Inertial Navigation using Atom Interferometry

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A dissertation submitted for ...

Abstract

This thesis describes work I did during my PhD...

Declaration

This dissertation is the result of my own work...

Jimmy Stammers



Acknowledgements

Some people worth thanking...



Preface

This thesis describes my research on various aspects of...

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Chapter 1.

Introduction

- atom interferometry experiments for precision measurements of inertial forces
- inertial navigation suffers from long-term bias drift
- recent experiments have demonstrated measuring acceleration in environments of interest to navigation

Chapter 2.

Theory

- Describe general principles of light-matter interaction
- Specific cases for laser cooling (doppler/sub-doppler) and Raman transitions
- Lead into atom interferometry
- Perhaps split this into two shorter chapters

4 Theory

- 2.1. Overview
- 2.2. Light-Matter Interactions
- 2.3. Laser Cooling of Rubidium-87
- 2.4. Raman Transitions in Rubidium-87
- 2.5. Light Pulse Atom Interferometry

Chapter 3.

MOTMaster

3.1. Chapter Overview

The aim of this chapter is to provide a description of the MOTMaster software, which was developed from a pre-existing version during my PhD. The design of MOTMaster assumes very little about the particular experiment it is being used for, so much of the discussion in this chapter will be kept general. This chapter begins by motivating the need to extend MOTMaster by developing a graphical interface to simplify the creation of experimental sequences, as well as implementing new methods of controlling hardware. This is followed by a description of the National Instruments hardware used by MOTMaster to configure the inputs and outputs to an experiment.

- New MOTMaster software to control experiment
- Control system diagram
- Interfacing with muquans/M Squared lasers
- Acquiring data from experiment axelsuite

3.2. Motivation

In the initial stages of my PhD, I decided to use Cicero Word Generator [1] to control the hardware for the experiment. This is a graphical-based control system developed by Wolfgang Ketterle's group at MIT, which was designed for controlling atomic physics experiments using National Instruments hardware. Over time, as the experiment became more complex, it started to become apparent that Cicero was not suited to meet all of our requirements for control software. This was most evident in the control of the M-Squared Raman laser system. Unlike the Muquans laser system, which can be controlled externally using analogue and digital voltages and serial messages, the M-Squared system is externally controlled by communicating JSON messages to a web server. Implementing such a drastically different scheme for controlling a specific component into Cicero was not deemed worthwhile. Around this time, I also realised that Cicero takes an appreciable amount of time (around 300ms) to re-calculate the experiment sequence between each shot. Since the design of Cicero was aimed at controlling experiments that take many seconds per cycle, this dead time between each cycle is not significant on those time scales. In contrast, each cycle of this experiment takes around 250ms. This unnecessary dead time needed to be addressed if we hoped to improve the repetition rate.

After it became clear that a potentially large amount of work would be needed to improve Cicero, I decided that it was worth moving to a new control system. Within Centre for Cold Matter (CCM), there is a collection of programs that are used to control a range of experiments. One of these, MOTMaster, was designed to control and acquire data from experiments investigating cold atoms trapped in a Magneto-optical Trap (MOT), as its name suggests.

Write more to motivate the switch to MOT-Master

3.3. Interfacing wth Hardware

The majority of the experimental hardware is controlled using analogue and digital voltages that are generated by Data Acquisition (DAQ) cards manufactured by National Instruments. MOTMaster is compatible with cards that use either the NI-DAQmx or NI-HSDIO device drivers. These are used to configure the generation or acquisition of digital or analogue voltage waveforms and are capable of precisely timing and synchronising their I/O across multiple devices. Most components in the experiment rely on this precise timing to function correctly. Other devices, where timing accuracy is less critical, are controlled by sending or receiving data using serial communication. This has the advantage of allowing more structured command beyond analogue or digital voltages, but the communication speed of the serial channel limits the accuracy of the execution time. An understanding of the low-level interface between control software and the experiment is very useful in both carrying out experiments and accurately interpreting the results.

3.3.1. Hardware Abstraction

When designing software, it is often useful to structure a program in such a way that modules which make use of other components do not need to know about their specific implementation in order to use them. In doing this, the submodule can be modified without harming the compatibility of these two components. In the context of experimental hardware, this is equivalent to requiring that changing specific components, for example the Voltage-Controlled Oscillator (VCO) that generates the RF power for an Acousto-optic Modulator (AOM), will not stop the experiment from working. This is done using abstract representations of the hardware, in the form of input and output channels that are used to communicate to each device.

3.3.2. Voltage Pattern Generation

All the analogue outputs controlled using MOTMaster are done using the NI-DAQmx software. Each output uses a Digital-to-Analogue Converter (DAC) to convert a floating-point number into an analogue voltage. To generate a sequence of voltages across multiple channels, the NI-DAQmx driver allocates a block of memory on the DAQ card for each output channel. This memory acts as a first-in first-out (FIFO) buffer for data streamed to it from a computer. The output of each channel is synchronised to a clock signal, so that every time a rising edge occurs on the clock, the voltage at each output transitions to the value corresponding to the next value in its corresponding buffer. Channels across multiple DAQ cards can be synchronised by sharing a clock signal, which can be done using the bus that connects cards in a PXIe chassis. Additional cards can also be configured to trigger the start of their output at the moment they receive the first clock pulse, rather than waiting for a command from the computer to start.

Digital outputs from NI-DAQmx cards are generated in much the same way as analogue voltages, except for the fact that they only take two values corresponding to either a low (0 V) or high (3.3/5 V) level. DAQ cards which operate using the NI-HSDIO driver function differently. These cards have faster on-board clocks than is usually available with NI-DAQmx hardware. For instance, NI-DAQmx PXI-6723 can operate with a maximum clock frequency of 200 kHz and typical sequence durations mean that a sample rate of 100 kHz is needed to fit the entire sequence into memory. However, the NI-HSDIO PXI-6541 card can generate digital voltages at sample rates up to 50 MHz and requires less memory to store a pattern. Rather than write the pattern as an array of values for each sample, the sequence is segmented into smaller patterns during which the state of each channel is constant, as illustrated in Figure 3.1.

Figure 3.1.: Scripted pattern generation for an NI-HSDIO digital output card. A pattern is split into segments which correspond to a duration for which all the channels output a constant value. Each of these smaller waveforms are written to the on-board memory, along with a script that instructs the card to output each pattern for the required number of times to reconstruct the original sequence. By reducing the amount of memory required to define the sequence, a faster clock frequency and hence timing resolution can be used to output digital control signals.

NI-HSDIO cards can be scripted to generate each of these patterns for the appropriate number of clock cycles.

3.3.3. Timed Serial Communication

Serial communication is used to control devices which require more complex control than is possible using analogue or digital voltages. This increase in complexity comes at the cost of slower response times, because it takes longer to communicate an array of bytes is longer than to change the voltage across an output terminal. Using the NI-VISA driver, the output of serial data can only be timed using a software clock on a computer, which is more prone to jitter than a hardware clock. One way to improve the synchronisation between serial data and hardware timed outputs is to use extra hardware to trigger the transmission of serial data. If the trigger is timed using the same clock as other outputs and the transmission delay is accounted for, then serial data can be output more synchronously. The scheme for timing serial messages is shown in Figure 3.2. Serial messages are stored as strings on the computer and a counter channel is configured so that every time it detects a rising edge, the computer outputs the next message. This counter is connected to a digital output channel, so that it acts as a trigger for the serial data output. Using this method, multiple serial messages can be sent to one device during a sequence even for devices which have no means of storing commands.

Figure 3.2.: Timing diagram for serial communication. A counter channel is configured to count edges from a digital output channel. Every time it sees a rising edge, it triggers the output of the next message on each serial channel from the computer. Multiple messages can be communicated during a single sequence without the need for software timing.

3.3.4. Voltage Acquisition

Analogue input channels are configured in a similar way to analogue output channels. A block of memory is allocated on the DAQ card for each input channel. Once the card is triggered to start acquiring, an Analogue-to-Digital Converter (ADC) converts the voltage across the input into a digital value at every rising edge of the clock signal. Once the sequence has finished, or the buffer has been filled, the card streams this data to the computer.

3.4. External Control

3.5. Building a Sequence

3.6. Experimental Sequence Structure

Chapter 4.

Cooling and Trapping in a MOT

- 4.1. Chapter Outline
- 4.2. The Navigator Vacuum Chamber
- 4.2.1. The 2D MOT system
- 4.2.2. The 3D MOT system
- 4.2.3. CCD Imaging
- 4.3. Generating MOT light
- 4.3.1. Muquans Laser Control

Frequency Control

Real-Time Communication

- 4.4. Controlling the MOTs
- 4.4.1. Optical Fibre Network

Chapter 5.

Preparing Atoms for Interferometry

This chapter presents the work that went towards the initial stages of the experimental sequence, where the main objective is to prepare a sufficiently cold ensemble of Rubidium-87 (⁸⁷Rb) in the same quantum state.

5.1. Chapter Outline

To-Do:

- Discuss loading atoms in 3D MOT from 2D
- Characterisation of the moving frame optical molasses
- Various schemes for preparing atoms into $|1,0\rangle$. *m*ention velocity selection here or in next chapter?

- **5.2.** Cooling in Optical Molasses
- 5.2.1. Real-time Frequency Control
- 5.2.2. Optimising the Temperature
- 5.3. State Preparation
- **5.3.1.** Schemes for Preparation
- 5.3.2. Optical Pumping Scheme
- **5.3.3. Including Microwave Transitions**

Wind-Freak Synthesiser

Chapter 6.

Acceleration-Sensitive Interference

This chapter describes the work towards realising an atom interferometer and subsequently measuring accelerations.

6.1. Chapter Outline

To-Do:

- Raman spectrum, identifying each transition
- Characterisation of velocity-selective pulse and each interferometer pulse using Rabi oscillations.
- Making a three-pulse atom interferometer
- Improving acceleration sensitivity and correlating vibrations using MEMS

6.2. Raman Optical Syster	Raman O	cal S	ysten
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- 6.2.1. Raman Beam Collimator
- 6.2.2. Retro-reflection Assembly
- 6.2.3. The MEMS Accelerometer
- 6.3. Driving Raman Transitions
- 6.3.1. Frequency and Phase Control
- 6.4. Atom Detection
- 6.4.1. Optical System
- 6.4.2. Measuring the Interferometer Phase
- 6.5. Individual Pulse Characterisation
- 6.5.1. Velocity-Selective Pulse
- **6.5.2.** Interferometer Pulses
- 6.6. Three-Pulse Atom Interference
- 6.7. Measuring Accelerations

Chapter 7.

Outlook

This final chapter describes some of the next steps and further work

7.1. Combining with classical accelerometers

- Discuss schemes for combining multiple sensors Kalman filtering
- Extend this to inertial navigation
- Steps towards overcoming sensitivity-bandwidth trade-off.

7.2. Extending to senstivity along three axes

- New chamber design
- Improvements to MSquared laser
- Required knowledge of gravitational axis for accurate navigation

Bibliography

- [1] Aviv Keshet and Wolfgang Ketterle. "A Distributed GUI-based Computer Control System for Atomic Physics Experiments". In: (2012). DOI: 10.1063/1.4773536. arXiv: 1208.2607.
- [2] μ Quans Laser System Specifications.

Acronyms

CCM Centre for Cold Matter

⁸⁷**Rb** Rubidium-87

⁸⁵**Rb** Rubidium-85

MOT Magneto-optical Trap

AOM Acousto-optic Modulator

EOM Electro-optic Modulator

PM Polarisation-Maintaining

QWP Quarter-wave Plate

HWP Half-wave Plate

MFD Mode Field Diameter

PPLN Periodically Poled Lithium Niobate

PLL Phase-Locked Loop

FPGA Field-Programmable Gate Array

EDFA Erbium-Doped Fibre Amplifier

ECDL External-Cavity Diode Laser

TTL Transistor-transistor Logic Circuit

NI National Instruments

DAQ Data Acquisition

VCO Voltage-Controlled Oscillator

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ADC Analogue-to-Digital Converter

DAC Digital-to-Analogue Converter

HAL Hardware Abstraction Layer

SPI Serial Programming Interface

DDS Direct Digital Synthesiser

PBS Polarising beam-splitter

DRO Dielectric Resonator Oscillator

Appendix A.

Laser Systems

This chapter provides a description of the hardware that makes up the experiment. Over the course of the project, the complexity of the experiment necessarily increased. The setup is presented in a bottom-up approach, starting from the most fundamental components, to provide a clear overview of the system.

To-Do:

- Figures describing each of the lasers
- Describe 3D and 2D MOT setups
- Imaging systems
- Microwave synthesisers
- Raman Assembly
- MOT light distribution

A.1. Chapter Overview

The first two sections describe the two commercial laser systems used in this experiment. The μ Quans laser system which generates the light used for cooling and repump in the 2D and 3D Magneto-optical Traps (MOTs), referred to as the MOT light. The design and operation of this laser is given in Section A.2. A secondary laser system, built by MSquared, is used to generate light to drive Raman transitions between

two hyperfine ground states in ⁸⁷Rb¹, otherwise referred to as Raman light. This is described in Section A.3. This is followed by a description of the vacuum chamber in Section ?? which contains both the 2D MOT (Section ??) and the 3D MOT (Section ??).

A.2. The μ Quans Laser System

To-Do:

- Laser Schematic
- Plots of lock signals
- DDS Serial communication
- Power output, stability
- Ref for error signal generation by current modulation
- Move some of this to appendix

All the MOT light in this experiment was generated by the μ Quans laser [2]. μ Quans is a French laser company that is a spin-off from the Institut d'Optique and Observatoire de Paris. Consequently, their technology has been developed over a long history of performing experiments into atom interferometry using Rubidium. A schematic of this laser system is shown in Figure A.1. The μ Quans laser is comprised of four 1560nm ecdls!s (ecdls!s) which are frequency-doubled to produce light at wavelengths close to 780nm. The telecommunications industry, which relies heavily on light in the 1530-1565nm wavelength band for optical communications, has motivated a rapid development in low-noise, robust lasers. In particular, this has enabled a design which does not require free-space optics and is much more resilient to effects such as temperature changes and vibrations, when compared to more conventional 780nm laser systems. The μ Quans laser contains one master laser ², which is locked to the $F = 3 \rightarrow F' = 3.4$ crossover point in Rubidium-85 (85Rb), and serves as an absolute frequency reference. The other three slave lasers are used for output. The first one is used to provide lightAll the MOT light in this experiment was generated by the μ Quans laser [2]. μ Quans is a French laser company that is a spin-off from the

¹The μ Quans laser also has a pair of lasers designed for driving Raman transitions, but these are not used in this experiment. Section A.3 gives an explanation for this.

²see Section A.2.1 for more details

Institut d'Optique and Observatoire de Paris. Consequently, their technology has been developed over a long history of performing experiments into atom interferometry using Rubidium. A schematic of this laser system is shown in Figure A.1. The μQuans laser is comprised of four 1560nm ECDLs!s (ECDLs!s) which are frequencydoubled to produce light at wavelengths close to 780nm. The telecommunications industry, which relies heavily on light in the 1530-1565nm wavelength band for optical communications, has motivated a rapid development in low-noise, robust lasers. In particular, this has enabled a design which does not require free-space optics and is much more resilient to effects such as temperature changes and vibrations, when compared to more conventional 780nm laser systems. The μ Quans laser contains one master laser ³, which is locked to the $F = 3 \rightarrow F' = 3.4$ crossover point in ⁸⁵Rb, and serves as an absolute frequency reference. The other three slave lasers are used for output. The first one is used to provide light for cooling, as well as repump light by modulating the phase of this laser using an Electro-optic Modulator (EOM). The other two make up a pair of lasers for driving Raman transitions. One laser is frequencyoffset locked to the master and the other is phase-locked to the first, to ensure that the relative phase between the two lasers is constant. It should be noted that this Raman laser was not used in this experiment, so will not be discussed in great detail. Each of these slave lasers is amplified in an Erbium-Doped Fibre Amplifier (EDFA) before being frequency doubled in a Periodically Poled Lithium Niobate (PPLN) and passed through an AOM which is used to control the output power during the experiment. for cooling, as well as repump light by modulating the phase of this laser using an **EOM**. The other two make up a pair of lasers for driving Raman transitions. One laser is frequency-offset locked to the master and the other is phase-locked to the first, to ensure that the relative phase between the two lasers is constant. It should be noted that this Raman laser was not used in this experiment, so will not be discussed in great detail. Each of these slave lasers is amplified in an EDFA before being frequency doubled in a PPLN and passed through an AOM which is used to control the output power during the experiment.

A.2.1. Absolute Frequency Reference

The purpose of the master laser is to provide an absolute frequency reference so that the frequency of the output lasers can be controlled by comparing the difference frequency

³see Section A.2.1 for more details

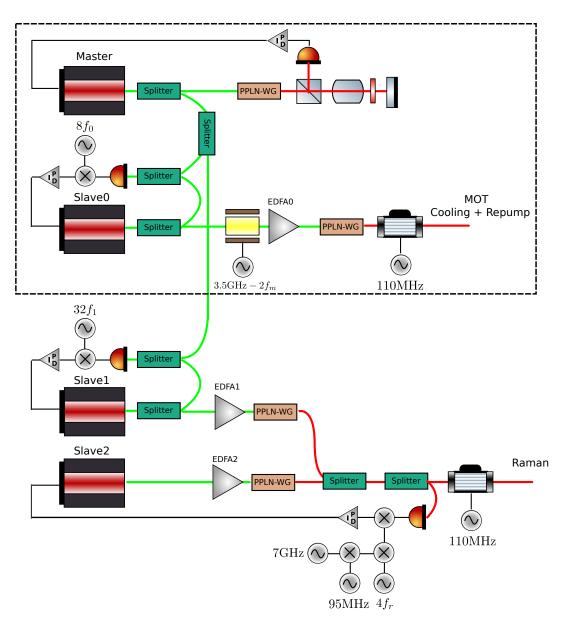


Figure A.1.: Schematic of the μ Quans laser system. Each output laser is derived from a 1560nm ECDL (shown in green) which is amplified using an EDFA and then frequency-doubled to 780nm using a PPLN crystal. A master laser is locked to the 3,4 crossover in ⁸⁵Rb and the output lasers are offset-locked to their corresponding frequencies. The dashed region indicates the components used for generating light for the MOTs, which was the only function of this laser for this experiment.

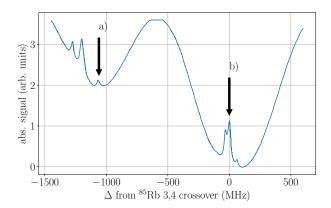


Figure A.2.: Saturated absorption spectroscopy using the Rubidium vapour cell in the μ Quans laser. The absorption features indicated are a): the $F=2 \rightarrow F'=3$ transition in ⁸⁷Rb and b): the crossover resonance between the $F=3 \rightarrow F'=3$ and $F=3 \rightarrow F'=4$ transitions in ⁸⁵Rb which is used to lock the frequency of the master laser.

between them and the master. Lasers with linewidths narrower than their natural linewidth can be achieved by using a servo to stabilise their frequency and is essential for any experiment that requires laser light of a precise frequency. The frequency reference for the master is obtained using saturated absorption spectroscopy inside a Rubidium vapor cell. The sub-Doppler features in this spectrum are insensitive to temperature changes, and under sufficiently weak laser power have linewidths close to the natural linewidth of Rubidium ($\Gamma \sim 2\pi \times 6 \text{MHz}$). Figure A.2 shows the saturated absorption spectrum using the μ Quans master laser. This is obtained by fine adjustment of the temperature of the master External-Cavity Diode Laser (ECDL). The laser is set to lock to the crossover resonance between the $F = 3 \rightarrow F' = 3$ and $F = 3 \rightarrow F' = 4$ transitions in ⁸⁵Rb (indicated as b)), which is the strongest feature in the spectrum as well as being relatively close the the cooling transition in ⁸⁷Rb (indicated as a)). Some form of feed-back onto the master laser is required to keep its frequency fixed. The simplest way to achieve this is to use a signal that is linearly proportional to the deviation in frequency from the set-point, if one exists. The frequency of the laser is modulated by weakly modulating the current to the master ECDL. add more detail about the error signal lineshape The error signal shown in Figure ?? is obtained by demodulating the absorption signal using a lock-in amplifier. In fact, this current modulation is always present on the master laser and the saturated absorption spectrum shown previously has been processed to average out the effects from this fast frequency modulation. In addition to proportional feed-back

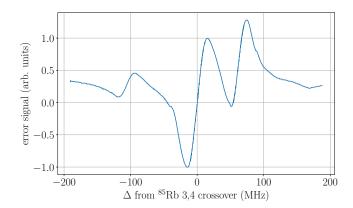


Figure A.3.: Error signal obtained by modulating the laser current. Close to the lock point, the signal is approximately linear. This signal is used in a feed-back loop to correct for frequency changes of the master laser.

from the error signal, the servo that controls the master frequency also contains an integrator to compensate for long-term drifts. Typically, these arise from external temperature changes and if unaccounted for, they could cause the laser to unlock. In the conditions of our laboratory, where the temperature is externally controlled, this has never occurred.

A.2.2. Generating MOT light

A.2.3. Raman light

A.2.4. Real-time Frequency Control

A.3. The M-Squared Laser System

To-Do:

- Schematic
- Raman PLL phase-noise
- Laser Control
- DCS module

- A.3.1. Laser Specifications
- A.3.2. The DCS Control Module
- A.3.3. Frequency Control of the Raman Lasers
- A.3.4. Controlling the Phase Difference