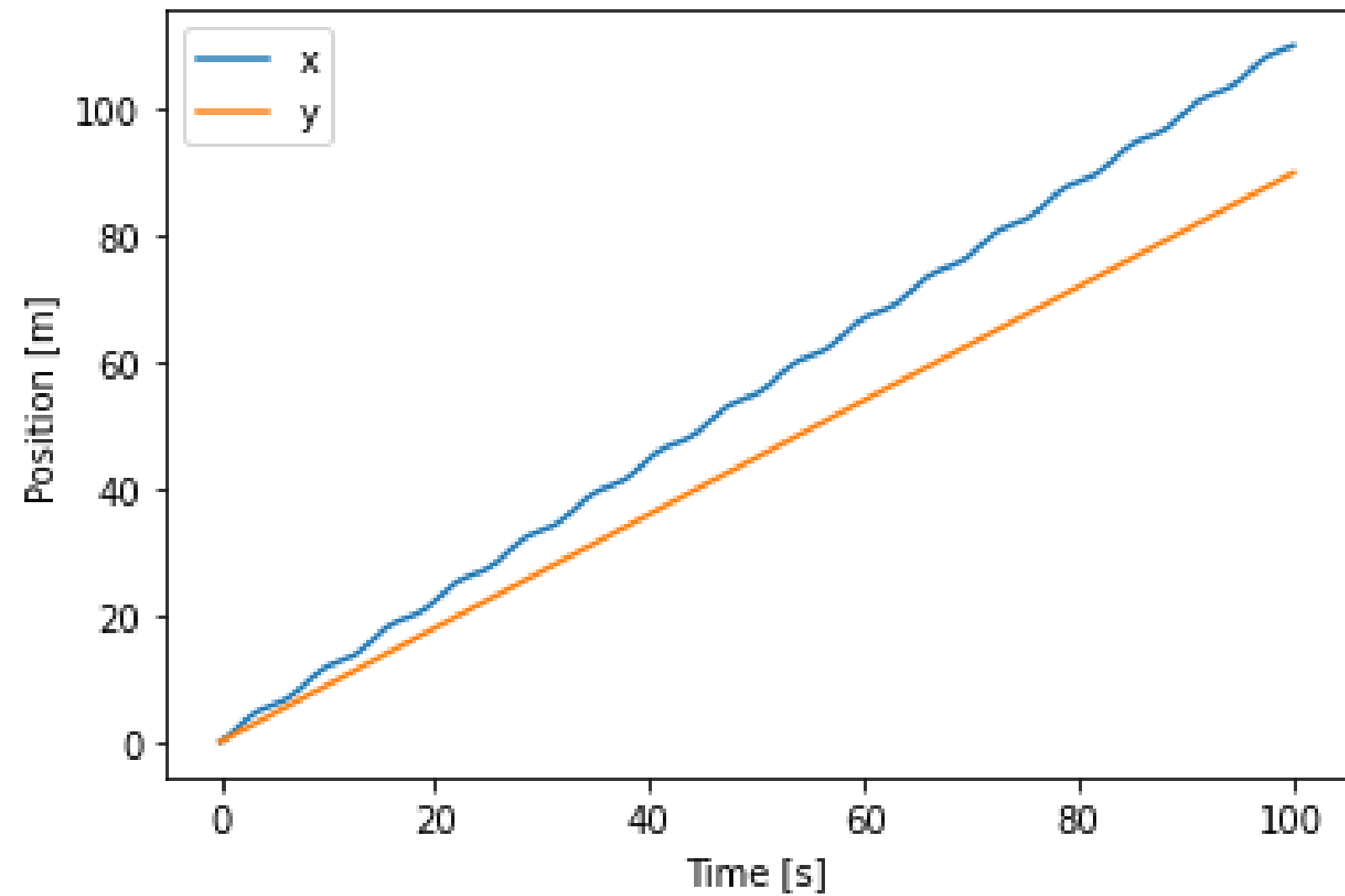
$$\omega_c = \frac{qB_0}{m} = 1 \text{ in code}$$

if  $\omega \ll \omega_c$ , the perturbations happen so quickly that it looks like the trajectory is constant like for constant  $E$  field.

if  $\omega \sim \omega_c$ , the local perturbations are much more noticeable & the  $x$  &  $y$  positions move about a guiding center w/ sinusoidal perturbations about the trajectory of constant  $E$

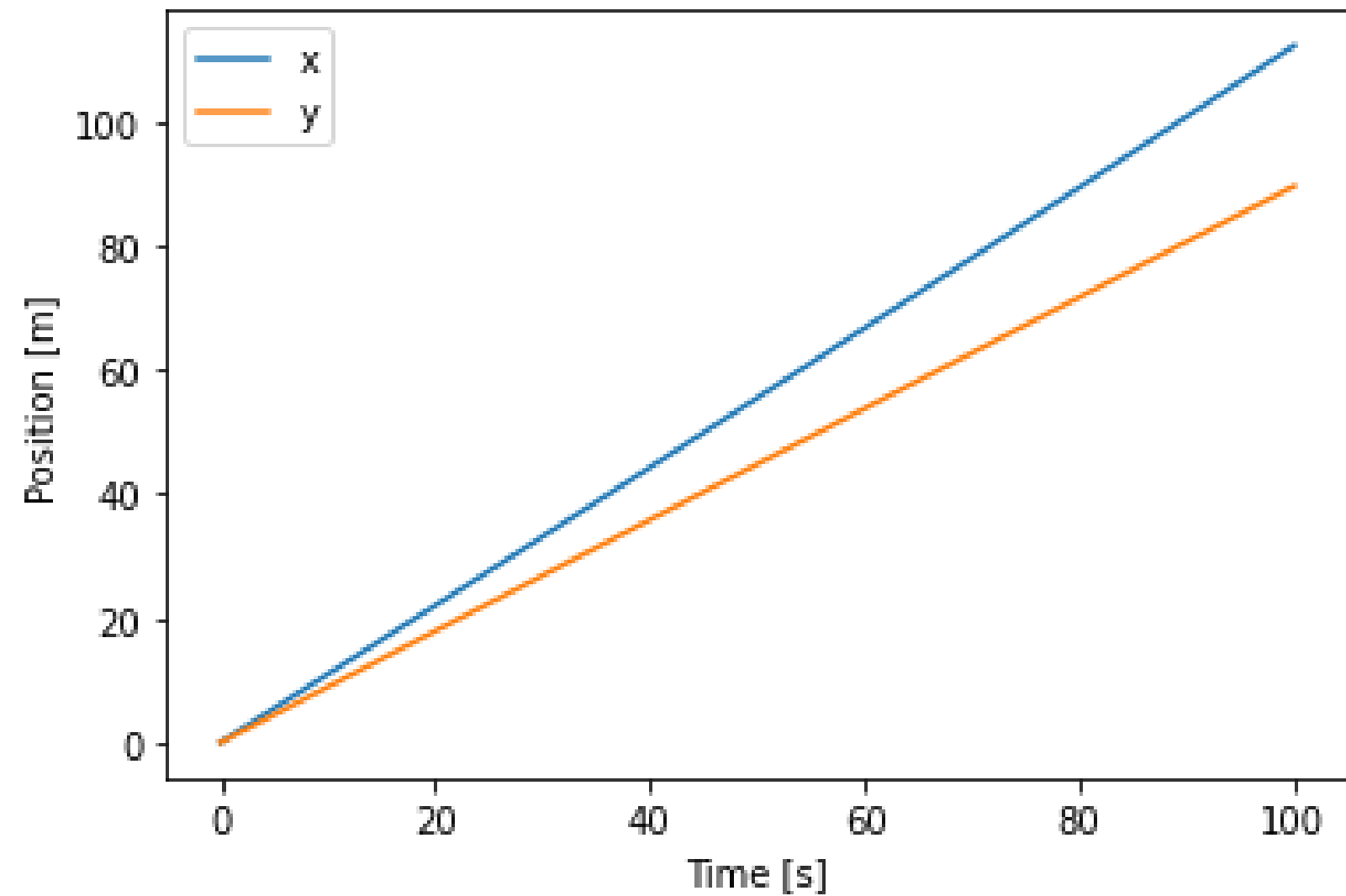
For  $\omega \ll \omega_c$  &  $\omega \sim \omega_c$ , the mean behavior is roughly the same. Although the trajectory for  $y$  is the same in all cases, the sinusoidally oscillating  $E$  fields have a lower final  $x$  position than a constant  $E$  field. This makes sense as the varying  $E$  will exert less force when averaged over a period of oscillation than a constant field over the same period.

4.3,  $w=1$





4.3,  $w=0.001$





### 4.3, Constant E

