



LEIDEN UNIVERSITY

Study of BCG-Subtracted Images of Nearby Clusters

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“Not only is the Universe stranger than we think, it is stranger than we can think..”

Werner Heisenberg

Abstract

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(This is just a simple draft taken from the original idea) We have obtained deep imaging data for a sample of low redshift massive clusters. The light from the BGC overwhelms the images from background galaxies and faint cluster members in the cluster core, and needs to be carefully subtracted. This is expected to reveal background galaxies that are strongly lensed. Identifying such systems allows for unique follow-up studies. Also the number density of faint cluster members may tell us something about the dynamical state of the cluster and how BGCs form. The aim of this project is to model the BCG light and search for strong lensing candidates and study the properties of faint cluster members in the core..

Acknowledgements

I would like to thank ...

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*Dedicated to my parents, whose love and support are my biggest
motivation. . .*

Chapter 1

Introduction

Stellar mass to light ratio and Stellar populations in the BCGs (ac'a es donde se deben hacer las citaciones mas importantes)

-Galaxy Clusters-

IMF is a very fundamental and important quantity in the study of stellar systems because it constraints the physics of star formation but also because it allows us to infer stellar masses through observed luminosities.

the correct use of an IMF in the context of gravitational lensing on massive objects like early type galaxies in galaxy clusters can help us constraint the amount of stellar mass and thus also infer the amount of dark matter in these systems.

Studying the amount of dark matter contribution, one could in principle make a good estimation of the stellar mass to light ratio.

For galaxies that are far away, it is impossible to make star counts, for this reason, the mass to light ratio of the stellar population provides a simple constraint on the IMF (Russell J. Smith and John R. Lucey)

Strong gravitational lensing of background galaxies provides a useful method to determine masses in elliptical galaxies, since it is difficult to constraint the IMF via M/L

massive galaxies - salpeter is a good IMF

A Kroupa IMF finds a value of gamma of around 4 for the mass to light ratio. (R. J. Smith 2014)

DM fraction in comparison with the IMF

Studying the matter distribution given by strong gravitational lensing can give us information about the IMF of the BCGs

percentage of dark matter will allow me to define the IMF more precisely. I want to see what fraction of the mass, what fraction of the surface density is stars

strong lensing at different radii is useful.

if I got to certain radius I will have more dark matter, because light drops quickly.

basically find how much dark matter and how many stars are there in the profile

Chapter 2

Theoretical Framework

Tyter.

2.1 Galaxy Clusters

Glas.

dwarf stars contribute very little to the integrated light from an old stellar population (Smith 2015)

Galaxy clusters contain a population of stars gravitationally unbound to individual galaxies, yet still bound to the clusters overall gravitational potential, created by the stripping of stars from galaxies during interactions and mergers

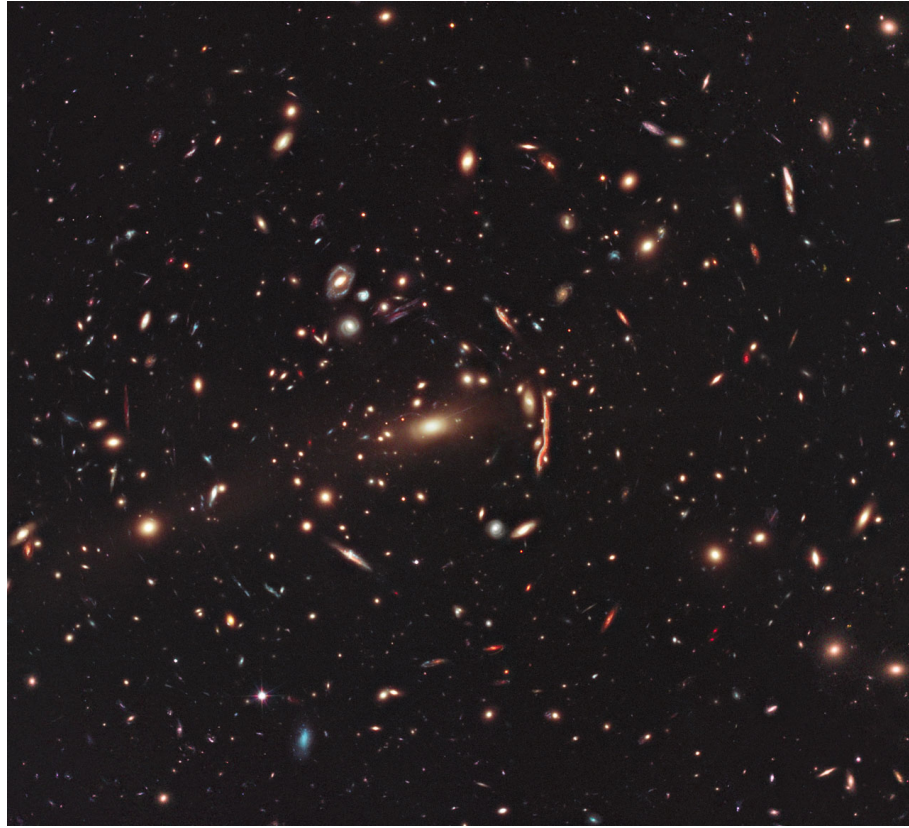


FIGURE 2.1: G

Quoted (need to change this):The image of galaxy cluster MACS J1206.2-0847 (or MACS 1206) is part of a broad survey with NASA Hubble Space Telescope. The distorted shapes in the cluster are distant galaxies from which the light is bent by the gravitational pull of an invisible material called dark matter within the cluster of galaxies. This cluster is an early target in a survey that will allow astronomers to construct the most detailed dark matter maps of more galaxy clusters than ever before. These maps are being used to test previous, but surprising, results that suggest that dark matter is more densely packed inside clusters than some models predict. This might mean that galaxy cluster assembly began earlier than commonly thought.

Scientists are planning to observe a total of 25 galaxy clusters under a project called CLASH (Cluster Lensing and Supernova survey with Hubble). One of the first objects observed for the new census is the galaxy cluster MACS J1206.2-0847. This conglomeration of galaxies is one of the most massive structures in the universe, and its gigantic gravitational pull causes stunning gravitational lensing. MACS 1206 lies 4 billion light-years from Earth. In addition to curving of light, gravitational lensing often produces double images of the same galaxy. In the new observation of cluster MACS J1206.2-0847, astronomers counted 47 multiple images of 12 newly identified galaxies. The era when the first clusters formed is not precisely known, but is estimated to be at least 9

billion years ago and possibly as far back as 12 billion years ago. If most of the clusters in the CLASH survey are found to have excessively high accumulations of dark matter in their central cores, then it may yield new clues to the early stages in the origin of structure in the universe.

2.2 Gravitational Lensing

Galaxies and clusters of galaxies that act as gravitational lenses can be approximated by single isothermal spheres. It is easy to relate an angular scaling parameter ξ_E , referred to as the Einstein radius, to the mass inside the corresponding light cone. The Einstein radius corresponds to the ring image of a point source aligned exactly on the axis of the lens.

Summary of isothermal sphere:

$$\rho(r) = \frac{\sigma^2}{2\pi G r^2} \quad (2.1)$$

$$\Sigma(\xi) = \frac{\sigma^2}{2G\xi} \quad (2.2)$$

$$\xi_E = 4\pi \left(\frac{\sigma}{c}\right)^2 \frac{D_{ds}}{D_s} \quad (2.3)$$

At small radii, stars dominate the lensing mass, so that lensing provides a direct probe of the stellar mass to light ratio, with only small corrections needed for dark matter.

In the paper of Russell Smith (a giant elliptical galaxy with a lightweight initial mass function) they find a stellar mass to light ratio of 3.01 plus minus 0.25

Modelling the lensing configuration provides the total projection mass within an aperture.

NFW profile for dark matter that is basically the dynamical mass of the cluster

let's take the case of ABELL1068, it's magnitude in U is 21.94, in I is 18.46, in g is 20.09, in r is 19.5, also $M_{200} = 4.3 \times 10^{14} M_{\odot}$

The bolometric luminosity of Abell1068 is 10^4 erg/s that in solar luminosities is $1.9 \times 10^{12} L_{\odot}$, this gives an effective brightness of $2^2 M_{\odot}/\text{kpc}^2$.

the distance to the galaxy is 591.42857 Mpc

salpeter mass function is $n(M) \propto M^{-2.3}$

de Vaucouleurs law for the surface brightness distribution in giant elliptical galaxies is:

$$I(R) = I_e e^{-b[(R/R_e)^{1/4} - 1]} \quad (2.4)$$

where $b = 7.67$ and I_e is the effective brightness which is basically the brightness at the effective radius R_e

From the paper of Lokas and Mamon, for constant mass-light-ratio we have $\Sigma_M(R) = \Gamma I(R)$ where I is the surface brightness.

the NFW density profile is

$$\rho(r) = \frac{\delta_c \rho_c}{(r/r_s)(1 + r/r_s)^2} \quad (2.5)$$

where the characteristic overdensity is:

$$\delta_c = \frac{200}{3} \frac{c^3}{\ln(1+c) - c/(1+c)} \quad (2.6)$$

from Munoz quartas et. al. we see that the concentration parameter depends on the mass and the redshift as we see in the following plot:

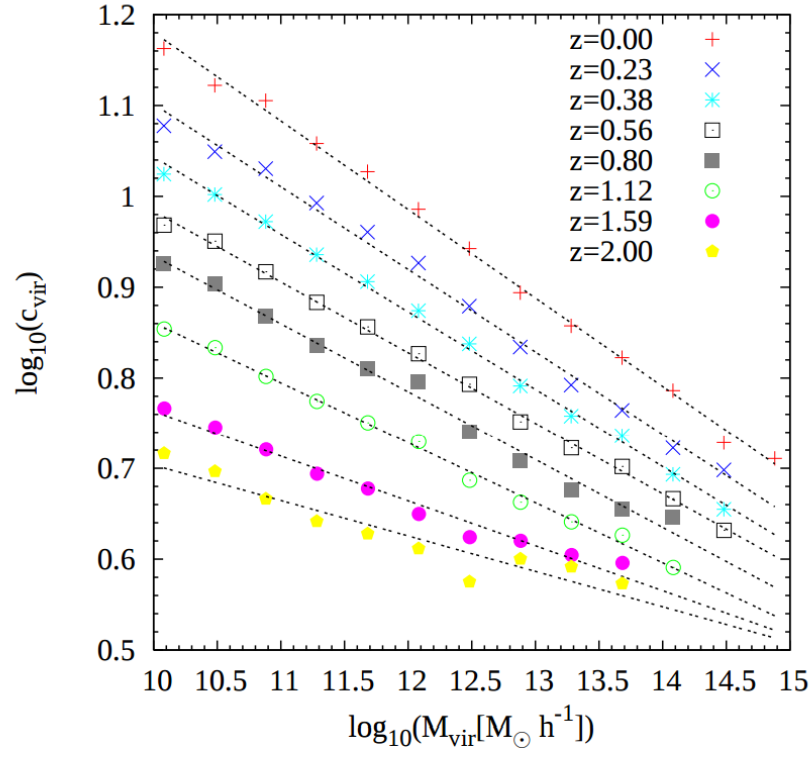


FIGURE 2.2: G

$$c = r_{200}/r_s$$

The concentration parameter c is strongly correlated with Hubble type, $c=2.6$ separating early from late-type galaxies. Those galaxies with concentration incides $c_i > 2.6$ are early-type galaxies reflecting the fact that the light is more concentrated towards their centres

From Dutton and Maccio 2014 we get the following image for the concentration parameter:

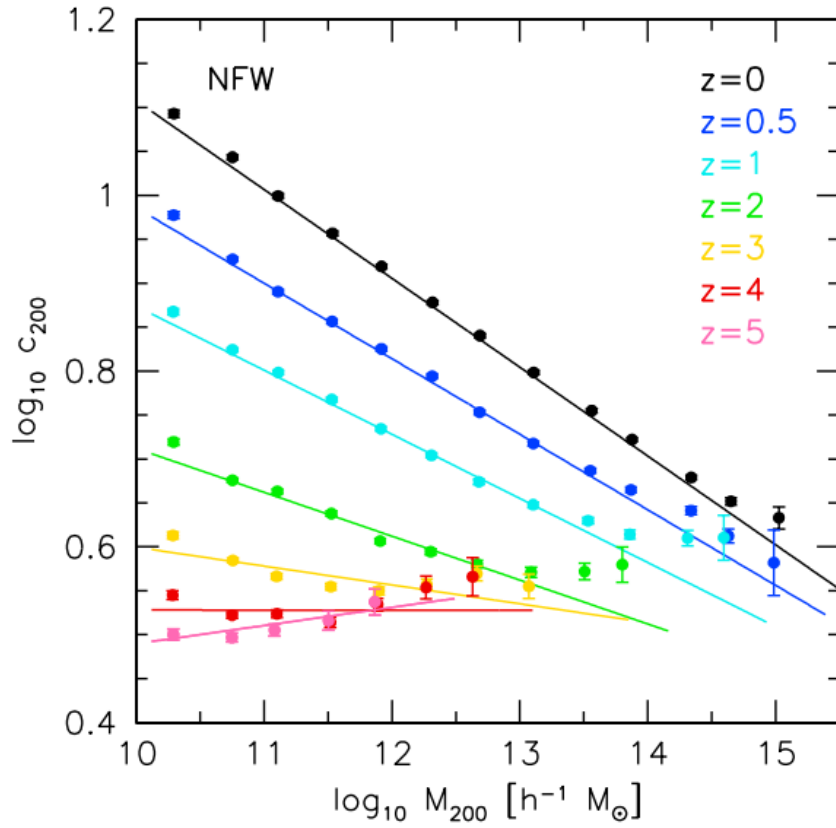


FIGURE 2.3: G

The surface mass density is given by:

$$\Sigma_{\text{NFW}}(x) = \begin{cases} \frac{2r_s\delta_c\rho_c}{(x^2-1)} \left[1 - \frac{2}{\sqrt{1-x^2}} \operatorname{arctanh} \sqrt{\frac{1-x}{1+x}} \right] (x < 1) & (x < 1) \\ \frac{2r_s\delta_c\rho_c}{3} (x = 1) & (x = 1) \\ \frac{2r_s\delta_c\rho_c}{(x^2-1)} \left[1 - \frac{2}{\sqrt{x^2-1}} \operatorname{arctan} \sqrt{\frac{x-1}{1+x}} \right] (x > 1) & (x > 1) \end{cases} \quad (2.7)$$

then the concentration parameter for ABELL1068 is about 7.9 supposing a mass of the galaxy of $10^{12.5} M_\odot$

so from the critical density:

$$\rho_c = \frac{3H^2(z)}{8\pi G} \quad (2.8)$$

the critical density would be: 2×10^{-26} in SI units so in Msol/pc^3 it is 2.9×10^{-7}

$$H(z) = H_0(1 + \Omega z)^{3/2}$$

the Hubble parameter at $z=0.138$ is $H(z)=85.6$

delta c is 25315 (dimensionless)

The characteristic radius is given by $r_{1/2} = 1.34R_e$

The mass to light ratio is $\gamma = 4$

Here the surface mass density for both the stellar content and the NFW profile:

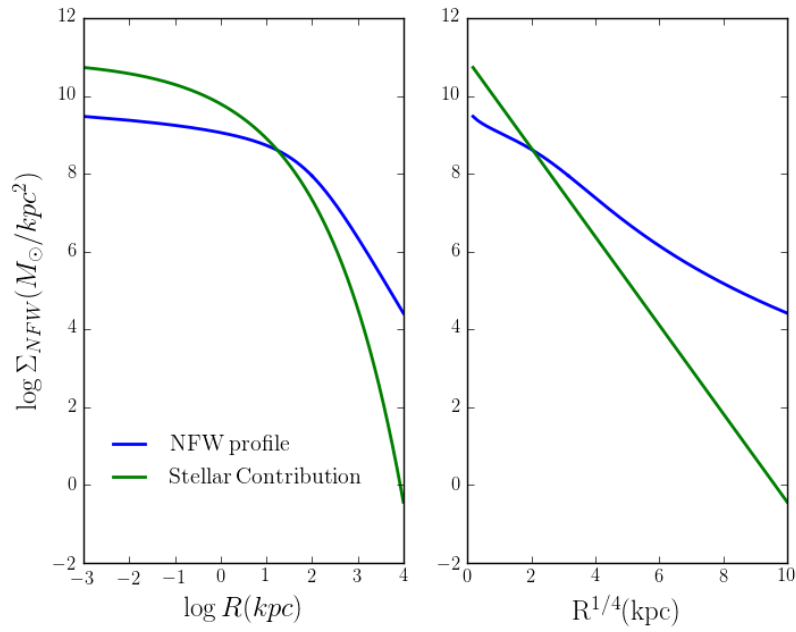


FIGURE 2.4: G

But we are more interested in the enclosed mass which can be done by integrating the surface mass density:

$$M(R) = \int_0^R 2\pi R \Sigma(R) dR \quad (2.9)$$

And we can recover our luminosity by integrating the surface brightness profile like this:

$$L = \int_0^R 2\pi R I(R) dR \quad (2.10)$$

That should give us a value comparable to the one found using Faber-Jackson relation:

$$L = C \times \sigma^4$$

And the plot for the enclosed mass is:

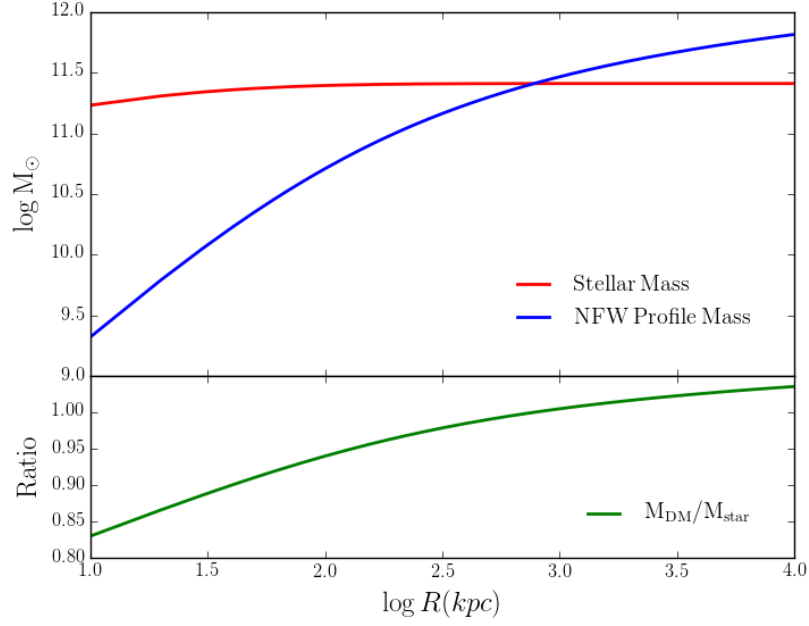


FIGURE 2.5: G

The radial dependence on the shear is:

$$\gamma_{\text{NFW}}(x) = \begin{cases} \frac{r_s \delta_c \rho_c}{\Sigma_c} g_{<}(x) & (x < 1) \\ \frac{r_s \delta_c \rho_c}{\Sigma_c} \left[\frac{10}{3} + 4 \ln \left(\frac{1}{2} \right) \right] & (x = 1) \\ \frac{r_s \delta_c \rho_c}{\Sigma_c} g_{>}(x) & (x > 1) \end{cases} \quad (2.11)$$

where:

$$g_{<}(x) = \frac{8 \operatorname{arctanh} \sqrt{\frac{1-x}{1+x}}}{x^2 \sqrt{1-x^2}} + \frac{4}{x^2} \ln \left(\frac{x}{2} \right) - \frac{2}{(x^2-1)} + \frac{4 \operatorname{arctanh} \sqrt{\frac{1-x}{1+x}}}{(x^2-1)(1-x^2)^{1/2}} \quad (2.12)$$

$$g_{>}(x) = \frac{8 \operatorname{arctan} \sqrt{\frac{x-1}{1+x}}}{x^2 \sqrt{x^2-1}} + \frac{4}{x^2} \ln \left(\frac{x}{2} \right) - \frac{2}{(x^2-1)} + \frac{4 \operatorname{arctan} \sqrt{\frac{x-1}{1+x}}}{(x^2-1)^{3/2}} \quad (2.13)$$

and with the critical surface mass density:

$$\Sigma_c \equiv \frac{c^2}{4\pi G} \frac{D_s}{D_d D_{ds}} \quad (2.14)$$

these equations come from the paper Wright and Brainerd 1999

the plot of the shear dependence on the radius is:

dependence on radius.png dependence on radius.png dependence on radius.png
dependence on radius.png dependence on radius.png dependence on radius.png
dependence on radius.png dependence on radius.png dependence on radius.png

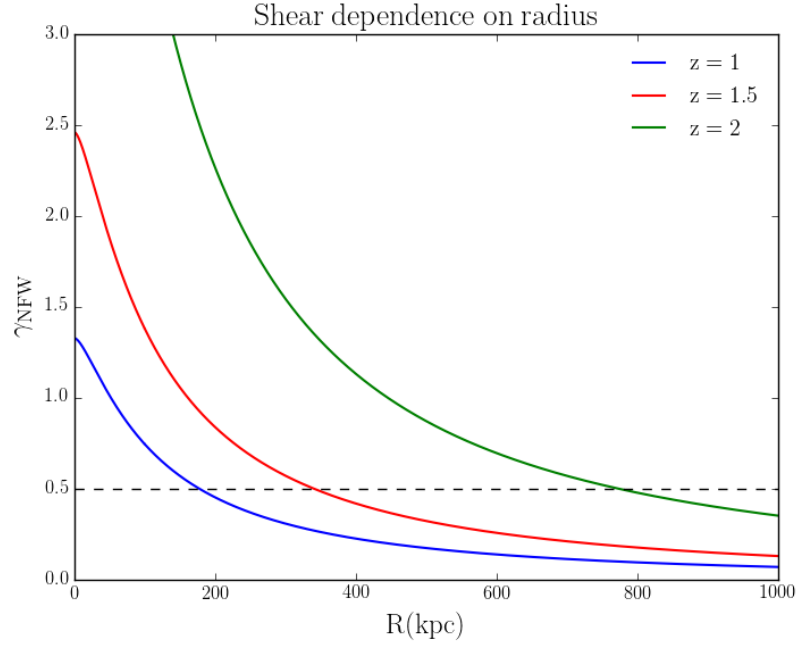


FIGURE 2.6: G

The magnification tensor is:

$$\frac{\partial \beta}{\partial \theta} = \delta_{ij} - \frac{\partial^2 \psi}{\partial \theta_i \partial \theta_j} = \begin{pmatrix} 1 - \kappa - \gamma_1 & -\gamma_2 \\ -\gamma_2 & 1 - \kappa + \gamma_1 \end{pmatrix} \quad (2.15)$$

The total magnification μ is given by the determinant of the magnification tensor:

$$\mu = \frac{1}{(1 - \kappa)^2 - \gamma_1^2 - \gamma_2^2} \quad (2.16)$$

Where κ is the convergence that determines the magnification and γ_1 and γ_2 are the shear components that determine the distortion of the background objects.

2.3 IMF in BCGs

number of stars / unit mass

Kroupa, Chabrier, Salpeter,

Heavyweight

It's difficult to see how much of the faint stars contribute to the mass of the system. We only see the new bright ones

BCG - giant ellipticals

For stars, measurements of the luminosity function can be used to derive the Initial Mass Function (IMF). For galaxies, this is more difficult because Mass to light ratio (M/L) of the stellar population depends upon the star formation history of the galaxy.

bulges have heavier IMFs than disks

Several recent studies have presented evidence for "heavyweight" IMFs in giant ellipticals, with a mass-to-light-ratio twice that of a Milky Way like IMF.

Chapter 3

Observational Procedures

the full description of the survey is in: D. J. Sand et. al. 2011

MegaCam wide field imager on the CFHT (Canada-France-Hawaii Telescope). The cluster sample consisted of 101 clusters within the range of redshifts from $0.05 < z < 0.55$

58 clusters from the MENEACs (Multi-Epoch nearby cluster survey)

The meneacs clusters represent all clusters in the BAX X-ray cluster database that are observable for the CFHT

About 60 clusters, but we used only 30 for the final studies and paid special attention to 10, marked with *

G, U, I and R images

The original omages have dimesions of [11000:11000] pixels but since our relevant region is the center of the cluster where the BCG is located, we cut the images with dimension of [1000,1000] for the color analysis and [4000:4000] to characterize the colors and discriminate between cluster and non-cluster members.

The INT images were multiple exposures so it was neccesary to make a mosaic of them using SWARP.

3.1 Sextractor

Segmentation image that will be used as a mask image (bad pixels) for galfit

We need to discriminate between field stars and the galaxies of the cluster so in order to do this, we used some of the parameters found by sextractor that allow us to constraint

Cluster	z	$\sigma(km/s)$	$d(Mpc)$	$\theta_E(^{\circ})$
A1033	0.126	762		14.6155
A1068*	0.138	740		13.5945
A1132	0.136	727		13.1515
A119*	0.044	875		21.0798
A1413*	0.143	881		19.1569
A1650	0.084	720		13.6758
A1651	0.085	903		21.4876
A1795	0.062	778		16.3514
A2029*	0.077	1152		35.2776
A2050	0.118	854		18.5258
A2055	0.102	697		12.5642
A2064	0.108	675		11.7048
A2065*	0.073	1095		32.0110
A2069	0.116	966		23.7574
A2142*	0.091	1086		30.8756
A2319*	0.056	1101		32.9563
A2420	0.085	800		16.8653
A2440	0.091	766		15.3608
A2597	0.085	682		12.2569
A2627	0.126	800		16.1096
A2703	0.114	800		16.3307
A399	0.072	800		17.1049
A553	0.066	800		17.2155
A655*	0.127	800		16.0911
A754*	0.054	800		17.4367
A763	0.085	800		16.8653
A795	0.136	800		15.9252
A85*	0.055	800		17.4182
A961	0.124	800		16.1464
A990	0.144	800		15.7778

TABLE 3.1: the redshifts of the clusters as given by C. Bildfell et. al. 2012

the fitted data. These are class_star, flux_radius, and FWHM (full width half maximum). Class_star uses the neural network star/galaxy of sExtractor that will give values close to 1 for stars and 0 for galaxies. flux_radius, and FWHM are closely related to each other and give the radius which contains half of the light of the object so it will be small for stars and bigger for extended objects.

Color magnitude diagram for A1068

we used a zero point magnitude of 30

mag.png mag.png mag.png mag.png mag.png mag.png mag.png mag.png mag.png

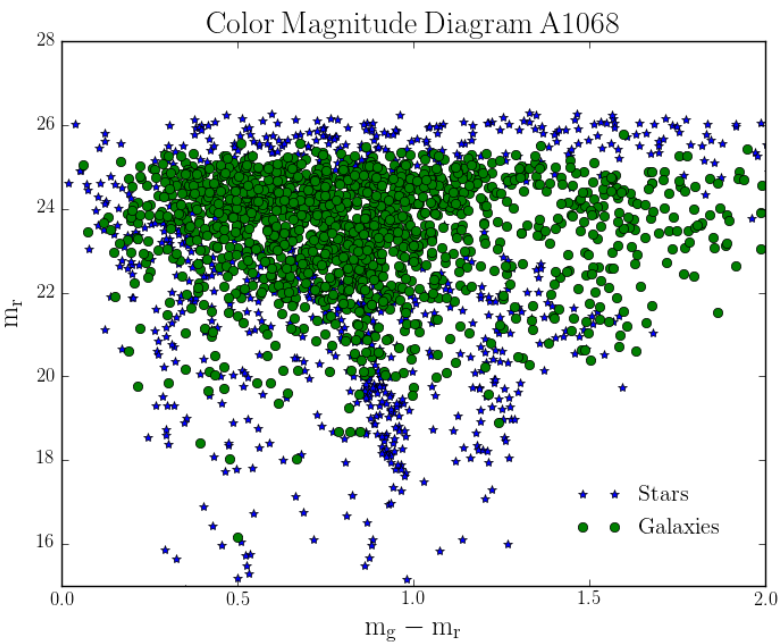


FIGURE 3.1: G

Mag vs flux rad to discriminate
vs flux rad.png vs flux rad.png vs flux rad.png vs flux rad.png vs flux rad.png vs flux
rad.png vs flux rad.png vs flux rad.png vs flux rad.png

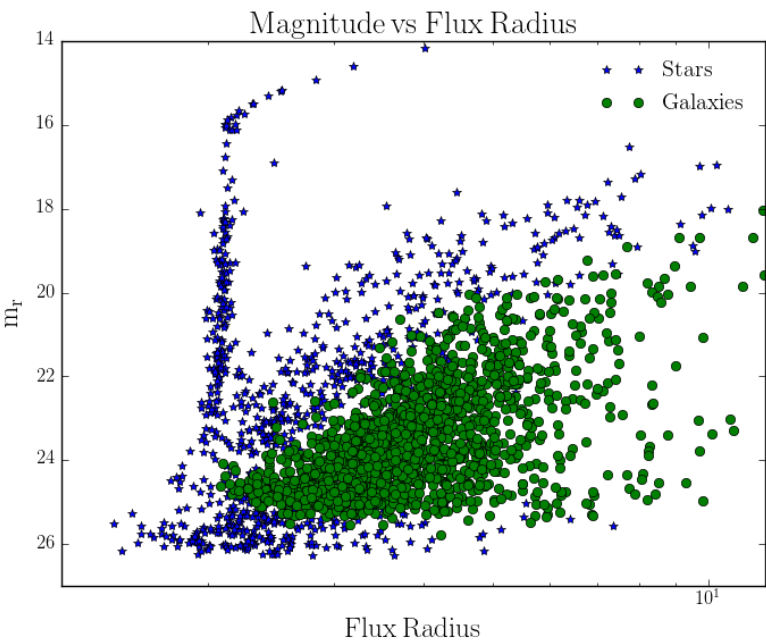


FIGURE 3.2: G

3.2 Galfit

galfit fits two dimensional profiles so it is a useful tool to remove the light from the BCG and allow us to observe background objects

Fit sersic profiles with $n=4$ which is de vacouleurs profile.

A first run gives us a rough idea of the true position of the center of the BCG so we can set this values in a second run for each cluster. We needed to combine different sersic parameters, as well as Fourier and bending modes for some of the BCGs.

We use the segmentation masks given by sextractor to mask bright objects in the fitting of the BCG

the fitting of many objects (not only the BCG)

the best results were given when we masked the innermost region of the BCG so the fitting will put more weight in the rest of the profile, thus reducing most of the light that hides the background objets.

we have to take into account the magnification bias

The parameters C0, B1, B2, F1, F2, etc. listed below are hidden from the user unless he/she explicitly requests them. These can be tagged on to the end of any previous components except, of course, the PSF and the sky – although galfit won't bar you from doing so, and will just ignore them. Note that a Fourier or Bending mode amplitude of exactly 0 will cause GALFIT to crash because the derivative image GALFIT computes internally will be entirely 0. If a Fourier or Bending amplitude is set to 0 initially GALFIT will reset it to a value of 0.01. To prevent GALFIT from doing so, one can set it to any other value.

Bending modes B1) 0.07 1 Bending mode 1 (shear) B2) 0.01 1 Bending mode 2 (banana shape) B3) 0.03 1 Bending mode 3 (S-shape)

Azimuthal fourier modes F1) 0.07 30.1 1 1 Az. Fourier mode 1, amplitude and phase angle F2) 0.01 10.5 1 1 Az. Fourier mode 2, amplitude and phase angle F6) 0.03 10.5 1 1 Az. Fourier mode 6, amplitude and phase angle F10) 0.08 20.5 1 1 Az. Fourier mode 10, amplitude and phase angle F20) 0.01 23.5 1 1 Az. Fourier mode 20, amplitude and phase angle

Traditional Diskyness/Boxyness parameter c C0) 0.1 0 traditional diskyness(-)/boxyness(+)

3.3 Color images

In er.

Here we take an isothermal sphere to model the Einstein ring in a distance of background objects of $z=1$

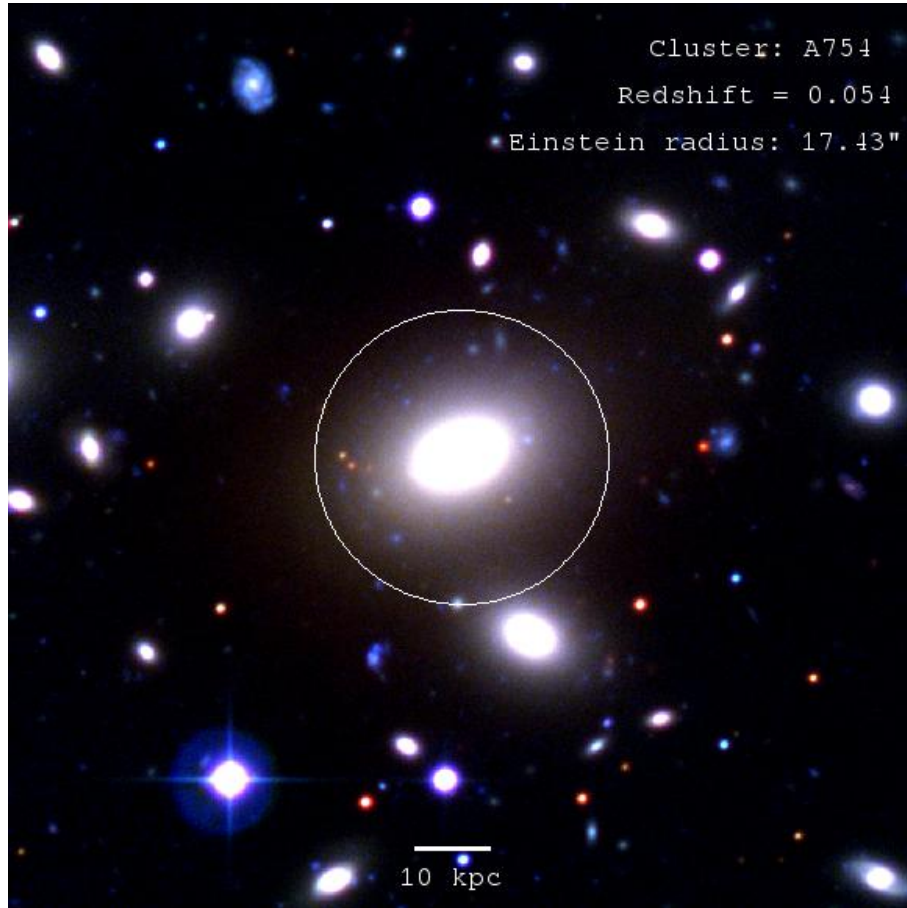


FIGURE 3.3: G

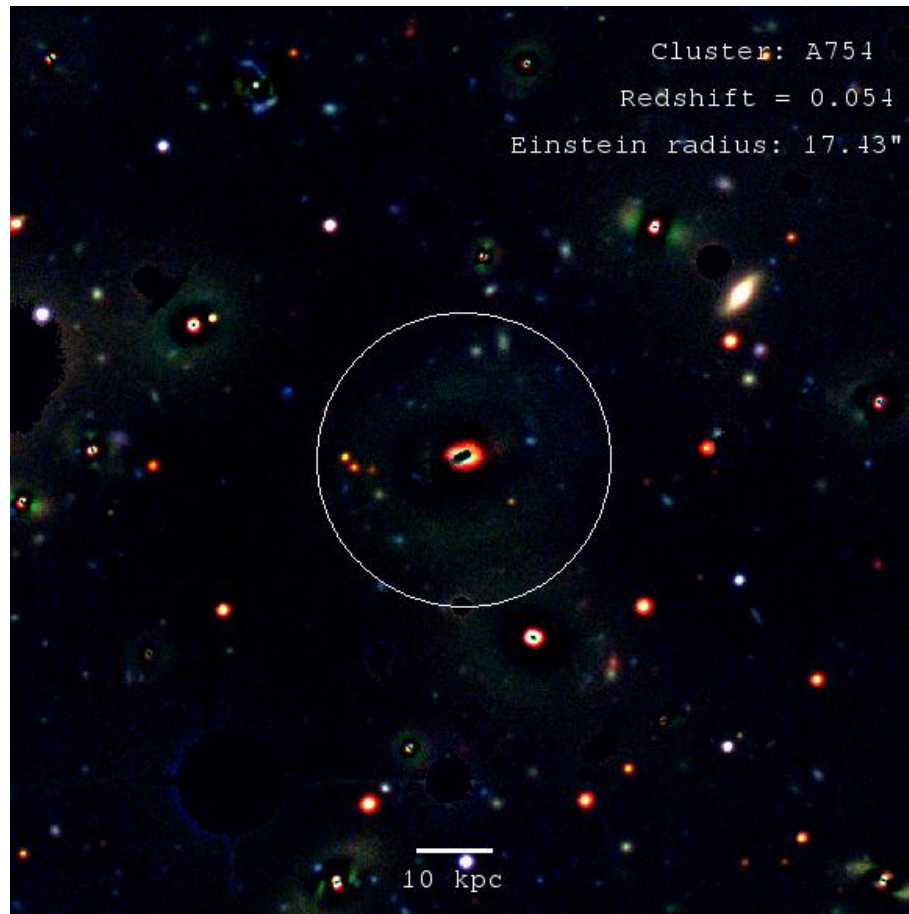


FIGURE 3.4: G

3.4 Photometric Redshift

(using as reference Benitez, Narciso 2000)

Chapter 4

Study of images

We ter.

Chapter 5

Conclusions

Thes.

Bibliography

- [1] Treu, Tommaso. 2010 *Strong Lensing by Galaxies*. Annu. Rev. Astron. Astrophysics. 2010. 48:87-125.
- [2] R. F. J. Van der Burg et. al. 2015 *Evidence for the inside-out growth of the stellar mass distribution in galaxy clusters since $z\tilde{1}$* . preprint arXiv:1412.2137v2.
- [3] Binney J., Tremaine S. *Galactic Dynamics*. Princeton University Press, 1994.
- [4] C. O. Wright & Teresa G. Brainerd, Teresa. 1999 *Gravitational Lensing by NFW halos*. preprint arXiv:astro-ph/9908213v1.
- [5] Smith, Russell. 2014 *Variations in the initial mass function in early-type galaxies: a critical comparison between dynamical and spectroscopic results*. MNRASL 443, L69-L73 (2014).
- [6] C. Bildfell et. al. 2012 *Evolution of the red sequence giant to dwarf ratio in galaxy clusters out to $z\tilde{0.5}$* . MNRAS 425, 204-221 (2012).
- [7] Smith, Russell & Lucey, John. 2013 *A giant elliptical galaxy with a lightweight initial mass function*. MNRAS 000, 1-14 (2013).
- [8] R. J. Smith et. al. 2015 *The IMF-sensitive $1.14\text{-}\mu\text{m}$ Na I doublet in early-type galaxies*. MNRAS 000, 1-14 (2013).
- [9] C. Sifon et. al. 2015 *Constraints on the alignment of galaxies in galaxy clusters from $\tilde{14000}$ spectroscopic members*. A&A 575, A48 (2015).
- [10] S. M. Adams et. al. 2012 *The environmental dependence of the incidence of galactic tidal features*. The Astrophysical Journal, 144:128(11pp) (2012).
- [11] D. J. Sand et. al. 2011 *Intracuster supernovae in the multi-epoch nearby cluster survey*. The Astrophysical Journal, 729:142 (13pp) (2011).
- [12] J. C. Muñoz Cuartas et. al. 2010 *The redshift evolution of Λ CDM halo parameters: concentration, spin and shape*. MNRAS, 000,1-11 (2010).