Index notation in Lean 4

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Abstract

The physics community relies on index notation to effectively manipulate tensors of specific types. This paper introduces the first formally verified implementation of index notation in the interactive theorem prover Lean 4. By integrating index notation into Lean, we bridge the gap between traditional physics notation and formal verification tools, making it more accessible for physicists to write and prove results within Lean. In the background, our implementation leverages a novel application of category theory.

1. INTRODUCTION

HepLean In previous work, the author initiated the digitalization (or formalization) of high energy physics results using the interactive theorem prover Lean 4 in a project called HepLean. Lean is an interactive theorem prover and programming language with syntax resembling traditional pen-and-paper mathematics. Lean allows users to write definitions, theorems, and proofs which are then automatically checked for correctness, using its foundation of dependent-type theory. HepLean has four main motivations: A linear storage of information makes look-up of results easier; Allows for the creation and proof of new results using automated tactics like AI; Makes it easier to check the correctness of results; And allows for new ways high-energy physics and computer science can be taught.

Other works in Lean HepLean is part of a broader movement of projects to formalize parts, or all, of mathematics and science. The largest of these projects is Mathlib, which aims to formalize mathematics. Indeed, HepLean is built downstream of Mathlib, meaning it has Mathlib as a dependency and uses many of the definitions and theorems from there. Other projects include the ongoing effort led by Kevin Buzzard to formalize the proof of Fermat's Last Theorem into Lean. In the realm of the sciences, js: sorry

Notation in physics and index notation: Physicists heavily rely on specialized notation to express complex concepts succinctly. Among these, index notation is particularly prevalent, as it provides a compact and readable way to represent specfic types of tensors and operations between them.

Challange: Formalizing index notation into an interactive theorem prover like Lean is chalanging due to the complexity and implicitness of the notation and the need for flexiable and rigour in it's formalization. This paper presents such an implementation of index notation within the HepLean project.

Motivation: The primary motivations for implementing index notation into Lean are:

1. To make easier to write and prove results from high energy physics into Lean.
2. To make the syntax of Lean more familiar to high energy physicists use to index notation.

We hope that the implementation of index notation into Lean not only enhances usability but also promotes the adoption of formal methods in the physics community.

Conclusion: The conclusion of the implentation of index notation into Lean 4 is that we can write results like the following:

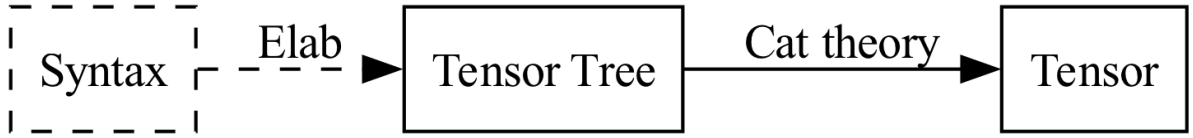
$$js : Addexample \tag{1}$$

and lean will correctly intpret this result as a formal mathematical object which proves can be applied to.

Other papers: Previous attempts to formalize index notation have been made in programming languages like Haskell [**?**]. However, these implementations do not provide the formal verification capabilities inherent in Lean. The formal verification requirement of Lean introduces unique challenges in implementing index notation, necessitating (what we beleive is) a novel solution.

Outline: This paper is split into two main sections. Section js: ref dicusses the implementation of index notation into Lean. Section js: ref gives two examples of theorems and proofs using index notation. A more minor section, section js: ref disucsses the future work related to this project.

2. IMPLEMENTATION OF INDEX NOTATION INTO LEAN 4



The implementation of index notation in Lean can be broken down into three main components, illustrated in Figure js: ref. The first component is the *syntax for tensor expressions*, which is what users interact with when writing results in Lean. This syntax closely mirrors the notation familiar to physicists, making it intuitive and accessible. It appears directly in the Lean files and can be thought of as an informal string that represents the tensor expressions.

The second component involves transforming this syntax into a *tensor tree*. The tensor tree is a formal mathematical representation of the tensor expression. By parsing the informal syntax into a structured tree, we establish a rigorous foundation that captures the tensor expression. This formal representation allows us to easily manipulate tensor expressions and prove results related to them in a way that Lean accepts as formal.

The third and final component is the conversion of the tensor tree into an actual *tensor*. This process utilizes properties of the symmetric-monoidal category of representations to translate the abstract tensor tree into a concrete tensor.

These three steps—syntax, tensor tree, and tensor—are illustrated in Figure js: ref. Although Lean processes information from left to right, starting with the syntax and proceeding to the tensor tree, and then finally (when we ask it to) to the tensors, it is more effective to discuss the implementation from right to left. This reverse approach is advantageous because the tensors, which are the primary objects of interest, are located on the right side of the diagram. The left and middle parts of the diagram represent intermediate stages that facilitate the manipulation and understanding of these tensors and their expressions.

2.1.1 Building blocks of tensors and color

Tensors of a species, such as complex Lorentz tensors, are constructed from a set of building block representations of a group G over a field k. For complex Lorentz tensors, the group G is $SL(2,\mathbb{C})$, the field k is the field of complex numbers.

There are six building block representations for complex Lorentz tensors. These are

- `Fermion.leftHanded` is the representation of $SL(2,\mathbb{C})$ corresponding $v \mapsto Mv$, corresponding to Left-handed Weyl fermions.
- `Fermion.altLeftHanded` is the representation of $SL(2,\mathbb{C})$ corresponding $v \mapsto M^{-1T}v$, corresponding to alternative Left-handed Weyl fermions (as we will call them).
- `Fermion.rightHanded` is the representation of $SL(2,\mathbb{C})$ corresponding $v \mapsto M^{*}v$, corresponding Right-handed Weyl fermions.
- `Fermion.altRightHanded` is the representation of $SL(2,\mathbb{C})$ corresponding $v \mapsto M^{-1\dagger}v$, corresponding to alternative Right-handed Weyl fermions.
- `Lorentz.complexContr` is the representation of $SL(2,\mathbb{C})$ induced by the homomorphism of $SL(2,\mathbb{C})$ into the Lorentz group and the contravariant action of the Lorentz group on four-vectors
- `Lorentz.complexCo` is the representation of $SL(2,\mathbb{C})$ induced by the homomorphism of $SL(2,\mathbb{C})$ into the Lorentz group and the covariant action of the Lorentz group on four-vectors.

As an example the representation `Fermion.leftHanded` is defined in Lean as follows:

```
def leftHanded : Rep ℂ SL(2,ℂ) := Rep.of {
    toFun := fun M => {
      toFun := fun (ψ : LeftHandedModule) =>
        LeftHandedModule.toFin2ℂEquiv.symm (M.1 * ψ.toFin2ℂ),
      map_add' := by
        intro ψ ψ'
        simp [mulVec_add]
      map_smul' := by
        intro r ψ
        simp [mulVec_smul]}
    map_one' := by
      ext i
      simp
    map_mul' := fun M N => by
      simp only [SpecialLinearGroup.coe_mul]
      ext1 x
      simp only [LinearMap.coe_mk, AddHom.coe_mk, LinearMap.mul_apply,
    LinearEquiv.apply_symm_apply,
        mulVec_mulVec]}
```

js: Not exact definiton as could not type $*_v$ Here the outer `toFun` argument is defining a map from $SL(2,\mathbb{C})$ to linear maps from `LeftHandedModule` (equivalent to $\mathbb{C}^2$) to itself. The inner `toFun`, `map_add'`

and `map_smul'` are defining defining respectively, the underlying function of this linear map, and proving it it linear with respect to addition and scalar multiplication. The `map_one'` and `map_mul'` are proving that action of the identity is trivial, and that the action of the product of two elements is the product of the actions of the elements.

We assign a unique label, which we refer to as a *color*, to each of the building block representations. We denote the type of colors for a given species of tensors as C. For complex Lorentz tensors C is defined as

```
inductive Color
  | upL : Color
  | downL : Color
  | upR : Color
  | downR : Color
  | up : Color
  | down : Color
```

Here `Color` is the name of our type and `upL`, `downL`, etc. are the colors.

To formally associate each color with its corresponding representation, we define a discrete functor from the set C to the category of k-representations of G, denoted RepkG, that is, a functor

$$F_D : C \Rightarrow \text{Rep}kG. \tag{2}$$

For complex Lorentz tensors this functor in Lean as follows:

```
FDiscrete := Discrete.functor fun c =>
  match c with
  | Color.upL => Fermion.leftHanded
  | Color.downL => Fermion.altLeftHanded
  | Color.upR => Fermion.rightHanded
  | Color.downR => Fermion.altRightHanded
  | Color.up => Lorentz.complexContr
  | Color.down => Lorentz.complexCo
```

The reason why we have used a functor here, rather than just a function, will become apparent in what follows.

2.1.2 A general tensor

For a given tensor species, given the symmetric monoidal structure on RepkG, we can take the tensor product of the building block representations. Elements of such tensor products form the general notion of a tensor for that given species.

To formalize this in Lean, we consider the category $\mathscr{S}^{\times}_{/C}$. The objects of $\mathscr{S}^{\times}_{/C}$ are functions $f : X \to C$ for some type X. A morphism from $f : X \to C$ to $g : Y \to C$ is a bijection $\phi : X \to Y$ such that $f = g \circ \phi$. This category is equivalent to the core of the category of types ($\int$) over C, hence our notation.

Any functor $H : C \Rightarrow \text{Rep}kG$ can be lifted to a symmetric monoidal functor $\mathscr{S}^{\times}_{/C} \Rightarrow \text{Rep}kG$ which takes f to the tensor product $\bigotimes_{x \in X} H(f(x))$ and morphisms to the linear maps of representations corresponding to reindexings of tensor products.

This construction is itself functorial, giving a functor:

$$\text{lift} : \text{Fun}(C, \text{Rep}kG) \Rightarrow \text{SymmMonFun}(\mathscr{S}^{\times}_{/C}, \text{Rep}kG) \tag{3}$$

In the previous subsection we defined the functor F_D which associates to each color its corresponding representation. We can then define a functor F by the image of F_D under the lift functor.

A tensor of a given species can then be thought of as a vector in the representations in the of F. Defining a tensor in this way allows us to utalize the structur of monodial categories and functors in a useful way.

In HepLean we formally define lift and F.

For most purposes in physics X will be a type Fin$n$, corresponding to the type (set) of natural numbers (including 0) less n. Later on we will restrict to these types.

2.1.3 Basic operations

Now we have discussed how tensors of a given species can be formally defined, we can define basic operations on these tensors.

The simipalist of these operations are addition, scalar multiplication and group action. Addition and scalar multiplication are given to us for free from the vector space structure on $F(f)$. The action of the group G on tensors also comes for free from the fact that $F(f)$ lives in RepkG.

The next simplest operation is permutation of indices (or building block representations). Such permutations arise from applying F to morphisms in $\mathscr{S}^{\times}_{/C}$. Theses sorts of permutations are often implicit in the notation physicists use, but arise in tensor expressions such as $T_{\mu\nu} = T_{\nu\mu}$.

The next operation is the tensor product of two tensors. Given $f : X \to C$ and $g : Y \to C$, and tensors $t \in F(f)$ and $s \in F(g)$ we can form a vector in $F(f) \otimes F(g)$. Formally this can be done by taking the tensor product of the morphisms $\mathbb{1} \to F(f)$ and $\mathbb{1} \to F(g)$ in the category of vector spaces over k. We can then use the tensorate of F to get a tensor in $F(f \otimes g)$.

We now turn to the slightly more complicated operation of contraction of indices, and the related metric and unit. To define the contraction of introduce an involution $\tau : C \to C$. For complex Lorentz tensors this is defined as follows

```
τ := fun c =>
  match c with
  | Color.upL => Color.downL
  | Color.downL => Color.upL
  | Color.upR => Color.downR
  | Color.downR => Color.upR
  | Color.up => Color.down
  | Color.down => Color.up
```

We say that τ takes a color to it's dual. With τ and F_D we can define the functor $F_\tau : C \Rightarrow RepkG$ which takes c to $F(c) \otimes F(\tau c)$. Letting $\mathbb{1}$ be the constant functor from C landing on the, we basic data for contraction is contained in a natural transformation from F_τ to $\mathbb{1}$. That is, for each $c \in C$ an equivariant map from $F(c) \otimes F(\tau c)$ to the trivial representation.

To see how this is can be used to contract indices, consider a tensor $t \in F(f)$, for $f : \text{Fin}\, n.succ.succ \to C$. Here Fin $n.succ.succ$ is the type of all number $0, 1, 2, \ldots, n+1$ less than $n.succ.succ = n+2$. Choosing

two distinct indices $i, j : \mathrm{Fin}\ n.succ.succ$ we can contract them using the following chain of maps

$$F f \equiv F_\tau(fi) \otimes F(f'i) \to \mathbb{1} \otimes F(f'i) \to F(f'i) \tag{4}$$

Here $f' : \mathrm{Fin}\ n \to C$ is the function defived from f by removing the indices i and j. The first equivalence is somewhat complicated to define formally, but essentially involves extracting i and j.

The metric and unit are defined in a similar way. The metric is a natural transformation from $\mathbb{1}$ to $F_D \otimes F_D$. That is, for each $c \in C$ an equivariant map from the trivial representation to $F_D(c) \otimes F_D(c)$. The unit is a natural transformation from $\mathbb{1}$ to F'_τ which takes c to $F_D(\tau c) \otimes F_D(c)$. That is, for each $c \in C$ an equivariant map from the trivial representation to $F_D(\tau c) \otimes F_D(c)$.

The choice of contraction, metric and unit are subject to a number of conditions.

- Contraction must be symmetric.

$$\begin{array}{ccc} F_D(c) \otimes F(\tau c) & \longrightarrow & F_D(\tau c) \otimes F_D(c) \\ \downarrow & & \downarrow \\ \mathbb{I} & \longrightarrow & F_D(\tau c) \otimes F_D(\tau\tau c) \end{array} \tag{5}$$

- Contraction with the unit does nothing.

$$\begin{array}{ccccc} F_D(c) & \longrightarrow & F_D(c) \otimes \mathbb{I} & \longrightarrow & F_D(c) \otimes (F_D(\tau c) \otimes F_D(c)) \\ \uparrow & & & & \downarrow \\ \mathbb{I} \otimes F_D(c) & & \longleftarrow & & (F_D(c) \otimes F_D(\tau c)) \otimes F_D(c) \end{array} \tag{6}$$

- The unit is symmetric

$$\begin{array}{ccc} \mathbb{I} & \longrightarrow & F_D(\tau c) \otimes F_D(c) \\ \downarrow & & \uparrow \\ F_D(\tau\tau c) \otimes F_D(\tau c) & \longrightarrow & F_D(\tau c) \otimes F_D(\tau\tau c) \end{array} \tag{7}$$

- Contraction with the metric its dual returns the unit

$$\begin{array}{ccccc} \mathbb{I} & & \longrightarrow & & F_D(\tau c) \otimes F_D(c) \\ \downarrow & & & & \uparrow \\ (F_D c \otimes F_d c) \otimes (F_D(\tau c) \otimes F_D(\tau c)) & & & & F_D(c) \otimes F_D(\tau c) \\ \downarrow & & & & \uparrow \\ F_D c \otimes (F_d c \otimes (F_D(\tau c) \otimes F_D(\tau c))) & & \longrightarrow & & F_D c \otimes ((F_d c \otimes F_D(\tau c)) \otimes F_D(\tau c)) \end{array} \tag{8}$$

In Lean we quite write these conditions in this way, instead we give them in a simpiler equivalent form in which we apply the maps in the above diagram to pure tensors. As an example, the condition is written in Lean as

The last operation we want to talk about is evaluation. Where one specifies the exact value of one of the indices of a tensor e.g. η_{0i}. All of the operations disucsed thus far have related to the category RepkG. As such, they have respect the group action. Evaluation, on the other hand, does not respect the group

action. For example, taking the 0th component of a four-vector is not Lorentz invariant. Nevertheless, we can define evaluation in the categor of vector spaces. To do this, we need a basis for each $F_D(c)$, the coordinates of a vector in $F_D(c)$ in that basis is discrbed by a linear map of vector spaces from $F_D(c)$ to k. This can be used to evaluate the index of a tensor in the following way:

$$Ff \equiv F_D(fi) \otimes F(f'i) \to \nVdash \otimes F(f'i) \to F(f'i) \tag{9}$$

where f' is f with the ith index removed, and again the first equivalence is somewhat complicated to define but involves extracting the ith index.

2.1.4 Tensor Species

- The data we have given so far consitutes what we have losely being calling a Tensor species.
- In Lean we make this more precise

2.2. TENSOR TREES

2.2.1 Structure

- A tensor expression consists of a series of tensors and operations performed on them.
- We can represent such an expression using a tensor tree, similar to the notion of a syntax tree.
- A tensor tree has different types of nodes either representing a tensor or a operation on or between tensors.
- Since we really only care about tensors with $X = Finn$, tensor trees in Lean are implemented only for these.
- Let us give the definition of a tensor tree and then dicuss in turn each of the nodes.
- js: Lean defn of tensor tree
- The basic node is a tensor node. The definition in Lean tells us how this is defined.
- Then for each operation addition, scalar mult, group action etc. we get a node.
- For example let us look at contraction
- js: sorry
- How we turn these tensor trees into tensors will be dicussed in the next section

2.2.2 To a tensor

- The notion of a tensor tree is defined without reference to the category theory we have been discussing.
- We can however turn a tensor tree into a tensor using said constructions.
- The definition of how we do this is defined recursively.

2.2.3 Using Tensor trees in proofs

- js: Discuss fact about 'tensor_eq'

2.3. ELBORATION

- We have dicussed the implementation of tensor species into Lean, and how to write tensor expressions using tensor trees.
- Really this is all one needs to effectively.
- However, we want our theorems to look like they do on paper.
- This is the role of the elaborator, which here takes a string written in lean code and turns it into a tensor tree.
- It is perhpase easist to give examples:
- The basic notation for a tensor nodes is js: sorry. Note that ... are free indices, it does not matter what we call them the expression is the same.
- The product of two tensors is written as js: sorry.
- The contraction of two tensors is written as js: sorry. Again note that the indices are free so we can call them anything without chaning how lean reads the expression.
- We also define a special notation of equality and addition. which takes account of permutation.
- This part of the Lean code is not formally verified, it is just telling lean how to read the notation. Once the tensor tree is created and we start using that, things are formally verifeid.

3. EXAMPLES

- We give two examples of theorems and proves related to index notation.
- Our first example is related to symmetric and anti-symmetric tensors.
- For this example we will go into explicit detail.
- Our second example is a more detailed one, where we prove that ...
- For this example we will only give a sketch of the prove, and discuss how things are done.

3.1. EXAMPLE 1: SYMMETRIC AND ANTI-SYMMETRIC TENSOR

3.2. EXAMPLE 2: CONTRACTING PAULI MATRICES

3.3. EXAMPLE 3: DEFINITIONS OF BISPINORS

4. FUTURE WORK

- Informal lemmas and definitions.
- Improvement of tactics.

- spinor-hleicity formalism.
- Tensor fields and derivatives.

REFERENCES