Index notation in Lean 4

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Abstract

The physics community relies on index notation to effectively manipulate types of tensors. This paper introduces the first formally verified implementation of index notation in the interactive theorem prover Lean 4. By integrating index notation into Lean, we bridge the gap between traditional physics notation and formal verification tools, making it more accessible for physicists to write and prove results within Lean. In the background, our implementation leverages a novel application of category theory.

The file you are currently reading is a draft. There will be typos, mistakenesses, notation that is not defined etc.

1. INTRODUCTION

In previous work [1], the author initiated the digitalization (or formalization) of high energy physics results using the interactive theorem prover Lean 4 in a project called HepLean. Lean is a programming language with syntax resembling traditional pen-and-paper mathematics. Users can write definitions, theorems, and proofs in Lean, which are then automatically checked for correctness using dependent-type theory. The HepLean project is driven by four primary motivations: to facilitate easier look-up of results through a linear storage of information; to support the creation and proof of new results using automated tactics and AI tools; to facilitate correctness checking of result in high energy physics; and to introduce new pedagogical methods for high-energy physics and computer science.

HepLean is part of a broader movement of projects to formalize parts, or all, of mathematics and science. The largest of these projects is Mathlib [2], which aims to formalize mathematics. Indeed, HepLean uses many results from Mathlib, and is built downstream thereof. Other projects in mathematics include the ongoing effort led by Kevin Buzzard to formalize the proof of Fermat's Last Theorem into Lean [3]. In the realm of the sciences, the paper [4] looks at absorption theory, thermodynamics, and kinematics in Lean, whilst the package SciLean [5], is a push in the direction of scientific computing within Lean.

Physicists rely heavily on specialized notation to express mathematical concepts succinctly. Among these index notation is particularly prevalent, as it provides a compact and readable way to represent specific types of tensors and operations between them. Such tensors form a backbone of modern physics.

Having a way to use index notation in Lean is crucial for the digitalisation of results from high energy physics. As well as, making results from high energy physics easier to write and prove in Lean, it will make the syntax more familiar to high energy physicists. Such an implementation is the subject of this paper. It is a challenge because, we need an implementation that is nice and flexible to use, but which is also formally well-defined and rigorous. We hope that the implementation presented here will



Figure 1: Overview of the implementation of index notation in Lean. The solid lines represent formally verified parts of the implementation.

not only enhances usability of Lean but also promotes the adoption of formal methods in the physics community.

As a taster of what is to come, in our implementation the result regarding Pauli matrices that $\sigma^{\nu\alpha\dot{\beta}}\sigma^{\alpha'\dot{\beta}'}_{\nu}=2\varepsilon^{\alpha\alpha'}\varepsilon^{\beta\beta'}$ is written in Lean as follows

$$\{\text{pauliCo} \mid v \ \alpha \ \beta \ \otimes \ \text{pauliContr} \mid v \ \alpha' \ \beta' = 2 \cdot_t \varepsilon \text{L} \mid \alpha \ \alpha' \ \otimes \varepsilon \text{R} \mid \beta \ \beta'\}^\text{T}$$

Lean will correctly interpret this result as a tensor expression with the correct contraction of indices and permutation of indices between the sides of the expression.

Previous implementations of index notation have been made in programming languages like Haskell [6]. However, the programs they appear in do not provide the formal verification capabilities inherent in Lean. The formal verification requirement of Lean introduces unique challenges in implementing index notation, necessitating (what we believe is) a novel solution.

In Section 2 of this paper, we will discuss the details of our implementation. In Section 3 we will give two examples of definitions, theorems and proofs using index notation, to give the reader an idea of how our implementation works in practice. The first of these examples involves a lemma regarding the contraction of indices of symmetric and antisymmetric tensors. The second involves examples related to the Pauli matrices and bispinors. We finish this paper in Section 4 by discussing potential future work related to this project.

NOTATION

In this paper, we will follow Lean's notation for types and terms. For example if C is a type (which can be thought of as similar to a set), then c: C is an element of the type C. In addition, instead of having two streams of notation for mathematical objects, one from the Lean code and one in LaTeX, we will use Lean code throughout represent mathematical objects.

2. IMPLEMENTATION OF INDEX NOTATION INTO LEAN 4

Our implementation of index notation can be thought of as three different representations of tensor expressions and maps between them. This is illustrated in Figure 1.

The first representation is *syntax*. This can (roughly) be thought of as the informal string that represents the tensor expression. It is what the user interacts with when writing results in Lean, and what appears in raw Lean files. The code snippet above, for Pauli matrices is an example of syntax.¹

The second representation of a tensor expression is a *tensor tree*. This representation is mathematically formal, but easy to use and manipulate. It is a structured tree which has different types of nodes for each of the main operations that one can perform on tensors.

¹In practice there is a representation before syntax, which is as a sequence of tokens in file which is parsed by Lean into a more structured syntax. For simplicity, we think of these representations as the same here.

The process of going from syntax to a tensor tree is done via an elaborator which follows a number of (informally defined) rules.

The third and final representation is a bona-fide *tensor*. This representation is the mathematical object that we are actually interested in. However, in going to this representation we lose the structure of the tensor expression itself, making it difficult to work with.

The process of going from a tensor tree to a tensor involves properties from a symmetric monoidal category of representations.

Before we discuss these processes in more detail we do some set-up by defining formally a 'tensor species'.

2.1. TENSOR SPECIES

A tensor species is a formalization of the data needed to define e.g., complex Lorentz tensors, real Lorentz tensors or Einstein tensors. The word 'species' is a nod to 'graphical species' defined in [7, 8], from which our construction is inspired.

We will start by giving the complete definition of a tensor species in Lean, and then dissect this definition, discussing each of the components in turn.

In Lean a tensor species is defined as follows:

```
/-- The structure of a type of tensors e.g., Lorantz tensors, Einstein tensors,
complex Lorentz tensors. -/
structure TensorSpecies where
  /-- The commutative ring over which w
                                                     to consider the tensors to live in,
    usually '\mathbb{R}' or '\mathbb{C}'. -/
  k : Type
  /-- An instance of 'k' as a commut
  k_commRing : CommRing k
  /-- The symmetry group acting on these tensor e.g. the Lorentz group or SL(2,\mathbb{C}).
    -/
  G : Type
  /-- An instance of 'G' as a group. -/
  G_group : Group G
  /-- The colors of indices e.g. up or down. -/
  C : Type
  /-- A functor from 'C' to 'Rep k G' giving our building block representations.
    Equivalently a function 'C 
ightarrow Re k G'. -/
  \mathtt{FD} : \mathtt{Discrete}\ \mathtt{C} \Rightarrow \mathtt{Rep}\ \mathtt{k}\ \mathtt{G}
  /-- A specification of the dimension of each color in C. This will be used for
    explicit
    evaluation of tensors. -/
  \mathtt{repDim}\,:\,\mathtt{C}\,\to\,\mathbb{N}
  /-- repDim is not zero for any color. This allows casting of 'N' to 'Fin
    (S.repDim c)'. -/
  repDim_neZero (c : C) : NeZero (repDim c)
  /-- A basis for each Module, determined by the evaluation map. -/
  basis : (c : C) \rightarrow Basis (Fin (repDim c)) k (FD.obj (Discrete.mk c)).V
  /-- A map from 'C' to 'C'. An involution. -/
  \tau : C \rightarrow C
```

```
/-- The condition that '	au' is an involution. -/
	au_{	ext{-involution}} : Function.Involutive 	au
/-- The natural transformation describing contraction. -/
contr : OverColor.Discrete.pair	au FD 	au 	o \mathbb{1}_{\_} (Discrete C \Rightarrow Rep k G)
/-- Contraction is symmetric with respect to duals. -/
contr_tmul_symm (c : C) (x : FD.obj (Discrete.mk c))
     (y : FD.obj (Discrete.mk (\tau c))) :
  (contr.app (Discrete.mk c)).hom (x \otimes_t [k] y) = (contr.app (Discrete_mk (1))
  c))).hom
  (y \otimes_t (FD.map (Discrete.eqToHom (	au_involution c).symm)).hom x
/-- The natural transformation describing the unit. -/
unit : \mathbb{1}_{\_} (Discrete C \Rightarrow Rep k G) \rightarrow OverColor.Discrete.	auPai
/-- The unit is symmetric. -/
unit_symm (c : C) :
  ((unit.app (Discrete.mk c)).hom (1 : k)) =
  ((FD.obj (Discrete.mk (\tau (c)))) \triangleleft
     (FD.map (Discrete.eqToHom (\tau_involution c)))).hom
  ((\beta_{-} (FD.obj (Discrete.mk (\tau (\tau c)))) (FD.obj (Discrete.mk (\tau (c))))).hom.hom
  ((unit.app (Discrete.mk (\tau c))).hom (1 : k)))
/-- Contraction with unit leaves invariant. -/
contr_unit (c : C) (x : FD.obj (Discrete.mk (c))) :
  (\lambda_{-} (FD.obj (Discrete.mk (c)))).hom.hom
  (((contr.app (Discrete.mk c)) ▷ (FD obj (Discrete.mk (c)))).hom
  ((\alpha_{-} - (FD.obj (Discrete.mk (c)))).inv.hom
  (x \otimes_t [k] \text{ (unit.app (Discrete.mk c)).hom } (1 ; k)))
/-- The natural transformation describit the metric. -/
\texttt{metric} \; : \; \mathbb{1}_{\_} \; (\texttt{Discrete} \; C \; \Rightarrow \; \texttt{Rep} \; \; k \; \; \texttt{G}) \; \rightarrow \; \texttt{OverColor}. \\ \texttt{Discrete}. \\ \texttt{pair} \; \; \texttt{FD}
/-- On contracting metrics we get back the unit. -/
contr_metric (c : C)
  (\beta_{-}(FD.obj(Discrete.mk c))(FD.obj(Discrete.mk(\tau c)))).hom.hom
  (((\mathbf{FD.obj}(\mathbf{Discrete.mk} c)) \triangleleft (\lambda_{-}(\mathbf{FD.obj}(\mathbf{Discrete.mk} (\tau c)))).hom).hom)
  (((FD.obj (Discrete.mk c)) \triangleleft ((contr.app (Discrete.mk c)) \triangleright
  (FD.obj (Discrete.mk (\tau c)))).hom
  (((FD.obj (Discrete.mk c)) \triangleleft (\alpha_{-} (FD.obj (Discrete.mk (c)))
     (FD.obj (Discrete.mk (	au c))) (FD.obj (Discrete.mk (	au c)))).inv).hom
  ((\alpha_{(FD.obj (Discrete.mk (c)))} (FD.obj (Discrete.mk (c)))
     (FD.obj (Discrete.mk (\tau c)) \otimes FD.obj (Discrete.mk (\tau c))).hom.hom
  ((metric.app (Discrete.mk c)).hom (1 : k) \otimes_t[k]
     (metric.app (Discrete.mk (\tau c))).hom (1 : k)))))
  = (unit.app (Discrete.mk c)).hom (1 : k)
```

Let us go through this definition piece by piece. The first part of the definition of a tensor species is a type k which is a commutative ring via k_commRing. For a given S: TensorSpecies, we retrieve k and all of the components of TensorSpecies via S.k etc. For the tensor species of complex Lorentz tensors, denoted complexLorentzTensor, the ring complexLorentzTensor.k is the ring of complex numbers.

The second part of the definition of a tensor species is a type G which is a group via G_group. For complex Lorentz tensors, complexLorentzTensor.G is the group $SL(2,\mathbb{C})$.

The next part of the definition is a type C which we call the type of colors. For complex Lorentz tensors, complexLorentzTensor.C is set equal to the type

which contains six colors. Colors can be thought of as labels for each of the building block representations making up the tensor species. This is made formal by the next part of the definition of a tensor species, a functor FD, which assigns to each color a representation of G over k. Note that for c: C to apply FD we have to write FD.obj (Discrete.mk c) in Lean. This is somewhat cumbersome, so in what follows we will abbreviate this to FD c. For complex Lorentz tensors, the functor complexLorentzTensor.FD is defined as:

The representations appearing here are:

- The representation of left-handed Weyl fermions, denoted in Lean as Fermion.leftHanded, and corresponding to the representation of $SL(2,\mathbb{C})$ taking $v \mapsto Mv$ for $M \in SL(2,\mathbb{C})$.
- The representation of 'alternative' left-handed Weyl fermions, denoted in Lean as Fermion.altLeftHanded, and corresponding to the representation of $SL(2,\mathbb{C})$ taking $v\mapsto M^{-1T}v$ for $M\in SL(2,\mathbb{C})$.
- The representation of right-handed Weyl fermions, denoted in Lean as Fermion.rightHanded, and corresponding to the representation of $SL(2,\mathbb{C})$ taking $v \mapsto M^*v$ for $M \in SL(2,\mathbb{C})$.
- The representation of 'alternative' right-handed Weyl fermions, denoted in Lean as Fermion.altRightHanded, and corresponding to the representation of $SL(2,\mathbb{C})$ taking $v \mapsto M^{-1\dagger}v$ for $M \in SL(2,\mathbb{C})$.
- The representation of contravariant Lorentz tensors, denoted in Lean as Lorentz.complexContr, and corresponding to the representation of $SL(2,\mathbb{C})$ induced by the homomorphism of $SL(2,\mathbb{C})$ into the Lorentz group and the contravariant action of the Lorentz group on four-vectors.
- The representation of covariant Lorentz tensors, denoted in Lean as Lorentz.complexCo, and corresponding to the representation of $SL(2,\mathbb{C})$ induced by the homomorphism of $SL(2,\mathbb{C})$ into the Lorentz group and the covariant action of the Lorentz group on four-vectors.

As an example of how these are defined in Lean, the representation of left-handed Weyl fermions Fermion.leftHanded is given by:

```
/-- The vector space \mathbb{C}^2 carrying the fundamental representation of \mathrm{SL}(2,\mathbb{C}).
  In index notation corresponds to a Weyl fermion with indices \psi^a. -/
\operatorname{def} leftHanded : Rep \mathbb{C} SL(2,\mathbb{C}) := Rep.of {
  /- The function from \mathrm{SL}(2,\mathbb{C}) to endomorphisms of LeftHandedModule
    (which corresponds to the vector space \mathbb{C}^2). -/
  toFun := fun M => {
    /- Start of the definition of the linear map. -/
    /- The function underlying the linear map. Defined as the dot pro
    toFun := fun (\psi : LeftHandedModule) =>
      LeftHandedModule.toFin2\mathbb{C}Equiv.symm (M.1 *, \psi.toFin2\mathbb{C}),
    /- Proof that the function is linear with respect to add
    map_add' := by
      intro \psi \psi,
      simp [mulVec_add]
    /- Proof that the function is linear with res
                                                                scalar mul
    map_smul' := by
      intro r \psi
      simp [mulVec_smul]
    /- End of the definition of the linear map. -/}
  /- Proof that (the outer) to Fun goes the identity map on the
    SL(2,\mathbb{C}). -/
  map_one' := by
    ext i
    simp
  /- Proof that the action of the product
                                                        elements :
    the product of the actions of the element
  map_mul' := fup M N => b
    simp only [SpecialLinearGroup.coe_mul]
    simp only [LinearMap.coe_mk, AddHom.coe_mk, LinearMap.mul_apply,
    LinearEquiv.apply_symm_apply
      mulVec_mulVec]}
```

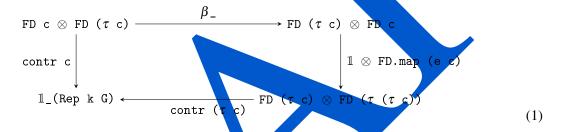
where we have added some explanatory comments to this code, not seen in the actual Lean code, to give the reader an idea of what each part does.

The next part of the definition of a tensor species is the map repDim, which assigns to each color a natural number, which corresponds to the dimension of the representation associated to that color. The condition is placed on the representations that they are non-empty, i.e., that the dimension is not equal to zero repDim_neZero. The part of the definition of a tensor species basis gives a basis indexed by Fin (repDim c) (numbers from 0 to repDim c-1) of each representation for each $c \in C$. We will use this basis in the definition of evaluation of tensor indices. For complex Lorentz tensors, these are the standard basis for Lorentz vectors and Weyl-fermions.

Next in the definition of a tensor species is the map τ , which is an involution via $\tau_{involution}$. This assigns to each color its 'dual' corresponding to the color it can be contracted with. So, for complex Lorentz tensors the map τ is given by

```
τ := fun c =>
match c with
| Color.upL => Color.downL
| Color.downL => Color.upL
| Color.upR => Color.downR
| Color.downR => Color.upR
| Color.up => Color.down
| Color.down => Color.down
```

The contraction itself is defined in a tensor species by contr, which is a natural transformation from the functor FD $_{-} \otimes$ FD $(\tau$ $_{-})$, to the functor $\mathbb{1}_{-}$ (Discrete C \Longrightarrow Rep k G) takes every object in C to the trivial representation $\mathbb{1}_{-}$ (Rep k G). This natural transformation is simply the assignment to each color c a linear-map from FD c \otimes FD $(\tau$ c) to k which is equivariant with respect to the group action. This contraction cannot be defined arbitrarily, but must satisfy the symmetry condition contr_tmul_symm, which corresponds to the commutative diagram.



where β_{-} is the braiding of the symmetric monoidal category, and e is the isomorphism in Discrete S.C between c and τ (τ c), which exists since τ is an involution.

Along with the contraction the definition of a tensor species includes the unit, along with its own symmetry condition unit_symm. The unit is a natural transformation from the functor $\mathbb{1}$ _ (Discrete C \Rightarrow Rep k G) to FD $(\tau_{-}) \otimes$ FD $_{-}$. This is the assignment to each color c an object of FD $(\tau_{-}) \otimes$ FD c which is invariant with respect to the group action. The symmetry condition unit_symm is represented by the commutative diagram

The unit is actually a unit in the sense that contraction with it does nothing. This is made formal with

the condition contr_unit, which is corresponds to the diagram

where ρ_{-} is the right-unitor, λ_{-} is the left-unitor, and α_{-} is the associator.

The final part of the definition of a tensor species is the metric, metric and it's interaction with the contraction and unit contr_metric. The metric is a natural transformation from the functor $\mathbb{1}_{-}$ (Discrete $C \Rightarrow \text{Rep } k$ G) to the functor $\text{FD }_{-} \otimes \text{FD }_{-}$. It thus represents the assignment to each color c an object of $\text{FD }_{-} \otimes \text{FD }_{-}$ c which is invariant with respect to the group action. The metric can be used to change an index into a dual index. The condition contr_metric corresponds to the diagram

For complex Lorentz tensors, the contraction is defined through the dot product, in the sense that e.g., the contraction of ψ^{μ} and ϕ_{μ} is via dot product of the underlying vectors. The metric is defined through the Minkowski metric and the metric tensors e.g. $\varepsilon^{\alpha\dot{\alpha}}$ for Weyl fermions. Lastly ,the units are defined through identity matrices.

2.2. TENSORS

Given any type C, we define the category OverColor C as follows. Objects are functions $f: X \to C$ for some type X. A morphism from $f: X \to C$ to $g: Y \to C$ is a bijection $\varphi: X \to Y$ such that $f = g \circ \varphi$. This category is equivalent to the core of the category of types sliced over C.

The category 0verColor C is not any old category, it carries a symmetric monoidal structure, which we will denote \otimes . The structure such that $f \otimes g$ for objects f and g is the induced map $X \oplus Y \to C$ where \oplus denotes the disjoint union of types. In Lean this is denoted Sum.elim f g.

For a given tensor species S, the functor FD can be lifted to a functor to a symmetric monoidal functor from OverColor S.C to Rep k G. This functor takes $f: X \to S.C$ to the tensor product, \otimes [k] x, S.FD (f x), of all FD (f x). This construction is general and functorial, allowing us to define

the functor

```
def OverColor.lift : (Discrete S.C \Rightarrow Rep S.k S.G) \Rightarrow BraidedFunctor (OverColor S.C) (Rep S.k S.G) where ...
```

where Lean calls symmetric monoidal functors BraidedFunctor.

We denote the lift of S.FD by F. Namely, we have

```
def F : BraidedFunctor (OverColor S.C) (Rep S.k S.G) := (OverColor.lift).obj S.FD
```

We can think of an object $f: X \to S.C$ of OverColor S.C as a type of indices X, and a specification of what color or representation each index is associated to. We can apply the functor S.F to $f: X \to S.C$ as follows S.F.obj (OverColor.mk f), which we will abbreviate to S.F f. This representation, S.F f, is the tensor product of each of these representations, so that vectors v: S.F f can be thought of as tensors with indices indexed by X of color C.

With this in mind, we define a general tensor of a species S as a vector in a representation S.F f.

In physics, we typically focus on objects $f: X \to S.C$ of OverColor S.C where X is a finite type of the form Fin n for some $n: \mathbb{N}$. When appropriate in what follows, we will restrict to these objects.

2.3. TENSOR TREES AND THEIR MAP TO TENSORS

Tensor trees are tress with a node for each of the basic operations one can perform on a tensor. Namely, tensor trees have nodes for addition of tensors, permutation of tensor indices, negation of tensors, scalar multiplication of tensors, group action on a tensor, tensor product of tensors, contraction of tensor indices, and evaluation of tensor indices. They also have nodes for tensors themselves.

Given a species S, we have a type of tensor tree for each map of the form $c: Fin n \to S.C$ in OverColor S.C. This restriction to Fin n is done for convience.

Tensor trees are defined inductively through a number of constructors:

```
inductive TensorTree (S : TensorSpecies) : \{n : \mathbb{N}\} \to (\text{Fin } n \to S.C) \to \text{Type where}
  /-- A general tensor node. -/
   \mid tensorNode {n : N} {c : Fin n \rightarrow S.C} (T : S.F.obj (OverColor.mk c)) :
     TensorTree S c
  /-- A node correpsonding to the scalar multiple of a tensor by a element of the
     field. -/
  \mid smul \{ \texttt{n} : \mathbb{N} \} \{ \texttt{c} : \texttt{Fin n} 	o \texttt{S.C} \} : \texttt{S.k} 	o \texttt{TensorTree S c} 	o \texttt{TensorTree S c}
  /-- A node corresponding to negation of a tensor. -/
  \mid neg {n : \mathbb{N}} {c : Fin n \rightarrow S.C} : TensorTree S c \rightarrow TensorTree S c
  /-- A node corresponding to the addition of two tensors. -/
  | \  \, \text{add} \  \, \{ \text{n} : \, \mathbb{N} \} \  \, \{ \text{c} : \, \text{Fin} \  \, \text{n} \, \rightarrow \, \text{S.C} \} \, : \, \text{TensorTree S c} \, \rightarrow \, \text{TensorTree S c} \, \rightarrow \, \text{TensorTree}
     Sc
  /-- A node correpsonding to the action of a group element on a tensor. -/
  \mid action \{ \texttt{n} : \mathbb{N} \} \ \{ \texttt{c} : \texttt{Fin} \ \texttt{n} 	o \texttt{S.C} \} : \texttt{S.G} 	o \texttt{TensorTree} \ \texttt{S} \ \texttt{c} 	o \texttt{TensorTree} \ \texttt{S} \ \texttt{c}
  /-- A node corresponding to the permutation of indices of a tensor. -/
  | \text{ perm } \{ \text{n m} : \mathbb{N} \} \ \{ \text{c} : \text{Fin n} \to \text{S.C} \} \ \{ \text{c1} : \text{Fin m} \to \text{S.C} \}
         (\sigma : (OverColor.mk c) 
ightarrow (OverColor.mk c1)) (t : TensorTree S c) :
     TensorTree S c1
  /-- A node corresponding to the product of two tensors. -/
```

```
| prod {n m : N} {c : Fin n \rightarrow S.C} {c1 : Fin m \rightarrow S.C} (t : TensorTree S c) (t1 : TensorTree S c1) : TensorTree S (Sum.elim c c1 o finSumFinEquiv.symm)

| --- A node corresponding to the contraction of indices of a tensor. -/
| contr {n : N} {c : Fin n.succ.succ \rightarrow S.C} : (i : Fin n.succ.succ) \rightarrow (j : Fin n.succ) \rightarrow (h : c (i.succAbove j) = S.\tau (c i)) \rightarrow TensorTree S (c o Fin.succAbove i o Fin.succAbove j)

| --- A node corresponding to the evaluation of an index of a tensor. |
| eval {n : N} {c : Fin n.succ \rightarrow S.C} : (i : Fin n.succ) \rightarrow (x : N) \rightarrow TensorTree S c \rightarrow TensorTree S (c o Fin.succAbove i)
```

Each constructor here, e.g. tensorNode, smul, neg, etc., can be thought of a forming a different type of node in a tensor tree.

Since the interpretation of each of the constructors is down to how we turn them into a tensor, we discuss this before discussing each of the constructors in turn. The process of going from a tensor tree to a tensor is proscribed by a function TensorTree S c \rightarrow S.F.obj (OverColor.mk c), which is defined recursively as follows:

```
/-- The underlying tensor a tensor unit corresponds to. -/

def TensorTree.tensor {n : N} {c : Fin n \rightarrow S.C} : TensorTree S c \rightarrow S.F.obj
   (OverColor.mk c) := fun
   | tensorNode t => t
   | smul a t => a · t.tensor
   | neg t => - t.tensor
   | add t1 t2 => t1.tensor + t2.tensor
   | action g t => (S.F.obj (OverColor.mk _)).\rho g t.tensor
   | perm \sigma t => (S.F.map \sigma) hom t.tensor
   | prod t1 t2 => (S.F.map (OverColor.equivToIso finSumFinEquiv).hom).hom
   ((S.F.\mu _ _).hom (t1.tensor \otimes_t t2.tensor))
   | contr i j h t => (S.contrMap _ i j h).hom t.tensor
   | eval i e t => (S.evalMap i (Fin.ofNat' _ e)) t.tensor
```

Let us now dicuss each of the constructors in turn, after which we will give some examples tensor trees.

tensorNode: The constructor tensorNode based on c creates a tensor tree from a tensor t in S.F.obj (OverColor.mk c). This tensor tree consists of a single node that directly represents the tensor. Since all other tensor tree constructors require an existing tensor tree as input, tensorNode serves as the foundational base case for building more complex trees. Diagramatically, such a tree is



Naturally, the tensor associated with this node is exactly the tensor provided during construction.

smul: The constructor smul takes a scalar a and an existing tensor tree t based on c and constructs a new tensor tree also based on c. Conceptually, this new tree has a root node labeled smul a, with the tensor tree t as its child. Diagramatically, such a tree is



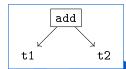
The tensor associated with this new tree is obtained by multiplying the tensor associated with t by the scalar a.

neg: The constructor neg takes an existing tensor tree t based on c and constructs a new tensor tree also based on c. This new tree has a root node labeled neg, with the tensor tree t as its child. Diagramatically, we have



The tensor associated with this new tree is obtained by negating the tensor associated with t.

add: The constructor add takes two existing tensor trees t1 and t2, based on the same c, and constructs a new tensor tree. This new tree has a root node labeled add, with the tensor trees t1 and t2 as its children. Diagramatically, this corresponds to



The tensor associated with this new tree is obtained by adding the tensors associated with t1 and t2.

action: The constructor action takes a group element g of S.G, an existing tensor tree t based on c, and constructs a new tensor tree also based on c. This new tree has a root node labeled action g, with the tensor tree t as its child. Diagramatically,



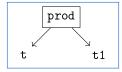
The tensor associated with this new tree is obtained by acting on the tensor associated with t with the group element g.

perm: The constructor perm takes a morphism σ from c to c1 in OverColor S.C and an existing tensor tree t based on c, and constructs a new tensor tree based on c1. This new tree has a root node labeled perm σ , with the tensor tree t as its child. Diagramatically, this is given by:



The tensor associated with this new tree is obtained by applying the image of the morphism σ under the functor S. F to the tensor associated with t.

proof: The constructor proof takes two existing tensor trees t based on c and t1 based on c1 and constructs a new tensor tree based on Sum.elim c c1 o finSumFinEquiv.symm which is the map from Fin (n + n1) acting via c i on i = 0, ..., n-1 and via c1 (i - n) on i = n, ..., n+n1-1. This new tree has a root node labeled proof, with the tensor trees t and t1 as its children. Diagramatically, this is given by



The tensor associated with this new tree is obtained by taking the tensor product of the tensors associated with t and t1, giving a vector in S.F.obj (OverColor.mk c) \otimes S.F.obj (OverColor.mk c1), using the tensorate of S.F to map this vector into a vector in S.F.obj (OverColor.mk c \otimes OverColor.mk c1), and finally using an isomorphism between OverColor.mk c \otimes OverColor.mk c1 and OverColor.mk (Sum.elim c c1 o finSumFinEquiv.symm) to map this vector into S.F.obj (OverColor.mk (Sum.elim c c1 o finSumFinEquiv.symm)).

contr: The constructor contr firstly, takes an existing tensor tree t based on a c. Fin n. succ. succ. -> S.C. Here n. succ. succ is n+1+1 with succ meaning the successor of a natural number. It also takes, an i of type Fin n. succ. succ, j of type Fin n. succ and a proof that c (i. succabove j) = S. τ (c i), where i. succabove is the map from Fin n. succ to Fin n. succ. succ with a hole at i. The proof says that the color of the index i. succabove j is the dual of the color of the index i, and thus that these two indices can be contracted. Note that we use a j in Fin n. succ and i. succabove j as the index to be contracted, instead of another index in Fin n. succ. succ to ensure the two indices to be contracted are not the same. The constructor outputs a new tensor tree based on c o i. succabove o j. succabove. This new tree has a root node labeled contraction j, with the tensor tree t as its child. Diagramatically,



The tensor associated with the new tensor tree is construced as follows. We start with a tensor in S.F.obj (OverColor.mk c). This is mapped into a vector in (S.FD.obj (c i) \otimes S.FD.obj (S. τ (c i))) \otimes S.F.obj (OverColor.mk (c \circ Fin.succAbove i \circ Fin.succAbove j)) via an equivalence from S.F.obj (OverColor.mk c). The equivalence is constructed using an equivalence in OverColor C to extract i and i.succAbove j from c, then using the tensorate of S.F, and the fact that c (i.succAbove j) = S. τ (c i). Using the contraction for c i, we then get a vector in 1 \otimes S.F.obj (OverColor.mk (c \circ Fin.succAbove i \circ Fin.succAbove j)) which is then mapped to a tensor in S.F.obj (OverColor.mk (c \circ Fin.succAbove i \circ Fin.succAbove j)) using the left-unitor of Rep S.G S.k. This is all contained within the S.contrMap appearing in the definition of function TensorTree.tensor.

eval: The final constructor eval takes an existing tensor tree t based on a c: Fin n.succ -> S.C, an i of type Fin n.succ, and a natural number x. The constructor outputs a new tensor tree based on c o Fin.succAbove i. This new tree has a root node labeled eval i x, with the tensor tree t as its child. Diagramatically, this corresponds to



The tensor associated with the new tensor tree is constructed as follows. We start with a tensor in S.F.obj (OverColor.mk c). This is mapped into a vector in S.F.obj (c i) \otimes S.F.obj (OverColor.mk (c o Fin.succAbove i)) via an equivalence from S.F.obj (OverColor.mk c). The equivalence is constructed using an equivalence in OverColor C to extract i from c, and then using the tensorate of S.F. Using the evaluation for c i at the basis element indicated by x (if x is too big a natural number for

the number of basis elements it defaults to 0), we then get a vector in 1 \otimes S.F.obj (OverColor.mk (c \circ Fin.succAbove i)). This mapping should be thought of as occuring in the category of modules over S.k, rather then in Rep S.G S.k, since the evaluation is not invariant under the group action. We then fininally map into S.F.obj (OverColor.mk (c \circ Fin.succAbove i)) using the left-unitor. This is all contained within the S.evalMap appearing in the definition of function TensorTree.tensor

The main reason tensor trees are nice to work with is that they have the following property. Define a subtree of a tensor trees to be a node and all child nodes of that node. If T is a tensor tree and S a subtree of T, we can replace S in T with another tensor tree S' to get a new overall tensor-tree T'. If S and S' have the same underlying tensor, then T and T' will.

In Lean this property manifests in a series of lemmas. For instance, for instance for the control constructor we have the lemma:

```
lemma contr_tensor_eq {n : N} {c : Fin n.succ.succ \rightarrow S.C} {T1 T2 : TensorTree S c} (h : T1.tensor = T2.tensor) {i : Fin n.succ.succ} {j : Fin n.succ} {h' : c (i.succAbove j) = S.\tau (c i)} : (contr i j h' T1).tensor = (contr i j h' T2).tensor := by simp only [Nat.succ_eq_add_one, contr_tensor] rw [h]
```

These lemmas allow us to navigate to the certain places in tensor trees and replace subtrees with other subtrees. We will see this used exensively in the examples.

2.4. SYNTAX AND THEIR MAP TO TENSOR TREES

We now dicuss how we make the Lean code look similar to what we would use on pen-and-paper physics.

This is done using a two-step process. Firstly, we define syntax for tensor expressions. Then we write code, called elaboration, to turn this syntax into a tensor tree. This process is not formally defined or verified, Lean takes the outputed tensor tree as the formal object to work with.

Instead of delving into the nitty-gritty details of how this process works under the hood, we give some examples to see how it works.

In what follows we will assume that T, T1, etc are tensors defined as elements of S.F c, S.F c1, etc for some tensor species S and some $c : Fin n \to S.C, c1 : Fin n1 \to S.C,$ etc for which the below expressions make sense.

The syntax allows us to write the following

```
\{T \mid \mu \ v\}^T tensorNode T
```

for a tensor node. Here the μ and v are free variables and it does not matter what we call them - Lean will elaborate the expression in the same way. The elaborator also knows how many indices to expect for a tensor T and will raise an error if the wrong number are given. The $\{_\}^T$ notation is used to tell Lean that the syntax is to be treated as a tensor expression. Throughout this section we will use the two-sided boxes given above, which on the left denote the syntax and on the right the expression it is elaborated to.

We can write e.g.,

```
\{\mathtt{T} \mid \mu \ v\}^{\mathrm{T}}.\mathtt{tensor} (tensorNode T).tensor
```

to get the underlying tensor.

Note that we have not lowered or risen the indices, as one would expect from pen-and-paper notation. There is one primary reason for this, whether an index is lower or risen does not carry any information, since this information comes from the tensor itself. Also, for something like complex Lorentz tensors, there are at three different types of upper-index.

If we want to evaluate an index we can put an explicit index in place of μ or ν above, for example

```
\{T \mid 1 \mid v\}^T eval 0 1 (tensorNode T)
```

The syntax and elaboration for negation, scalar multiplication and the group action are fairly similar. For negation we have:

$$\{\mathtt{T} \mid \mu \ v\}^{\mathtt{T}} \mid \mathtt{neg} \ (\mathtt{tensorNode} \ \mathtt{T})$$

For scalar multiplication by $a \in k$ we have:

$$\{\mathtt{a}\cdot_t\mathtt{T}\mid \mu\ v\}^{\mathrm{T}}$$
 smul a (tensorNode T)

For the group action of $g \in G$ on a tensor T we have:

$$\{g\cdot_a T\mid \mu\ v\}^T \mid ext{action } g ext{ (tensorNode T)}$$

The product of two tensors is also fairly simple, with us having:

```
\{ 	extstyle 	e
```

The syntax for contraction is as one expect

```
\{T \mid \mu \ v \otimes T2 \mid v \ \sigma\}^T \mid \text{contr 1 1 rfl (prod (tensorNode T) (tensorNode T2))}
```

On the RHS here the first argument (1) of contr is the index of the first v on the LHS, the second argument (also 1) is the second index. The rfl is a proof that the colors of the two contracted indices are actually dual to one another. If they are not, this proof will fail and the elaborator will complain. It will also complain if more than two indices are trying to be contracted. Although, this depends on where exactly the indices sit in the expression. For example

```
\{T \mid \mu \ v \otimes T2 \mid v \ v\}^T (prod (tensorNode T) (contr 0 0 rfl (tensorNode T2)))
```

works fine because the inner contraction is done before the product.

We now turn to addition. Our syntax allows for e.g., $\{T \mid \mu \ v + T2 \mid \mu \ v\}^T$ and also $\{T \mid \mu \ v + T2 \mid \nu \ \mu\}^T$, provided of course that the indices are of the correct color (which Lean will check). The elaborator handles both these cases, and generalizations thereof by adding a permutation node. Thus we have

```
\{T \mid \mu \mid v + T2 \mid \mu \mid v\}^T \mid \text{add (tensorNode T) (perm _ (tensorNode T2))}
```

where here the _ is a placeholder for the permutation. For this case will the permutation will be the identity, but for

$$\{T \mid \mu \mid v + T2 \mid v \mid \mu\}^T \mid \text{add (tensorNode T) (perm _ (tensorNode T2))}$$

it will be the permutation for the two identities.

Despite not forming part of a node in our tensor tree, we also give syntax for equality. This is done in a very similar way to addition, with the addition of a permutation node to account for e.g., the fact that $T_{\mu\nu} = T_{\nu\mu}$.

```
\{T \mid \mu \ v = T2 \mid v \ \mu\}^T (tensorNode T).tensor = (perm _ (tensorNode T2)).tensor
```

Note the use of the .tensor to extract the tensor from the tensor tree, it does not really mean much to ask for equality of the tensor trees themselves.

With this syntax we can write complicated tensor expressions in a way close to pen-and-paper index notation. We will see examples of this in the next section.

3. EXAMPLES

We give two examples in this section. The first example is a simple theorem involving index notation and tensor trees. We will show here, in rather explicit detail, how we can manipulate tensor trees to solve such theorems. The second example we shall give will show a number of definitions in HepLean concerning index notation. Here we will not give so much detail, the point rather being to show the reader the broad use of our construction.

3.1. EXAMPLE 1: SYMMETRIC AND ANTISYMMETRIC TENSOR

If $A^{\mu\nu}$ is an antisymmetric tensor and $S_{\mu\nu}$ and S is a symmetric tensor, then it is true that $A^{\mu\nu}S_{\mu\nu} = -A^{\mu\nu}S_{\mu\nu}$. In Lean this result, and it's proof are written as follows:

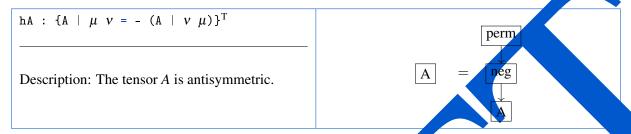
```
lemma antiSymm_contr_symm
    {A : complexLorentzTensor.F.obj (QverColor.mk ![Color.up, Color.up])}
    {S : complexLorentzTensor.F.obj (OverColor.mk ! [Color.down, Color.down])}
    (hA : \{A \mid \mu \ v = - \ (A \mid v \ \mu)\}^T) (hs
                                               {S | 🖊
    \{A \mid \mu \ v \otimes S \mid \mu \ v = - \ A \mid \mu \ v \otimes S
  conv =>
    lhs
    rw [contr_tensor_eq < contr_tensor_eq < prod_tensor_eq_fst < | hA]
    rw [contr_tensor_eq < contr_tensor_eq < | prod_tensor_eq_snd < | hs]</pre>
        [contr_tensor_eq < contr_tensor_eq < prod_perm_left _ _ _ _]</pre>
       [contr_tensor_eq < | contr_tensor_eq < | perm_tensor_eq < | prod_perm_right _ _
    rw [contr_tensor_eq <| contr_tensor_eq <| perm_perm _ _ _ _]</pre>
        [contr_tensor_eq < | perm_contr_congr 1 2]</pre>
    rw [perm_contr_congr 0 0]
    rw [perm_tensor_eq < contr_contr _ _ _]</pre>
    rw [perm_perm]
    rw [perm_tensor_eq < | contr_tensor_eq < | contr_tensor_eq < | neg_fst_prod _ _]
    rw [perm_tensor_eq <| contr_tensor_eq <| neg_contr _]</pre>
    rw [perm_tensor_eq <| neg_contr _]</pre>
  apply perm_congr _ rfl
  decide
```

Let us break this down. The statements

```
{A : complexLorentzTensor.F.obj (OverColor.mk ![Color.up, Color.up])}
{S : complexLorentzTensor.F.obj (OverColor.mk ![Color.down, Color.down])}
```

are simply defining A and S to be tensors of type $A^{\mu\nu}$ and $S_{\mu\nu}$ respectively. Here ![Color.up, Color.up] is shorthand for the map Fin 2 \rightarrow complexLorentzTensor.C that sends 0 to Color.up and 1 to Color.up.

The parameter hA is stating that A is antisymmetric. Expanded in terms of tree diagrams we have

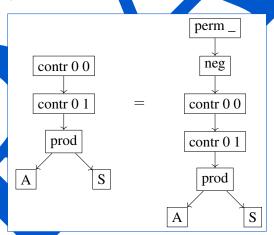


where by equality of tensor trees in the above diagram, we really mean equality of the underlying tensors of those tensor trees. This will be left implicit throughout.

Similarly, the parameter hs is stating that S is symmetric. Expanded in terms of tree diagrams



The line $\{A \mid \mu \ v \otimes S \mid \mu \ v = -A \mid \mu \ v \otimes S \mid \mu \ v\}^T$ is the statement we are trying to prove. In terms of tree diagrams it says that

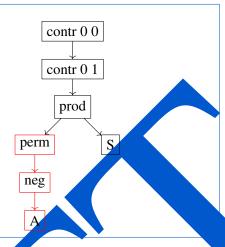


The perm here actually does nothing, but is included by Lean.

The lines of the proof in the conv block are manipulations of the tensor tree on the LHS of the equation. The rw tactic is used to rewrite the tensor tree using the various lemmas. We go through each step in turn.

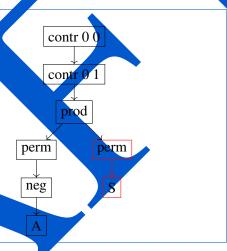
```
rw [contr_tensor_eq <| contr_tensor_eq
<| prod_tensor_eq_fst <| hA]</pre>
```

Description: Here contr_tensor_eq and prod_tensor_eq_fst navigate to the correct place in the tensor tree, whilst hA replaces the node A with the RHS of hA.



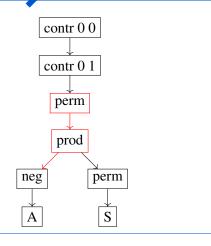
```
rw [contr_tensor_eq <| contr_tensor_eq
<| prod_tensor_eq_snd <| hs]</pre>
```

Description: Here contr_tensor_eq and prod_tensor_eq_fst navigate to the correct place in the tensor tree, whilst hS replaces the node S with the RHS of hs.



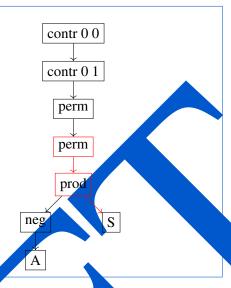
```
rw [contr_tensor_eq <| contr_tensor_eq
<| prod_perm_left _ _ _]</pre>
```

Description: Here contr_tensor_eq navigates to the correct place in the tensor tree, whilst prod_perm_left moves the permutation on the left through the product.



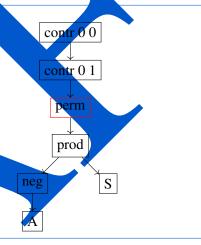
```
rw [contr_tensor_eq <| contr_tensor_eq
<| perm_tensor_eq <| prod_perm_right _ _
_ _]</pre>
```

Description: Here contr_tensor_eq and perm_tensor_eq navigate to the correct place in the tensor tree, whilst prod_perm_right moves the permutation on the right through the product.



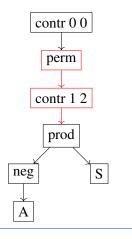
```
rw [contr_tensor_eq <| contr_tensor_eq
<| perm_perm _ _ _]</pre>
```

Description: Here contr_tensor_eq navigates to the correct place in the tensor tree, whilst perm_perm uses functoriality to combine the two permutations.



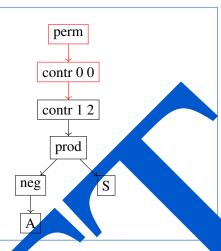
```
rw [contr_tensor_eq <| perm_contr_congr
1 2]</pre>
```

Description: Here contr_tensor_eq navigates to the correct place in the tensor tree, whilst perm_contr_congr moves the permutation through the contraction, and simplifies the contraction indices to 1 and 2 (Lean will check if this is correct).



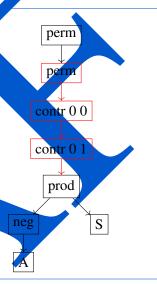
```
rw [perm_contr_congr 0 0]
```

Description: Here perm_contr_congr moves the permutation through the contraction, and simplifies the contraction indices to 0 and 0 (Lean will check if this is correct).



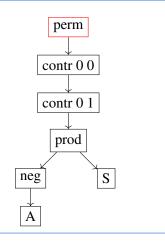
```
rw [perm_tensor_eq <| contr_contr _ _ _]</pre>
```

Description: Here perm_tensor_eq navigates to the correct place in the tensor tree, whilst contr_contr swaps the two contractions, in the meantime inducing a permutation.



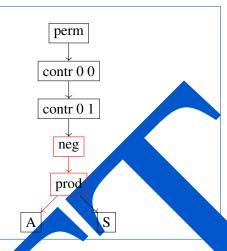
rw [perm_perm]

Description: Here perm_perm uses functoriality to combine the two permutations.



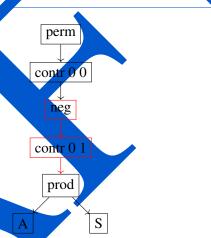
```
rw [perm_tensor_eq <| contr_tensor_eq <|
contr_tensor_eq <| neg_fst_prod _ _]</pre>
```

Description: Here perm_tensor_eq and contr_tensor_eq navigate to the correct place in the tensor tree, whilst neg_fst_prod moves the negation through the product.

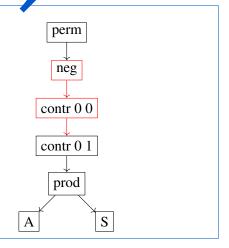


```
rw [perm_tensor_eq <| contr_tensor_eq <|
neg_contr _]</pre>
```

Description: Here perm_tensor_eq and contr_tensor_eq navigate to the correct place in the tensor tree, whilst neg_contr moves the negation through the first contraction.



Description: Here perm_tensor_eq navigates to the correct place in the tensor tree, whilst neg_contr moves the negation through the second contraction.



The remainder of the proof

```
apply perm_congr _ rfl
decide
```

helps Lean to understand that the two sides of the equation are equal.

3.2. EXAMPLE 2: PAULI MATRICES AND BISPINORS

Using the formlism we have set up thus far it is possible to define Pauli matrices and bispinors as complex Lorentz tensors.

The Pauli matrices appear in HepLean as follows:

```
/-- The Pauli matrices as the complex Lorentz tensor '\sigma^{\mu}\alpha^{\alpha}{dot \beta}'. -/ def pauliContr := {PauliMatrix.asConsTensor | v \alpha \beta}<sup>T</sup>.tensor /-- The Pauli matrices as the complex Lorentz tensor '\sigma_{\mu}\alpha^{\alpha}{dot \beta}'. -/ def pauliCo := {\eta' | \mu v \otimes \text{pauliContr} | v \alpha \beta}<sup>T</sup>.tensor /-- The Pauli matrices as the complex Lorentz tensor '\sigma_{\mu}\alpha^{\alpha}{dot \beta}' -/ def pauliCoDown := {pauliCo | \mu \alpha \beta \otimes \epsilon L' | \alpha \alpha' \otimes \epsilon R' | \beta \beta'}<sup>T</sup>.tensor /-- The Pauli matrices as the complex Lorentz tensor '\sigma^{\mu}\alpha^{\alpha} \epsilon R' | \epsilon R" | \epsilon R' | \epsilon R
```

The first of these definitions depends on PauliMatrix.asConsTensor which is defined using an explcit basis expansion as a map $\mathbb{1}_{-}$ (Rep \mathbb{C} SL(2, \mathbb{C})) \to complexContr \otimes Fermion.leftHanded \otimes Fermion.rightHanded, which Lean knows how to treat as a tensor. Here complexContr is the representation of complex Lorentz vectors under $SL(2,\mathbb{C})$.

In these expressions we have the appearence of metrics η , is the metric usually written as $\eta_{\mu\nu}$, η the metric $\eta^{\mu\nu}$, εL , the metric $\varepsilon_{\alpha\alpha'}$, and εR , the metric $\varepsilon_{\dot{R}\dot{B}'}$.

With these we can also define bispinors

```
/-- A bispinor 'p<sup>aa</sup>' created from a lorentz vector 'p'\mu'. -/ def contrBispinorUp (p : complexContr) := {pauliCo | \mu \alpha \beta \otimes p | \mu}<sup>T</sup>.tensor

/-- A bispinor '\alpha' created from a lorentz vector 'p°\mu'. -/ def contrBispinorDown (p : complexContr) := {\epsilonL' | \alpha \alpha' \otimes \epsilonR' | \beta \beta' \otimes contrBispinorUp p | \alpha \beta}<sup>T</sup>.tensor

/-- A bispinor \alpha<sup>aa</sup>' created from a lorentz vector 'p_{-}\mu'. -/ def coBispinorUp (p : complexCo) := {pauliContr | \mu \alpha \beta \otimes p | \mu}<sup>T</sup>.tensor

# Siminor '\alpha<sup>aa</sup>' complexCo := {\alpha<sup>a</sup> complexCo := {\alpha
```

Here complexCo corresponds to the representation of covariant complex Lorentz vectors.

Using these definitions we can start to prove results about the pauli matrices and bispinors. These proofs rely on essentially the sorts of manipulations in the last section, although in some cases we expand tensors in terms of a basis and use rules about how the basis interacts with the operations in a tensor tree.

As an example, we prove the following lemmas

```
lemma coBispinorDown_eq_pauliContrDown_contr (p : complexCo) : {coBispinorDown p | \alpha \beta = pauliContrDown | \mu \alpha \beta \otimes p | \mu}^T := by
```

```
/-- The statement that '\eta_{\mu\nu} of \mu and of \beta of \nu a' dot \beta' = 2 efaa' efact \beta dot \beta'. -/

theorem pauliCo_contr_pauliContr:

{pauliCo | \nu a \beta \otimes pauliContr | \nu a' \beta' = 2 · \nu EL | a a' \nu ER | \nu by
```

We do not give the proofs of these lemmas explicitly. The former, however, is a fairly simple application of associativity of the tensor product, and appropriatly shuffling around of the contractions.

4. FUTURE WORK

The scale of formalizing all results regarding index notation is a task that surpasses the capacity of any single individual. Inspired by the Lean community's blueprint projects, we have added to HepLean informal lemmas related to index notation and tensors. An example of such is

```
informal_lemma coBispinorUp_eq_metric_contr_coBispinorDown where math := "{coBispinorUp p | \alpha \beta = \varepsilonL | \alpha \alpha' \otimes \varepsilonR | \beta \beta' \otimes coBispinorDown p | \alpha , \beta', }<sup>T</sup>" proof := "Expand 'coBispinorDown' and use fact that metrics contract to the identity." deps := [''coBispinorUp, ''coBispinorDown, ''leftMetric, ''rightMetric]
```

These informal lemmas are written in strings, and are not type checked. They also include dependencies indicating what definitions and lemmas we expect to be used in their statement or proof. They are intended to be a guide for future formalization efforts, either by humans or automated.

As demonstrated in our earlier examples, manipulating tensor expressions can involve tedious calculations, especially when dealing with directly with tensor trees. In the future, we intend to automate many of these routine steps by developing suitable tactics within Lean. We are optimistic that the structured nature of tensor trees will lend itself well to such automation, thereby streamlining computations and enhancing the efficiency of formal proofs involving index notation and tensor species.

There are two primary directions in which we can extend the concepts presented in this work. First, we could incorporate the spinor-helicity formalism, which is used in the study of scattering amplitudes. Second, we could extend our approach to encompass tensor *fields*, their derivatives etc. We do not anticipate any insurmountable challenges in pursuing these extensions. They represent promising avenues for future research and have the potential to significantly enhance the utility of formal methods in physics. sdfljk

ACKNOWLEDGMENTS

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