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A Numerical Study of Unsteady Natural Convection in a Rectangular Enclosure -- The Effect of Variable Thermodynamic and Transport Properties

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A Numerical Study of Unsteady Natural Convection in a Rectangular Enclosure – The Effect of Variable Thermodynamic and Transport Properties

A Thesis

Submitted to the Graduate faculty of
University of New Orleans
in partial fulfillment of the
requirements for the degree of

Master of Science
in
Mechanical Engineering
Thermal Fluids Science

by

Manohar Chidurala

B.Tech., JNTU College of Engineering, 2006
M.S., University of New Orleans, 2009

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To my Mom, Dad, Srinivas, Premalatha and Ambika

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Nomenclature

Symbols

AR	Aspect ratio (H/L)
c_p	Specific heat at constant pressure
\bar{c}_p	Non-dimensional specific heat at constant pressure
c_v	Specific heat at constant volume
\bar{c}_v	Non-dimensional specific heat at constant volume
E_t	Total energy per unit volume
e	Internal energy per unit mass
\vec{V}	Velocity vector
\vec{f}	Force per unit mass
$\vec{\Pi}$	Stress tensor
\vec{q}	Heat transfer by conduction
Q	Heat transfer by external agencies
Fr	Froude number (u_{ref}^2/gH)
g	Gravitational acceleration
\vec{g}	Body force vector per unit volume
Gr	Grashof number
H	Height of the cavity
\bar{H}	Non-dimensional height of the cavity
L	Characteristic length

\bar{L}	Non-dimensional width of the cavity
k	Thermal conductivity
p	Thermodynamic pressure
\bar{p}	Non-dimensional thermodynamic pressure
Pr	Prandtl number
R	Gas constant
\bar{R}	Non-dimensional gas constant
Ra	Rayleigh number ($g \beta \Delta T H^3 / \nu \alpha$)
h	Heat transfer coefficient
Nu	Nusselt number
T	Temperature
\bar{T}	Non-dimensional temperature
t	Time
\bar{t}	Non-dimensional time
u	Horizontal velocity
\bar{u}	Non-dimensional horizontal velocity
\tilde{u}	Non-dimensional horizontal velocity ($u H / \alpha$)
u_{ref}	Reference velocity ($Gr^{1/2} \nu / H$)
v	Vertical velocity
\bar{v}	Non-dimensional vertical velocity
\tilde{v}	Non-dimensional vertical velocity ($v H / \alpha$)
x	Horizontal spatial coordinate
\bar{x}	Non-dimensional horizontal spatial coordinate

y	Vertical spatial coordinate
\bar{y}	Non-dimensional vertical spatial coordinate

Greek Symbols

α	Thermal diffusivity; co-efficient of auxiliary matrix, P
β	Coefficient of thermal expansion
γ	Isentropic constant
$\Delta\bar{t}$	Non-dimensional time increment
ΔT	Temperature difference between hot and cold walls (T_h-T_c)
$\Delta\bar{x}$	Non-dimensional increment in the horizontal direction
$\Delta\bar{y}$	Non-dimensional increment in the vertical direction
ε	Convergence criterion for SIP algorithm
ε_p	Convergence criterion for pressure
μ	Absolute viscosity
$\bar{\mu}$	Non-dimensional absolute viscosity
ν	Kinematic viscosity
$\bar{\nu}$	Non-dimensional kinematic viscosity
ρ	Density
δ	Kronecker delta function
θ	Angle of inclination of the enclosure

Subscripts

c	Cold
h	Hot
w	Wall

i Index for space increment in x direction

j Index for space increment in y direction

ref Reference

Superscripts

k Index for number of iterations (k^{th} iteration in the CMSIP)

n Index for time increment

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Abstract

A two-dimensional mathematical model is adopted to investigate the development of buoyancy driven circulation patterns and temperature contours inside a rectangular enclosure filled with a compressible fluid where one of the vertical walls of the enclosure is kept at a higher temperature than the opposite one. Fluid thermodynamic and transport properties are assumed to be functions of temperature. The governing equations are discretized using second order accurate differencing for spatial and temporal derivatives and then linearized using Newton's linearization method. The resulting set of algebraic equations is solved using the Coupled Modified Strongly Implicit Procedure for the unknowns of the problem. The results of this study show that the variable property model predicts lower values for wall heat fluxes and Nu number than the constant property one for Rayleigh numbers between 10^4 and 10^5 .

Keywords: Natural convection, CMSIP, Buoyancy Driven flows with Variable Properties

Chapter 1

Introduction

Natural convection within enclosures has been widely studied by many authors. Some of the applications of natural convection within enclosures include heat removal from electronic components, cooling of nuclear reactors, climate control in rooms, and crystal growth in liquids. Numerical studies of buoyancy driven flows have been carried out using finite difference, finite volume, and finite element techniques. Experimental studies have also been carried out utilizing PIV (Particle Image Velocimetry), LDA (Laser Doppler Anemometry), smoke visualization, and interferometry. A comprehensive literature survey is provided in Chapter 2.

The main objective of this thesis is to study the steady state and transient natural convection phenomenon of a compressible fluid (an ideal gas with $Pr=0.72$) inside a square enclosure with differentially heated vertical walls and to determine the effect of variable thermodynamic and transport properties on circulation patterns, wall heat flux, and Nusselt number. A more detailed description of the physical problem is given in Chapter 3.

A physics based mathematical model is developed to obtain numerical solutions for the above problem by assuming the thermodynamic and transport properties of the fluid to be functions of temperature. The mathematical model comprises of conservation of mass, momentum, and energy equations. Formulation of conservation equations in terms of the primitive variables of the problem (u , v , p , and T) together with the appropriate boundary conditions and assumptions are provided in Chapter.4. All the dimensional variables are

transformed into non-dimensional form. These variables are then used to transform the conservation equations and boundary conditions into their corresponding non-dimensional forms.

The discretization and linearization of the non-dimensional conservation equations are given in Chapter.5. This chapter also includes the numerical solution and algorithm employed to predict the values of the primitive variables of the problem. A detailed description of the Coupled Modified Strongly Implicit Procedure (CMSIP) used in the solution algorithm, which is second order accurate in time differencing is also given in this chapter.

The accuracy of the numerical code is verified by applying the developed code to predict the results of a benchmark case study. Grid independence and time convergence studies were carried to verify the solution algorithm is presented in Chapter 6. The study of the second order accurate in time model is also provided including the grid independence and time convergence studies and then compared to the first order accurate in time model in Chapter 6. A parametric study was carried out to determine the effect of various parameters such as Rayleigh number and wall temperature difference and the size of the enclosure on circulation patterns and heat transfer in Chapter 7.

The accuracy of the results achieved by variable properties is compared to those results of constant properties in Chapter 8. The orientation of the enclosure is also studied in the Chapter 9. In Chapter 10, the results of various case studies to verify the accuracy of the solution algorithm for second order accurate in time model are presented. Transient development of circulations and stratification patterns are provided in Chapter 11. The wall heat flux and Nusselt number calculations are presented in Chapter 12.

The conclusions drawn from the results of the present study are presented in Chapter 13. Recommendations to improve the mathematical and solution procedure proposed in this study

are given in Chapter 14. A list of references used in the literature survey is provided in the “List of References” section of the thesis.

Chapter 2

Literature Survey

Natural convection within enclosures has been widely studied by many researchers. An excellent review of the developments in understanding and modeling of these studies natural convection in enclosures can be found in the review paper by Ostrach [1]. This article discussed the complexities of the natural convection phenomenon that arise due to the inherent interaction between the boundary layer near the walls and the outer core region. Suggestions for different variable scaling are also made. Use of experimental results to support numerical predictions was strongly advocated. Moreover, enclosures with different geometries such as horizontal cylinders, high and low aspect ratio rectangular enclosures, and annuli were also reviewed.

Kimura and Bejan [4] have studied natural convection in a rectangular cavity with uniform heat flux imposed at the sides of the cavity. They specifically studied the heating of the cavity from the top and from the sides. Their results showed that the convection driven by the horizontal temperature gradients persists even when the vertical stabilizing gradient is larger than the horizontal gradient. Buoyancy induced flows subjected to partially heated flows studied by T. H. Chen and L. Y. Chen [5] and also by Nithyadevi [6]. Cormack et al. [7] have considered the effect of the upper surface conditions on the buoyancy driven flow in a shallow cavity.

The natural convective heat transfer with varying boundary conditions studied by many authors such as Sathiyamoorthy et al. [8], Calcagni [9], Ganzarolli [10], Le Peutrec and Lauriat [11], Mariani [12], Aydin[13], and Basak [14]. They have concluded that the surface boundary conditions have an important influence on the temperature and flow structure within a cavity.

Nonlinear variations in liquid density with respect to temperature and its effect on buoyancy driven flow has been studied by Lin and Nansteel [15]. They have identified a dimensionless density distribution parameter which characterizes the distribution of the buoyancy force and how it determines the possible variations in flow structure. The effect of thermal stratification on natural convection in a vertical porous insulation layer has been studied by Rees and Lage [16].

Previous research indicates that the circulation patterns and thermal stratification in buoyancy driven flows are dependent on two-dimensional parameters: Prandtl and Rayleigh. Many researchers have carried out their studies for air ($Pr=0.7$). The results for low Pr number fluids (liquid metals) are reported by Viskanta et al. [17]. Fluid with high Pr numbers was studied by Hiller et al. [18]. Moh [19] studied the simulations of two-dimensional, low Prandtl number natural convection in harmonically oscillated, differentially heated enclosures. Buoyancy driven flows can be laminar or turbulent depending on the value of the Rayleigh number. The mathematical models proposed by Henkes et al. [20] can handle both types of flow regimes. Simulation of Laminar Buoyancy driven Flows in an Enclosure carried out by Evern Selamet et al. [21]. Numerical investigations have been conducted of flow transitions in Deep Cavities by Xia and Murthy [22]. Numerical investigation of turbulent natural convection in a square enclosure with localized heating from below and symmetrical cooling from the vertical side walls were carried out by Anil kumar Sharma et al. [23]. A numerical procedure to solve turbulent flow which makes use of k- ϵ model has been developed by Elkaim et al [24]. Laminar and turbulent natural convection in irregular shaped enclosures has been investigated by Coulter and Guceri [25]. An experimental and numerical investigation is presented concerning the natural convection of water near its density in a differentially heated rectangular enclosure at high Rayleigh numbers by Ho and Tu [26], in which an oscillatory convection regime may arise.

The unsteady, compressible effects were studied numerically by Akyuzlu et al. [27]. Their model included incompressible flow equations together with Boussinesq approximation and it was used to predict the transient circulation patterns and resulting thermal stratification of the propellant in a cryogenic storage tank for the constant heat flux boundary conditions. Steady, compressible effects were studied in closed cavities by Mazumder [28]. However, most of the researchers of natural convection in cavities and enclosures assume the density to be function of temperature only in the buoyancy term in the momentum equation (Boussinesq approximation.) According to Gray [29] a temperature difference less than 28.6 deg C between the hot and cold walls of the enclosure ensures the validity of this approximation.

There are many natural convection processes in various fields, and it is still a hot topic to investigate the fluid dynamics and heat transfer of natural convection. The analytical solutions are meaningful in both theoretical investigation and practical applications. Specially, they are very useful to computational fluid dynamics and heat transfer as benchmark solutions to check the numerical solutions and to develop numerical differencing schemes, and grid generation methods and so forth. An analytical study has been conducted for natural convection in a cavity of different aspect ratios with uniform volumetric heat generation by Joshi et al. [30]. Two explicit analytical solutions of 2-D steady laminar natural convection along a vertical porous plate and between two vertical plates were derived by Cai and Zhang [31]. The numerical study of natural convection in a differentially heated cavity with internal heat generation for different aspect ratios has been carried out by Fusegi et al. [32].

Most of the researchers have carried out their studies in two-dimension as it is done in the present study. This approach is being justified when the depth aspect ratio of the enclosure is greater than 5 [33]. Also, a three-dimensional approach may be the right choice when

multicellular, transient flow characteristics are expected as in the case of high Grashof number flows where three-dimensional effects are important. Researchers like Viskanta [17], Hiller [18], Fusegi [33], and Wakitani [34] have addressed this problem by proposing three-dimensional mathematical models which required a new method of solution than the previously adopted ones. (See publications by de Vahl Davis [35] and Kublbeck [36]). The effect of wall thermal boundary conditions on the development of three-dimensional and unsteady natural convective flows was studied by Oosthuizen et al. [37]. In 1997, multicellular solutions for air filled enclosure were developed by Wakitani [38]. A high aspect ratio ($AR = 16$) enclosure was considered in his study. A numerical study is conducted to investigate the effect of subcooling on natural convection in a densified cryogenic propellant by Akyuzlu et al. [39].

Validation of a proposed model can be done by comparing the predicted values of velocity and temperature fields to those obtained in experiments. Although limited, there are published experimental data in the literature. (See work done by Viskanta [17], Nicolette [40], Kirkpatrick [41], Hiller [18], and Linthorst [42]). The experimental analysis through 2D-PIV system was provided by Corvaro and Paroncini in 2009 [43]. The same authors developed interferometry techniques in 2007 [44]. Natural convection of air in a tall vertical cavity was studied using a smoke patterns and interferometry by Wright et al. [45]. An experimental study of high Rayleigh number natural convection in an enclosure was conducted by Giel and Schmidt [46]. The experimental benchmark data for turbulent natural convection to validate a proposed CFD code is presented in a publication by Ampofo et al. [47]. Akyuzlu, Nemani, and Chakravarthy [48] used the Particle Velocimetry (PIV) system to experimentally investigate the effect of different heat transfer boundary conditions on natural convection inside a rectangular storage tank. Water was used as the working fluid. Another method of validation of a newly proposed

mathematical model or solver is to predict the results of the benchmark case (heating of a square enclosure by isothermal walls.) In 1983, De Vahl Davis [2] numerically investigated the natural convection of air in a square enclosure. This study has widely been accepted as the benchmark case. The study used the non-dimensional stream function-vorticity formulation of the governing equations. The well known Boussinesq approximation was also incorporated.

Natural convection flows were also studied numerically inside rectangular cavities with inclination. Two dimensional natural convection flows in titled cavities for porous media and homogeneous fluids has been studied by Baez and Nicolas [49]. Natural convection in two dimensional enclosures with three flat and one wavy wall is numerically investigated by Dalal and Das [50]. In this study, one wall is having a sinusoidal temperature profile and other three walls including the wavy wall are maintained at constant cold temperature. A numerical study of the effect of a hot wavy wall of laminar natural convection in an inclined square cavity, differentially heated, was carried out by Adjilout et al. [51]. Flow visualization observations are described for natural convection flows in rectangular inclined enclosures by Linthorst et al. [42].

Variable property mathematical models were also developed to study natural convection in cavities and enclosures. The transient version of the classical differentially heated square cavity problem, considering fluid properties as functions of temperature has been studied by De Souza et al. [52]. Same problem with constant and variable fluid properties has been also studied by Leal et al. [53]. Variable property effects in laminar natural convection in a square enclosure were studied by Zhong et al. [54] and Emery et al. [55]. They indicate that the variable properties produced observable changes in the temperature and velocities, however, the overall heat transfer was found to be unaffected by it. The effects of combined temperature dependent thermodynamics and transport fluid properties on the heat transfer rate, heat function fields and

profiles in a fluid filled square enclosure was quantitatively studied by Waheed [56]. The variations in fluid properties were also considered in the numerical studies conducted by Markatos and Pericleous [57].

Present study is a continuation of an effort undertaken in the Cryogenics Laboratory at the University of New Orleans in collaboration with the Advanced Programs division of LMMO in New Orleans and NASA Stennis Space Center to study the transient heating of densified liquid propellants in a cryogenic storage tank. Previous studies conducted by one of the authors of this paper [58 and 59] has indicated the importance of predicting, accurately, the circulation patterns that develop inside the propellant due to natural convection. This process dominates the heating of the propellant and is driven by the heat leak at the walls of the storage tank which leads to unwanted thermal stratification in the tank. This fact led our research group in the Cryogenics lab to adopt a compressible flow model (without Boussinesq assumption) with constant properties to study natural convection in square enclosures [27]. i.e., a two-dimensional rectangular enclosure with insulated top and bottom and with vertical walls kept at a prescribed temperature difference.

To solve the governing equations of the proposed mathematical model, an in-house numerical solver based on an implicit finite difference technique, a modified version of the SIP algorithm [60] called Coupled Modified Strongly Implicit Procedure (CMSIP) [61], was developed by our group. (See reference 35 for the details of this algorithm.) In the first set of numerical experiments using this solver, the fluid was assumed to be an ideal gas with $Pr = 1$ [27]. In the present study, the mathematical model has been improved by assuming the thermodynamic and transport properties of the fluid to be functions of temperature. The solution algorithm had to be modified accordingly. Furthermore, the present study uses a second order accurate time differencing for the CMSIP solution algorithm (see references [76] and [77]).

Chapter 3

Description of the Physical Model

A two dimensional square enclosure is considered for this present study. The left wall of enclosure is maintained at a constant but higher temperature than the right wall. The top and bottom walls of the enclosure are assumed to be perfectly insulated. The enclosure is filled with air, which is considered to be a compressible ideal gas. Initially, the fluid inside the tank is assumed to be at an average temperature of the hot and cold walls. This temperature is called the reference temperature. The aspect ratio (ratio of height to width) is determined by the geometry of the enclosure. In this study, since the length of the enclosure is taken to be equal to the height, the aspect ratio is computed to be unity. No slip conditions are assumed on all the walls of the enclosure and the walls are presumed to be impermeable. Orientation of the enclosure is depends on direction of gravitational force. Refer to Figure 1 for more details of the physical model.

The driving mechanism behind the natural convection phenomenon is buoyancy. In this case, the buoyancy is associated with gravity. Heat is transferred from the hot wall to the adjacent air particles resulting in a decrease in the density of the air near the hot wall. The less dense air begins to rise while the heavier (more dense) air starts falling downward near the cold wall consequently introducing a clockwise circulation. This type of buoyancy induced flow can be seen in heating or cooling of rooms and in cooling of transistors and transformers.

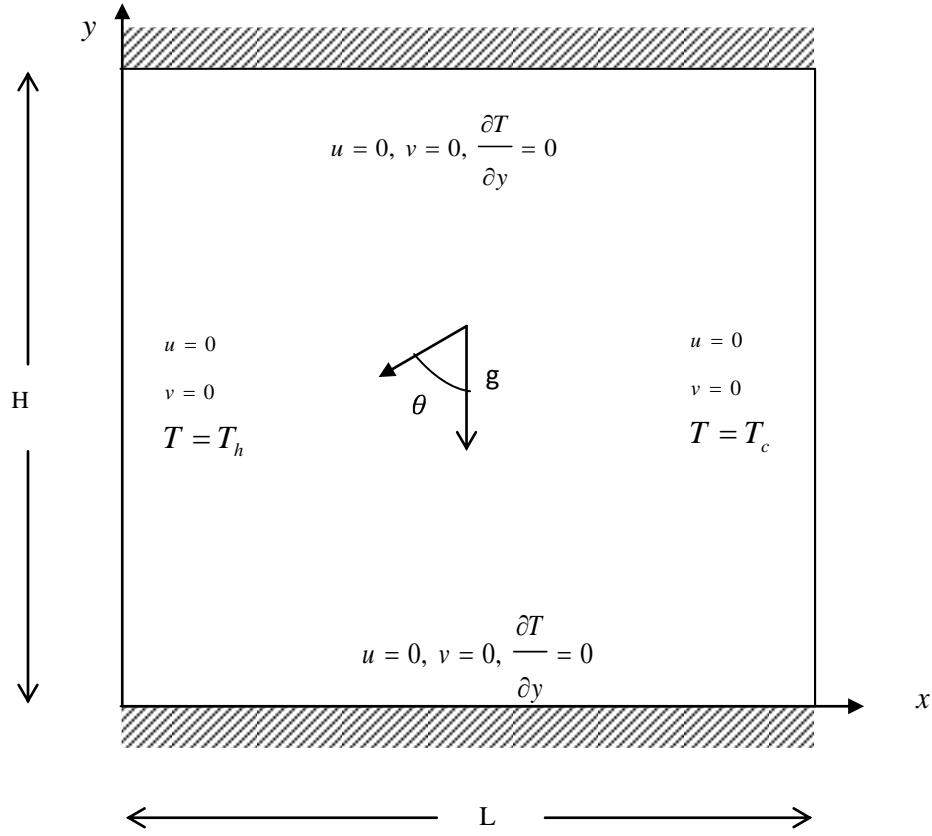


Figure 1- Schematic of the Physical Domain (Dimensional) and the Boundary Conditions

Depending on the Rayleigh number, the flow within the enclosure can be categorized as laminar, transitional, or turbulent. In this study, only laminar flows are considered. The Rayleigh numbers are kept below the transition limit (less than 1×10^7). The Rayleigh number is dependent on the properties of the working fluid, temperature differential between the vertical walls, and the characteristic length of the enclosure.

All the dimensional variables are transformed into a non-dimensional domain. The non-dimensional physical model is shown in Figure 2.

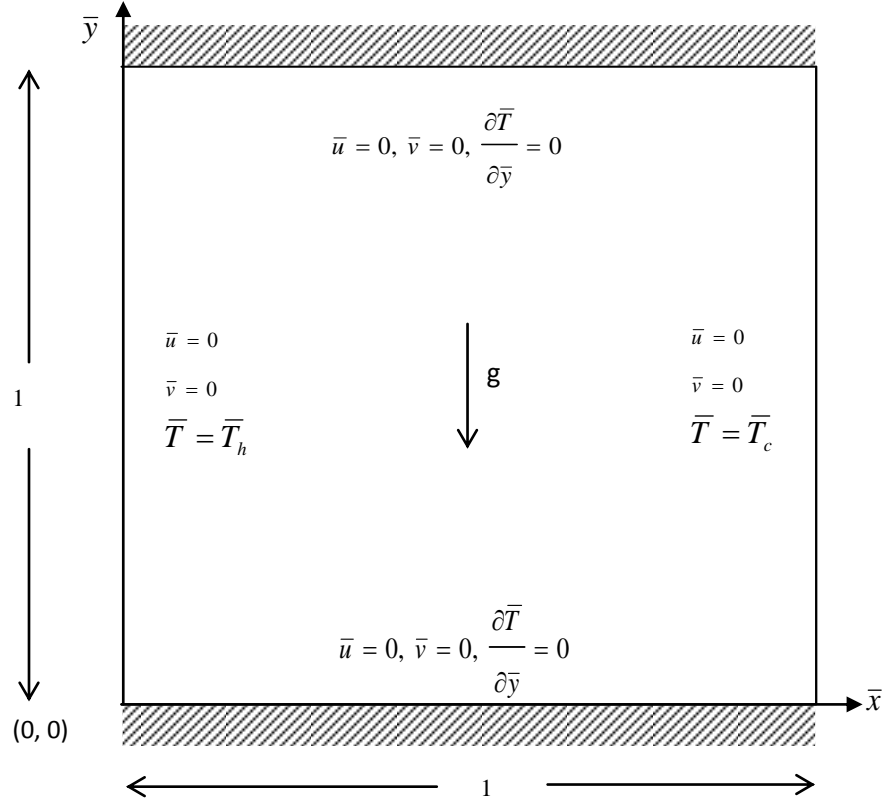


Figure 2- Schematic of the Physical Domain (Non- dimensional) and the Boundary Conditions

Chapter 4

Description of Mathematical Model

The mathematical formulation of the conservation equations together with the initial and boundary conditions in first order and second order accurate in time models are given below. The dimensional governing equations were derived from the respective vector form (refer to Appendix I) and together with the initial and boundary conditions were then transformed into non-dimensional form. The non-dimensional forms of these equations as well as the assumptions made in the derivation of these equations are presented in this chapter.

4.1. Assumptions of Mathematical Model

The following assumptions were made for the present study.

1. The physical domain is two-dimensional and the equations are in Cartesian co-ordinates
2. The working fluid forms a continuum
3. The flow is unsteady, laminar, and viscous
4. The working fluid (air) is compressible (the density of the fluid is a function of temperature and pressure) and can be treated as an ideal gas.
5. The working fluid behaves like a Newtonian fluid with stokes assumptions
6. Pressure work term is negligible in the energy equation
7. Effects of viscous dissipation in the energy equation are insignificant
8. The kinetic and potential energy terms in the energy equations are neglected
9. Radiation heat transfer is ignored

10. There are no internal heat sources
11. Heat conduction within the fluid follows Fourier's law
12. The physical properties of the fluid are assumed to be constant
13. Thermodynamic and transport properties are assumed to be function of temperature

4.2. Dimensional Formulation

After making the necessary assumptions given in section 4.1, each individual conservation equation was then derived from the respective vector form of that equation. The dimensional form of governing differential equations in conservative form can be referred to Appendix IIA. Then the equation of state was substituted in the conservational equations. The dimensional formulations of the governing equations for the four primitive variables (u , v , p and T) of the present problem are given as follows

Continuity equation

$$\frac{\partial}{\partial t} \left(\frac{p}{T} \right) + \frac{\partial}{\partial x} \left(\frac{pu}{T} \right) + \frac{\partial}{\partial y} \left(\frac{pv}{T} \right) = 0 \quad (4.1)$$

Momentum equation in the x-direction

$$\begin{aligned} \frac{\partial}{\partial t} \left(\frac{pu}{T} \right) + \frac{\partial}{\partial x} \left(\frac{puu}{T} \right) + \frac{\partial}{\partial y} \left(\frac{puv}{T} \right) + \frac{\partial(Rp)}{\partial x} - \frac{\partial}{\partial x} \left[\frac{2}{3} R\mu \left(2 \frac{\partial u}{\partial x} - \frac{\partial v}{\partial y} \right) \right] \\ - \frac{\partial}{\partial y} \left[R\mu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \right] + \left(\frac{p}{T} \right) g \sin \theta = 0 \end{aligned} \quad (4.2)$$

Momentum equation in the y-direction

$$\begin{aligned} \frac{\partial}{\partial t} \left(\frac{pv}{T} \right) + \frac{\partial}{\partial x} \left(\frac{puv}{T} \right) + \frac{\partial}{\partial y} \left(\frac{pvv}{T} \right) + \frac{\partial(Rp)}{\partial y} \\ - \frac{\partial}{\partial x} \left[R\mu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \right] - \frac{\partial}{\partial y} \left[\frac{2}{3} R\mu \left(2 \frac{\partial v}{\partial y} - \frac{\partial u}{\partial x} \right) \right] + \left(\frac{p}{T} \right) g \cos \theta = 0 \end{aligned} \quad (4.3)$$

Energy equation

$$\frac{\partial}{\partial t}(c_p p) + \frac{\partial}{\partial x}(c_p pu) + \frac{\partial}{\partial y}(c_p pv) - \frac{\partial}{\partial x}(Rk \frac{\partial T}{\partial x}) - \frac{\partial}{\partial y}(Rk \frac{\partial T}{\partial y}) = 0 \quad (4.4)$$

In both momentum equations, θ assumed to be zero for square enclosure with variable thermodynamic and transport properties, and also orientation of the enclosure was studied at different angles of θ .

4.3. Sub Models for Variable Thermodynamic and Transport Properties

Specific heat at constant pressure and viscosity are assumed to be functions of temperature in the present study of natural convection. That is,

$$C_p = f(T) \text{ and } \mu = f(T)$$

The specific heat at constant pressure is function of temperature for an ideal gas, C_{p0} , can be given as (refer to appendix III)

$$C_{p0} = C_0 + C_1\phi + C_2\phi^2 + C_3\phi^3 \quad (4.5)$$

The viscosity function of temperature for an ideal gas can be given as (refer to appendix III)

$$\frac{\mu}{\mu_0} = \left(\frac{T}{T_0} \right)^{\frac{3}{2}} \frac{T_0 + S}{T + S} \quad (4.6)$$

4.4. Non-dimensional Formulation

The non-dimensional conservation equations were transformed from the dimensional form of the equations given in the previous section 4.2 for present study of natural convection. The definitions of non-dimensional variables used in these governing differential equations are given below. In this study ρ_{ref} was determined at T_{∞} and P_{∞} . Reference characteristic length, L_{ref} , is then to be the initial height of the cavity, H ; i.e. $L_{\text{ref}} = H$ and also $T_{\text{ref}} = T_{\infty}$

$$Ra = \frac{g\beta\Delta T H^3}{\nu\alpha}$$

$$Gr = \frac{g\beta\Delta TH^3}{\nu^2}$$

$$u_{ref} = \sqrt{Gr} \frac{\nu}{H}$$

$$\bar{t} = \frac{t}{L_{ref}/u_{ref}}$$

$$\bar{x} = \frac{x}{L_{ref}}$$

$$\bar{y} = \frac{y}{L_{ref}}$$

$$\bar{u} = \frac{u}{u_{ref}}$$

$$\bar{v} = \frac{v}{u_{ref}}$$

$$\bar{p} = \frac{p}{\rho_{ref}u_{ref}}$$

$$\bar{T} = \frac{T}{T_{ref}}$$

$$\bar{\rho} = \frac{\rho}{\rho_{ref}}$$

$$\bar{\mu} = \frac{\mu}{\mu_{ref}}$$

$$\bar{R} = \frac{R}{u_{ref}^2/T_{ref}}$$

$$\bar{c}_p = \frac{c_p}{u_{ref}^2 / T_{ref}}$$

$$Pr = \frac{c_p \mu}{k}$$

$$Fr = \frac{u_{ref}}{\sqrt{gH}}$$

Using the above parameters, the dimensional forms of equations were transformed into their non-dimensional forms.

Continuity Equation

$$\frac{\partial}{\partial \bar{t}} \left(\frac{\bar{p}}{\bar{T}} \right) + \frac{\partial}{\partial \bar{x}} \left(\frac{\bar{p}\bar{u}}{\bar{T}} \right) + \frac{\partial}{\partial \bar{y}} \left(\frac{\bar{p}\bar{v}}{\bar{T}} \right) = 0 \quad (4.7)$$

Momentum equation in x-direction

$$\begin{aligned} \frac{\partial}{\partial \bar{t}} \left(\frac{\bar{p}\bar{u}}{\bar{T}} \right) + \frac{\partial}{\partial \bar{x}} \left(\frac{\bar{p}\bar{u}\bar{u}}{\bar{T}} \right) + \frac{\partial}{\partial \bar{y}} \left(\frac{\bar{p}\bar{u}\bar{v}}{\bar{T}} \right) + \frac{\partial}{\partial \bar{x}} (\bar{R}\bar{p}) \\ + \frac{1}{Gr^{1/2}} \frac{\partial}{\partial \bar{x}} \left[\frac{2}{3} \bar{R}\bar{\mu} \left(2 \frac{\partial \bar{u}}{\partial \bar{x}} - \frac{\partial \bar{v}}{\partial \bar{y}} \right) \right] - \frac{1}{Gr^{1/2}} \frac{\partial}{\partial \bar{y}} \left[\bar{R}\bar{\mu} \left(\frac{\partial \bar{u}}{\partial \bar{y}} + \frac{\partial \bar{v}}{\partial \bar{x}} \right) \right] + \frac{1}{Fr} \left(\frac{\bar{p}}{\bar{T}} \right) \sin \theta = 0 \end{aligned} \quad (4.8)$$

Momentum equation in y-direction

$$\begin{aligned} \frac{\partial}{\partial \bar{t}} \left(\frac{\bar{p}\bar{v}}{\bar{T}} \right) + \frac{\partial}{\partial \bar{x}} \left(\frac{\bar{p}\bar{u}\bar{v}}{\bar{T}} \right) + \frac{\partial}{\partial \bar{y}} \left(\frac{\bar{p}\bar{v}\bar{v}}{\bar{T}} \right) + \frac{\partial}{\partial \bar{y}} (\bar{R}\bar{p}) \\ - \frac{1}{Gr^{1/2}} \frac{\partial}{\partial \bar{x}} \left[\bar{R}\bar{\mu} \left(\frac{\partial \bar{v}}{\partial \bar{x}} + \frac{\partial \bar{u}}{\partial \bar{y}} \right) \right] - \frac{1}{Gr^{1/2}} \frac{\partial}{\partial \bar{y}} \left[\frac{2}{3} \bar{R}\bar{\mu} \left(2 \frac{\partial \bar{v}}{\partial \bar{y}} - \frac{\partial \bar{u}}{\partial \bar{x}} \right) \right] + \frac{1}{Fr} \left(\frac{\bar{p}}{\bar{T}} \right) \cos \theta = 0 \end{aligned} \quad (4.9)$$

Energy equation

$$\begin{aligned} \frac{\partial}{\partial \bar{t}} (\bar{c}_p \bar{p}) + \frac{\partial}{\partial \bar{x}} (\bar{c}_p \bar{p} \bar{u}) + \frac{\partial}{\partial \bar{y}} (\bar{c}_p \bar{p} \bar{v}) - \frac{1}{\text{Pr} Gr^{1/2}} \frac{\partial}{\partial \bar{x}} \left(\bar{R} \bar{c}_p \bar{\mu} \frac{\partial \bar{T}}{\partial \bar{x}} \right) \\ - \frac{1}{\text{Pr} Gr^{1/2}} \frac{\partial}{\partial \bar{y}} \left(\bar{R} \bar{c}_p \bar{\mu} \frac{\partial \bar{T}}{\partial \bar{y}} \right) = 0 \end{aligned} \quad (4.10)$$

4.4.1. Boundary Conditions

To complete the mathematical formulation, the following boundary conditions were used.

- a. No slip condition was imposed on all the walls of the cavity

$$\bar{v}(0, \bar{y}) = 0 \quad \text{at} \quad \bar{x} = 0$$

$$\bar{v}(1, \bar{y}) = 0 \quad \text{at} \quad \bar{x} = 1$$

$$\bar{u}(\bar{x}, 0) = 0 \quad \text{at} \quad \bar{y} = 0$$

$$\bar{u}(\bar{x}, 1) = 0 \quad \text{at} \quad \bar{y} = 1$$

- b. All the walls of the cavity were considered to be impermeable

$$\bar{u}(0, \bar{y}) = 0 \quad \text{at} \quad \bar{x} = 0$$

$$\bar{u}(1, \bar{y}) = 0 \quad \text{at} \quad \bar{x} = 1$$

$$\bar{v}(\bar{x}, 0) = 0 \quad \text{at} \quad \bar{y} = 0$$

$$\bar{v}(\bar{x}, 1) = 0 \quad \text{at} \quad \bar{y} = 1$$

- c. All the walls of the cavity were assumed to be perfectly insulated. The walls were considered to be adiabatic i.e., no heat flux

$$\bar{T}(0, \bar{y}) = \bar{T}_h \quad \text{at} \quad \bar{x} = 0$$

$$\bar{T}(1, \bar{y}) = \bar{T}_c \quad \text{at} \quad \bar{x} = 1$$

$$\frac{\partial \bar{T}}{\partial \bar{y}}(\bar{x}, 0) = 0 \quad at \quad \bar{y} = 0$$

$$\frac{\partial \bar{T}}{\partial \bar{y}}(\bar{x}, 1) = 0 \quad at \quad \bar{y} = 1$$

4.4.2. Initial Conditions

Initially, the fluid within the enclosure was assumed to be motionless and isothermal. The initial temperature was assumed to be the reference temperature. The initial pressure distribution was calculated such that the steady state momentum equations were satisfied everywhere in the computational domain.

$$\bar{u}(\bar{x}, \bar{y}) = 0 \quad at \quad \bar{t} = 0$$

$$\bar{u}(\bar{x}, \bar{y}) = 0 \quad at \quad \bar{t} = 0$$

$$\bar{T}(\bar{x}, \bar{y}) = \bar{T}_{ref} \quad at \quad \bar{t} = 0$$

Chapter 5

Numerical Formulation and Solution Technique

The governing differential equations for the comprehensive mathematical model are coupled therefore the Coupled Modified Strongly Implicit Procedure (CMSIP) was used to solve them in terms of the primitive variables.

The discretization method used for the inner points of the computational domains, which are represented by a uniform orthogonal structured mesh, will be given in the next section. The discretization method used for boundary conditions are different and described in section 5.2 separately. The Newton's Linearization technique was developed to linearize the resulting non-linear discretized governing differential equations in the section 5.3. The computational cell used for the discretization is shown below in figure 3.

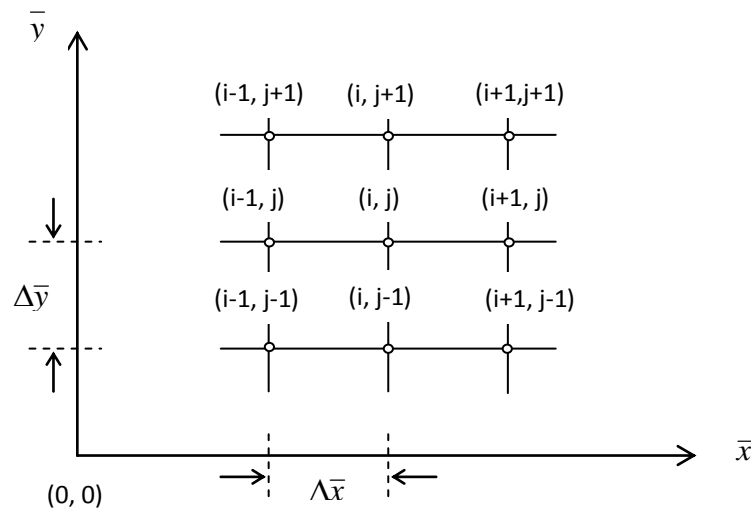


Figure 3- Computational Cell used for the Discretization of the Governing Differential Equations

5.1 Discretization of the Governing Differential Equations

The governing differential equations are discretized using second order accurate central differencing for spatial derivatives and second order (based on Taylor Expansion) finite differencing for time derivatives.

Second order accurate approximation for the time term in the inertial term becomes

$$\left(\frac{\bar{p}\bar{u}}{\bar{T}}\right)_{i,j}^{n+1} = \left(\frac{\bar{p}\bar{u}}{\bar{T}}\right)_{i,j}^n + \frac{\Delta t}{2} \left[\frac{\partial}{\partial t} \left(\frac{\bar{p}\bar{u}}{\bar{T}}\right)_{i,j}^n + \frac{\partial}{\partial t} \left(\frac{\bar{p}\bar{u}}{\bar{T}}\right)_{i,j}^{n+1} \right] + O(\Delta t^2, \Delta x^2, \Delta y^2) \quad (5.1)$$

$$\frac{\partial}{\partial t} \left(\frac{\bar{p}\bar{u}}{\bar{T}}\right)_{i,j}^{n+1} = \frac{2}{\Delta t} \left[\left(\frac{\bar{p}\bar{u}}{\bar{T}}\right)_{i,j}^{n+1} - \left(\frac{\bar{p}\bar{u}}{\bar{T}}\right)_{i,j}^n \right] - \frac{\partial}{\partial t} \left(\frac{\bar{p}\bar{u}}{\bar{T}}\right)_{i,j}^n \quad (5.2)$$

$\frac{\partial}{\partial t} \left(\frac{\bar{p}\bar{u}}{\bar{T}}\right)_{i,j}^n$ was derived from the Navier-Stokes equation at n^{th} time.

A central differencing was used for the first derivative of convection term in the axial direction:

$$\frac{\partial}{\partial \bar{x}} \left(\frac{\bar{p}\bar{u}\bar{u}}{\bar{T}}\right) = \left[\left(\frac{\bar{p}\bar{u}\bar{u}}{\bar{T}}\right)_{i+1,j}^{n+1} - \left(\frac{\bar{p}\bar{u}\bar{u}}{\bar{T}}\right)_{i-1,j}^{n+1} \right] / 2\Delta \bar{x} \quad (5.3)$$

The diffusion term in the vertical direction was discretized using central differencing as given below:

$$\frac{\partial}{\partial \bar{y}} \left(\bar{R}\bar{\mu} \frac{\partial \bar{v}}{\partial \bar{y}} \right)_{i,j}^{n+1} = \left[\left(\bar{R}\bar{\mu} \frac{\partial \bar{v}}{\partial \bar{y}} \right)_{i+\frac{1}{2},j}^{n+1} - \left(\bar{R}\bar{\mu} \frac{\partial \bar{v}}{\partial \bar{y}} \right)_{i-\frac{1}{2},j}^{n+1} \right] / \Delta \bar{y} \quad (5.4)$$

The term evaluated at the $i + \frac{1}{2}$ nodal point is evaluated by taking the average value at $i+1$ and i nodes.

$$(\bar{R}\bar{\mu})_{i+\frac{1}{2}} = \left[(\bar{R}\bar{\mu})_{i,j} + (\bar{R}\bar{\mu})_{i+1,j} \right] / 2 \quad (5.5)$$

The term evaluated at the $i - \frac{1}{2}$ nodal point is evaluated by taking the average value at $i-1$ and i nodes.

$$(\bar{R}\bar{\mu})_{i-\frac{1}{2}} = [(\bar{R}\bar{\mu})_{i,j} + (\bar{R}\bar{\mu})_{i-1,j}] / 2 \quad (5.6)$$

The first order derivative at plus half nodal point is evaluated as it shows below

$$\left(\frac{\partial \bar{v}}{\partial \bar{y}} \right)_{i+\frac{1}{2},j}^{n+1} = [\bar{v}_{i+1,j}^{n+1} - \bar{v}_{i,j}^{n+1}] / \Delta \bar{y} \quad (5.7)$$

and the same way the first order derivative for the minus half nodal point was evaluated

$$\left(\frac{\partial \bar{v}}{\partial \bar{y}} \right)_{i-\frac{1}{2},j}^{n+1} = [\bar{v}_{i,j}^{n+1} - \bar{v}_{i-1,j}^{n+1}] / \Delta \bar{y} \quad (5.8)$$

The second order mixed derivative was evaluated using central differencing as follows

$$\frac{\partial}{\partial \bar{y}} \left(\bar{R}\bar{\mu} \frac{\partial \bar{u}}{\partial \bar{x}} \right) = [\left(\bar{R}\bar{\mu} \frac{\partial \bar{u}}{\partial \bar{x}} \right)_{i,j+1}^{n+1} - \left(\bar{R}\bar{\mu} \frac{\partial \bar{u}}{\partial \bar{x}} \right)_{i,j-1}^{n+1}] / (2\Delta \bar{y}) \quad (5.9)$$

and

$$\left(\frac{\partial \bar{u}}{\partial \bar{x}} \right)_{i,j+1}^{n+1} = [\bar{u}_{i+1,j+1}^{n+1} - \bar{u}_{i-1,j+1}^{n+1}] / (2\Delta \bar{x}) \quad (5.10)$$

$$\left(\frac{\partial \bar{u}}{\partial \bar{x}} \right)_{i,j-1}^{n+1} = [\bar{u}_{i+1,j-1}^{n+1} - \bar{u}_{i-1,j-1}^{n+1}] / (2\Delta \bar{x}) \quad (5.11)$$

5.1.1 Finite Difference Approximation for Boundary Conditions

The first order derivative for the temperature gradient at the bottom and top boundaries was approximated by the four point forward and backward difference formulas, respectively. The elevation of the temperature gradient at the bottom wall is presented below as an example.

$$\frac{\partial \bar{T}}{\partial \bar{y}} = 0 \quad (5.12)$$

The approximation of the non-dimensional temperature at the all as given by the four point formula can be given by:

$$\bar{T}_1 = \frac{1}{11}(18\bar{T}_2 - 9\bar{T}_3 + 2\bar{T}_4) \quad (5.13)$$

5.2 Linearization of the Governing Differential Equations

After discretization the non-linear equations are linearized for the unknown variables by Newton's method using the following general formula

$$[uv\dots]^{k+1} = [uv\dots]^k + \frac{\partial[uv\dots]^k}{\partial u}(u^{k+1} - u^k) + \frac{\partial[uv\dots]^k}{\partial v}(v^{k+1} - v^k) + \dots \quad (5.14)$$

As an example the application of the above formula to various terms in the transform equation is presented below.

Consider the non-linear term $(\frac{\bar{p}\bar{u}}{\bar{T}})_{i,j}^{n+1}$ from the discretization of the time term in the x-momentum equation, where the primitive variables are \bar{p} , \bar{u} , \bar{T} application of Equation 5.14 results in:

$$\begin{aligned} \frac{\bar{p}_{i,j}^{n+1} \bar{u}_{i,j}^{n+1}}{\bar{T}_{i,j}} &= \left[\frac{\bar{p}\bar{u}}{\bar{T}} \right]_{i,j}^{n+1,k} + \frac{\partial \left[\frac{\bar{p}\bar{u}}{\bar{T}} \right]_{i,j}^{n+1,k}}{\partial \bar{u}} (\bar{u}_{i,j}^{n+1,k+1} - \bar{u}_{i,j}^{n+1,k}) \\ &\quad + \frac{\partial \left[\frac{\bar{p}\bar{u}}{\bar{T}} \right]_{i,j}^{n+1,k}}{\partial \bar{v}} (\bar{v}_{i,j}^{n+1,k+1} - \bar{v}_{i,j}^{n+1,k}) + \frac{\partial \left[\frac{\bar{p}\bar{u}}{\bar{T}} \right]_{i,j}^{n+1,k}}{\partial \bar{T}} (\bar{T}_{i,j}^{n+1,k+1} - \bar{T}_{i,j}^{n+1,k}) \\ &= \left[\frac{\bar{p}\bar{u}}{\bar{T}} \right]_{i,j}^{n+1,k} + \left[\frac{\bar{p}}{\bar{T}} \right]_{i,j}^{n+1,k} (\bar{u}_{i,j}^{n+1,k+1} - \bar{u}_{i,j}^{n+1,k}) \\ &\quad + \left[\frac{\bar{u}}{\bar{T}} \right]_{i,j}^{n+1,k} (\bar{p}_{i,j}^{n+1,k+1} - \bar{p}_{i,j}^{n+1,k}) + \left[-\frac{\bar{p}\bar{u}}{\bar{T}^2} \right]_{i,j}^{n+1,k} (\bar{T}_{i,j}^{n+1,k+1} - \bar{T}_{i,j}^{n+1,k}) \\ &= \left[\frac{\bar{p}}{\bar{T}} \right]_{i,j}^{n+1,k} (\bar{u}_{i,j}^{n+1,k+1}) + \left[\frac{\bar{u}}{\bar{T}} \right]_{i,j}^{n+1,k} (\bar{p}_{i,j}^{n+1,k+1}) + \left[-\frac{\bar{p}\bar{u}}{\bar{T}^2} \right]_{i,j}^{n+1,k} (\bar{T}_{i,j}^{n+1,k+1}) \end{aligned} \quad (5.15)$$

where n stands for a discrete time step and k for the number of iteration. In Equation (5.15),

$\bar{p}_{i,j}^{n+1,k+1}$, $\bar{u}_{i,j}^{n+1,k+1}$, and $\bar{T}_{i,j}^{n+1,k+1}$ are the unknown and the rest are known coefficients from the

previous iteration.

Another example of anon-linear term with four variables is $(\frac{\bar{p}\bar{u}\bar{v}}{\bar{T}})^{n+1}_{i,j+1}$ which is one of the convective terms in the y-momentum equation. Application of the linearization formula as given by Equation 5.14 yields:

$$\begin{aligned}
\frac{\bar{p}_{i,j+1}^{n+1,k+1}\bar{u}_{i,j+1}^{n+1,k+1}\bar{v}_{i,j+1}^{n+1,k+1}}{\bar{T}_{i,j+1}^{n+1,k+1}} &= \left[\frac{\bar{p}\bar{u}\bar{v}}{\bar{T}} \right]_{i,j+1}^{n+1,k} + \frac{\partial \left[\frac{\bar{p}\bar{u}\bar{v}}{\bar{T}} \right]_{i,j+1}^{n+1,k}}{\partial \bar{u}} (\bar{p}_{i,j+1}^{n+1,k+1} - \bar{p}_{i,j+1}^{n+1,k}) \\
&+ \frac{\partial \left[\frac{\bar{p}\bar{u}\bar{v}}{\bar{T}} \right]_{i,j+1}^{n+1,k}}{\partial \bar{u}} (\bar{u}_{i,j+1}^{n+1,k+1} - \bar{u}_{i,j+1}^{n+1,k}) + \frac{\partial \left[\frac{\bar{p}\bar{u}\bar{v}}{\bar{T}} \right]_{i,j+1}^{n+1,k}}{\partial \bar{v}} (\bar{v}_{i,j+1}^{n+1,k+1} - \bar{v}_{i,j+1}^{n+1,k}) \\
&+ \frac{\partial \left[\frac{\bar{p}\bar{u}\bar{v}}{\bar{T}} \right]_{i,j+1}^{n+1,k}}{\partial \bar{T}} (\bar{T}_{i,j+1}^{n+1,k+1} - \bar{T}_{i,j+1}^{n+1,k}) \\
&= \left[\frac{\bar{p}\bar{u}\bar{v}}{\bar{T}} \right]_{i,j+1}^{n+1,k} + \left[\frac{\bar{u}\bar{v}}{\bar{T}} \right]_{i,j+1}^{n+1,k} (\bar{p}_{i,j+1}^{n+1,k+1} - \bar{p}_{i,j+1}^{n+1,k}) \left[\frac{\bar{p}\bar{v}}{\bar{T}} \right]_{i,j+1}^{n+1,k} (\bar{u}_{i,j+1}^{n+1,k+1} - \bar{u}_{i,j+1}^{n+1,k}) \\
&+ \left[\frac{\bar{p}\bar{u}}{\bar{T}} \right]_{i,j+1}^{n+1,k} (\bar{v}_{i,j+1}^{n+1,k+1} - \bar{v}_{i,j+1}^{n+1,k}) + \left[-\frac{\bar{p}\bar{u}\bar{v}}{\bar{T}^2} \right]_{i,j+1}^{n+1,k} (\bar{T}_{i,j+1}^{n+1,k+1} - \bar{T}_{i,j+1}^{n+1,k}) \\
&= \left[\frac{\bar{u}\bar{v}}{\bar{T}} \right]_{i,j+1}^{n+1,k} (\bar{p}_{i,j+1}^{n+1,k+1}) + \left[\frac{\bar{p}\bar{v}}{\bar{T}} \right]_{i,j+1}^{n+1,k} (\bar{u}_{i,j+1}^{n+1,k+1}) + \left[\frac{\bar{p}\bar{u}}{\bar{T}} \right]_{i,j+1}^{n+1,k} (\bar{v}_{i,j+1}^{n+1,k+1}) \\
&+ \left[-\frac{\bar{p}\bar{u}}{\bar{T}^2} \right]_{i,j+1}^{n+1,k} (\bar{T}_{i,j+1}^{n+1,k+1}) - \left[\frac{\bar{p}\bar{u}\bar{v}}{\bar{T}} \right]_{i,j+1}^{n+1,k}
\end{aligned} \tag{5.16}$$

5.3 Solution Technique

A modified version of the Coupled Strongly Implicit Procedure (CSIP) developed by Zedan and Schneider [60] was used to solve for the primitive variables of the non-dimensional conservation equations. This modified version was proposed by Chen and Pletcher [67] and is called the Coupled Modified Strongly Implicit Procedure (CMSIP).

The descritized non-dimensional conservation equations after linearized are put into the following form for any nodal point (i, j) of the computational domain

$$\begin{aligned} A_{i,j}^6 x_{i-1,j-1} + A_{i,j}^5 x_{i,j-1} + A_{i,j}^4 x_{i+1,j-1} + A_{i,j}^7 x_{i-1,j} + A_{i,j}^9 x_{i,j} \\ + A_{i,j}^3 x_{i+1,j} + A_{i,j}^8 x_{i-1,j+1} + A_{i,j}^1 x_{i,j+1} + A_{i,j}^2 x_{i+1,j+1} = b_{i,j} \end{aligned} \quad (5.17)$$

In the present study, where we have four primitive variables, the x in Equation 5.17 is 4 element vectors and A is a 4 by 4 matrix.

The computational molecule for A^1 , A_i^2 , A^3 , . . . and A^9 in Equation 5.17 is shown in Figure 4.

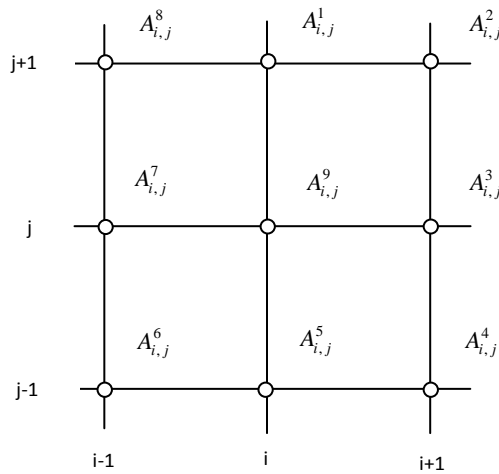


Figure 4- Two Dimensional Computational Molecules for the Elements of the A Matrix

Similar equations are generated for the rest of the inner nodal points of the computational domain. The resulting sets of algebraic equations (as many as the number of the unknowns) are then put into block matrix in the following form:

$$[A] x = b \quad (5.18)$$

where $[A]$ is the coefficient matrix, x is the unknown vector to solve for and b is the known vector. Except for boundaries with no slip condition, all boundary conditions were treated implicitly by writing the governing differential equations on the boundary points and then discretized and linearized before incorporating them into the block matrix form. The unknown vector has as many elements as the number of inner nodes in the computational domain. The $[A]$ matrix is a block matrix and is defined as shown below:

$$A = \begin{bmatrix} A_{1,1}^9 & A_{1,1}^3 & & & A_{1,1}^1 & A_{1,1}^2 & & & \\ * & * & * & & * & * & * & & \\ & * & * & * & & * & * & * & \\ & & & & & & & & - \\ & & & & & & & & - \\ * & * & & & * & * & * & * & * & * \\ A_{i,j}^6 & A_{i,j}^5 & A_{i,j}^4 & & A_{i,j}^7 & A_{i,j}^9 & A_{i,j}^3 & & A_{i,j}^8 & A_{i,j}^1 & A_{i,j}^2 \\ & * & * & * & & * & * & * & * & * & * \\ & & & & & & & & & & - \\ & & & & & & & & & & - \\ & & & & * & * & * & & * & * & * \\ & & & & & * & * & * & & * & * & * \\ & & & & & & & & A_{im,jm}^6 & A_{im,jm}^5 & & A_{im,jm}^7 & A_{im,jm}^9 \end{bmatrix}$$

The vectors \mathbf{x} and \mathbf{b} are also of block form are presented below:

$$\mathbf{x} = \begin{bmatrix} \begin{bmatrix} u \\ v \\ p \\ T \end{bmatrix}_{1,l} \\ \begin{bmatrix} u \\ v \\ p \\ T \end{bmatrix}_{2,l} \\ \cdot \\ \cdot \\ \begin{bmatrix} u \\ v \\ p \\ T \end{bmatrix}_{jn,kn-1} \\ \begin{bmatrix} u \\ v \\ p \\ T \end{bmatrix}_{jn,kn} \end{bmatrix} \quad \mathbf{b} = \begin{bmatrix} \begin{bmatrix} b_u \\ b_v \\ b_p \\ b_T \end{bmatrix}_{1,l} \\ \begin{bmatrix} b_u \\ b_v \\ b_p \\ b_T \end{bmatrix}_{2,l} \\ \cdot \\ \cdot \\ \begin{bmatrix} b_u \\ b_v \\ b_p \\ b_T \end{bmatrix}_{jn,kn-1} \\ \begin{bmatrix} b_u \\ b_v \\ b_p \\ b_T \end{bmatrix}_{jn,kn} \end{bmatrix}$$

The coefficient matrix $[A]$ is decomposed into upper and lower diagonal matrices $[U]$ and $[L]$ to solve the equation 5.18 and to make this possible an auxiliary matrix $[P]$ was added to both sides of the equation.

$$[A + P] x^{k+1} = b + [P]x^k \quad (5.19)$$

where transcript k is for iteration

setting

$$\delta^{k+1} = x^{k+1} - x^k$$

and

$$R^k = b - [A] x^k$$

in equation (5.19) becomes

$$[A + P] \delta^{k+1} = R^k \quad (5.20)$$

Replacing the matrix $[A+P]$ with the product of Lower-block triangular matrix $[L]$ and Upper-block triangular matrix $[U]$ in equation 5.20 (see appendix V), it becomes

$$[L][U] \delta^{k+1} = R^k \quad (5.21)$$

Defining the vector W by

$$W^{k+1} = [U] \delta^{k+1} \quad (5.22)$$

then equation (5.22) can be written as

$$[L] W^{k+1} = R^k \quad (5.23)$$

The solution of equation 5.19 is as follows:

The vector W^{k+1} is calculated from equation (5.23) by forward substitution procedure and then vector δ^{k+1} is computed from equation (5.22) by backward substitution. This procedure was repeated for the calculation of the new residual vector R followed by direct calculation of the W and δ until the convergence of X is attained.

The convergence depends upon the following criterion:

$$\left(\frac{\sum_{i,j=1}^{im,jm} \sum_{n=1}^{n=4} \left| \frac{\varphi_n^{k+1} - \varphi_n^k}{\varphi_{n,rms}^{k+1}} \right|^2}{4 \times im \times jm} \right)^{0.5} \leq \varepsilon \quad (5.24)$$

where,

k is the iteration level,

im is the number of grid points in the x-direction

jm is the number of grid points in the y direction

n is the variable index

φ_n is the component of the unknown vector φ

$\varphi_{n,rms}$ is the root mean square value of φ

The value of ε was set to 1×10^{-7} . This value of ε was chooses by numerical experimentation.

A separate criterion for pressure was also considered.

$$\left| p_n^{k+1} - p_n^k \right| \leq \varepsilon_p \quad (5.25)$$

Its convergence value was set to 1×10^{-7} .

Second order time accuracy approximation was stabilized by adding pressure smoothening function as shown below.

$$p_{i,j} = \sum_{i=2,j=2}^{im-1,jm-1} \left(p_{i,j} + \frac{w}{2\Delta x} (p_{i+1,j} - 2p_{i,j} + p_{i-1,j}) + \frac{w}{2\Delta y} (p_{i,j+1} - 2p_{i,j} + p_{i,j-1}) + \right) \quad (5.26)$$

where

w is small constant assigned as 0.0005

5.4. Solution Algorithm

The solution was developed to implement the solution procedure (presented in previous section) used to solve the governing equations of this study.

The steps of this algorithm are given below.

1. All the necessary variables were defined and the corresponding arrays were dimensioned

2. The input parameters such as non-dimensional computational, geometrical, and operational parameters were defined
3. The initial conditions were imposed, the velocities were set to zero and the temperature was assumed to be uniform with a value of unity. For natural convection, a hydrostatic pressure distribution was assumed in the vertical direction while the pressure distribution in the horizontal direction was assumed invariant.
4. The appropriate boundary conditions were imposed on all the walls of the enclosure. The velocities were set to zero. The temperatures on the left and right vertical walls were set to the corresponding non-dimensional hot and cold values.
5. The matrices for the corresponding inner nodes were generated
6. The Coupled Modified Strongly Implicit Procedure was implemented as explained in the previous section and counter for linearization is activated.
7. Values of the primitive variables were computed for the future time step. This was an iterative process. The iteration was stopped when the difference between the predicted values at iteration level $k+1$ and k was smaller than a given criteria. The separate criterion was also verified for pressure.
8. Steps 5 through 7 were repeated for each time step until the final simulation time was reached.
9. Steady state was verified for all the nodal points in the domain.

The above solution algorithm was implemented into a computer code in Fortran by Akyuzlu and Antoniou [62]. This code was executed to generate the simulations of natural convection of an incompressible fluid inside an enclosure with differentially heated side walls.

Chapter 6

The Validation of the Mathematical and Computer Codes

In this study, various numerical studies have been carried out to establish the validity of the natural convection phenomenon. When numerical methods are used to solve differential equations, getting a numerical stable solution does not immediately imply that this is the correct one. A time convergence study and grid independence study requires that there are not significant changes in the results when the computational time increment and grid size are varied. Furthermore, the computer code developed for the present study was tested for accuracy using benchmark studies. These case studies are presented in this chapter.

6.1. First Order Accurate in Time Model

Time convergence and grid independent studies were presented for first order accurate in time model in the following section.

6.1.1. Time Convergence Study

A time convergence study was carried out using three different non-dimensional time steps 1×10^{-5} , 5×10^{-6} and 1×10^{-6} for the present study. Once the grid size is set, the time computational time convergence is determined by satisfying the diagonal dominance condition of the coefficient matrix. For the present code validation study, the biggest non-dimensional computational time step is 1×10^{-5} . It was found that a time increment smaller than this did not result in any significant changes in qualitative and quantitative results. To illustrate this point, non-dimensional vertical velocity distribution along the horizontal centerline of the enclosure were compared for different time increments such as $\Delta \bar{t} = 1 \times 10^{-5}$, $\Delta \bar{t} = 5 \times 10^{-6}$ and $\Delta \bar{t} = 1 \times$

10^{-6} for $Ra = 10^5$ and 21 x 21 mesh points. The quantitative comparison is presented in Table 1 and the comparison of predicted non-dimensional vertical velocity profiles is presented in Figure

5.

Table 1- Steady State Non-dimensional values of Vertical Velocity along the Horizontal Centerline of the Enclosure for $\Delta\bar{t} = 1.0E-5$, $\Delta\bar{t} = 5.0E-6$ and $\Delta\bar{t} = 1.0E-6$ ($Ra = 10^5$ and 21 x 21)

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Vertical Velocity, \bar{v} ($\Delta\bar{t} = 1.0E-5$)	Non-dimensional Vertical Velocity, \bar{v} ($\Delta\bar{t} = 5.0E-6$)	Non-dimensional Vertical Velocity, \bar{v} ($\Delta\bar{t} = 1.0E-6$)
1.00	0.0000000	0.0000000	0.0000000
0.95	-0.2465031	-0.2405049	-0.2455756
0.90	-0.2011154	-0.2007765	-0.2008536
0.85	-0.1036450	-0.1085817	-0.1036823
0.80	-0.0343976	-0.1489494	-0.0347446
0.75	-0.0019368	-0.0411735	-0.0023072
0.70	0.0096405	0.0066313	0.0093005
0.65	0.0114227	0.0093207	0.0111182
0.60	0.0103105	0.0085196	0.0100377
0.55	0.0077418	0.0060819	0.0074892
0.50	0.0049455	0.0033997	0.0047312
0.45	0.0020850	0.0008425	0.0019073
0.40	0.0001912	-0.0005219	0.0000633
0.35	0.0011802	0.0012908	0.0010889
0.30	0.0084302	0.0096339	0.0083634
0.25	0.0281221	0.0306583	0.0280601
0.20	0.0667380	0.0706031	0.0666712
0.15	0.1275761	0.1322839	0.1275338
0.10	0.1938100	0.1977258	0.1938031
0.05	0.2009478	0.2019970	0.2009947
0.00	0.0000000	0.0000000	0.0000000

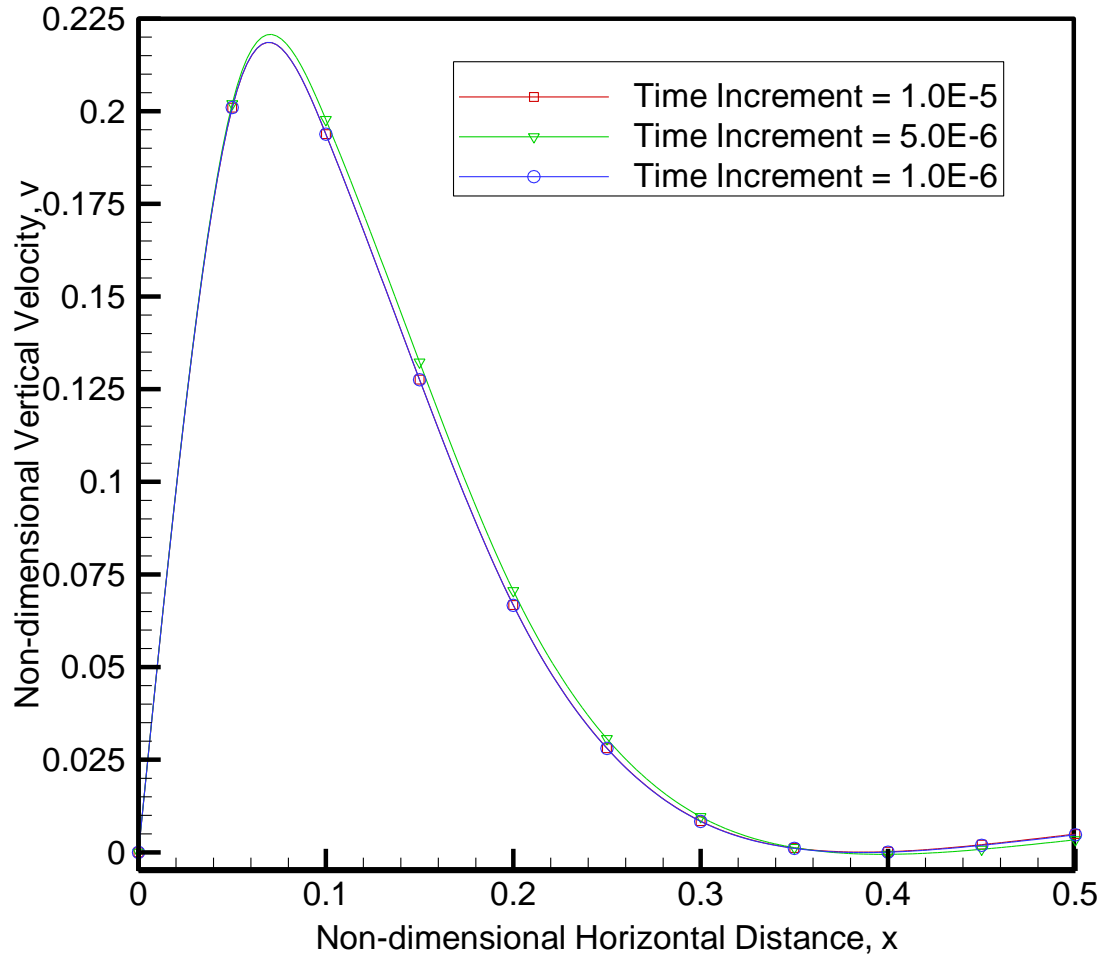


Figure 5- Comparison of Steady State Non-dimensional Vertical Velocity Distribution along the Half of the Horizontal Centerline for $\Delta \bar{t} = 1.0\text{E-}5$, $\Delta \bar{t} = 5.0\text{E-}6$ and $\Delta \bar{t} = 1.0\text{E-}6$ at $Ra = 10^5$

6.1.2. Grid Independence Study

In order to validate the accuracy and convergence of the computer code, a grid independence study for the present study ($Ra = 10^5$) was also conducted. The grid size chosen for the present study was 21×21 . To verify that the converged solutions were independent of the grid chosen one more study was carried out with grid size of 41×41 . Steady state results using uniform, orthogonal 21×21 and 41×41 meshes were obtained using the present computer code. The Comparison of the predicted non-dimensional vertical velocity profiles along the x-axis for grid independence study at $Ra = 10^5$ is presented in Figure 6 and the quantitative comparison is presented in Table 2.

Table 2- Steady State Non-dimensional values of Vertical Velocity along the Horizontal Centerline of the Enclosure for 21 X 21 and 41 X 41 ($Ra = 10^5$ and 21 x 21)

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Vertical Velocity, \bar{v} (21 x 21)	Non-dimensional Vertical Velocity, \bar{v} (41 x 41)	% Deviation in \bar{v}
1.00	0.0000000	0.0000000	0.0000000
0.95	-0.2440615	-0.2412806	-1.1525584
0.90	-0.2006345	-0.2090486	4.0249492
0.85	-0.1048553	-0.1148455	8.6988170
0.80	-0.0364389	-0.0460759	20.915360
0.75	-0.0036786	-0.0116106	68.316591
0.70	0.0081803	0.0014213	-475.55753
0.65	0.0103389	0.0043900	-135.50819
0.60	0.0093217	0.0034348	-171.38688
0.55	0.0069599	0.0011907	-484.48458
0.50	0.0042331	-0.0013476	414.12458
0.45	0.0016037	-0.0037290	143.00665
0.40	-0.0001891	-0.0053963	96.496226
0.35	0.0010085	-0.0050037	120.15436
0.30	0.0082728	0.0007297	-1033.7072
0.25	0.0281219	0.0187185	-50.235167
0.20	0.0668939	0.0602751	-10.981119
0.15	0.1284077	0.1349891	4.8755048
0.10	0.1954522	0.2255438	13.341798
0.05	0.2031731	0.2394256	15.141447
0.00	0.0000000	0.0000000	0.0000000

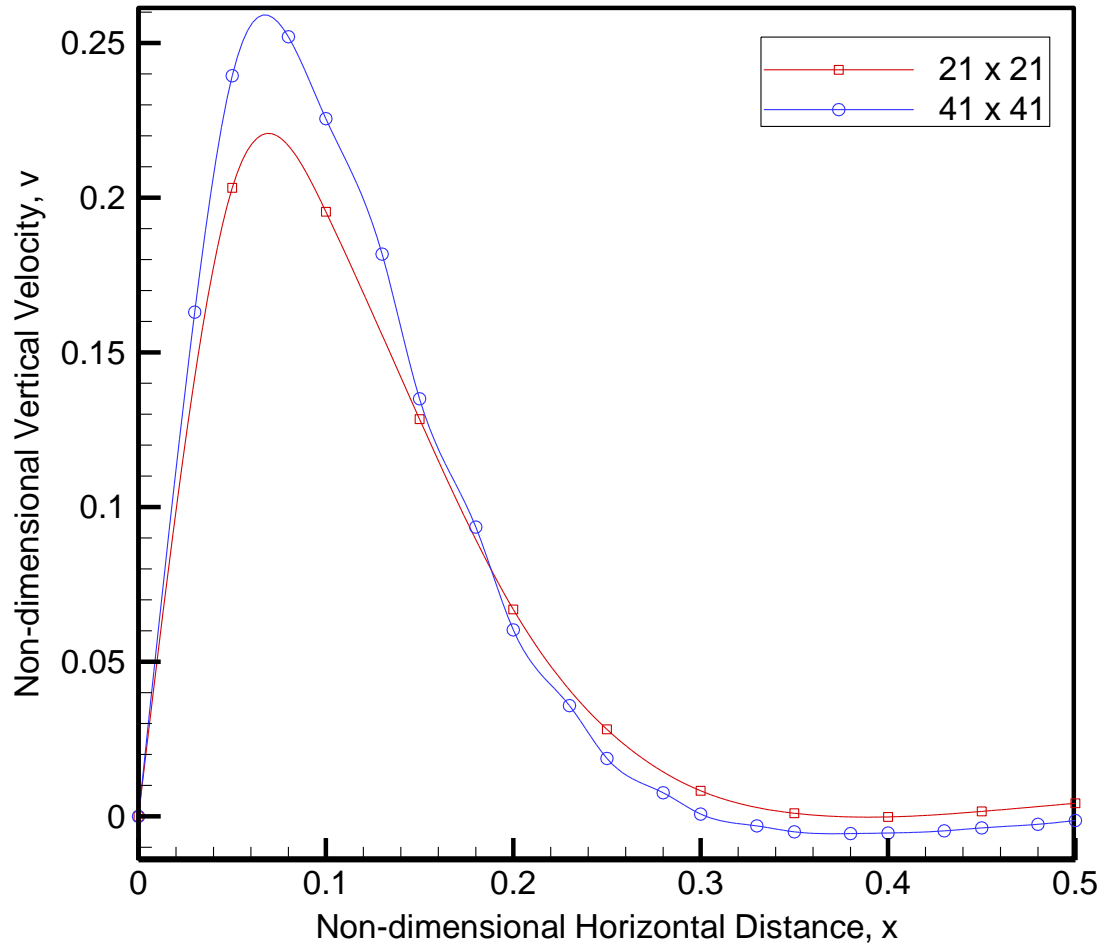


Figure 6- Comparison of Steady State Non-dimensional Vertical Velocity Distribution along the Half of the Horizontal Centerline for Grid Sizes of 21 x 21 and 41 x 41 at $Ra = 10^5$

It is evident in these comparisons that the accuracy of the results increased with finer grid sizes, as expected.

6.1.3. Benchmarking of the Developed Code

After verifying the time convergence and grid dependence studies, the computer code was developed for this study was tested for the accepted benchmark case provided by de Vahl Davis [2] (for incompressible flow with Boussinesq approximation). A quantitative comparison of maximum velocities in the horizontal and vertical directions as predicted by this study and de Vahl Davis is presented below in Table 3 where $\tilde{u} = u H/\alpha$ and $\tilde{v} = v H/\alpha$.

Table 3- Comparison of the Results of the Constant Properties [27] and the Benchmark Case (de Vahl Davis, [2]) for Different Rayleigh Numbers

Rayleigh Number	Constant Properties \tilde{u}_{max}	Benchmark Case \tilde{u}_{max}	% Deviation in \tilde{u}_{max}
10^4	15.772	16.178	- 2.509
10^5	37.241	34.730	+ 7.230

Rayleigh Number	Constant Properties \tilde{v}_{max}	Benchmark Case \tilde{v}_{max}	% Deviation in \tilde{v}_{max}
10^4	18.927	19.617	- 3.517
10^5	68.629	68.590	+ 0.057

A comparison of the results of this study (which is based on compressible flow with constant properties assumption) and that of de Vahl Davis's study [2] indicates observable differences. The absolute value of maximum computed horizontal velocity underestimated corresponding benchmark value by 2.5 % for $Ra = 10^4$. However, for $Ra = 10^5$, same value overestimated by 7.2 %. Similarly, the maximum difference between the benchmark vertical value and the present value is 3.5 % for $Ra = 10^4$, and almost negligible difference in the same value for $Ra = 10^5$.

6.2. Second Order Accurate in Time Model

A numerical study was carried out to determine the differences between the predicted unsteady velocity and temperature profiles, while the circulation patterns were developing after a step change in wall temperatures using the first order and second order accurate in time CMSIP algorithms. Time convergence and grid dependence studies were conducted to verify the accuracy of the second order accuracy in time model.

6.2.1. Time Convergence Study

A time convergence study was carried out for the present study which is second order accurate in time. As presented in section 6.1.1, once the grid size is set, the time computational time convergence is determined by satisfying the diagonal dominance condition of the coefficient matrix. For the present code validation study, a non-dimensional computational time increment of 1×10^{-5} was considered. It was found that a time increment smaller than this did not result in any significant changes in qualitative and quantitative results. To illustrate, non-dimensional vertical velocity distribution along the horizontal centerline of the enclosure were compared for different time increments such as $\Delta \bar{t} = 1 \times 10^{-5}$ and $\Delta \bar{t} = 5 \times 10^{-6}$ for $Ra = 10^5$ and 21×21 mesh points. The quantitative comparison is presented in Table 4 and qualitative comparison is presented in Figure 7.

Table 4- Steady State Non-dimensional values of Vertical Velocity along the Horizontal Centerline of the Enclosure for $\Delta\bar{t} = 1.0\text{E-}5$ and $\Delta\bar{t} = 5.0\text{E-}6$ at $Ra = 10^5$ in Second Order Time Accurate Approximation

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Vertical Velocity, \bar{v} ($\Delta\bar{t} = 1.0\text{E-}5$)	Non-dimensional Vertical Velocity, \bar{v} ($\Delta\bar{t} = 5.0\text{E-}6$)
1.00	0.0000000	0.0000000
0.95	-0.2440447	-0.2461275
0.90	-0.2002889	-0.2023205
0.85	-0.1043035	-0.1047269
0.80	-0.0357540	-0.0352146
0.75	-0.0031484	-0.0024742
0.70	0.0086785	0.0092366
0.65	0.0106349	0.0110582
0.60	0.0095869	0.0099664
0.55	0.0070514	0.0074107
0.50	0.0043028	0.0046463
0.45	0.0015101	0.0017887
0.40	-0.0003056	-0.0001617
0.35	0.0007269	0.0006464
0.30	0.0079743	0.0075860
0.25	0.0276678	0.0269726
0.20	0.0664488	0.0656514
0.15	0.1278468	0.1274489
0.10	0.1949693	0.1957688
0.05	0.2027110	0.2048181
0.00	0.0000000	0.0000000

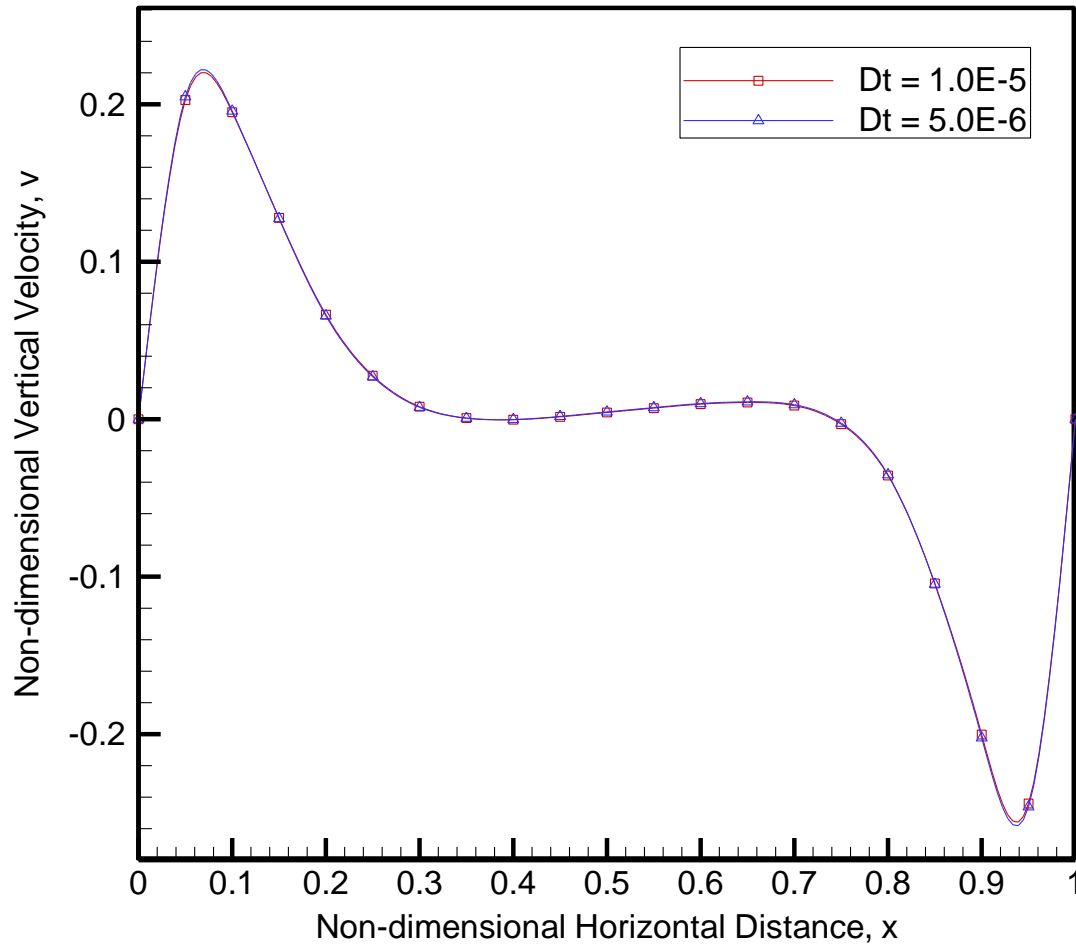


Figure 7- Steady State Non-dimensional Vertical Velocity Distribution along the Horizontal Centerline for $\Delta \bar{t} = 1.0\text{E-}5$, $\Delta \bar{t} = 5.0\text{E-}6$ and $\Delta \bar{t} = 1.0\text{E-}6$ at $Ra = 10^5$ in Second Order Time Accurate Approximation

6.2.2. Grid Independence Study

In order to validate the accuracy and convergence of the previous published results in the first order accurate in time, a grid independence study for the present study ($Ra = 10^5$) was also conducted. Steady state results using uniform, orthogonal 21×21 and 41×41 meshes were obtained using the present computer code. The centerline non-dimensional vertical velocity profiles for are presented in Figure 8 and in Table 5, quantitative comparison is presented. It is evident in these comparisons that the accuracy of the results increased with finer grid sizes.

Table 5- Steady State Non-dimensional values of Vertical Velocity along the Horizontal Centerline of the Enclosure for 21 X 21 and 41 X 41 at $Ra = 10^5$ in Second Order Time Accurate Approximation

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Vertical Velocity, \bar{v} (21 x 21)	Non-dimensional Vertical Velocity, \bar{v} (41 x 41)	% Deviation in \bar{v}
1.00	0.0000000	0.0000000	0.0000000
0.95	-0.2440447	-0.2408513	-1.3258803
0.90	-0.2002889	-0.2085564	3.9641554
0.85	-0.1043035	-0.1143165	8.7590155
0.80	-0.0357540	-0.0455260	21.464604
0.75	-0.0031484	-0.0110939	71.619711
0.70	0.0086785	0.001856	-367.35998
0.65	0.0106349	0.0047221	-125.21089
0.60	0.0095869	0.0036587	-162.02546
0.55	0.0070514	0.0013085	-438.88896
0.50	0.0043028	-0.0013337	422.59862
0.45	0.0015100	-0.0038197	139.53328
0.40	-0.0003056	-0.0055958	94.538590
0.35	0.0007269	-0.0053183	113.66860
0.30	0.0079743	0.0002984	-2571.8144
0.25	0.0276678	0.0181909	-52.097092
0.20	0.0664488	0.0597164	-11.273841
0.15	0.1278468	0.1345171	4.9587004
0.10	0.1949693	0.2252853	13.456714
0.05	0.2027110	0.2394137	15.330242
0.00	0.0000000	0.0000000	0.0000000

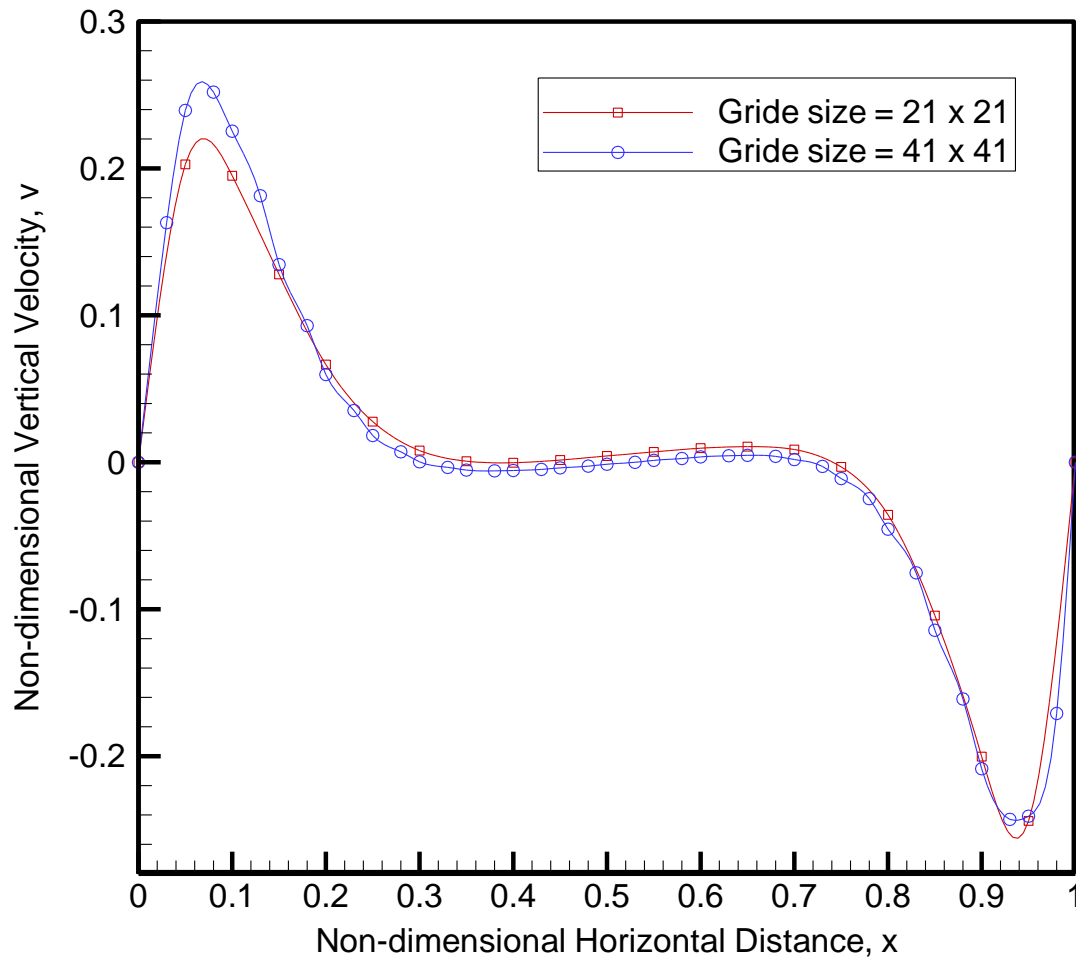


Figure 8- Steady State Non-dimensional Vertical Velocity Distribution along the Horizontal Centerline for Grid Sizes of 21 x 21, 41 x 41 and 81 x 81 at $Ra = 10^5$ in Second Order Time Accurate Approximation

6.2.3. Base Model Comparison with First Order Accurate in Time Model

The present study employs a second order accurate in time CMSIP procedure [74, 75, and 76] to solve the unsteady governing equations of the problem. This new algorithm uses Taylor expansion based second order accurate differencing of temporal terms and a second order accurate central difference approximations for all the spatial terms in the governing equations. Previous CMSIP algorithm used by Akyuzlu [27] and also by Chen and Pletcher [67] are first order accurate in time. The comparison of maximum values of horizontal and vertical velocities in first order and second order accurate in time models are presented in the Table 6.

Table 6- Comparison of the Results of the First Order and the Second Order Time Accurate CMSIP Models using Constant Properties for Different Rayleigh Numbers

Rayleigh Number	Second Order \bar{u}_{max}	First Order \bar{u}_{max}	% Deviation in \bar{u}_{max}
10^4 (Location)	0.154 ($\bar{y} = 0.782$)	0.155 ($\bar{y} = 0.780$)	- 0.649
10^5 (Location)	0.109 ($\bar{y} = 0.840$)	0.110 ($\bar{y} = 0.838$)	- 0.909

Rayleigh Number	Second Order \bar{v}_{max}	First Order \bar{v}_{max}	% Deviation in \bar{v}_{max}
10^4 (Location)	0.900 ($\bar{x} = 0.125$)	0.900 ($\bar{x} = 0.125$)	+0.000
10^5 (Location)	0.220 ($\bar{x} = 0.070$)	0.220 ($\bar{x} = 0.070$)	+ 0.000

The results of the second order accurate in time model has almost negligible difference with that of first order accurate in time model. The maximum computed values of horizontal velocity (non-dimensional), corresponding first order accurate in time model are deviated by 0.6 % for $Ra = 10^4$ and 0.9 % for $Ra = 10^5$. However, these differences in vertical velocities are negligible.

Chapter 7

Parametric Study of Natural Convection using the First Order Accurate in Time Model

A parametric study was conducted using the constant property model to investigate how the flow fields are affected by different Rayleigh numbers and wall temperature differences. In this chapter two case studies are presented.

7.1. The Effects of Variations in Rayleigh Numbers

In the first case of parametric study, the effects of variations in Rayleigh number was observed by changing the wall temperature difference at constant size of the enclosure. In this study, for different Rayleigh numbers (10^4 , 5×10^4 , 10^5 , and 5×10^5) the non dimensional vertical velocity distribution along the horizontal centerline of the enclosure was compared and it is concluded that there was a notable difference in magnitudes of velocities for higher Rayleigh numbers. As Rayleigh number increases the location of the maximum velocity point moves towards the hot side of the enclosure with increase in magnitude. The results are presented quantitatively and qualitatively in Table 7 and Figure 9, respectively.

Table 7-Steady State Non-dimensional values of Vertical Velocity along the Horizontal Centerline of the Enclosure for $Ra = 10^4$, $Ra = 5 \times 10^4$, $Ra = 10^5$, $Ra = 5 \times 10^5$ at $H=0.0254m$

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Vertical Velocity, \bar{v} ($Ra = 1.0E4$)	Non-dimensional Vertical Velocity, \bar{v} ($Ra = 5.0E4$)	Non-dimensional Vertical Velocity, \bar{v} ($Ra = 1.0E5$)	Non-dimensional Vertical Velocity, \bar{v} ($Ra = 5.0E5$)
1.00	0.0000000	0.0000000	0.0000000	0.0000000
0.95	-0.1453714	-0.2379227	-0.2465031	-0.2087094
0.90	-0.2075329	-0.2362243	-0.2011154	-0.0919643
0.85	-0.2082459	-0.1565491	-0.1036450	-0.0152543
0.80	-0.1759759	-0.0801576	-0.0343975	0.0045337
0.75	-0.1342380	-0.0316080	-0.0019367	0.0073687
0.70	-0.0945000	-0.0380077	0.0096405	0.0087798
0.65	-0.0614451	0.0038054	0.0114227	0.0097664
0.60	-0.0348230	0.0069824	0.0103105	0.0107504
0.55	-0.0129692	0.0067456	0.0077417	0.0109685
0.50	-0.0129692	0.0058315	0.0049454	0.0102460
0.45	0.0259200	0.0056771	0.0020850	0.0080060
0.40	0.0471586	0.0081480	0.0001911	0.0037837
0.35	0.0720149	0.0158039	0.0011802	-0.0022495
0.30	0.1011022	0.0318961	0.0084302	-0.0085098
0.25	0.1340680	0.0610267	0.0281221	-0.0087953
0.20	0.1668373	0.1055469	0.0667380	0.0085259
0.15	0.1915537	0.1621978	0.1275761	0.0582169
0.10	0.1922468	0.2081509	0.1938100	0.1449731
0.05	0.1453591	0.1878341	0.2009478	0.2161246
0.00	0.0000000	0.0000000	0.0000000	0.0000000

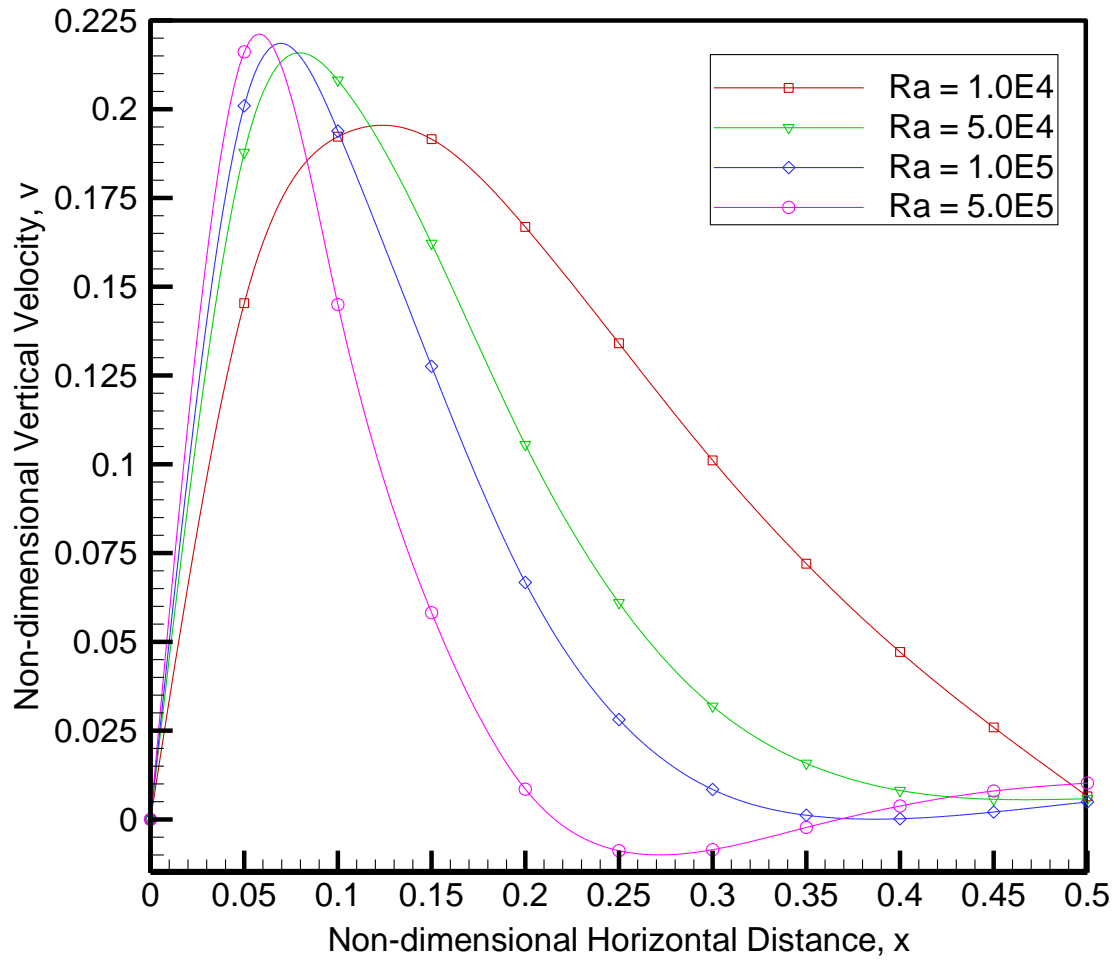


Figure 9- Comparison of Steady State Non-dimensional Vertical Velocity Distributions along the Horizontal Centerline of the Enclosure for $Ra= 10^4$, $Ra=5 \times 10^4$, $Ra= 10^5$, $Ra= 5 \times 10^5$ at $H=0.0254m$

7.2. The Effects of Variations in Wall Temperature Differences

In the second case of the parametric study, the effects of variations in wall temperature difference were observed by changing the size of the enclosure at constant Rayleigh number 10^5 . In this study, for different wall temperature difference (64.8 K, 356.6 K, and 648.4 K) the non dimensional vertical velocity distributions along the horizontal centerline of the enclosure were compared and it is concluded that there is notable difference in velocity magnitudes for higher temperature difference. As wall temperature difference increases the location of the maximum velocity moves towards the cold side of the enclosure with increase in magnitude. The results are presented quantitatively and qualitatively in Table 8 and Figure 10, respectively.

Table 8- Steady State Non-dimensional values of Vertical Velocity along the Horizontal Centerline of the Enclosure for $\Delta T= 64.8K$, $\Delta T= 356.6K$, $\Delta T= 648.4K$ at $Ra= 10^5$

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Vertical Velocity, \bar{v} ($\Delta T= 64.8K$)	Non-dimensional Vertical Velocity, \bar{v} ($\Delta T= 356.6K$)	Non-dimensional Vertical Velocity, \bar{v} ($\Delta T= 648.4K$)
1.00	0.0000000	0.0000000	0.0000000
0.95	-0.2465031	-0.2084725	-0.2020537
0.90	-0.2011154	-0.1727811	-0.1657409
0.85	-0.1036450	-0.0984031	-0.1009974
0.80	-0.0343975	-0.0445509	-0.0577585
0.75	-0.0019367	-0.0179709	-0.0366760
0.70	0.0096405	-0.0078904	-0.0281754
0.65	0.0114227	-0.0064076	-0.0281754
0.60	0.0103105	-0.0077848	-0.0229883
0.55	0.0077417	-0.0099892	-0.0192403
0.50	0.0049454	-0.0113427	-0.0120143
0.45	0.0020850	-0.0107910	0.0017215
0.40	0.0001911	-0.0066578	0.0244202
0.35	0.0011802	0.0042894	0.0606051
0.30	0.0084302	0.0266349	0.1129162
0.25	0.0281221	0.0673053	0.1836736
0.20	0.0667380	0.1305537	0.2656362
0.15	0.1275760	0.2118118	0.3355243
0.10	0.1938100	0.2762716	0.3393019
0.05	0.2009478	0.2377315	0.2282922
0.00	0.0000000	0.0000000	0.0000000

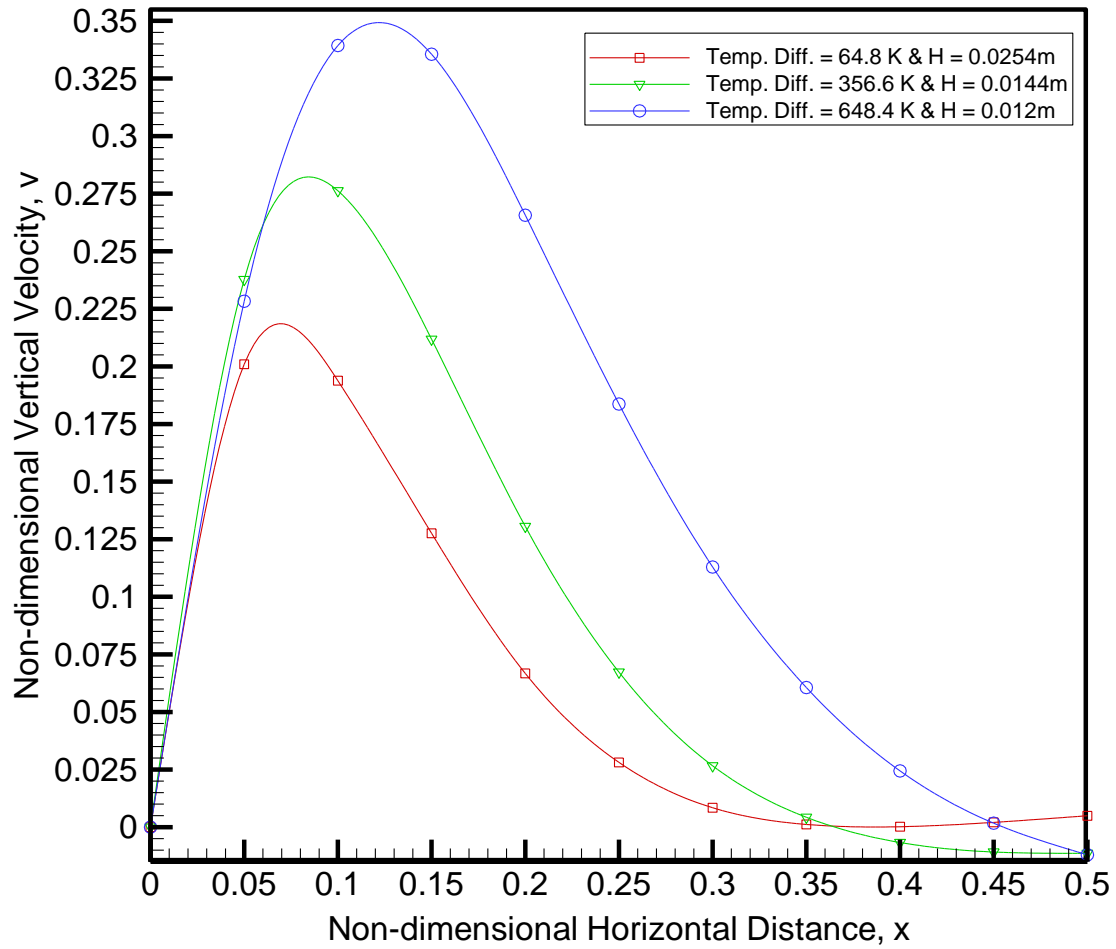


Figure 10- Comparison of Steady State Non-dimensional Vertical Velocity Distributions along the Horizontal Centerline of the Enclosure for $\Delta T= 64.8\text{K}$, $\Delta T= 356.6\text{K}$ and $\Delta T= 648.4\text{K}$ at $Ra= 10^5$

Chapter 8

The Study of Effects of Variable Fluid Properties in Natural Convection inside a Square Enclosure

In this chapter, the results of the constant property model are compared to that of variable property model. Then the variable property model was proposed as a base model in the present study. Also, the effects of variations in Rayleigh number for variable property model inside a square enclosure are presented in this chapter.

8.1. The Study of Base Model

In this study, the results of the constant property and variable property models are compared for $Ra = 10^5$ ($\Delta T = 64.8$ K). There is notable difference in the predicted values of velocity and temperature. These differences are presented quantitatively in Tables 9 through 11. The comparisons between these models are presented in terms of velocity and temperature distributions in Figure 11 through 13. From Figure 11, the predicted values of non-dimensional horizontal velocities along the vertical centerline of the enclosure are observed higher for variable property model at cold side of the enclosure and lower at hot side of the enclosure compared to constant property model. However, non-dimensional vertical velocity and temperature values along the horizontal centerline of the enclosure are predicted higher at hot side of the enclosure and lower at cold side of the enclosure for variable property model as shown in Figure 12 and 13. Horizontal and vertical velocity histograms are also compared in Figure 14 and 15, respectively. Horizontal velocities are predicted higher for variable property model and vertical velocities are predicted higher for constant property model.

The variable property model is proposed as the base model and is considered for $Ra\ 10^5$ with uniform, orthogonal mesh of 21×21 , in which wall temperature difference is 64.8 K. The predicted primitive variables along the horizontal centerline and vertical centerline of the enclosure are presented in Table 12 and Table 13, respectively. The computational domain used in the present study is shown in Figure 16. The velocity (u and v) histograms at a point closer to the high temperature wall where $\bar{x} = 0.05$ and $\bar{y} = 0.5$ are presented in Figures 17 and 18. The temperature stratification is shown in Figure 19. Temperature gradients are higher near the walls and are close to zero near the center. This fact is easily illustrated by the isotherms being almost horizontal in the center of the enclosure. The formation of the two vortices can be seen in Figure 20. The flow is not unicellular; instead, it separates into boundary layer and core flow. The core flow is more stagnant relative to the boundary regions. This characteristic is also observed in the streamlines shown in Figure 21. The non-dimensional horizontal velocity distribution along the vertical centerline, vertical velocity distribution along the horizontal centerline and temperature distribution along the horizontal centerline are presented in Figure 22, 23 and 24, respectively.

Table 9- Comparison of Predicted Non-dimensional values of Horizontal Velocity along the Vertical Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 10^5$ ($\Delta T = 64.8$ K, $H=0.0254$ m and 21×21)

Non-dimensional Vertical Distance, \bar{y}	Non-dimensional Horizontal Velocity, \bar{u} (Constant)	Non-dimensional Horizontal Velocity, \bar{u} (Variable)	% Deviation in \bar{u}
1.00	0.0000000	0.0000000	0.0000000
0.95	0.0535207	0.0571390	6.3324524
0.90	0.0810651	0.0876302	7.4917753
0.85	0.1004595	0.1072624	6.3422970
0.80	0.1009580	0.1071562	5.7842663
0.75	0.0929973	0.0981868	5.2853213
0.70	0.0789586	0.0821039	3.8308866
0.65	0.0645284	0.0656531	1.7130632
0.60	0.0495355	0.0484717	-2.1946804
0.55	0.0381889	0.0351714	-8.5795203
0.50	0.0241282	0.0201633	-19.663835
0.45	0.0125497	0.0079409	-58.039436
0.40	-0.0046300	-0.0088292	47.559387
0.35	-0.0228675	-0.0268213	14.741480
0.30	-0.0460588	-0.0497550	7.4286760
0.25	-0.0767948	-0.0800412	4.0558706
0.20	-0.1056613	-0.1087645	2.8531368
0.15	-0.1387613	-0.1398505	0.7788316
0.10	-0.1406555	-0.1405555	-0.0711462
0.05	-0.0878570	-0.0882356	0.4290444
0.00	0.0000000	0.0000000	0.0000000

Table 10- Comparison of Predicted Non-dimensional values of Vertical Velocity along the Horizontal Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 10^5$ ($\Delta T = 64.8$ K, $H=0.0254$ m and 21×21)

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Vertical Velocity, \bar{v} (Constant)	Non-dimensional Vertical Velocity, \bar{v} (Variable)	% Deviation in \bar{v}
1.00	0.0000000	0.0000000	0.0000000
0.95	-0.2502395	-0.2475826	-1.0731368
0.90	-0.1944804	-0.1971463	1.3522445
0.85	-0.0930081	-0.1008293	7.7567929
0.80	-0.0257422	0.0334961	176.85118
0.75	0.0031427	-0.0018893	266.33759
0.70	0.0122860	0.0094903	-29.458318
0.65	0.0130391	0.0111882	-16.543233
0.60	0.0117409	0.0099772	-17.677062
0.55	0.0092661	0.0073344	-26.337007
0.50	0.0065152	0.0045396	-43.518607
0.45	0.0034166	0.0018026	-89.540905
0.40	0.0008496	0.0001795	-373.20101
0.35	0.0005829	0.0016220	64.059846
0.30	0.0059338	0.0095263	37.710790
0.25	0.0231456	0.0300361	22.940638
0.20	0.0591808	0.0693695	14.687566
0.15	0.1184252	0.1299771	8.8876425
0.10	0.1861128	0.1934348	3.7852547
0.05	0.1988817	0.1948903	-2.0480239
0.00	0.0000000	0.0000000	0.0000000

Table 11- Comparison of Predicted Non-dimensional values of Temperature along the Horizontal Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 10^5$ ($\Delta T = 64.8$ K, $H=0.0254$ m and 21×21)

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Temperature, \bar{T} (Constant)	Non-dimensional Temperature, \bar{T} (Variable)	% Deviation in \bar{T}
1.00	0.8907647	0.8907647	0.0000000
0.95	0.9560220	0.9560723	0.0000005
0.90	1.0042400	1.0029240	-0.0000131
0.85	1.0271800	1.0256300	-0.0000113
0.80	1.0341910	1.0331320	-0.0000102
0.75	1.0347690	1.0340560	-0.0000068
0.70	1.0339310	1.0332200	-0.0000068
0.65	1.0333490	1.0324500	-0.0000087
0.60	1.0331680	1.0320330	-0.0000110
0.55	1.0332070	1.0318370	-0.0000132
0.50	1.0331700	1.0315720	-0.0000155
0.45	1.0327950	1.0309780	-0.0000176
0.40	1.0318170	1.0298150	-0.0000194
0.35	1.0300450	1.0279460	-0.0000204
0.30	1.0275110	1.0255040	-0.0000196
0.25	1.0247950	1.0232280	-0.0000153
0.20	1.0236580	1.0230580	-0.0000059
0.15	1.0276350	1.0285710	0.0000091
0.10	1.0420280	1.0445690	0.0000243
0.05	1.0709450	1.0736000	0.0000247
0.00	1.1092350	1.1092350	0.0000000

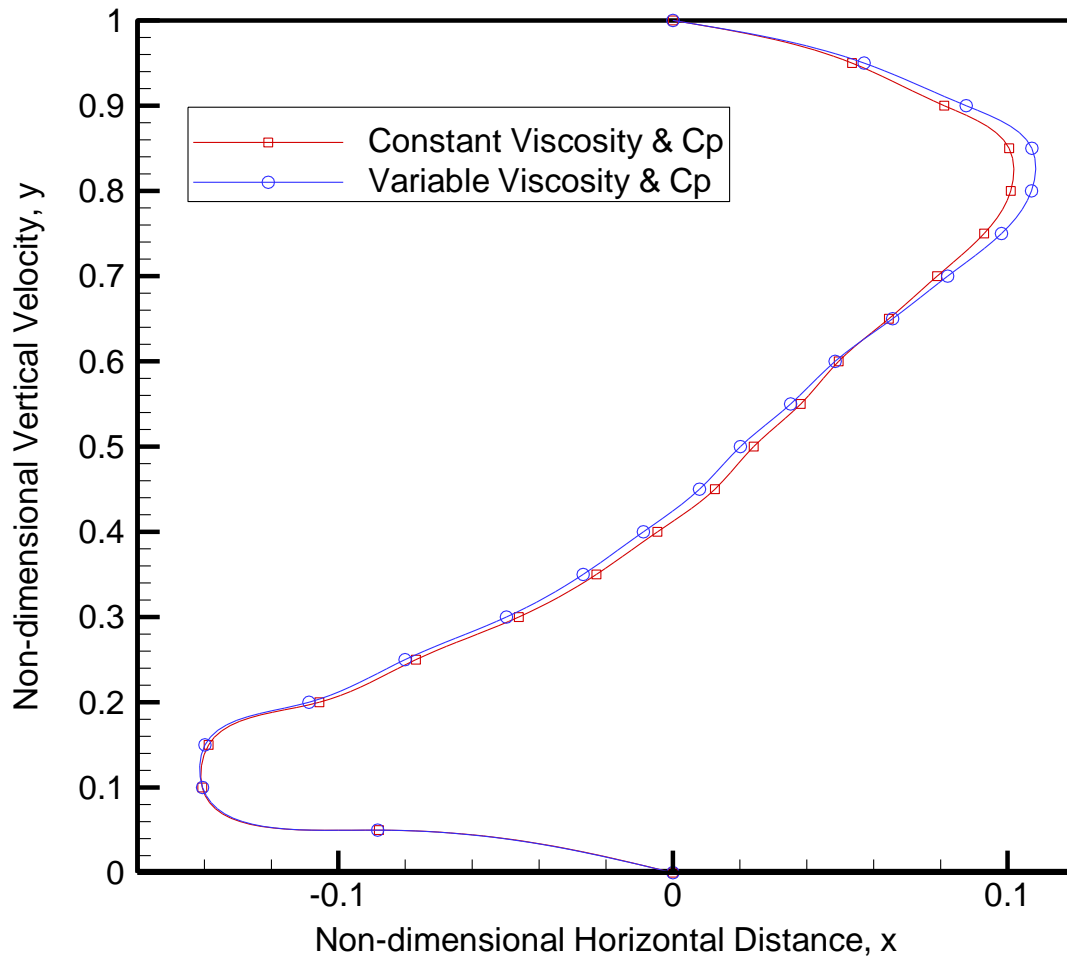


Figure 11- Comparison of Steady State Non-dimensional Horizontal Velocity Distributions along the Vertical Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 10^5$ ($\Delta T = 64.8$ K, $H = 0.0144$ m)

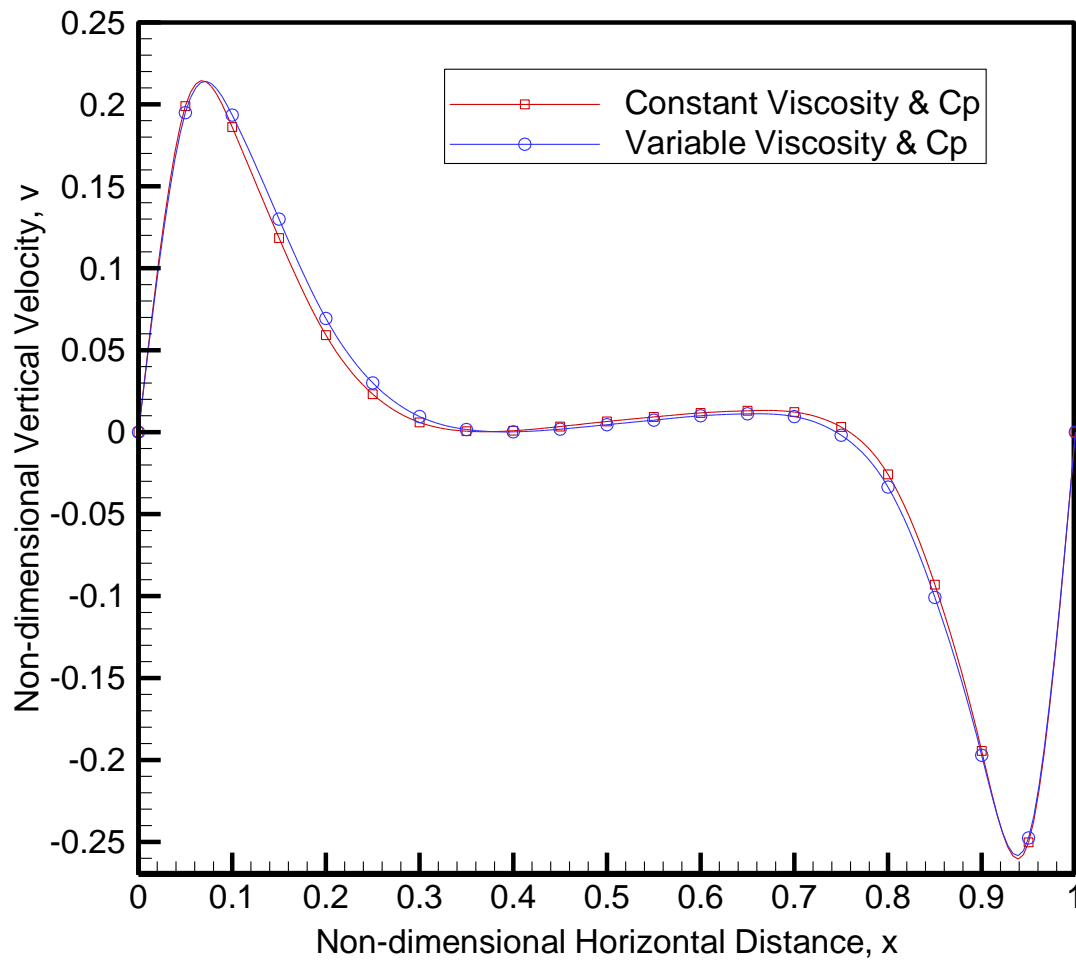


Figure 12- Comparison of Steady State Non-dimensional Vertical Velocity Distributions along the Horizontal Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 10^5$ ($\Delta T = 64.8$ K, $H = 0.0144$ m)

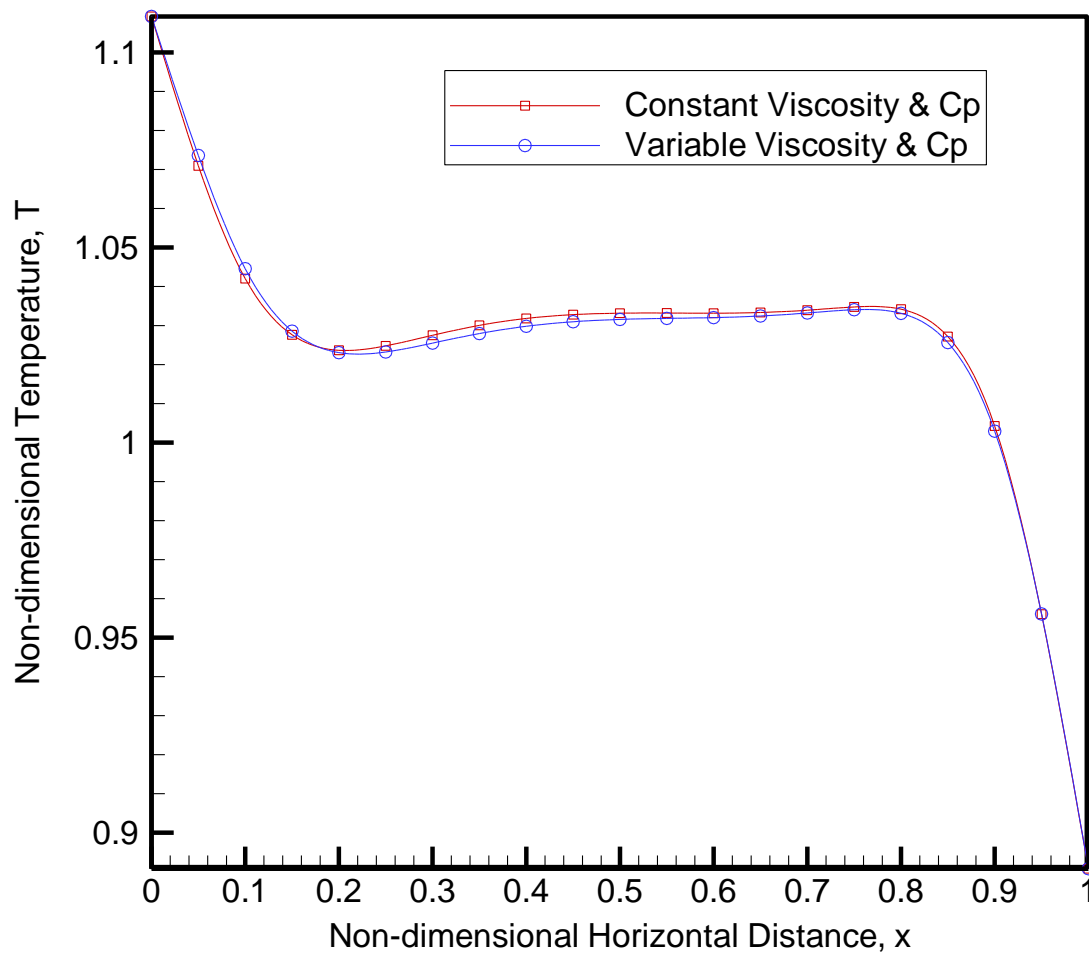


Figure 13- Comparison of Steady State Non-dimensional Temperature Distributions along the Horizontal Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 10^5$ ($\Delta T = 64.8$ K, $H = 0.0144$ m)

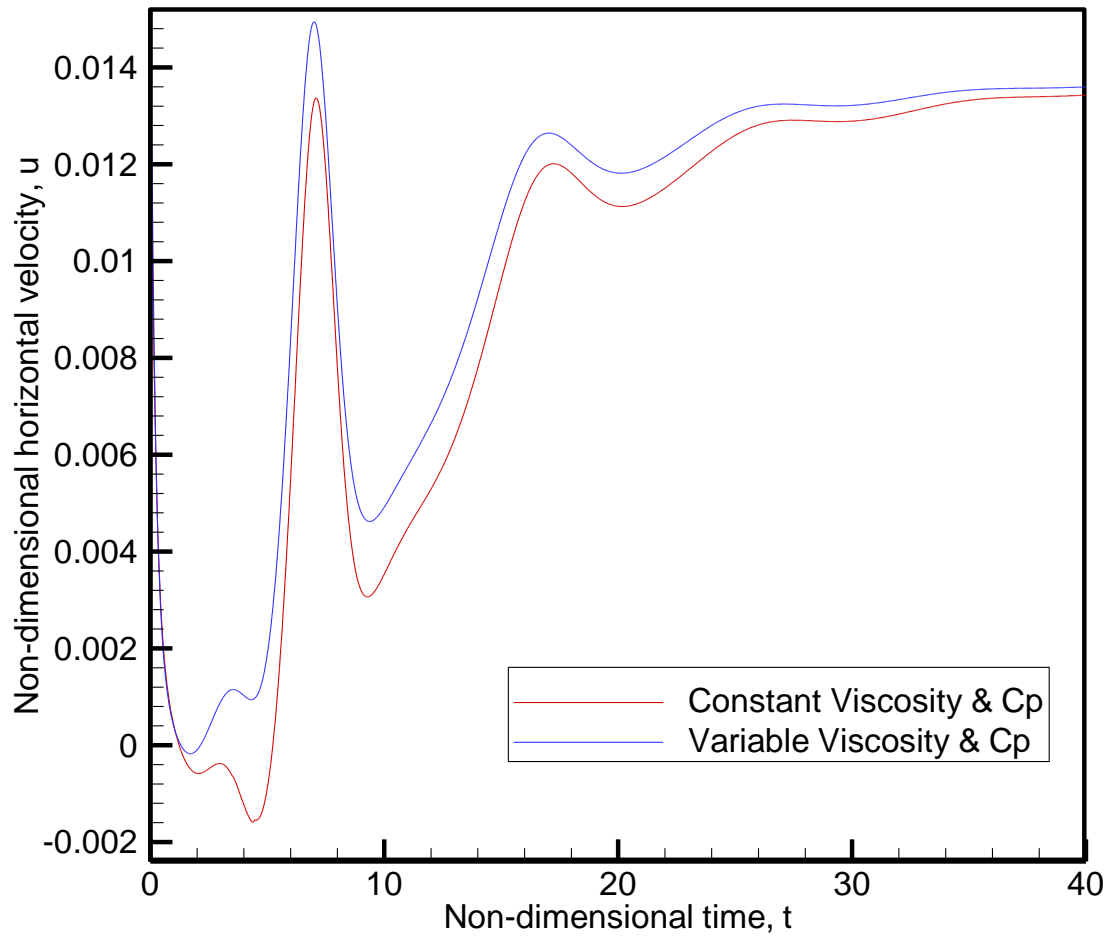


Figure 14- Comparison of Non-dimensional Horizontal Histograms of Constant and Variable Properties for $Ra= 10^5$ at $\bar{X}=0.05$ and $\bar{Y}=0.5$

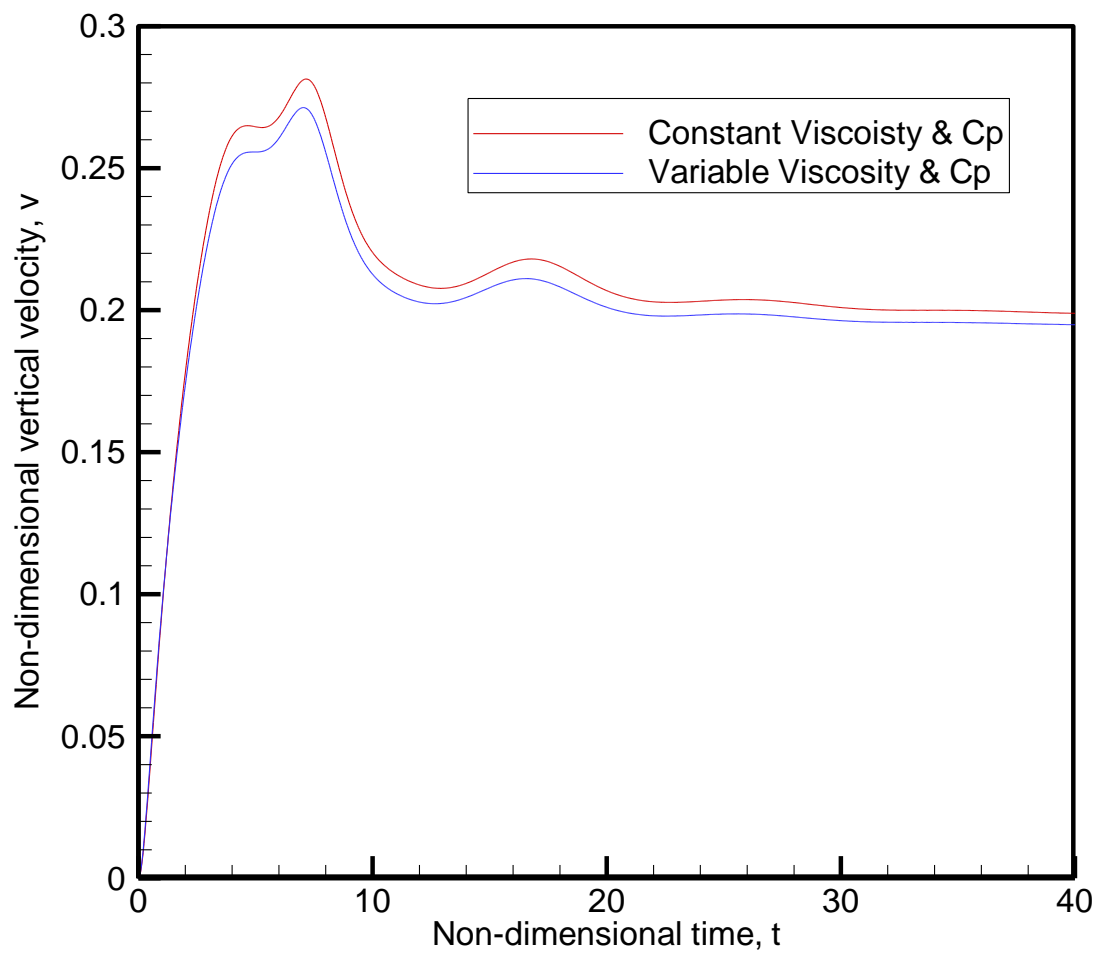


Figure 15- Comparison of Vertical Velocity Histograms of Constant and Variable Properties for $Ra=10^5$ at $\bar{X}=0.05$ and $\bar{Y}=0.5$

Table 12- Steady State Non-dimensional values of Primitive Variables along the Horizontal Centerline of the Enclosure for $Ra = 10^5$ ($\Delta\bar{t} = 1.0E-5$ and 21×21)

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Vertical Velocity, \bar{v}	Non-dimensional Pressure, \bar{p}	Non-dimensional Temperature, \bar{T}
1.00	0.0000000	1622.1422	0.8907647
0.95	-0.2475826	1622.1422	0.9560722
0.90	-0.1971463	1622.1422	1.0029243
0.85	-0.1008293	1622.1422	1.0256303
0.80	-0.0334962	1622.1422	1.0331324
0.75	-0.0018894	1622.1422	1.0340564
0.70	0.0094903	1622.1422	1.0332203
0.65	0.0111883	1622.1422	1.0324496
0.60	0.0099773	1622.1422	1.0320333
0.55	0.0073344	1622.1422	1.0318368
0.50	0.0045396	1622.1422	1.0315717
0.45	0.0018026	1622.1422	1.0309784
0.40	0.0001795	1622.1422	1.0298149
0.35	0.0016221	1622.1422	1.0279456
0.30	0.0095263	1622.1422	1.0255036
0.25	0.0300361	1622.1422	1.0232275
0.20	0.1299771	1622.1422	1.0230579
0.15	0.1299770	1622.1422	1.0285707
0.10	0.1934348	1622.1422	1.0445688
0.05	0.1948903	1622.1422	1.0736002
0.00	0.0000000	1622.1422	1.1092353

Table 13- Steady State Non-dimensional values of Primitive Variables along the Vertical Centerline of the Enclosure for $Ra = 10^5$ ($\Delta\bar{t} = 1.0E-5$ and 21×21)

Non-dimensional Vertical Distance, \bar{y}	Non-dimensional Horizontal Velocity, \bar{u}	Non-dimensional Pressure, \bar{p}	Non-dimensional Temperature, \bar{T}
1.00	0.0000000	1622.1400	1.0854359
0.95	0.0571390	1622.1401	1.0853203
0.90	0.0876302	1622.1404	1.0843379
0.85	0.1072624	1622.1406	1.0815360
0.80	0.1071562	1622.1409	1.0778722
0.75	0.0981868	1622.1410	1.0720498
0.70	0.0821039	1622.1413	1.0663142
0.65	0.0656531	1622.1415	1.0585184
0.60	0.0484717	1622.1417	1.0510841
0.55	0.0351714	1622.1420	1.0410634
0.50	0.0201634	1622.1422	1.0315717
0.45	0.0079409	1622.1424	1.0192259
0.40	-0.0088292	1622.1426	1.0082883
0.35	-0.0268214	1622.1429	0.9951756
0.30	-0.0497550	1622.1431	0.9846475
0.25	-0.0800413	1622.1434	0.9729092
0.20	-0.1087645	1622.1436	0.9638571
0.15	-0.1398505	1622.1439	0.9534625
0.10	-0.1405555	1622.1440	0.9461181
0.05	-0.0882356	1622.1444	0.9393139
0.00	0.0000000	1622.1446	0.9363195

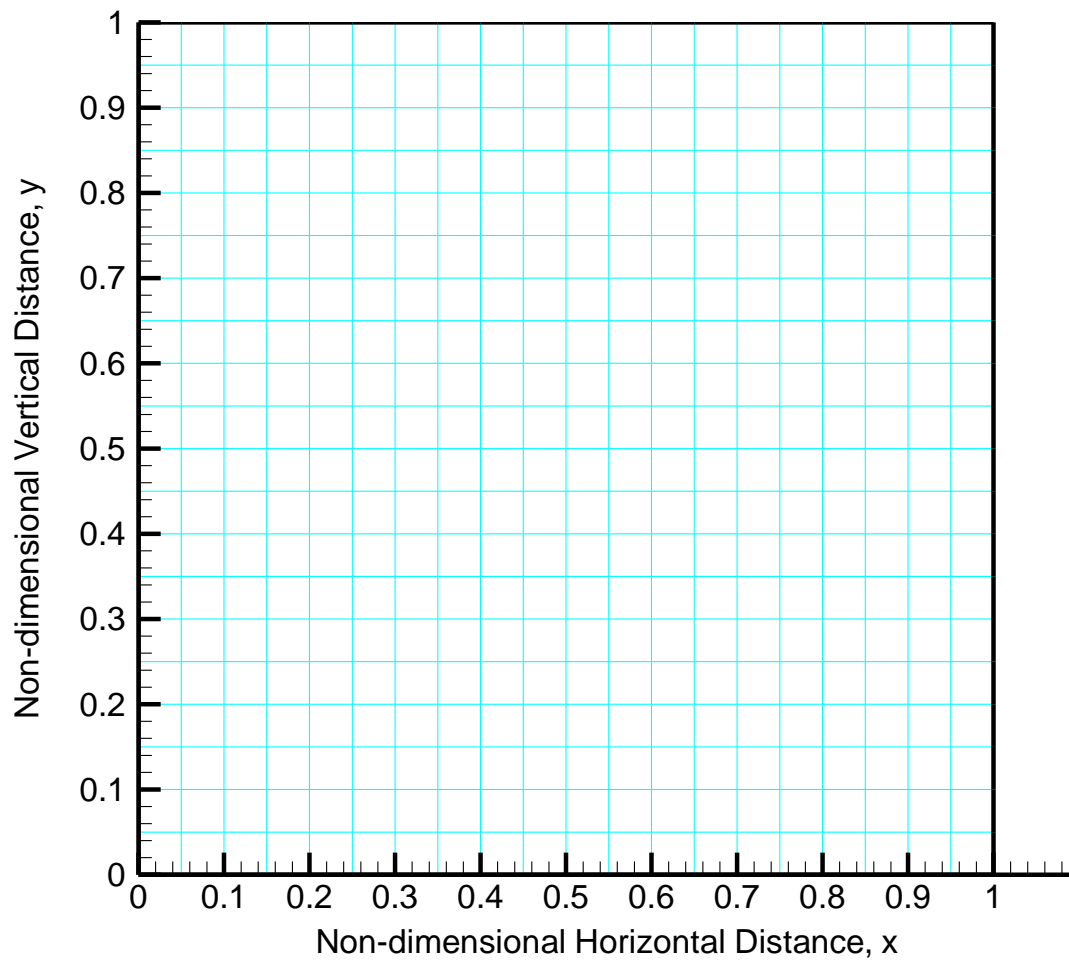


Figure 16- Computational Mesh of the Non-dimensional domain for the Natural Convection Case for Variable Properties

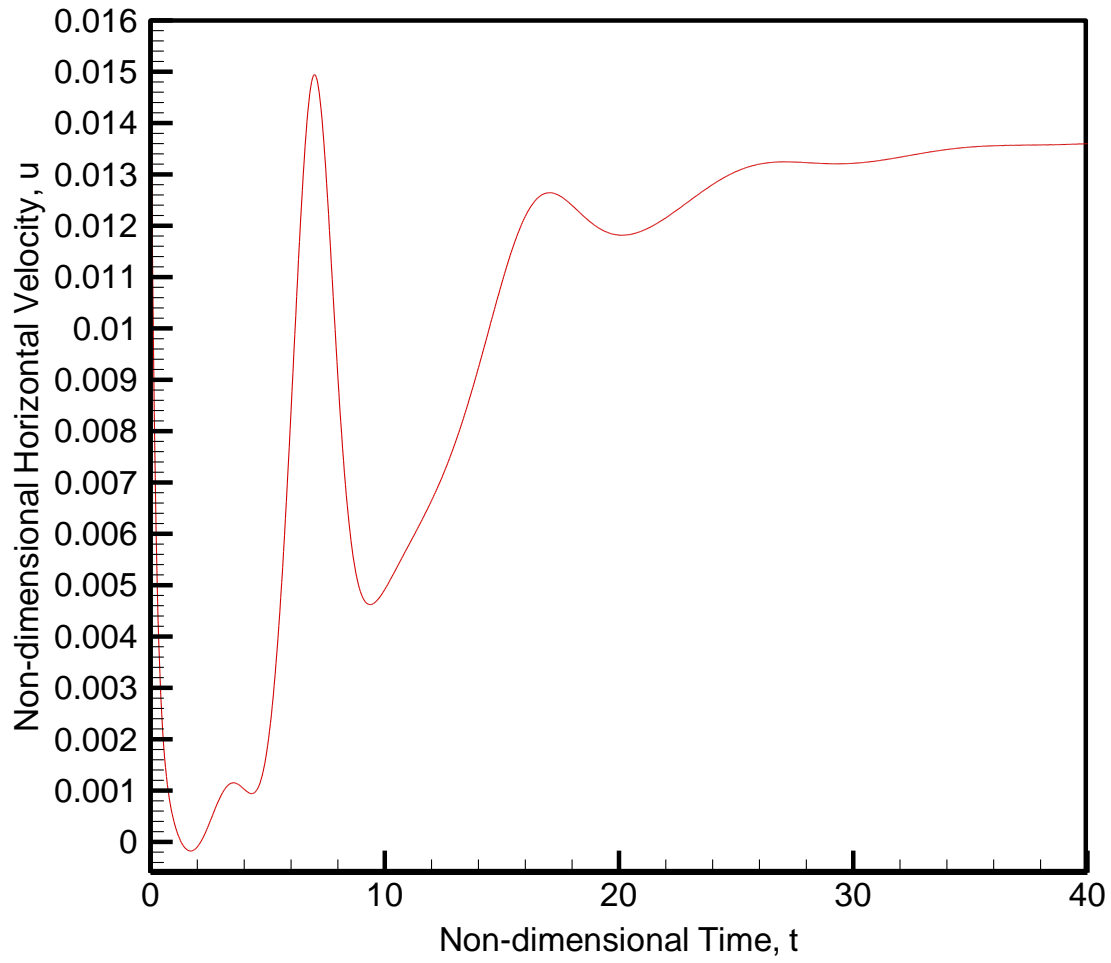


Figure 17- Non-dimensional Horizontal Velocity Histogram for $Ra=10^5$ at $\bar{X}=0.05$ and $\bar{Y}=0.5$ for Variable Properties (21 x 21, $\Delta \bar{t} = 1.0E-5$)

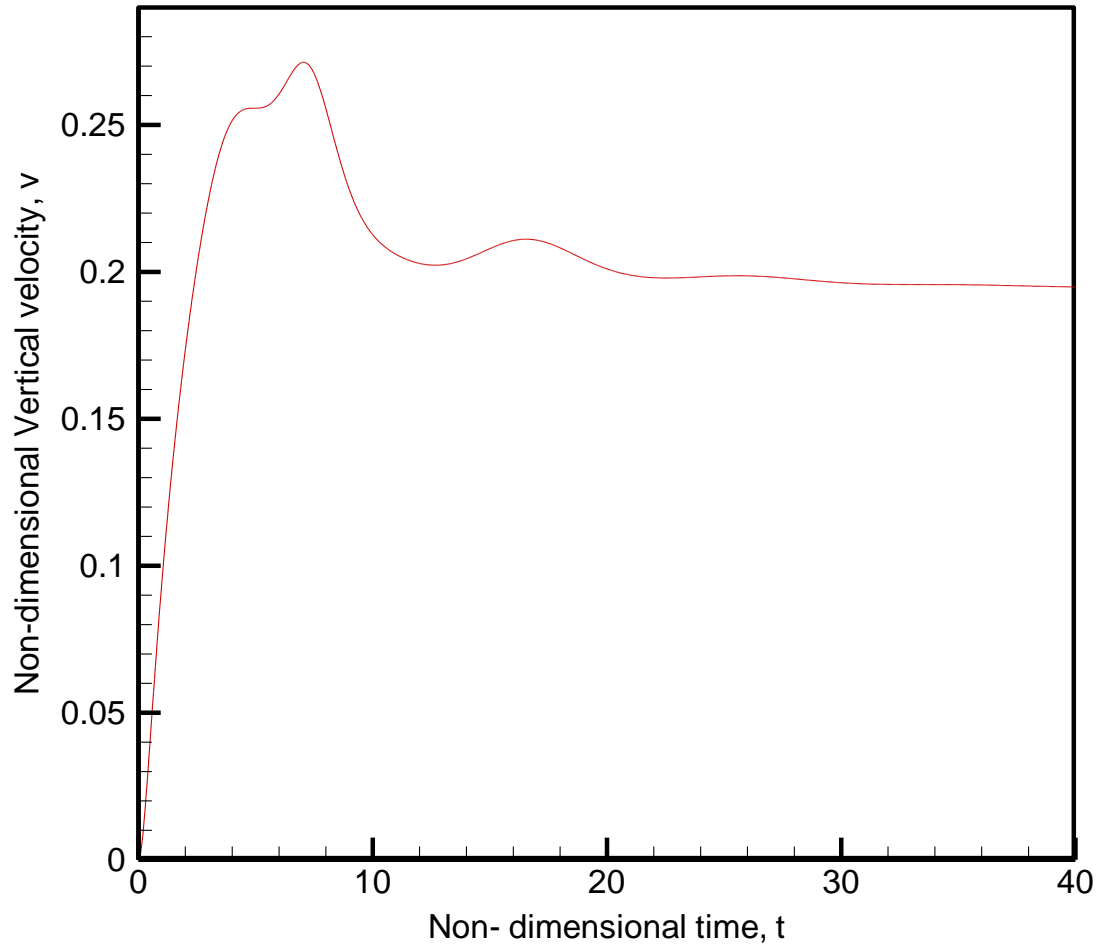


Figure 18- Non-dimensional Vertical Velocity Histogram for $Ra=10^5$ at $\bar{X}=0.05$ and $\bar{Y}=0.5$ for Variable Properties (21×21 , $\Delta \bar{t} = 1.0E-5$)

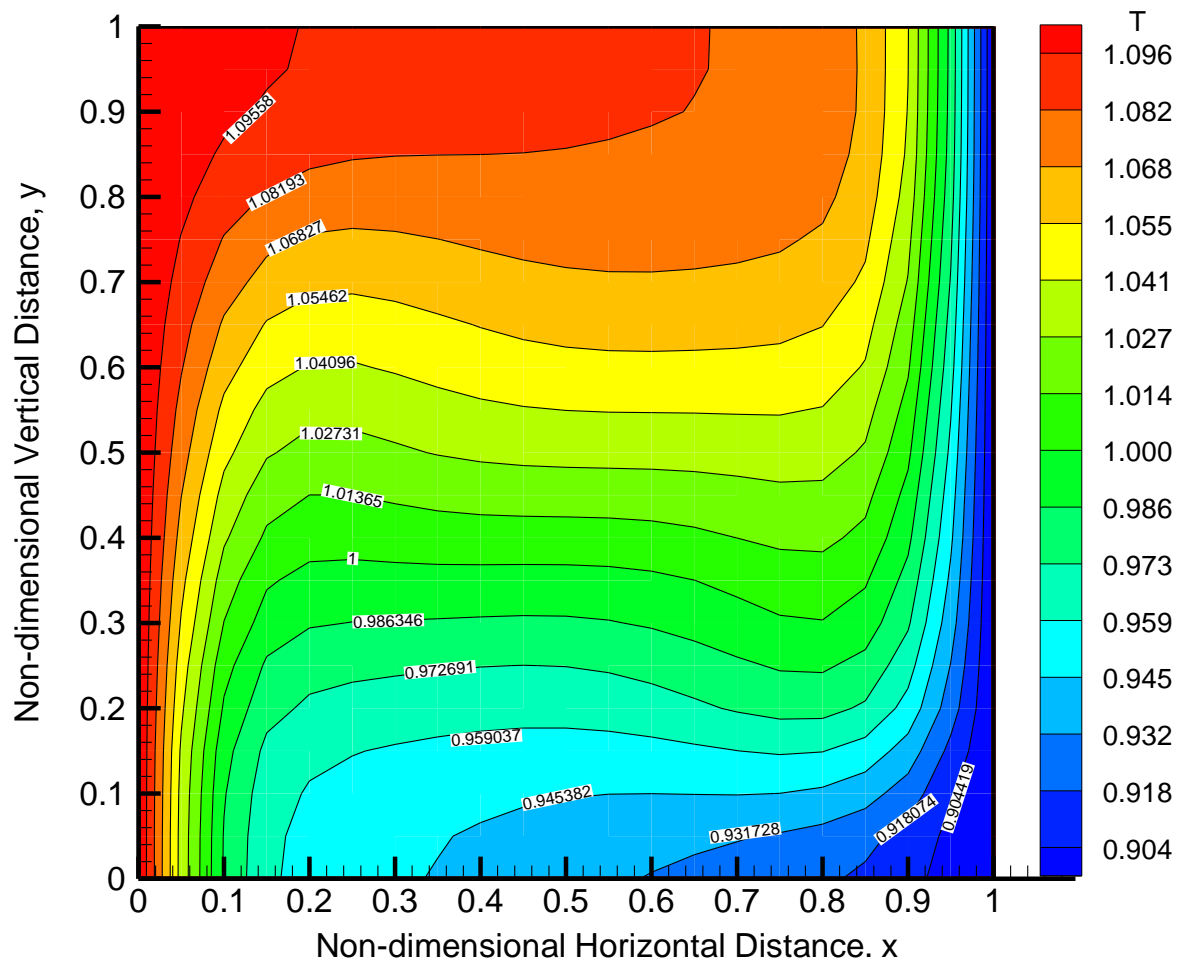


Figure 19- Steady State Non-dimensional Temperature Contours for $Ra=10^5$ for Variable Properties (21×21 , $\Delta \bar{t} = 1.0E-5$)

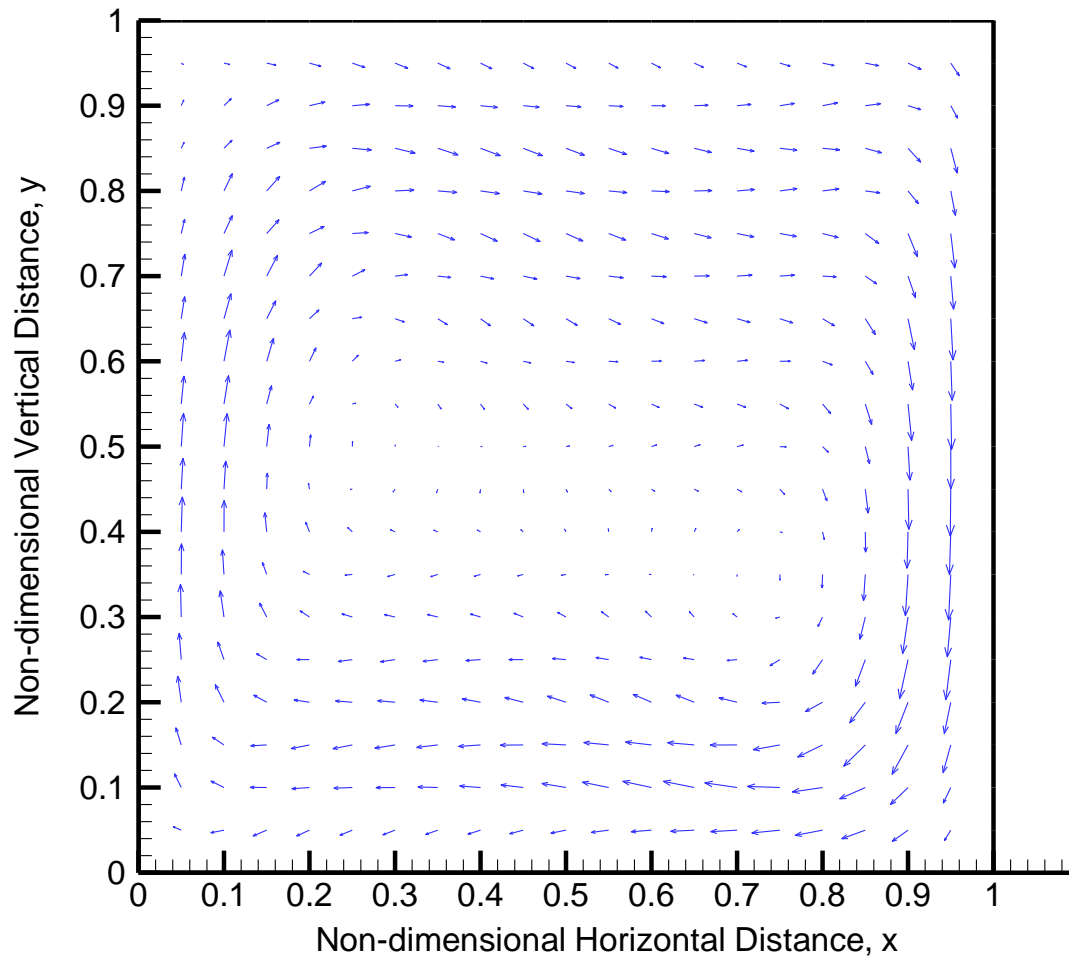


Figure 20- Steady State Non-dimensional Velocity Vectors for $Ra=10^5$ (21×21 , $\Delta \bar{t} = 1.0E-5$) for Variable Properties

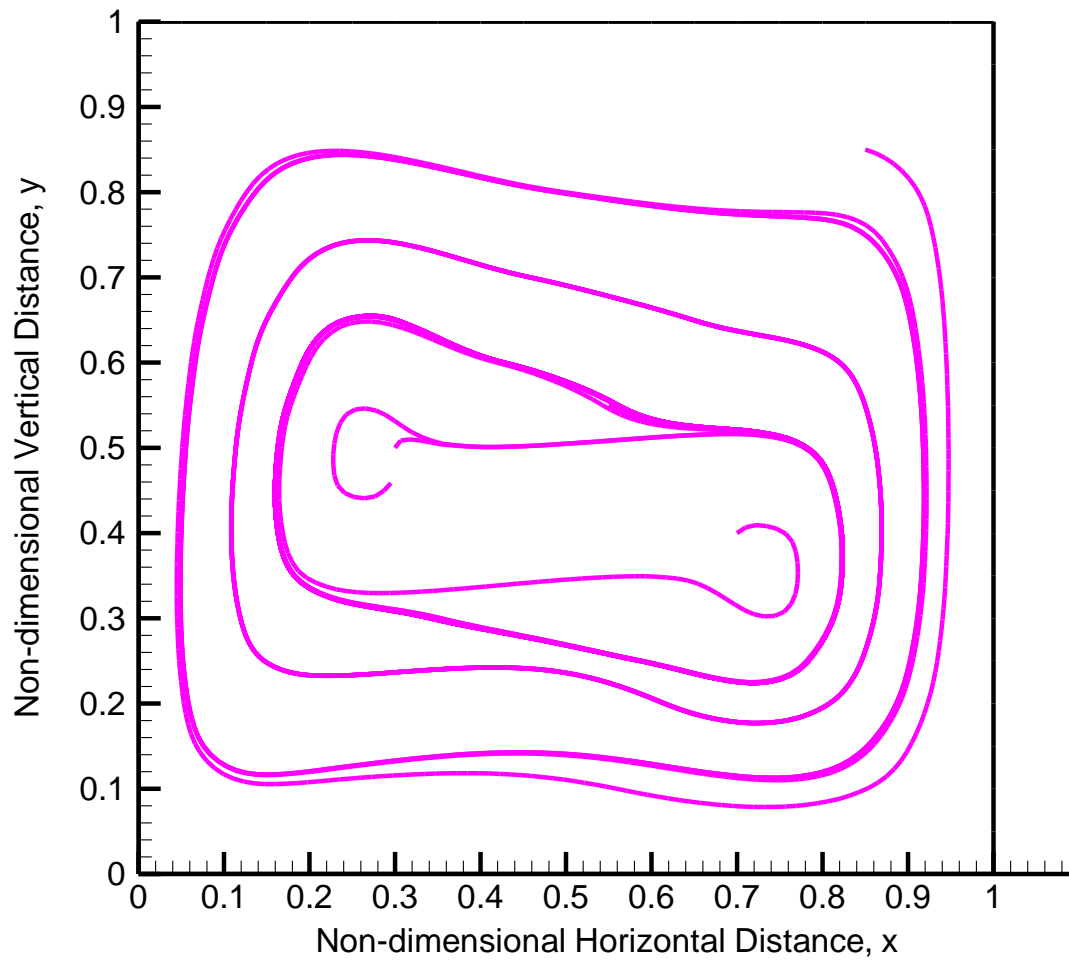


Figure 21- Steady State Non-dimensional Velocity Streamlines for $Ra=10^5$ for Variable Properties (21×21 , $\Delta \bar{t} = 1.0E-5$)

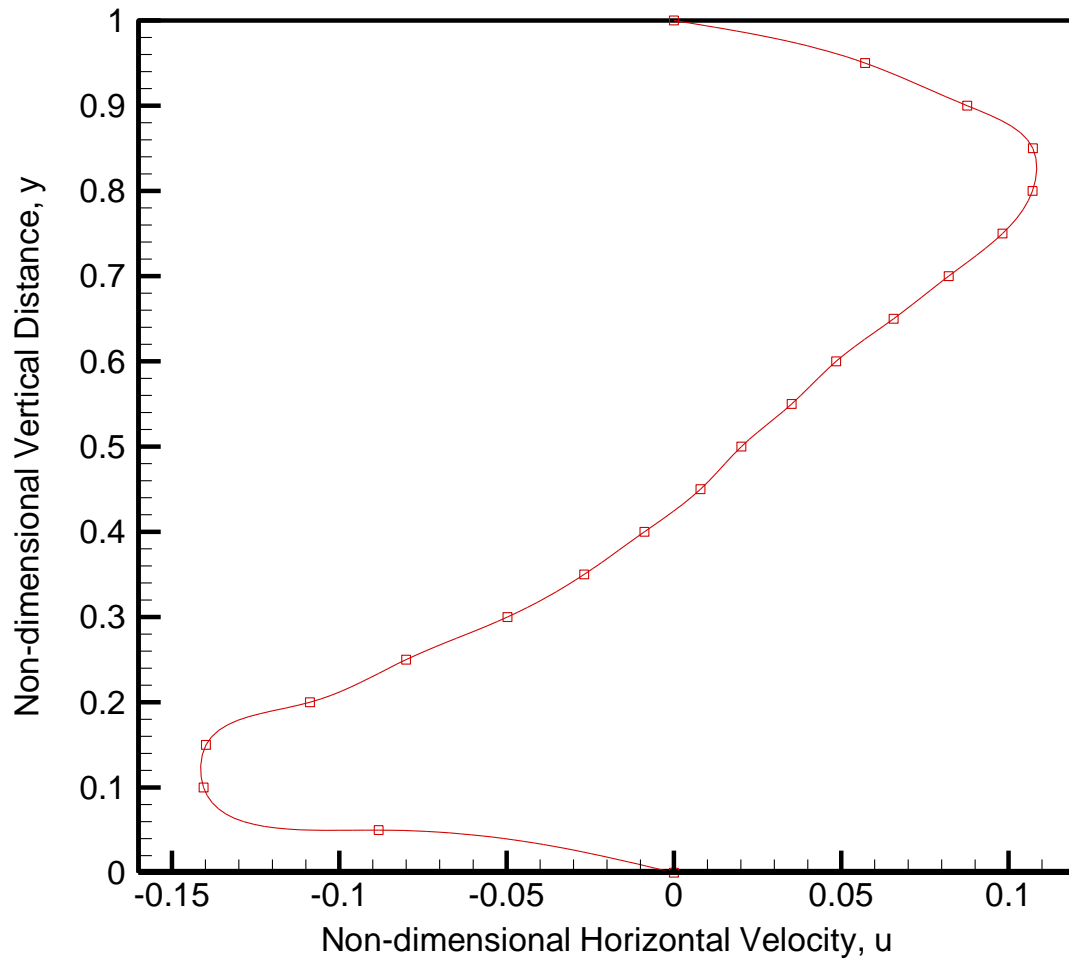


Figure 22- Steady State Non-dimensional Horizontal Velocity Distribution along the Vertical Centerline of the Enclosure for $Ra=10^5$ for Variable Properties
(21 x 21, $\Delta \bar{t} = 1.0E-5$)

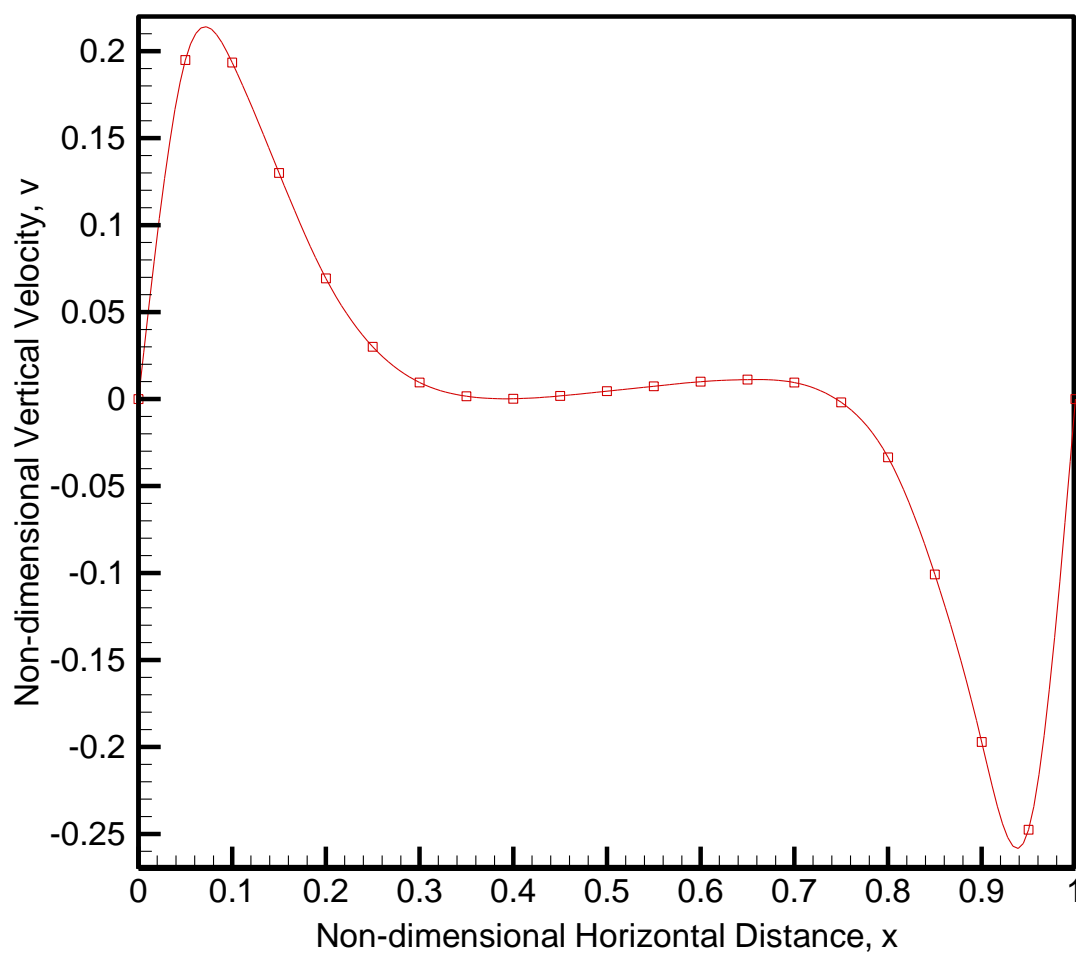


Figure 23- Steady State Non-dimensional Vertical Velocity Distribution along the Horizontal Centerline of the Enclosure for $Ra=10^5$ for Variable Properties
(21 x 21, $\Delta \bar{t} = 1.0E-5$)

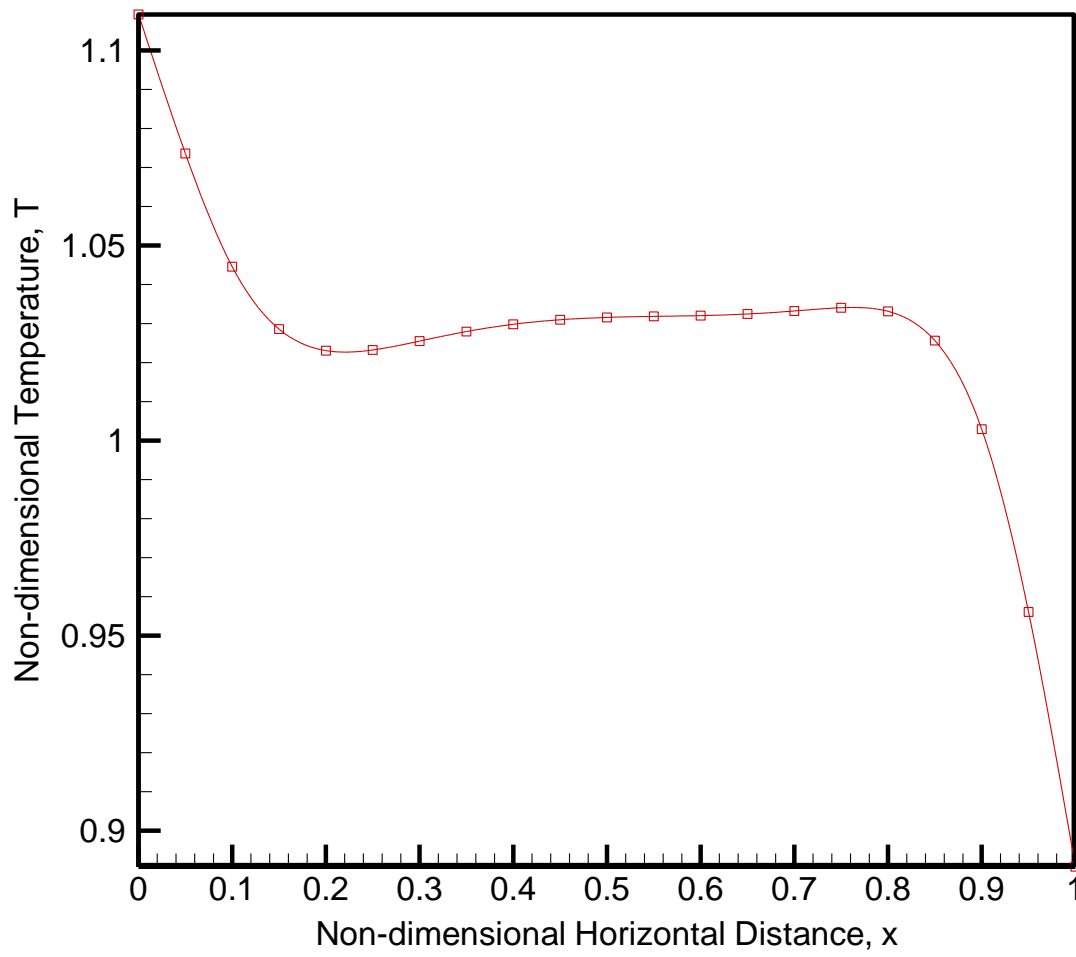


Figure 24- Steady State Non-dimensional Temperature Distribution along the Horizontal Centerline of the Enclosure for $Ra=10^5$ for Variable Properties
(21 x 21, $\Delta \bar{t} = 1.0E-5$)

8.2. The Effects of Variable Fluid Properties in Square Enclosures

The comparison of the present model with variable thermodynamic and transport properties to that with constant properties shows noticeable differences when the temperature difference between the enclosure walls is high. These differences can be seen when one compares the velocity and temperature profiles in horizontal and vertical directions. The quantitative comparisons of the results of constant and variable properties are presented for Rayleigh numbers 10^5 , 5×10^4 and 10^4 with $\Delta T = 356.6$ deg K in Table 14 through 22 respectively. These results are compared graphically as shown in Figures 25 through 33, respectively. The non-dimensional velocity values of the variable property model are predicted lower than the constant property model and these differences are increasing as the Rayleigh number decreases. However, the non-dimensional temperature values are predicted higher for variable property model than the constant property model and these differences are also increasing as the Rayleigh number decreases.

The comparison of the vertical velocity profiles for the different temperature differences ($\Delta T = 64.84$ deg K and $\Delta T = 356.6$ deg K) at Ra number 10^5 are also presented in this chapter. As wall temperature difference increases the maximum velocity in vertical direction increases in magnitude and moves towards the cold side of the enclosure as shown in Figure 34. For different Rayleigh numbers 10^4 , 5×10^4 and 10^5 at same wall temperature difference $\Delta T = 356.6$ deg K, the same comparison is presented in Figure 35. In this case, the maximum value of the vertical velocity higher as increases Rayleigh numbers and the location of this point moves towards the hot side of the enclosure.

Table 14- Comparison of Predicted Non-dimensional values of Horizontal Velocity along the Vertical Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 10^5$ ($\Delta T = 356.6$ K, $H=0.0144$ m and 21×21)

Non-dimensional Vertical Distance, \bar{y}	Non-dimensional Horizontal Velocity, \bar{u} (Constant)	Non-dimensional Horizontal Velocity, \bar{u} (Variable)	% Deviation in \bar{u}
1.00	0.0000000	0.0000000	0.0000000
0.95	0.0936828	0.0838684	-11.702097
0.90	0.1488197	0.1341210	-10.959283
0.85	0.1716902	0.1341210	-28.011422
0.80	0.1652194	0.1595315	-3.5653773
0.75	0.1426628	0.1461648	2.3959256
0.70	0.1111090	0.1217983	8.7762308
0.65	0.0794429	0.0934810	15.017077
0.60	0.0512550	0.0639186	19.811995
0.55	0.0290225	0.0372387	22.063447
0.50	0.0118728	0.0144168	17.646169
0.45	-0.0025051	-0.0047551	47.316622
0.40	-0.0157795	-0.0207039	23.784529
0.35	-0.0313873	-0.0365194	14.053032
0.30	-0.0475436	-0.0509674	6.7175475
0.25	-0.0669880	-0.0677428	1.1142725
0.20	-0.0828617	-0.0808987	-2.4264911
0.15	-0.0962331	-0.0929094	-3.5773448
0.10	-0.0923184	-0.0889011	-3.8439314
0.05	-0.0646962	-0.0638143	-1.3819460
0.00	0.0000000	0.0000000	0.0000000

Table 15- Comparison of Predicted Non-dimensional values of Vertical Velocity along the Horizontal Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 10^5$ ($\Delta T = 356.6$ K, $H=0.0144$ m and 21×21)

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Vertical Velocity, \bar{v} (Constant)	Non-dimensional Vertical Velocity, \bar{v} (Variable)	% Deviation in \bar{v}
1.00	0.0000000	0.0000000	0.0000000
0.95	-0.2093566	-0.2203230	4.9774195
0.90	-0.1630863	-0.1482642	-9.9970862
0.85	-0.0850964	-0.0773353	-10.035699
0.80	-0.0337917	-0.0353622	4.4411213
0.75	-0.0113363	-0.0141749	20.025142
0.70	-0.0043228	-0.0075802	42.972368
0.65	-0.0045029	-0.0060678	25.789552
0.60	-0.0066290	-0.0078408	15.455837
0.55	-0.0092771	-0.0088425	-4.9148134
0.50	-0.0112224	-0.0094607	-18.620428
0.45	-0.0117948	-0.0069627	-69.400468
0.40	-0.0096178	-0.0008446	-1038.7143
0.35	-0.0017328	0.0132710	113.05706
0.30	0.0165154	0.0388702	57.511265
0.25	0.0526970	0.0820472	35.772274
0.20	0.1129559	0.1426556	20.819161
0.15	0.1957559	0.2091187	6.3900550
0.10	0.2702196	0.2412191	-12.022472
0.05	0.2444604	0.1800874	-35.745421
0.00	0.0000000	0.0000000	0.0000000

Table 16- Comparison of Predicted Non-dimensional values of Temperature along the Horizontal Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 10^5$ ($\Delta T = 356.6$ K, $H=0.0144$ m and 21×21)

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Temperature, \bar{T} (Constant)	Non-dimensional Temperature, \bar{T} (Variable)	% Deviation in \bar{T}
1.00	0.4005175	0.4005175	0.0000000
0.95	0.6880814	0.8184327	15.926941
0.90	0.8781839	0.9827196	10.637388
0.85	0.9661263	1.0501240	7.9988363
0.80	0.9919955	1.0705230	7.3354332
0.75	0.9937108	1.0724990	7.3462259
0.70	0.9900559	1.0697560	7.4503064
0.65	0.9867782	1.0664140	7.4676251
0.60	0.9838625	1.0631580	7.4584868
0.55	0.9801099	1.0588480	7.4362042
0.50	0.9741970	1.0527800	7.4643325
0.45	0.9652914	1.0442670	7.5627784
0.40	0.9531033	1.0338590	7.8110941
0.35	0.9384271	1.0233100	8.2949350
0.30	0.9243982	1.0176280	9.1614814
0.25	0.9186205	1.0261480	10.478751
0.20	0.9359822	1.0637710	12.012811
0.15	0.9999081	1.1468870	12.815464
0.10	1.1384730	1.2812960	11.146760
0.05	1.3570710	1.4443730	6.0442835
0.00	1.5994830	1.5994830	0.0000000

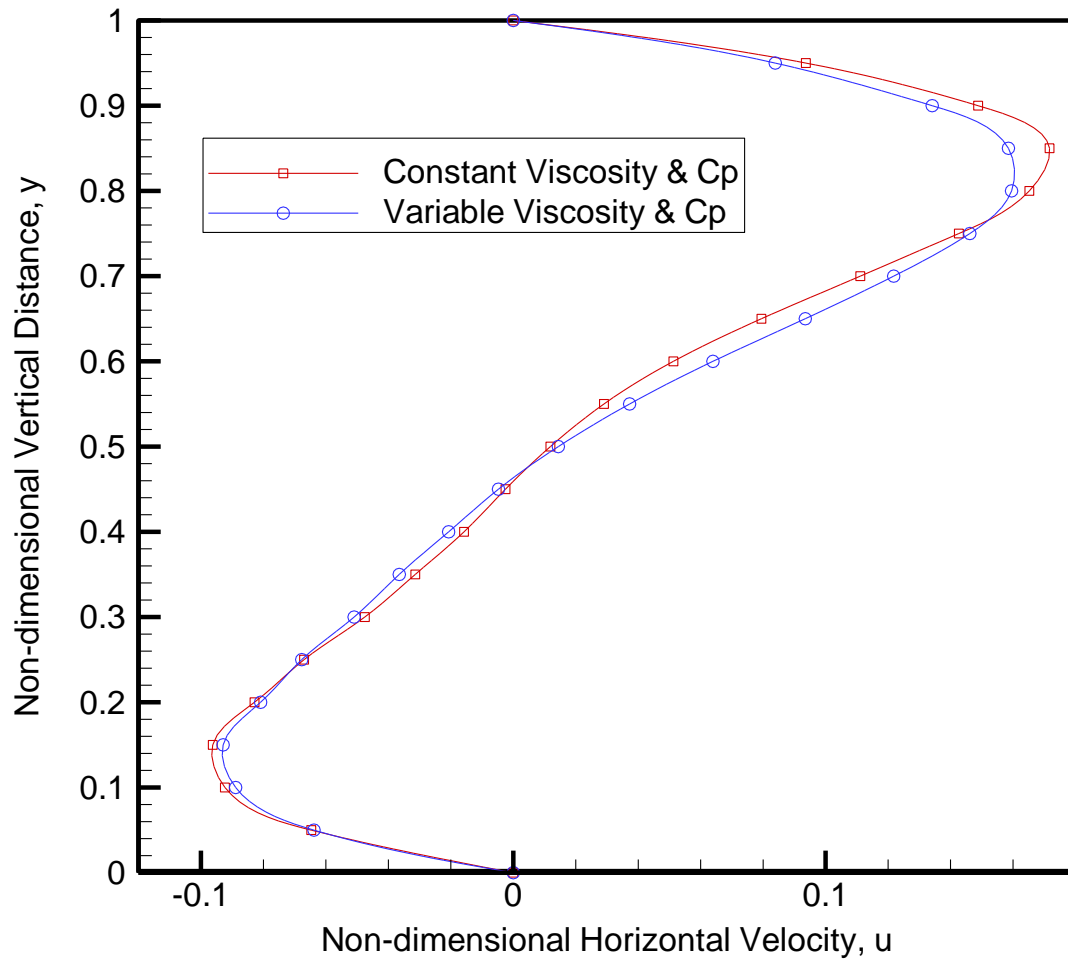


Figure 25- Comparison of Steady State Non-dimensional Horizontal Velocity Distributions along the Vertical Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 10^5$ ($\Delta T = 356.6$ K, $H = 0.0115$ m)

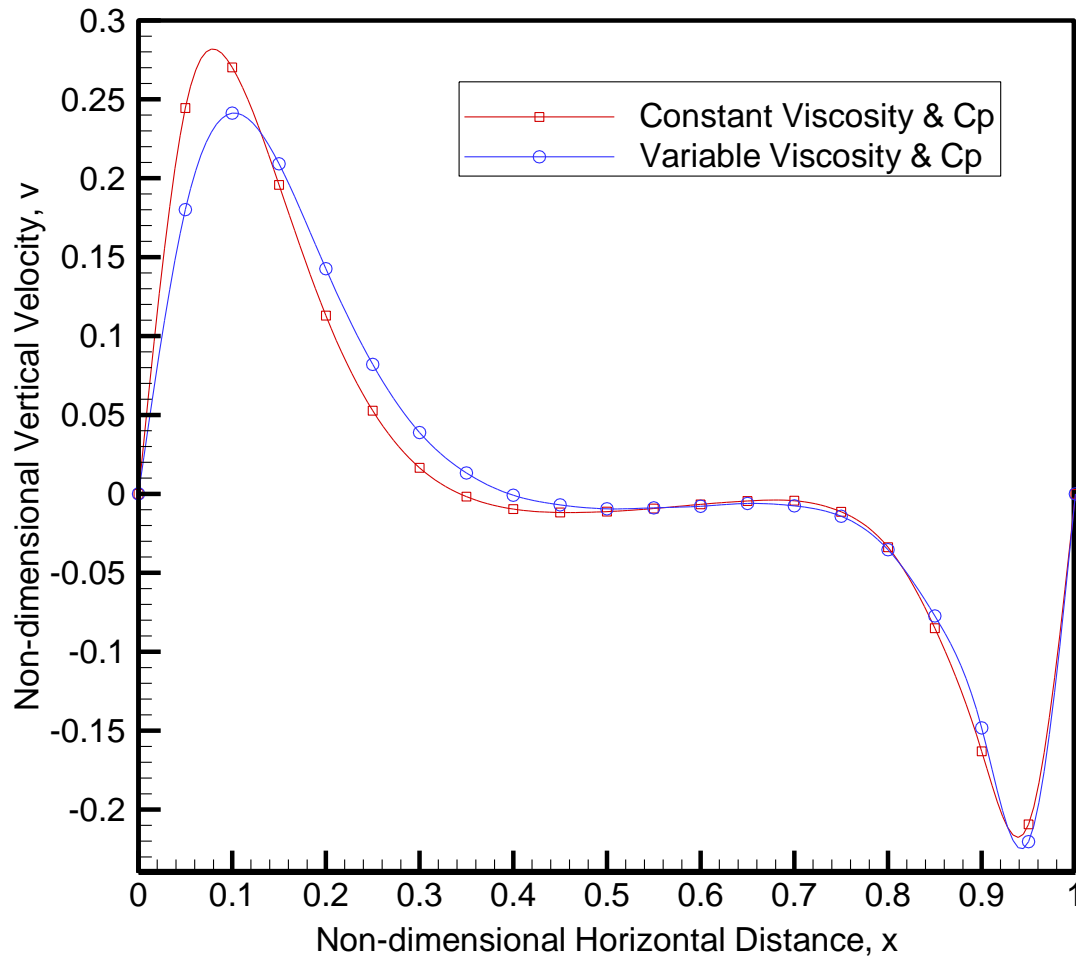


Figure 26- Comparison of Steady State Non-dimensional Vertical Velocity Distributions along the Horizontal Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 10^5$ ($\Delta T = 356.6$ K, $H = 0.0115$ m)

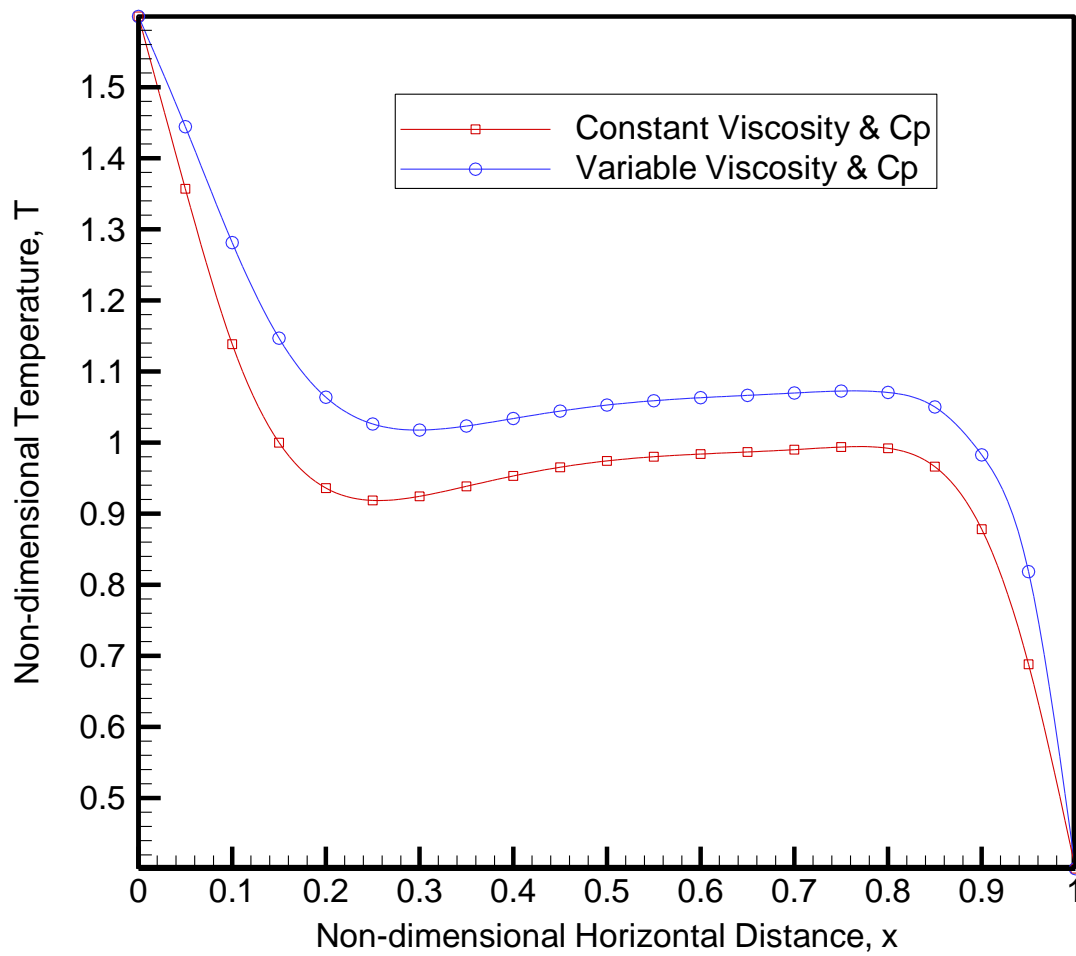


Figure 27- Comparison of Steady State Non-dimensional Temperature Distributions along the Horizontal Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 10^5$ ($\Delta T = 356.6$ K, $H = 0.0115$ m)

Table-17: Comparison of Predicted Non-dimensional values of Horizontal Velocity along the Vertical Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 5 \times 10^4$ ($\Delta T = 356.6$ K, $H=0.0115$ m and 21×21)

Non-dimensional Vertical Distance, \bar{y}	Non-dimensional Horizontal Velocity, \bar{u} (Constant)	Non-dimensional Horizontal Velocity, \bar{u} (Variable)	% Deviation in \bar{u}
1.00	0.0000000	0.0000000	0.0000000
0.95	0.1010880	0.0842762	-19.948369
0.90	0.1608237	0.1356474	-18.560105
0.85	0.1867887	0.1616705	-15.536662
0.80	0.1832999	0.1652628	-10.914192
0.75	0.1625545	0.1544248	-5.2645041
0.70	0.1307763	0.1320511	0.9653838
0.65	0.0965233	0.1042924	7.4492772
0.60	0.0635012	0.0735256	13.633982
0.55	0.0353098	0.0439724	19.700226
0.50	0.0122669	0.0172566	28.914988
0.45	-0.0068705	-0.0058550	-17.344015
0.40	-0.0233795	-0.0251376	6.9936982
0.35	-0.0407417	-0.0433361	5.9866850
0.30	-0.0581902	-0.0596543	2.4543729
0.25	-0.0785703	-0.0780983	-0.6044561
0.20	-0.0957358	-0.0932317	-2.6859203
0.15	-0.1085417	-0.1062663	-2.1412244
0.10	-0.1023104	-0.1009673	-1.3302326
0.05	-0.0692979	-0.0701084	1.1561509
0.00	0.0000000	0.0000000	0.0000000

Table 18- Comparison of Predicted Non-dimensional values of Vertical Velocity along the Horizontal Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 5 \times 10^4$ ($\Delta T = 356.6$ K, $H=0.0115$ m and 21×21)

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Vertical Velocity, \bar{v} (Constant)	Non-dimensional Vertical Velocity, \bar{v} (Variable)	% Deviation in \bar{v}
1.00	0.0000000	0.0000000	0.0000000
0.95	-0.2079249	-0.2181263	4.6768317
0.90	-0.1805588	-0.1735712	-4.0257830
0.85	-0.1105591	-0.1065565	-3.7563170
0.80	-0.0555643	-0.0564429	1.5565636
0.75	-0.0254680	-0.0274957	7.3747006
0.70	-0.0122137	-0.0132682	7.9472724
0.65	-0.0085119	-0.0080180	-6.1597252
0.60	-0.0085827	-0.0067362	-27.411106
0.55	-0.0098464	-0.0066076	-49.015896
0.50	-0.0102490	-0.0053554	-91.375590
0.45	-0.0082766	-0.0009993	-728.21051
0.40	-0.0019325	0.0087003	122.21279
0.35	0.0122559	0.0273276	55.151703
0.30	0.0386824	0.0584576	33.828290
0.25	0.0833502	0.1055321	21.019045
0.20	0.1483615	0.1650306	10.100611
0.15	0.2259767	0.221059	-2.2246097
0.10	0.2786205	0.2356686	-18.225550
0.05	0.2286864	0.1671887	-36.783407
0.00	0.0000000	0.0000000	0.0000000

Table 19- Comparison of Predicted Non-dimensional values of Temperature along the Horizontal Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 5 \times 10^4$ ($\Delta T = 356.6$ K, $H=0.0115$ m and 21×21)

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Temperature, \bar{T} (Constant)	Non-dimensional Temperature, \bar{T} (Variable)	% Deviation in \bar{T}
1.00	0.4115081	0.4115081	0.0000000
0.95	0.6617610	0.7668620	13.705334
0.90	0.8384050	0.9422544	11.021375
0.85	0.9328032	1.0257650	9.0626800
0.80	0.9688077	1.0578800	8.4198869
0.75	0.9757017	1.0652240	8.4040821
0.70	0.9720609	1.0628410	8.5412681
0.65	0.9660809	1.0577420	8.6657332
0.60	0.9598690	1.0521630	8.7718347
0.55	0.9532482	1.0462480	8.8888867
0.50	0.9453675	1.0394670	9.0526683
0.45	0.9356530	1.0315640	9.2976296
0.40	0.9243249	1.0232780	9.6702069
0.35	0.9131055	1.0171590	10.229816
0.30	0.9065261	1.0188430	11.023965
0.25	0.9133821	1.0377910	11.987856
0.20	0.9482343	1.0867650	12.747070
0.15	1.0301340	1.1763090	12.426581
0.10	1.1754250	1.3047870	9.9144151
0.05	1.3757300	1.4503020	5.1418256
0.00	1.5884920	1.5884920	0.0000000

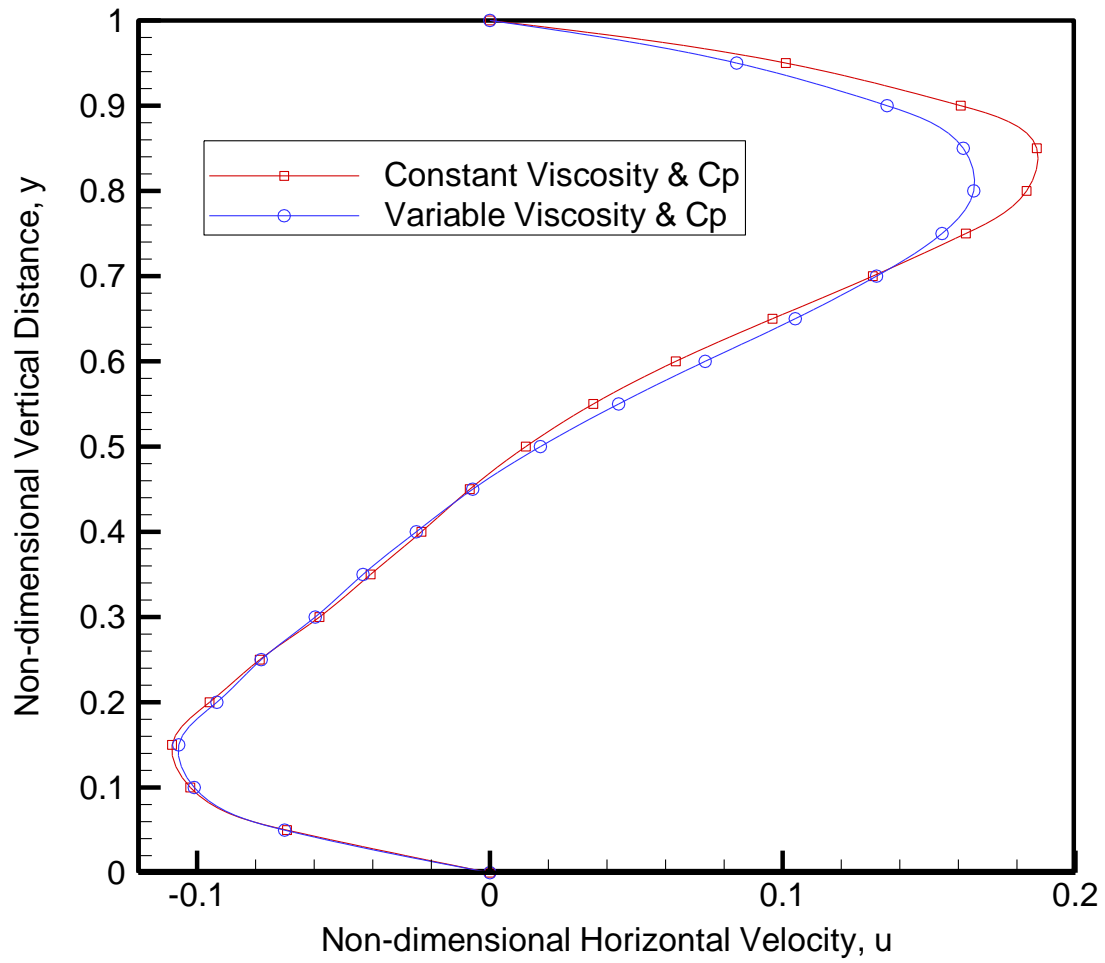


Figure 28- Comparison of Steady State Non-dimensional Horizontal Velocity Distributions along the Vertical Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 5 \times 10^4$ ($\Delta T = 356.6$ K, $H = 0.0115$ m)

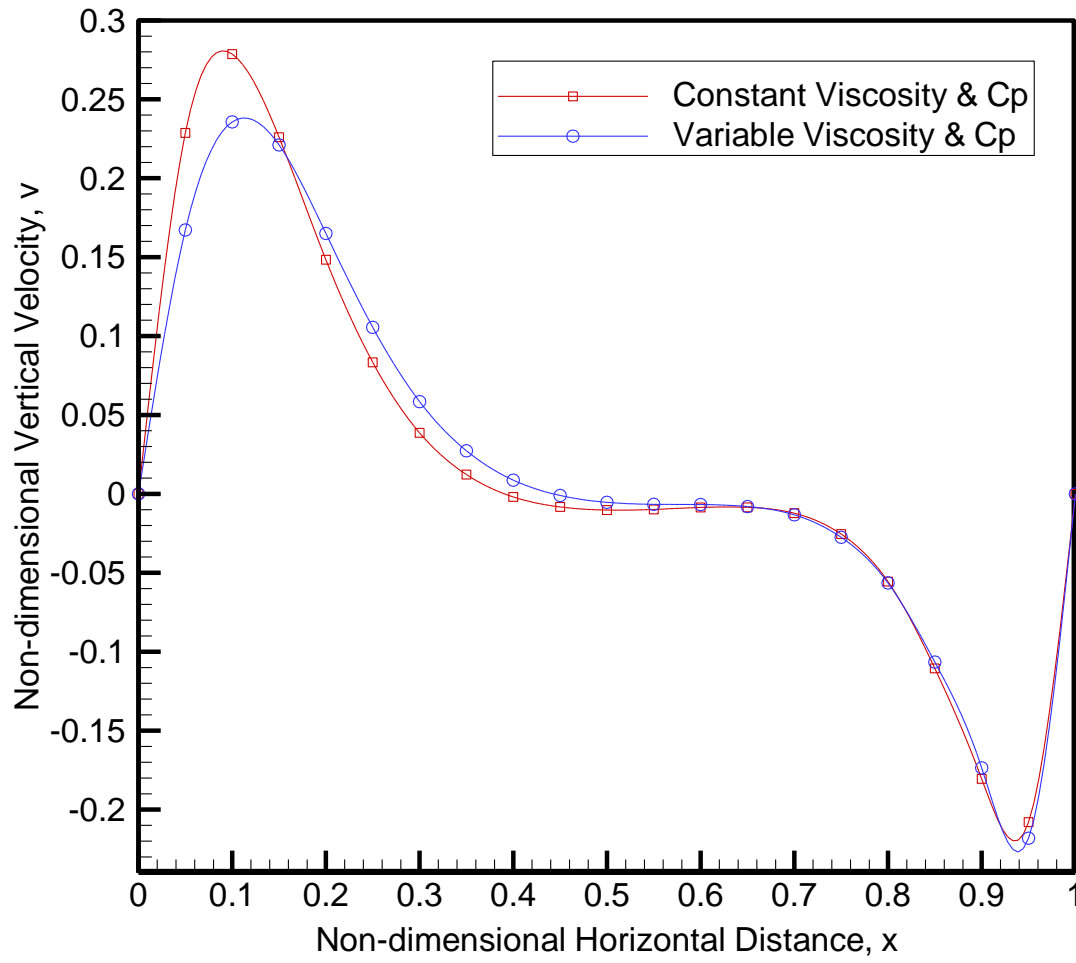


Figure 29- Comparison of Steady State Non-dimensional Vertical Velocity Distributions along the Horizontal Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra= 5 \times 10^4$ ($\Delta T = 356.6$ K, $H = 0.0115m$)

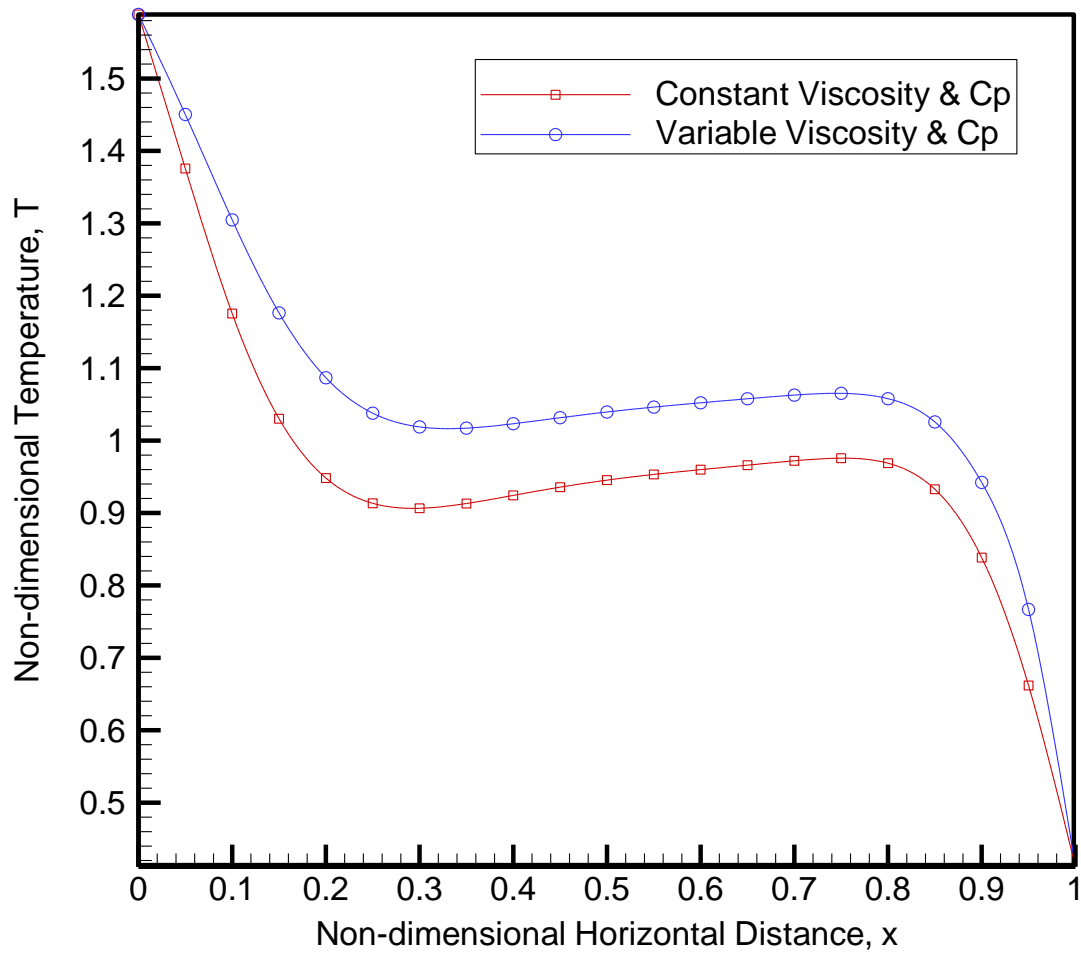


Figure 30- Comparison of Steady State Non-dimensional Temperature Distributions along the Horizontal Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 5 \times 10^4$ ($\Delta T = 356.6$ K, $H = 0.0115$ m)

Table 20- Comparison of Predicted Non-dimensional values of Horizontal Velocity along the Vertical Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 10^4$ ($\Delta T = 356.6$ K, $H=0.0067$ m and 21×21)

Non-dimensional Vertical Distance, \bar{y}	Non-dimensional Horizontal Velocity, \bar{u} (Constant)	Non-dimensional Horizontal Velocity, \bar{u} (Variable)	% Deviation in \bar{u}
1.00	0.0000000	0.0000000	0.0000000
0.95	0.1083175	0.0862691	-25.557658
0.90	0.1756615	0.1418459	-23.839673
0.85	0.2116735	0.1744832	-21.314544
0.80	0.2217043	0.1879760	-17.942875
0.75	0.2141084	0.1879882	-13.894595
0.70	0.1926927	0.1765157	-9.1646238
0.65	0.1629599	0.1569659	-3.8186637
0.60	0.1272570	0.1303389	2.3645281
0.55	0.0889332	0.0988132	9.9986813
0.50	0.0497051	0.0637028	21.973478
0.45	0.0108663	0.0265779	59.115146
0.40	-0.0264220	-0.0105711	-149.94404
0.35	-0.0627750	-0.0475411	-32.043747
0.30	-0.0961872	-0.0818469	-17.520782
0.25	-0.1268682	-0.1145173	-10.785182
0.20	-0.1486439	-0.1396459	-6.4434401
0.15	-0.1578665	-0.1546874	-2.0551770
0.10	-0.1419982	-0.1439025	1.3233265
0.05	-0.0931065	-0.0969537	3.9680692
0.00	0.0000000	0.0000000	0.0000000

Table 21- Comparison of Predicted Non-dimensional values of Vertical Velocity along the Horizontal Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 10^4$ ($\Delta T = 356.6$ K, $H=0.0067$ m and 21×21)

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Vertical Velocity, \bar{v} (Constant)	Non-dimensional Vertical Velocity, \bar{v} (Variable)	% Deviation in \bar{v}
1.00	0.0000000	0.0000000	0.0000000
0.95	-0.1823712	-0.2083991	12.489449
0.90	-0.2134350	-0.2101146	-1.5802804
0.85	-0.1835589	-0.1661266	-10.493382
0.80	-0.1361653	-0.1167885	-16.591359
0.75	-0.0923559	-0.0752924	-22.662900
0.70	-0.0569510	-0.0426731	-33.458736
0.65	-0.0304331	-0.0179429	-69.610304
0.60	-0.0094573	0.0021489	540.08904
0.55	0.0083900	0.0200227	58.097265
0.50	0.0263901	0.0386682	31.752292
0.45	0.0469496	0.0603234	22.170256
0.40	0.0727090	0.0869015	16.331734
0.35	0.1055913	0.1189967	11.265354
0.30	0.1459709	0.1546759	5.6278967
0.25	0.1914424	0.1888658	-1.3642491
0.20	0.2335142	0.2121595	-10.065398
0.15	0.2573774	0.2139577	-20.293590
0.10	0.2404851	0.1834930	-31.059549
0.05	0.1609832	0.1140888	-41.103421
0.00	0.0000000	0.0000000	0.0000000

Table 22- Comparison of Predicted Non-dimensional values of Temperature along the Horizontal Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 10^4$ ($\Delta T = 356.6$ K, $H=0.0067$ m and 21×21)

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Temperature, \bar{T} (Constant)	Non-dimensional Temperature, \bar{T} (Variable)	% Deviation in \bar{T}
1.00	0.4048320	0.4048320	0.0000000
0.95	0.5606439	0.6564321	14.592248
0.90	0.6961532	0.8109570	14.156582
0.85	0.7969056	0.9066408	12.103492
0.80	0.8610080	0.9619112	10.489866
0.75	0.8953147	0.9905590	9.6152071
0.70	0.9093859	1.0029400	9.3279857
0.65	0.9116908	1.0064310	9.4134818
0.60	0.9082556	1.0060100	9.7170405
0.55	0.9031278	1.0051570	10.150573
0.50	0.8991418	1.0065670	10.672434
0.45	0.8988663	1.0128870	11.257001
0.40	0.9052760	1.0271490	11.865172
0.35	0.9222708	1.0528780	12.404780
0.30	0.9548032	1.0936110	12.692611
0.25	1.0082650	1.1516130	12.447584
0.20	1.0870050	1.2264310	11.368434
0.15	1.1917150	1.3141070	9.3137012
0.10	1.3177410	1.4085590	6.4475822
0.05	1.4555180	1.5036310	3.1997877
0.00	1.5951680	1.5951680	0.0000000

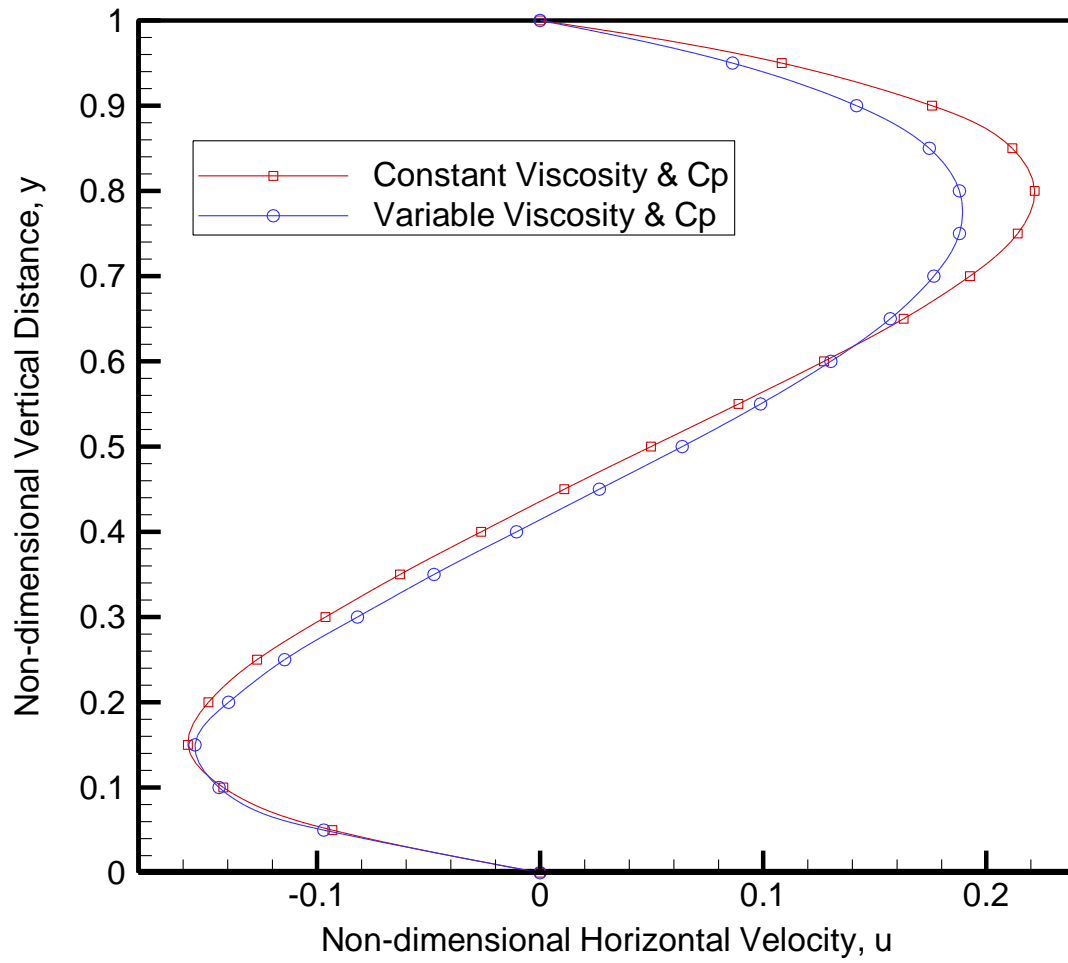


Figure 31- Comparison of Steady State Non-dimensional Horizontal Velocity Distributions along the Vertical Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra=10^4$ ($\Delta T = 356.6$ K, $H = 0.0067$ m)

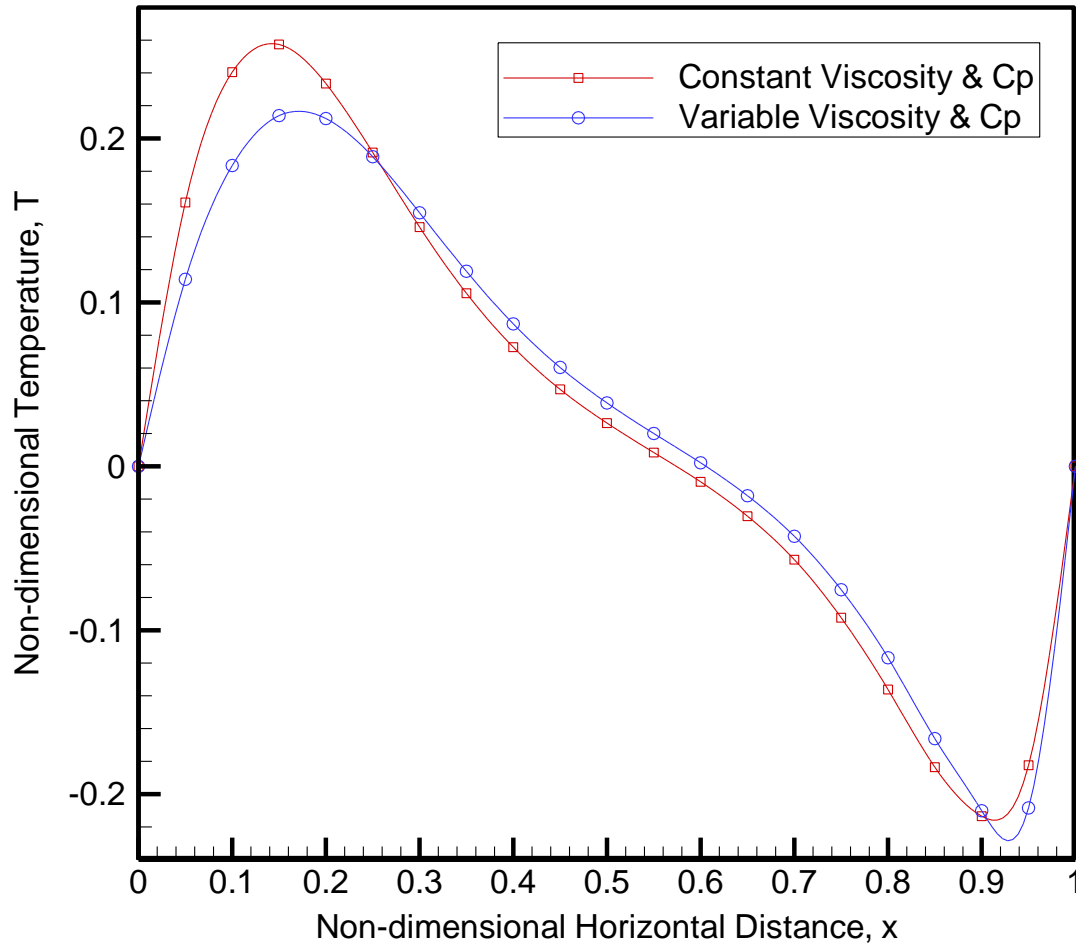


Figure 32- Comparison of Steady State Non-dimensional Vertical Velocity Distributions along the Horizontal Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 10^4$ ($\Delta T = 356.6$ K, $H = 0.0067$ m)

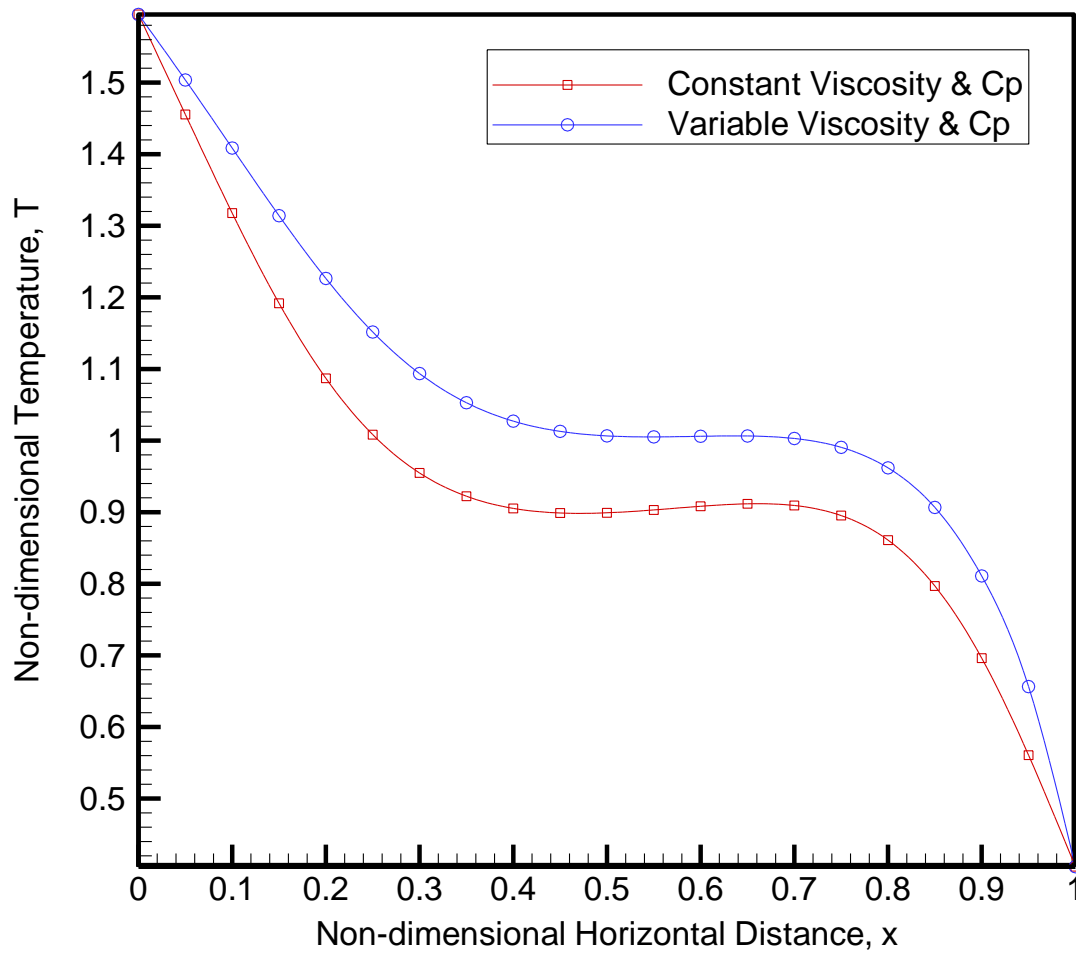


Figure 33- Comparison of Steady State Non-dimensional Temperature Distributions along the Horizontal Centerline of the Enclosure for Constant and Variable Fluid Properties at $Ra = 10^4$ ($\Delta T = 356.6$ K, $H = 0.0067$ m)

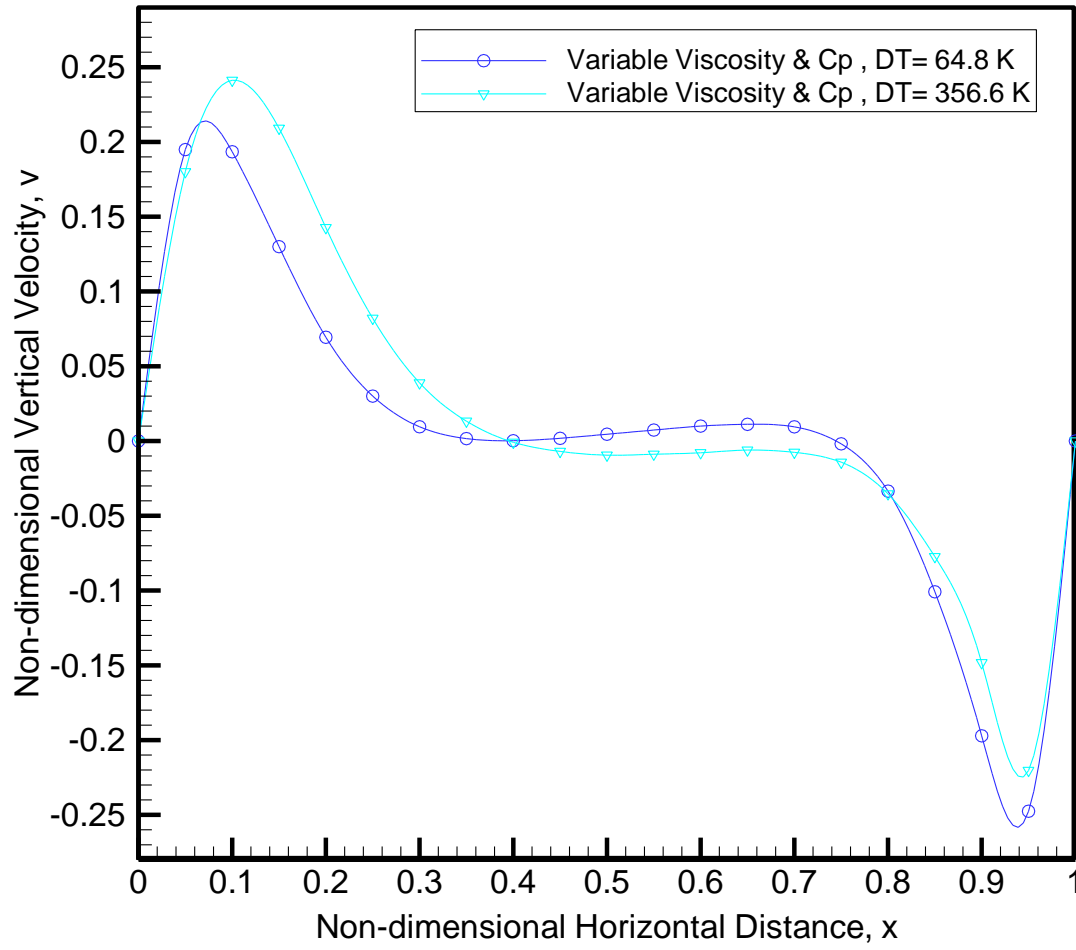


Figure 34- Comparison of Non-dimensional Vertical Velocity Profiles along the Horizontal Centerline of the Enclosure between $\Delta T = 64.8$ K and $\Delta T = 356.6$ K ($Ra = 10^5$) for Variable Properties.

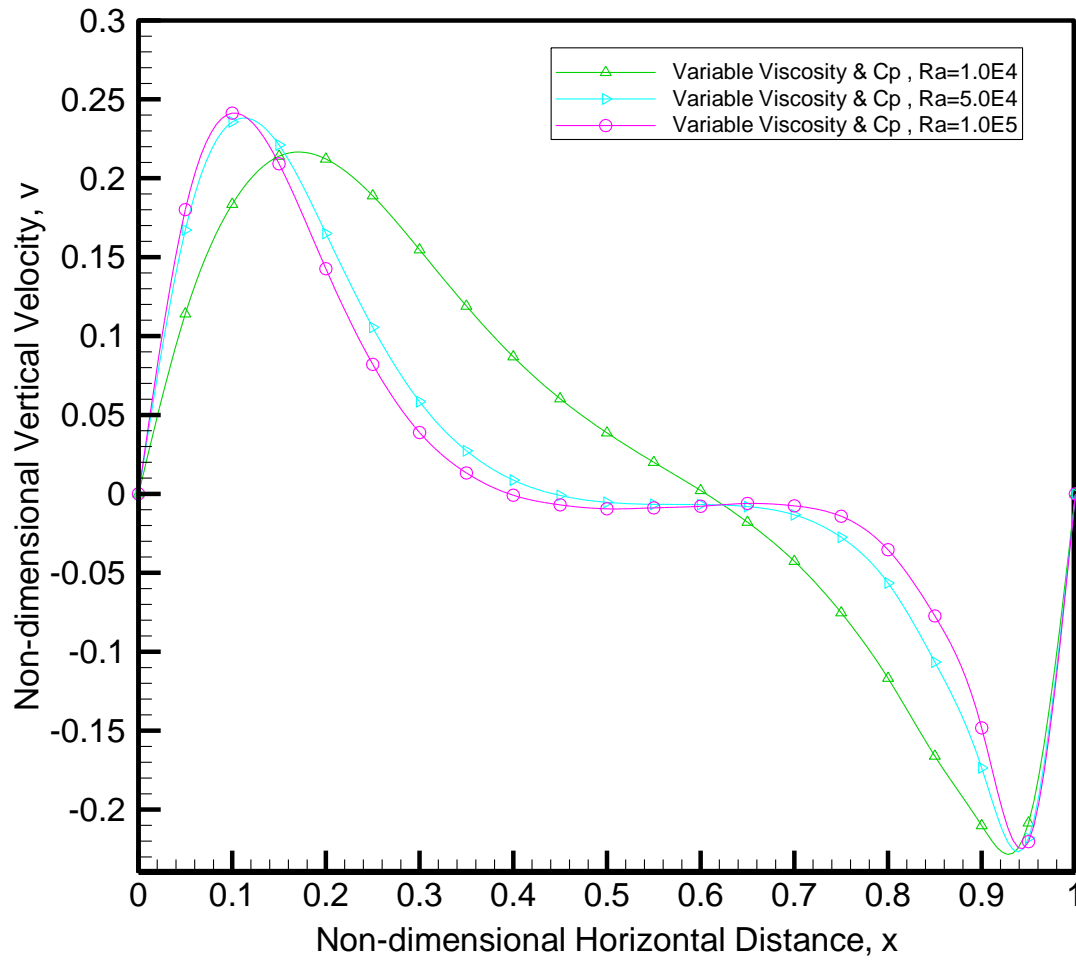


Figure 35- Comparison of Non-dimensional Vertical Velocity Profiles along the Horizontal Centerline of the Enclosure between $Ra = 10^4$, $Ra = 5 \times 10^4$, and $Ra = 10^5$ ($\Delta T = 356.6$ K) for Variable Properties.

Chapter 9

The Study of Effects of Inclination of the Enclosure on Natural Convection

The effects of inclination of the enclosure are presented for the laminar natural convection phenomenon in this chapter. The velocity fields and temperature distributions were observed at different angles of the enclosure.

In this study, primitive variables are predicted for different angles (0 deg, 30 deg, 45 deg, 60 deg and 90 deg) of the enclosure. The predicted non-dimensional values of horizontal and vertical velocities are presented for various angles of the enclosure in Table 23 and 24, respectively. As the inclination of the enclosure increases, the magnitudes of the horizontal and vertical velocities are increasing. When inclined angle becomes 90 deg, the velocity profile is observed in opposite direction to the previous cases as shown in Figure 36 and 37. The predicted temperature distribution is presented in Figure 38 for different angles of the enclosure. This shows that as angle of the enclosure increases, the magnitudes of the temperature are decreasing.

Table 23- Steady State Non-dimensional values of Vertical Velocity along the Horizontal Centerline of the Inclined Enclosure at Various Angles for $Ra = 10^5$ ($\Delta\bar{t} = 1.0E-5$ and 21×21)

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Vertical Velocity, \bar{v} (Angle 0)	Non-dimensional Vertical Velocity, \bar{v} (Angle 30)	Non-dimensional Vertical Velocity, \bar{v} (Angle 45)	Non-dimensional Vertical Velocity, \bar{v} (Angle 60)	Non-dimensional Vertical Velocity, \bar{v} (Angle 90)
1.00	0.00000	0.00000	0.00000	0.00000	0.00000
0.95	-0.24602	-0.26058	-0.25779	-0.24928	0.21951
0.90	-0.20141	-0.24630	-0.27096	-0.28980	0.30641
0.85	-0.10434	-0.16668	-0.21887	-0.26821	0.34053
0.8	-0.03502	-0.10694	-0.17004	-0.22956	0.31293
0.75	-0.00233	-0.07122	-0.13153	-0.18897	0.26161
0.70	0.00942	-0.04878	-0.09502	-0.14116	0.19542
0.65	0.01127	-0.03580	-0.06627	-0.09980	0.13691
0.60	0.01016	-0.02174	-0.03761	-0.05934	0.08114
0.55	0.00757	-0.01251	-0.01921	-0.03202	0.03820
0.50	0.00476	-0.00042	0.00002	-0.00343	-0.00849
0.45	0.00190	0.00849	0.01290	0.01678	-0.04752
0.40	0.00004	0.02114	0.03239	0.04633	-0.09612
0.35	0.00109	0.03349	0.05384	0.07773	-0.14396
0.30	0.00842	0.05287	0.08696	0.12320	-0.20211
0.25	0.02823	0.07945	0.12814	0.17658	-0.26020
0.20	0.06702	0.12094	0.17873	0.23593	-0.31358
0.15	0.12808	0.18484	0.23919	0.29394	-0.33307
0.10	0.19448	0.25241	0.29254	0.33381	-0.31983
0.05	0.20148	0.25950	0.27693	0.29038	-0.22554
0.00	0.00000	0.00000	0.00000	0.00000	0.00000

Table 24- Steady State Non-dimensional values of Horizontal Velocity along the Vertical Centerline of the Inclined Enclosure at Various Angles for $Ra = 10^5$ ($\Delta\bar{t} = 1.0E-5$ and 21×21)

Non-dimensional Vertical Distance, \bar{y}	Non-dimensional Horizontal Velocity, \bar{u} (Angle 0)	Non-dimensional Horizontal Velocity, \bar{u} (Angle 30)	Non-dimensional Horizontal Velocity, \bar{u} (Angle 45)	Non-dimensional Horizontal Velocity, \bar{u} (Angle 60)	Non-dimensional Horizontal Velocity, \bar{u} (Angle 90)
1.00	0.00000	0.00000	0.00000	0.00000	0.00000
0.95	-0.08853	-0.22133	-0.28419	-0.31367	0.31299
0.90	-0.14139	-0.28811	-0.34372	-0.37180	0.37906
0.85	-0.14126	-0.24311	-0.28114	-0.30926	0.33793
0.80	-0.11099	-0.15650	-0.18309	-0.21767	0.27140
0.75	-0.08158	-0.08556	-0.10562	-0.14397	0.21131
0.70	-0.05042	-0.03216	-0.05050	-0.09015	0.16028
0.65	-0.02672	-0.00095	-0.01738	-0.05259	0.11681
0.60	-0.00818	0.01021	-0.00296	-0.02824	0.07712
0.55	0.00879	0.01594	0.00743	-0.00613	0.03956
0.50	0.02091	0.01190	0.00873	0.00929	0.00203
0.45	0.03566	0.01725	0.01842	0.03015	-0.03673
0.40	0.04859	0.02441	0.02831	0.04963	-0.07771
0.35	0.06550	0.04901	0.05506	0.08079	-0.12216
0.30	0.08214	0.08378	0.09221	0.11902	-0.17037
0.25	0.09846	0.13280	0.14878	0.17484	-0.22378
0.20	0.10825	0.18676	0.21675	0.24286	-0.27964
0.15	0.10867	0.22557	0.27849	0.31166	-0.33096
0.10	0.08889	0.22107	0.29699	0.34198	-0.34849
0.05	0.05825	0.15283	0.22115	0.26664	-0.27077
0.00	0.00000	0.00000	0.00000	0.00000	0.00000

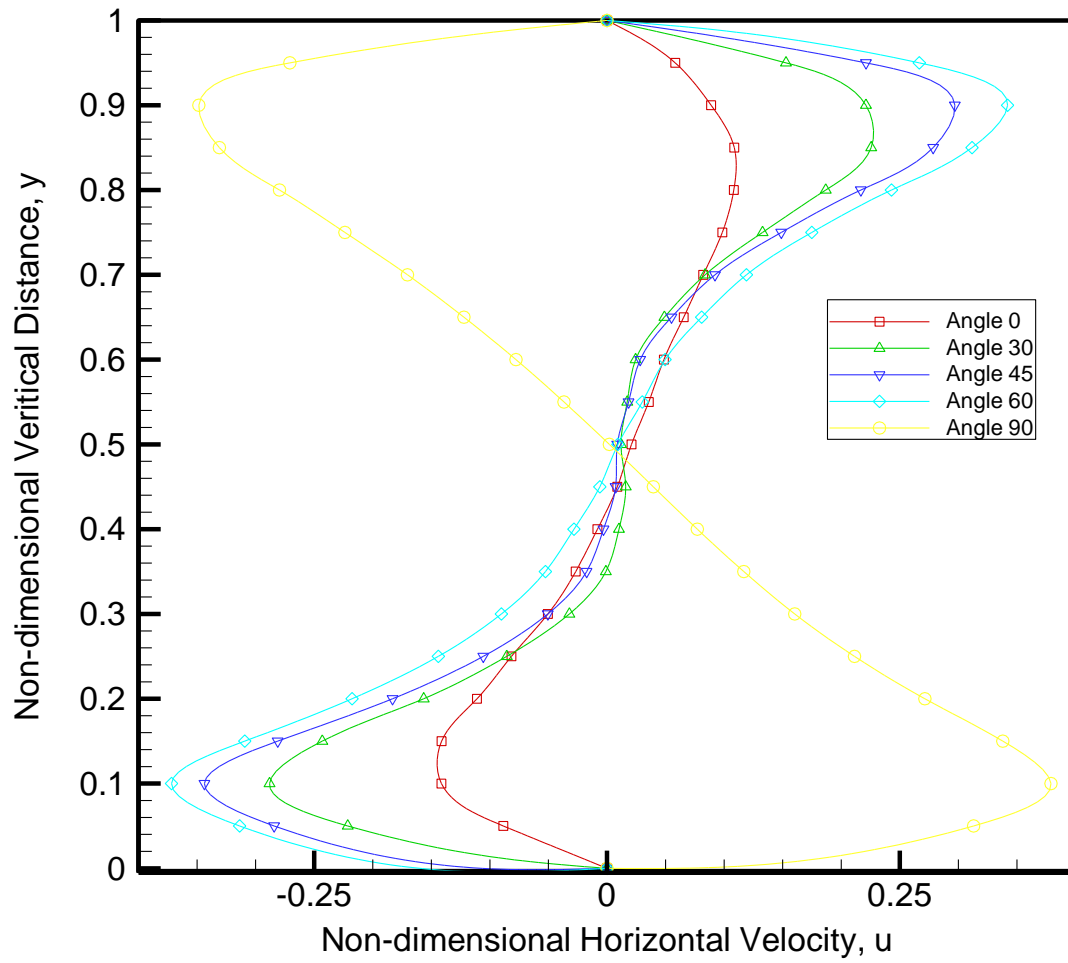


Figure 36- Steady State Non-dimensional Horizontal Velocity Distributions along the Vertical Centerline of the Inclined Enclosure for Constant Fluid Properties at $Ra= 10^5$

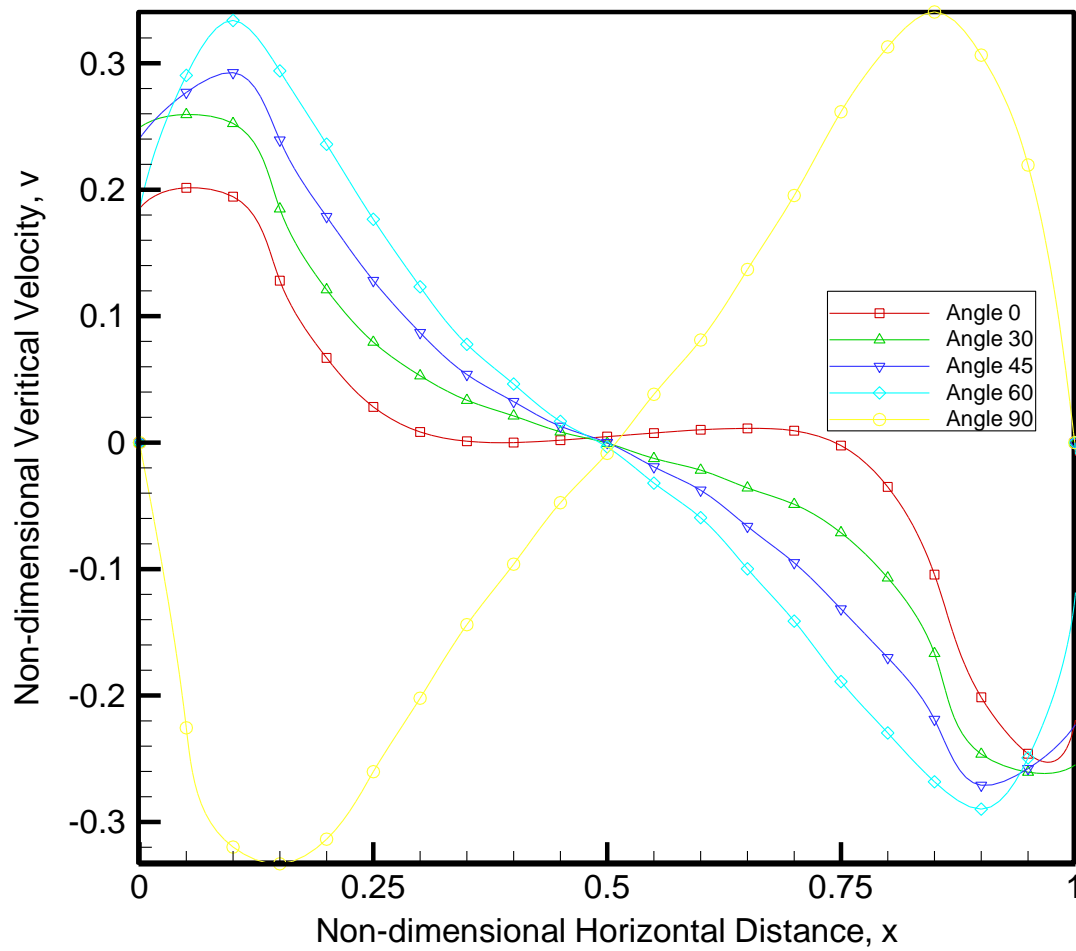


Figure 37- Steady State Non-dimensional Vertical Velocity Distributions along the Horizontal Centerline of the Inclined Enclosure for Constant Fluid Properties at $Ra= 10^5$

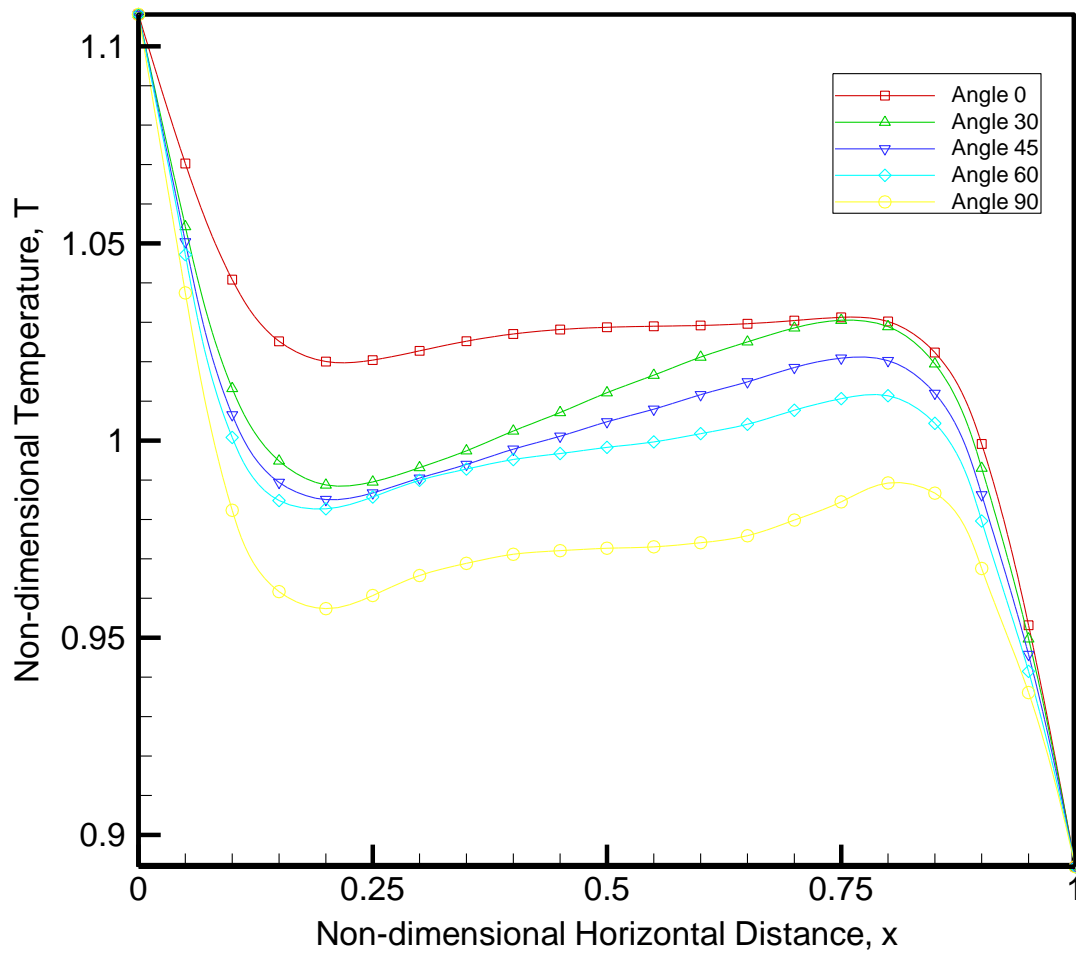


Figure 38- Steady State Non-dimensional Temperature Distributions along the Horizontal Centerline of the Inclined Enclosure for Constant Fluid Properties at $Ra = 10^5$

Chapter 10

The Study of Natural Convection Using the Second Order Accurate in Time Model

The study of natural convection phenomenon is also conducted using the second order accurate in time model. The results of the second order accurate in time model are presented in this chapter. Also, these results are compared to the first order accurate in time model to verify the accuracy of the second order accurate in time model for different Rayleigh numbers.

10.1. Results of the Second Order Accurate in Time Model

For, $Ra = 10^5$, the temperature between the vertical walls of the enclosure was taken to be 64.8 K. A uniform and orthogonal mesh of 21 x 21 was utilized. The predicted primitive variables along the horizontal centerline and vertical centerline of the enclosure are presented in Table 25 and 26 respectively at steady state. The qualitative results of general characteristics of computational mesh, vectors, streamlines, temperature contours are presented in Figure 39, 42, 43 and 44. The flow inside square enclosure is not unicellular. Two secondary vortices formation is evident from Figure 43. The core flow is more stagnant relative to the boundary regions. The velocity (u and v) histograms at a point closer to the high temperature wall where $\bar{x} = 0.05$ and $\bar{y} = 0.5$ are presented in Figures 40 and 41. The non-dimensional horizontal velocity distribution along the vertical centerline, vertical velocity distribution along the horizontal centerline and temperature distribution along the horizontal centerline are presented in Figure 45, 46 and 47 respectively. The maximum non-dimensional vertical velocity is observed as 0.25 at near the hot wall and the maximum non-dimensional horizontal velocity is observed as 0.11 at near the cold wall. The temperature gradients are higher near the walls and equal to zero at the center of the

enclosure. This fact is easily illustrated by isotherms being almost horizontal in the center of the enclosure. The qualitative results of general characteristics of computational mesh, vectors, streamlines, temperature contours are also presented in Figure 48, 49, 50 and 51 for grid size of 81 x 81.

Table 25- Steady State Non-dimensional values of Primitive Variables along the Horizontal Centerline of the Enclosure for $Ra = 10^5$ in Second Order Time Accurate Approximation ($\Delta \bar{t} = 1.0E-5$ and 21 X 21)

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Vertical Velocity, \bar{v}	Non-dimensional Pressure, \bar{p}	Non-dimensional Temperature, \bar{T}
1.00	0.0000000	1605.1374	0.8919274
0.95	-0.2440447	1605.1374	0.9541405
0.90	-0.2002889	1605.1374	0.9988405
0.85	-0.1043035	1605.1374	1.0214050
0.80	-0.0357540	1605.1374	1.0290520
0.75	-0.0031484	1605.1374	1.0300350
0.70	0.0086785	1605.1374	1.0291910
0.65	0.0106349	1605.1374	1.0283880
0.60	0.0095869	1605.1374	1.0279430
0.55	0.0070514	1605.1374	1.0277350
0.50	0.0043028	1605.1374	1.0274820
0.45	0.0015100	1605.1374	1.0269280
0.40	-0.0003056	1605.1374	1.0258250
0.35	0.0007269	1605.1374	1.0240220
0.30	0.0079743	1605.1374	1.0216240
0.25	0.0276678	1605.1374	1.0193210
0.20	0.0664488	1605.1374	1.0189950
0.15	0.1278468	1605.1374	1.0242100
0.10	0.1949693	1605.1374	1.0399920
0.05	0.2027110	1605.1374	1.0697490
0.00	0.0000000	1605.1374	1.1080730

Table 26- Steady State Non-dimensional values of Primitive Variables along the Vertical Centerline of the Enclosure for $Ra = 10^5$ in Second Order Time Accurate Approximation ($\Delta\bar{t}= 1.0E-5$ and 21×21)

Non-dimensional Vertical Distance, \bar{y}	Non-dimensional Horizontal Velocity, \bar{u}	Non-dimensional Pressure, \bar{p}	Non-dimensional Temperature, \bar{T}
1.00	0.0000000	1605.1353	1.0821941
0.95	0.0595109	1605.1354	1.0821070
0.90	0.0905230	1605.1357	1.0811453
0.85	0.1101803	1605.1358	1.0782586
0.80	0.1092251	1605.1361	1.0745235
0.75	0.0986764	1605.1363	1.0685488
0.70	0.0816850	1605.1365	1.0627430
0.65	0.0644987	1605.1367	1.0548044
0.60	0.0472117	1605.1370	1.0472641
0.55	0.0340867	1605.1372	1.0370585
0.50	0.0192216	1605.1374	1.0274822
0.45	0.0071675	1605.1376	1.0151296
0.40	-0.0096449	1605.1379	1.0043528
0.35	-0.0278360	1605.1381	0.9916311
0.30	-0.0510775	1605.1383	0.9815297
0.25	-0.0815223	1605.1386	0.9703938
0.20	-0.1101736	1605.1388	0.9618196
0.15	-0.1394019	1605.1391	0.9521419
0.10	-0.1390995	1605.1393	0.9454112
0.05	-0.0867933	1605.1396	0.9394930
0.00	0.0000000	1605.1398	0.9369507

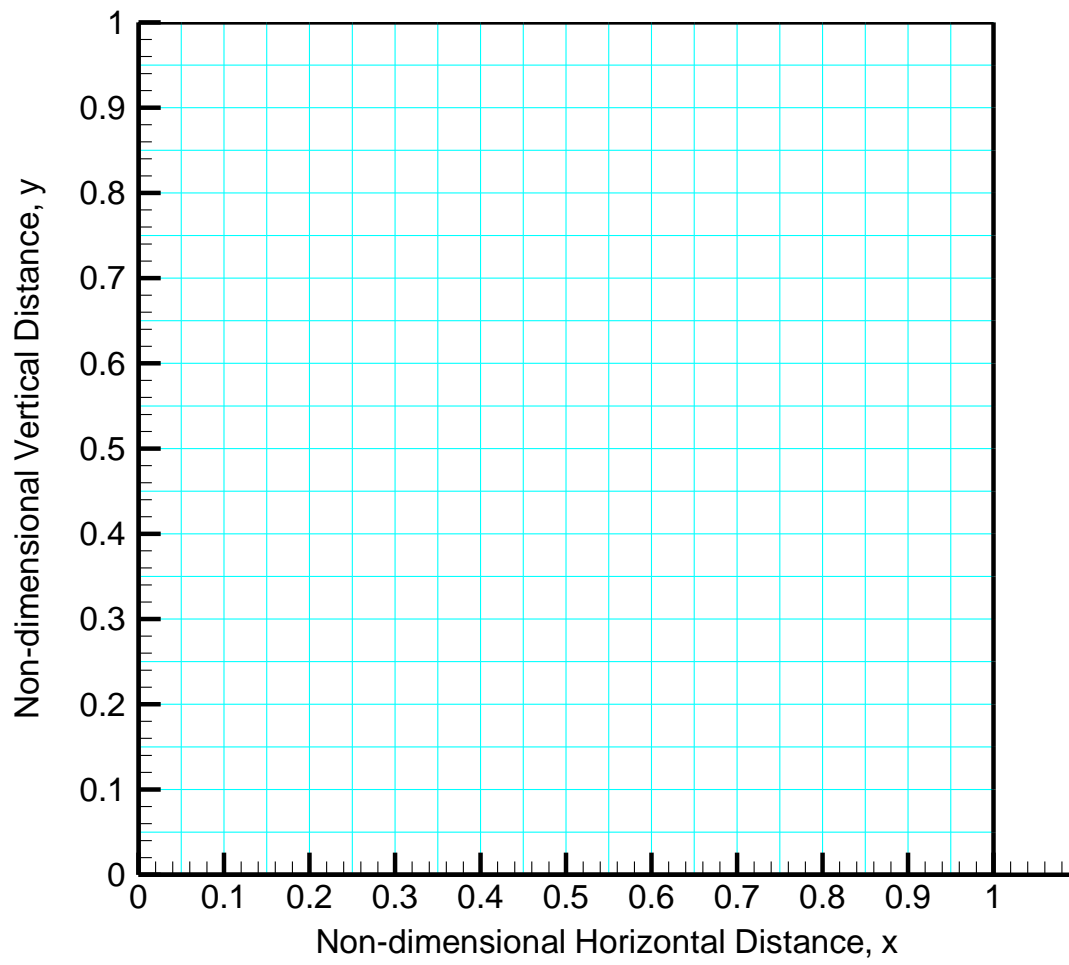


Figure 39- Computational Mesh of the Non-dimensional domain for the Natural Convection Case in Second Order Time Accurate Approximation (21×21 , $Ra = 10^5$)

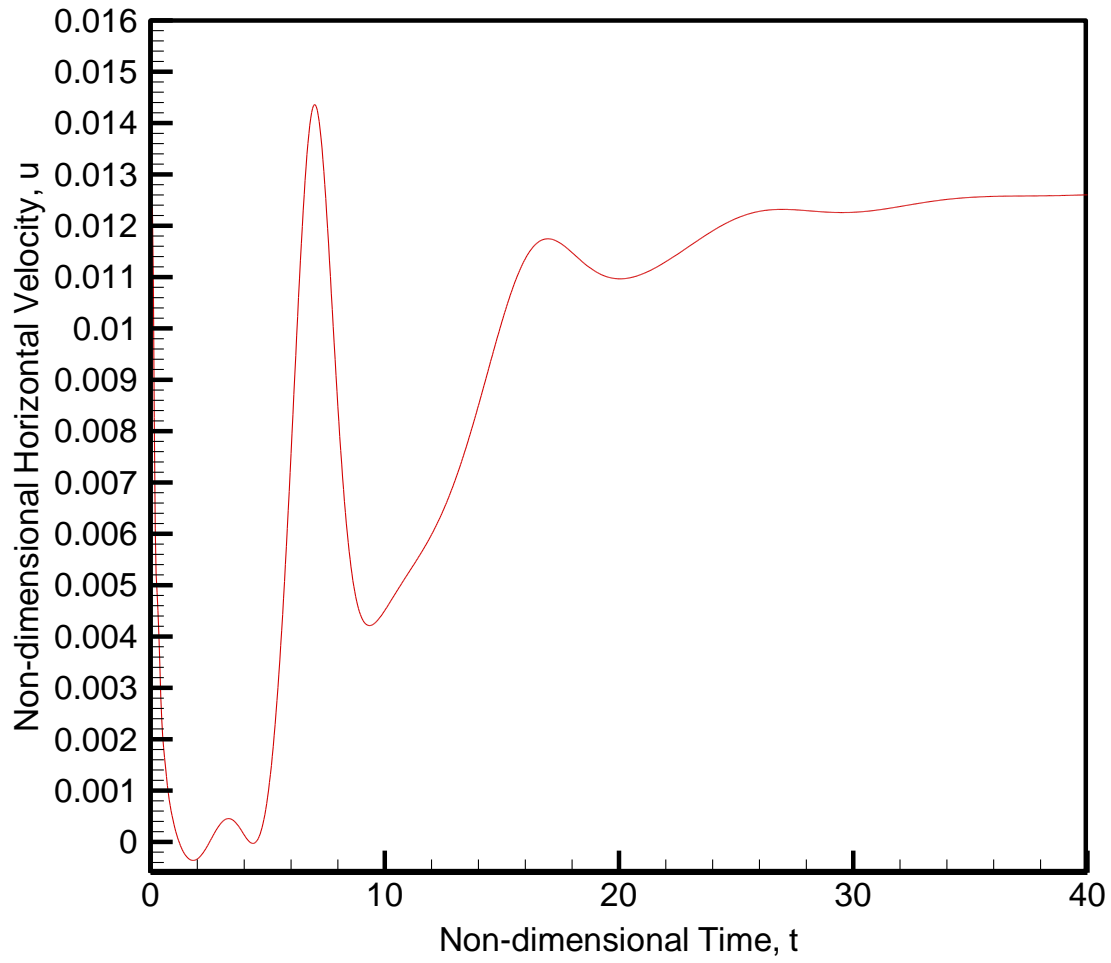


Figure 40- Non-dimensional Horizontal Velocity Histogram for $Ra=10^5$ at $\bar{X}=0.2$ and $\bar{Y}=0.5$ (21×21 , $\Delta \bar{t} = 1.0E-5$) in Second Order Time Accurate Approximation

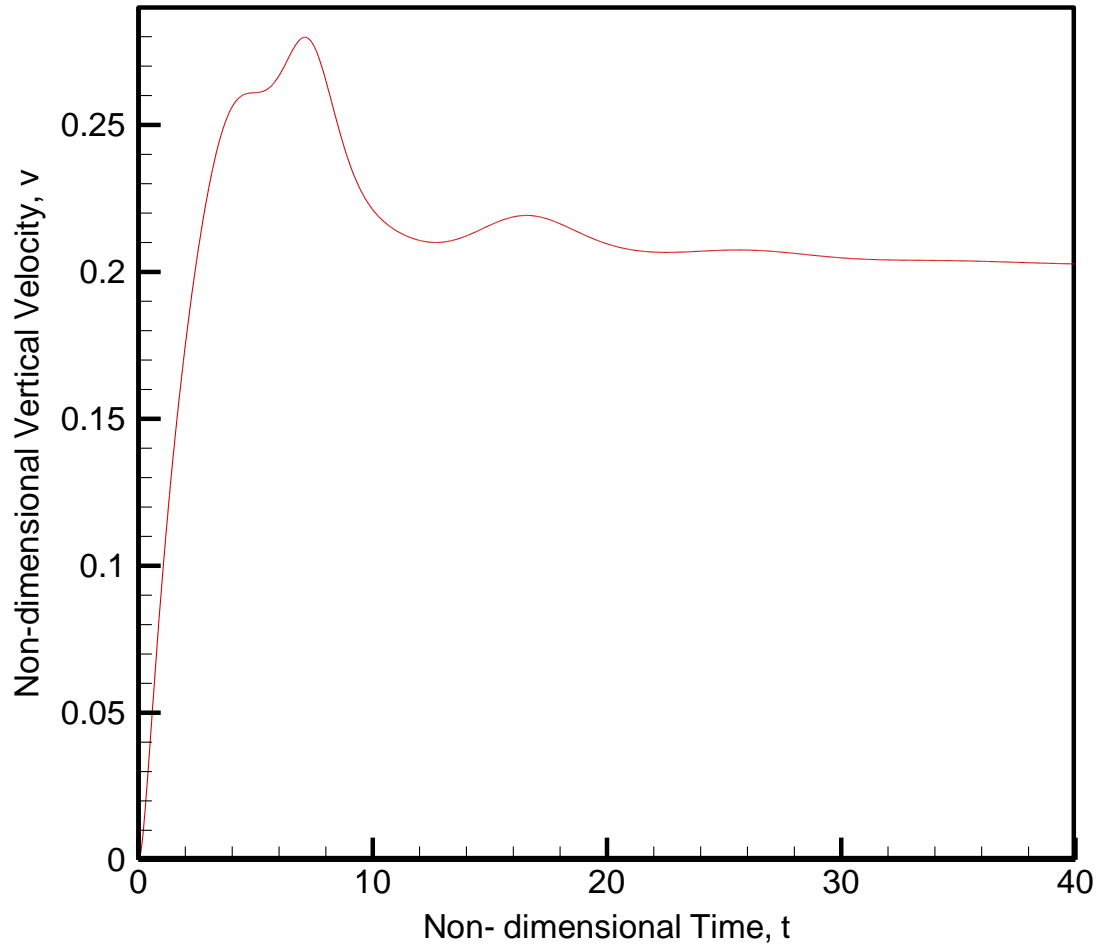


Figure 41- Non-dimensional Vertical Velocity Histogram for $Ra=10^5$ at $\bar{X}=0.2$ and $\bar{Y}=0.5$
 $(21 \times 21, \Delta \bar{t} = 1.0E-5)$ in Second Order Time Accurate Approximation

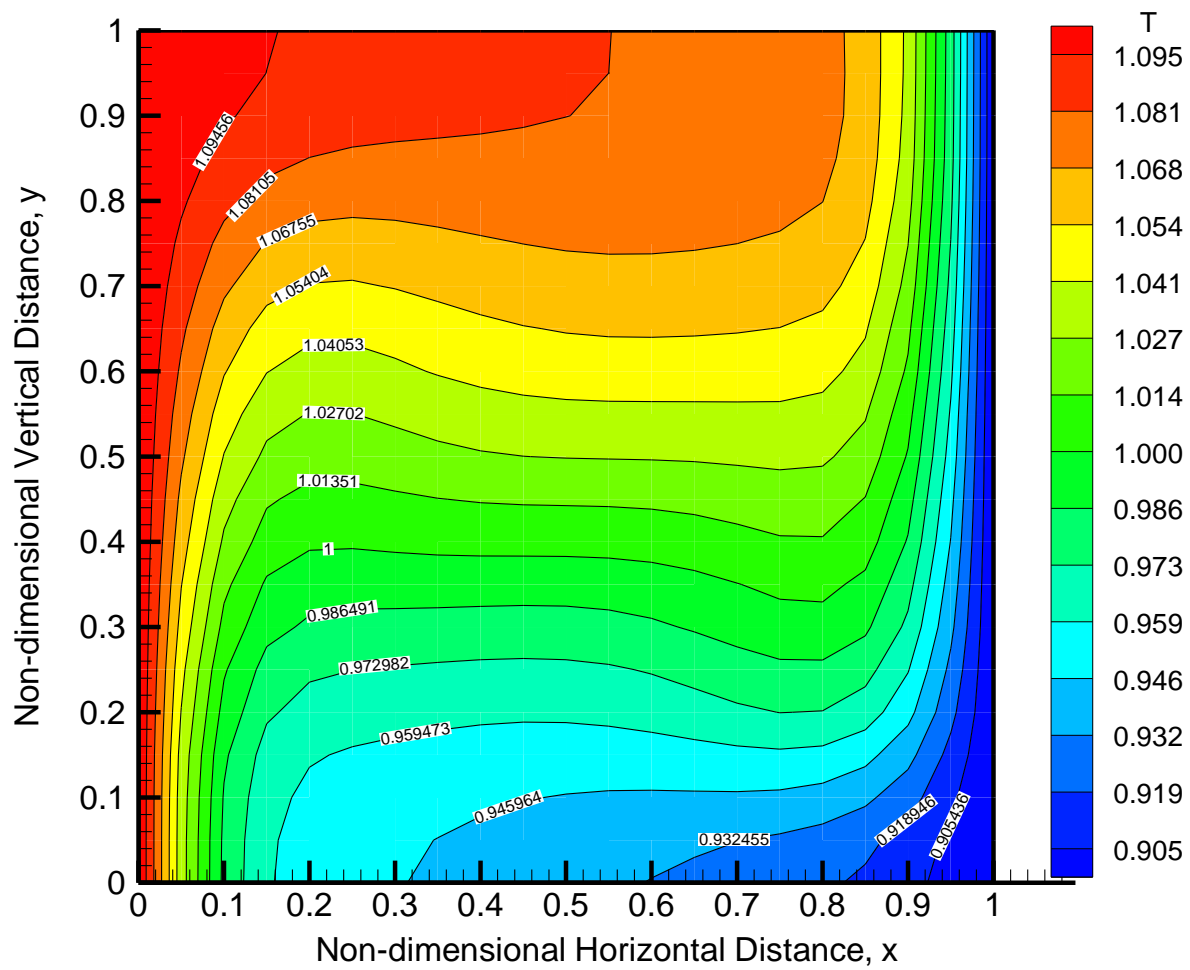


Figure 42- Steady State Non-dimensional Temperature Contours for $Ra=10^5$
 $(21 \times 21, \Delta \bar{t} = 1.0E-5)$ in Second Order Time Accurate Approximation

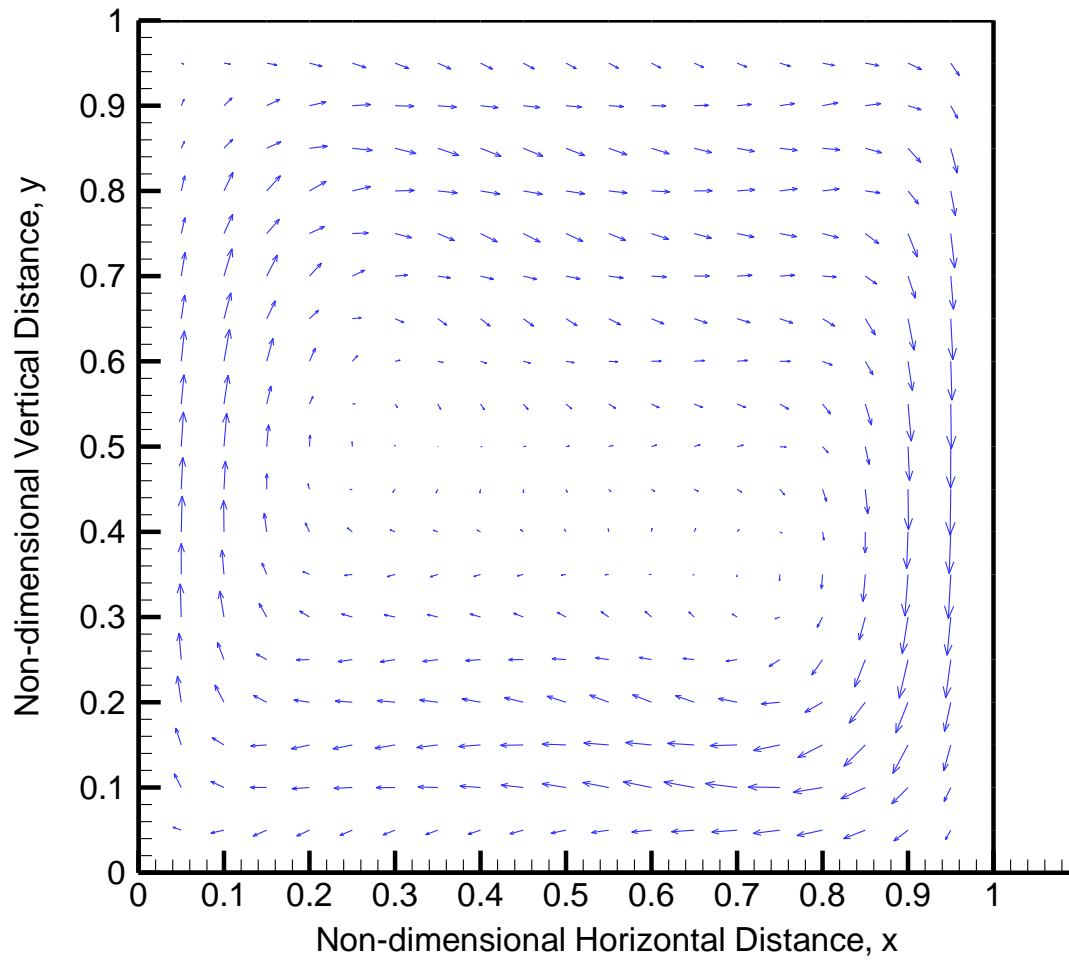


Figure 43- Steady State Non-dimensional Velocity Vectors for $Ra=10^5$ (21×21 , $\Delta \bar{t} = 1.0E-5$) in Second Order Time Accurate Approximation

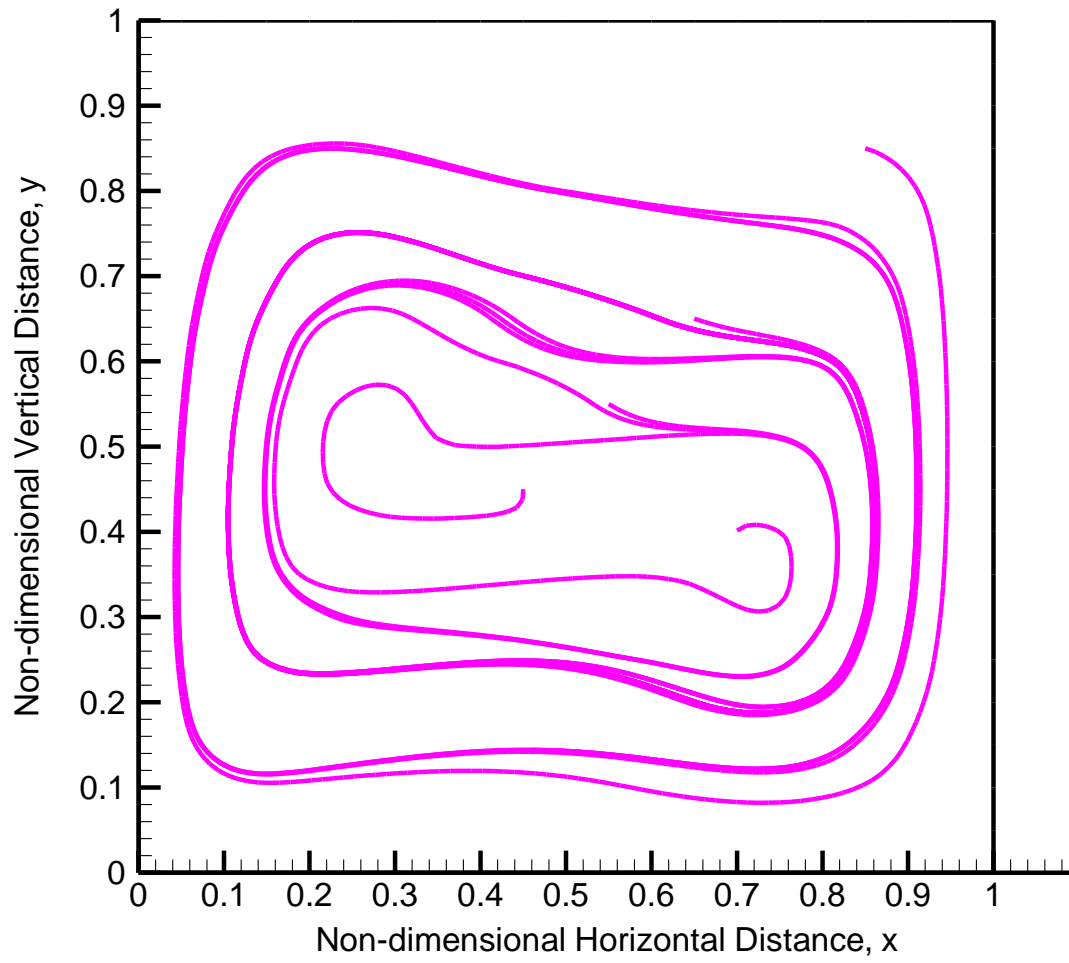


Figure 44- Steady State Non-dimensional Velocity Streamlines for $Ra=10^5$
 (21 x 21, $\Delta \bar{t} = 1.0E-5$) in Second Order Time Accurate Approximation

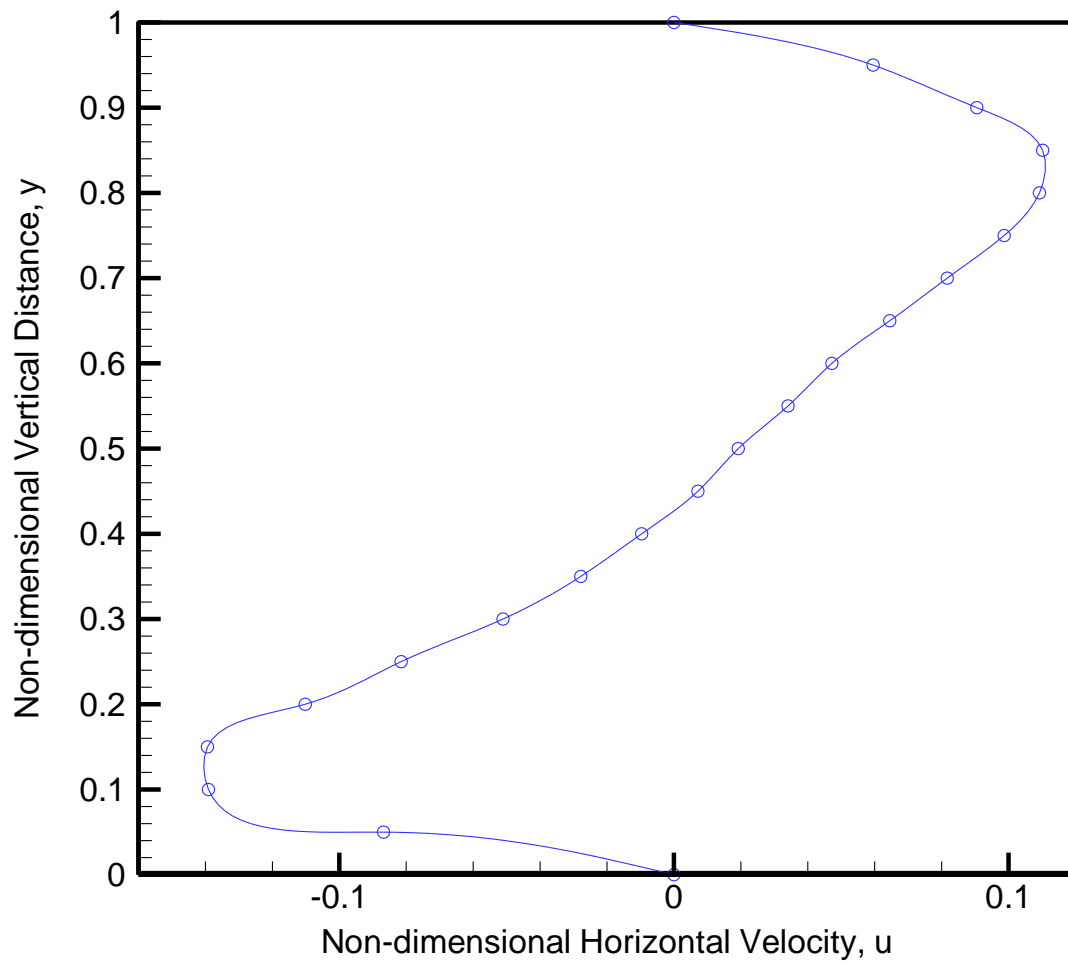


Figure 45- Steady State Non-dimensional Horizontal Velocity Distribution along the Vertical Centerline of the Enclosure for $Ra=10^5$ (21×21 , $\Delta \bar{t} = 1.0E-5$) in Second Order Time Accurate Approximation

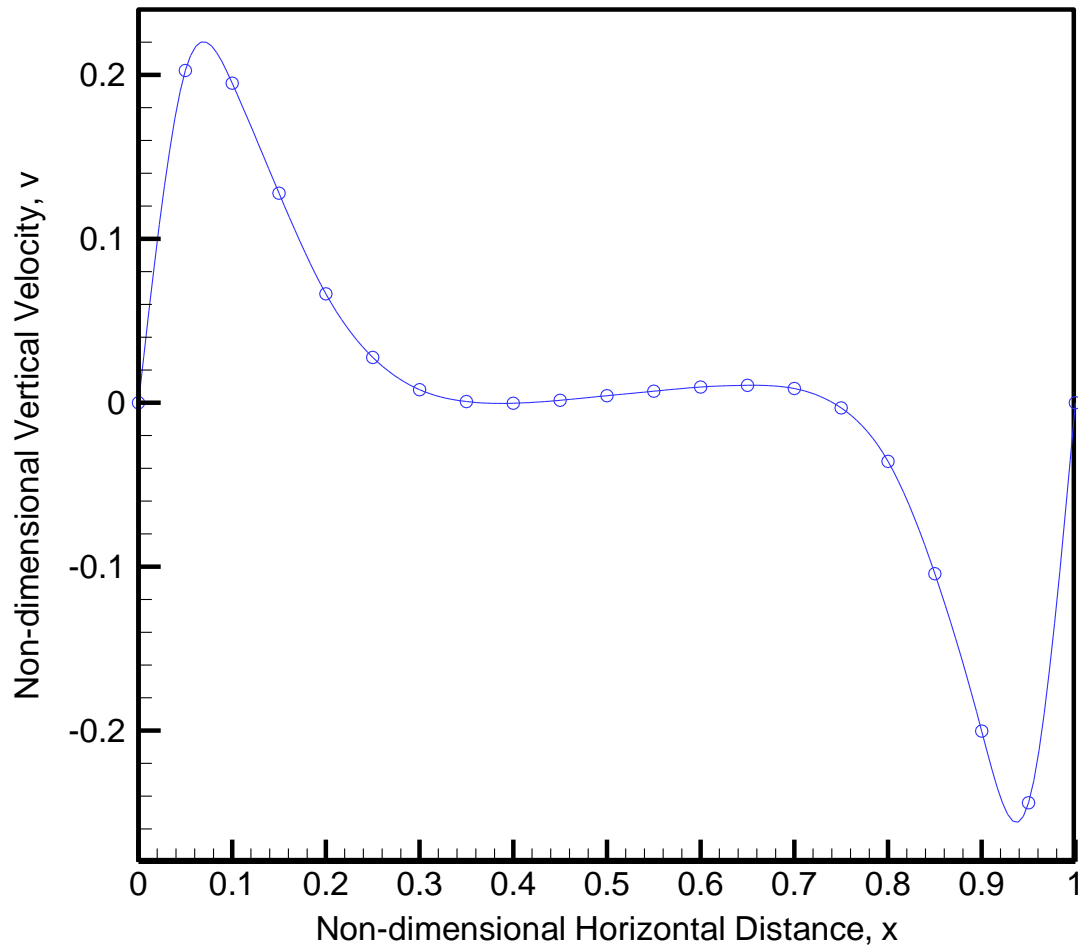


Figure 46- Steady State Non-dimensional Vertical Velocity Distribution along the Horizontal Centerline of the Enclosure for $Ra=10^5$ (21×21 , $\Delta \bar{t} = 1.0E-5$) in Second Order Time Accurate Approximation

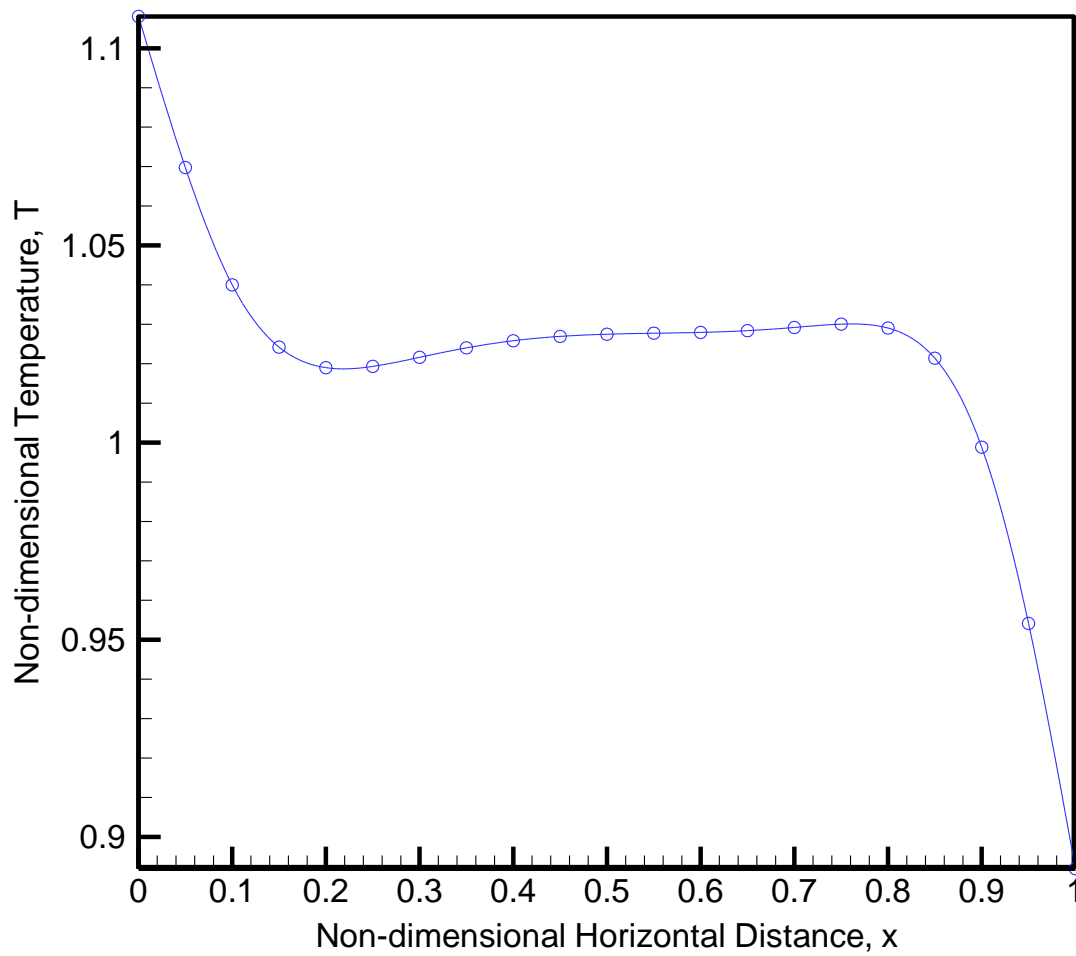


Figure 47- Steady State Non-dimensional Temperature Distribution along the Horizontal Centerline of the Enclosure for $Ra=10^5$ (21×21 , $\Delta \bar{t} = 1.0E-5$) in Second Order Time Accurate Approximation

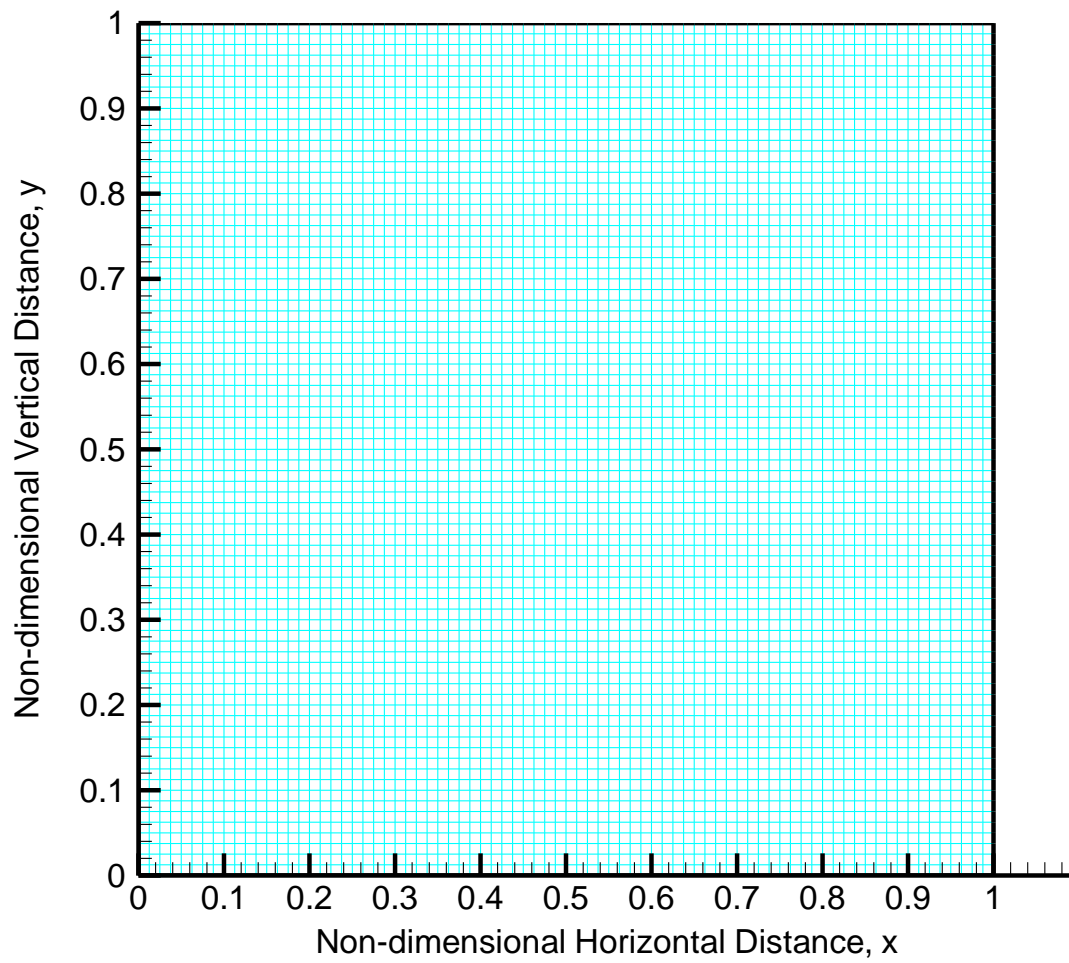


Figure 48- Computational Mesh of the Non-dimensional domain for the Natural Convection Case in Second Order Time Accurate Approximation (81×81 , $Ra = 10^5$)

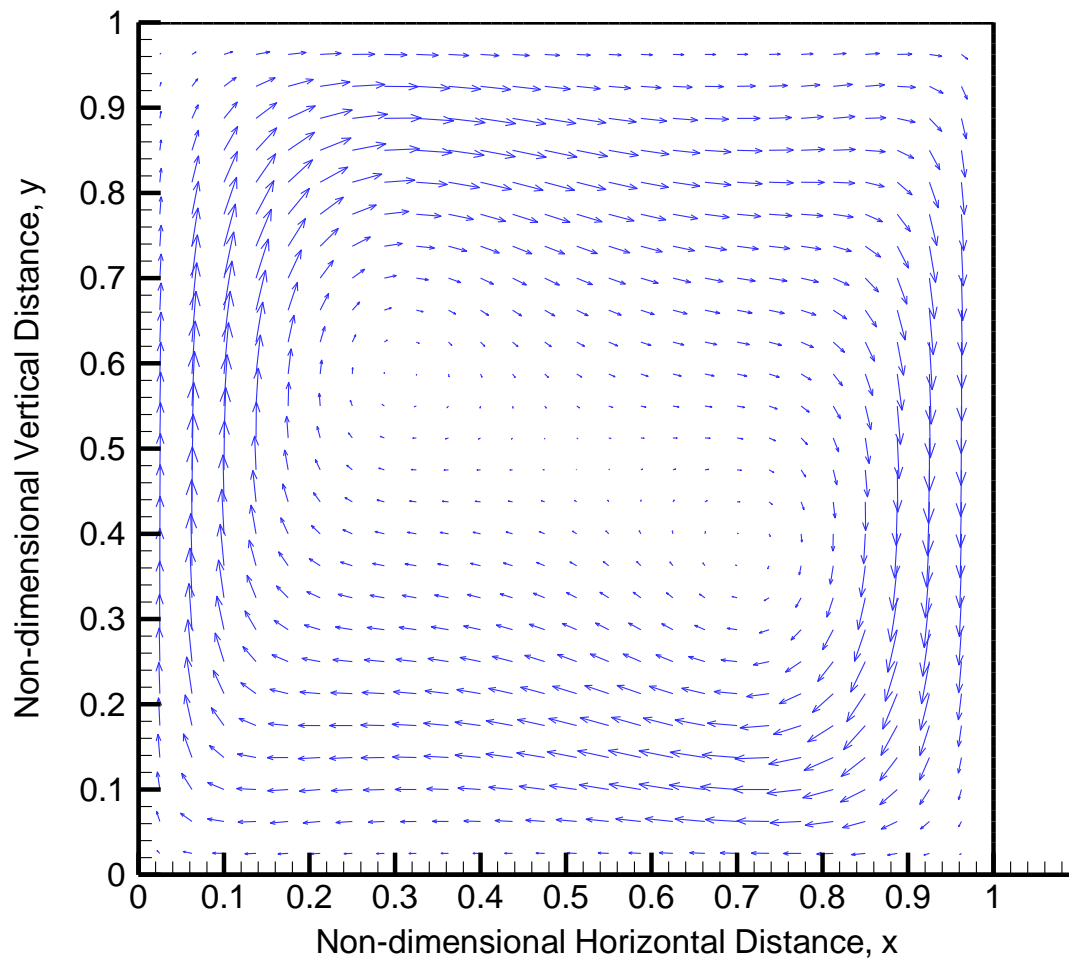


Figure 49- Steady State Non-dimensional Velocity Vectors for $Ra=10^5$ (81×81 , $\Delta \bar{t} = 1.0E-6$) in Second Order Time Accurate Approximation

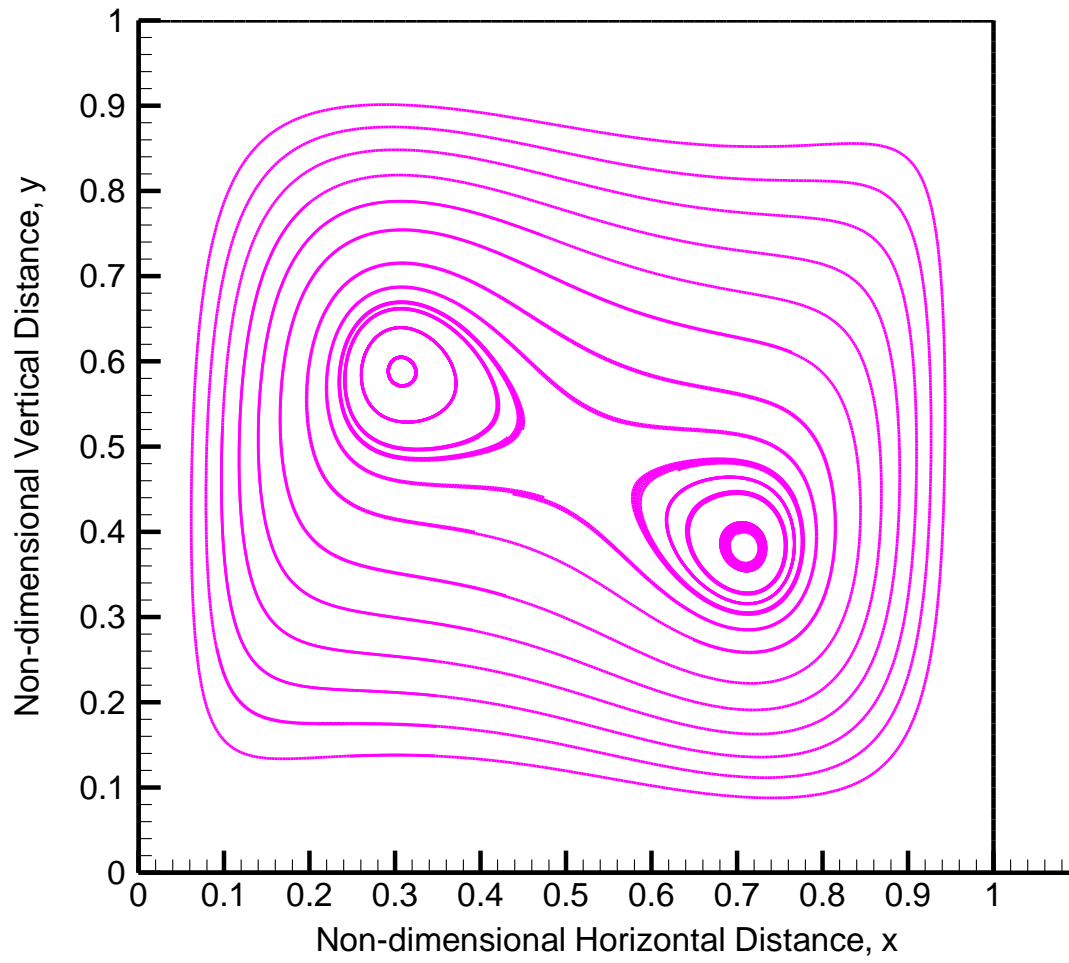


Figure 50- Steady State Non-dimensional Velocity Streamlines for $Ra=10^5$
(81 x 81, $\Delta \bar{t} = 1.0E-6$) in Second Order Time Accurate Approximation

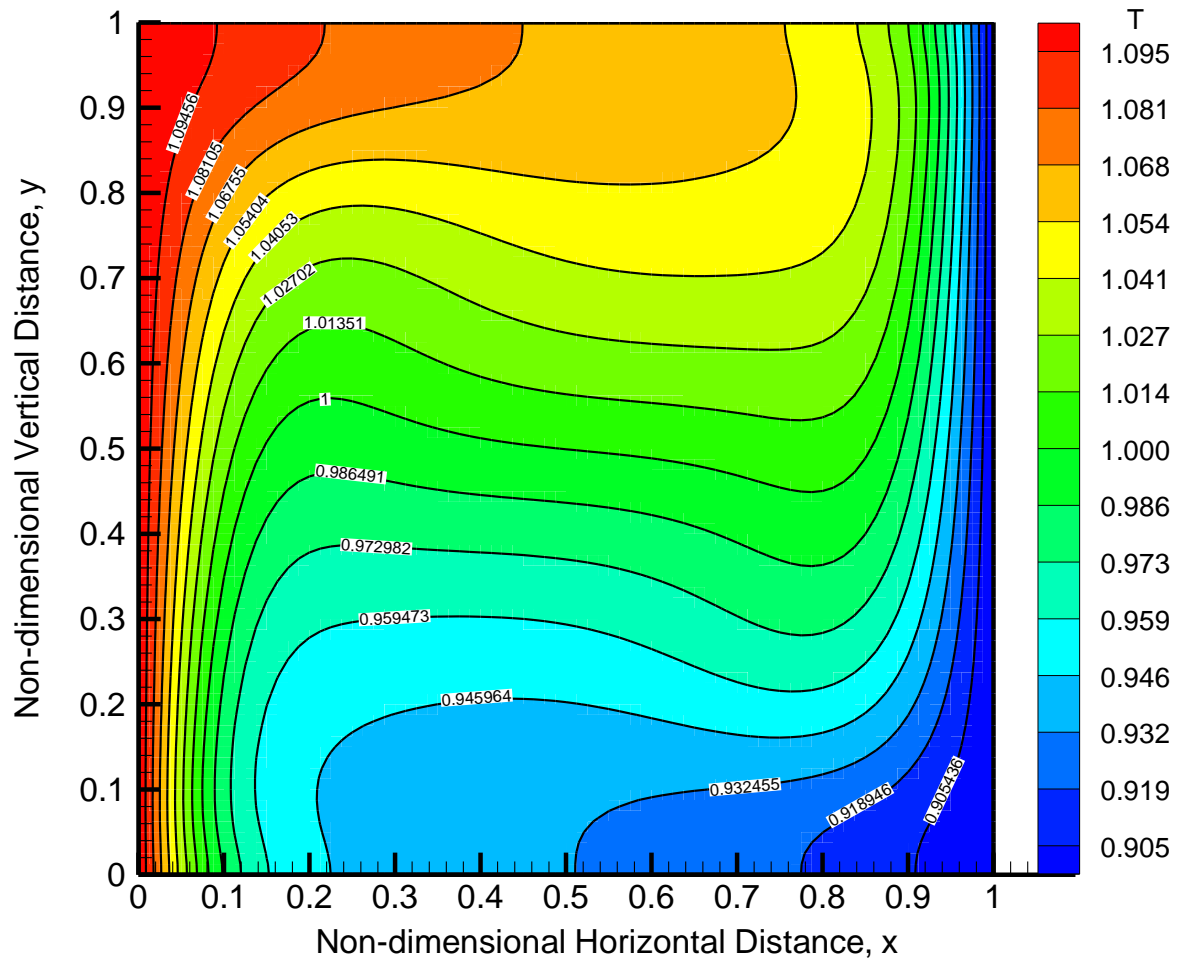


Figure 51- Steady State Non-dimensional Temperature Contours for $Ra=10^5$
 $(81 \times 81, \Delta \bar{t} = 1.0E-6)$ in Second Order Time Accurate Approximation

10.2. The Comparison between First Order and Second Order Accurate in Time Models

The comparison of the present model with first order accurate in time CMSIP model shows negligible differences. A quantitative comparisons of the first order and second order accurate in time CMSIP models for Rayleigh number 10^5 with $\Delta T = 64.84$ deg K are presented in Table 27 through 29. The qualitative comparison is presented in Figure 52 through 54. The results for Rayleigh number 10^4 by changing the temperature difference between the vertical walls of the enclosure with the same size are presented quantitatively in Table 30 through 32 and qualitatively in Figures 55 through 57. These results indicate that the accuracy of second order accurate in time model can be verified successfully with first order accurate in time model.

The differences can be seen when one compares the velocity and temperature profiles in horizontal and vertical directions for different Rayleigh numbers. The comparisons of the above results at constant wall temperature difference for different Rayleigh numbers 10^4 and 10^5 are presented. The qualitative comparisons of these results (velocity and temperature) are presented in Figure 58 through 60. As shown in Figure 58, the magnitudes of the horizontal velocities are predicted higher for Ra number 10^4 compared to Ra 10^5 . But the magnitudes of the vertical velocities and temperature are predicted higher for Ra 10^5 as depicted in Figure 59 and 60, because of the increase in temperature difference between the walls of the enclosure.

Table 27- Comparison of Non-dimensional values of Horizontal Velocity along the Vertical Centerline between First Order and Second Order Time Accurate Approximations ($Ra = 10^5$, $\Delta \bar{t} = 1.0E-5$ and 21×21)

Non-dimensional Vertical Distance, \bar{y}	Non-dimensional Horizontal Velocity, \bar{u} (1st Oder)	Non-dimensional Horizontal Velocity, \bar{u} (2nd Order)	% Deviation in \bar{u}
1.00	0.0000000	0.0000000	0.0000000
0.95	0.0592333	0.0595109	0.4665191
0.90	0.0900378	0.0905231	0.5360290
0.85	0.1097794	0.1101803	0.3638581
0.80	0.1089421	0.1092251	0.2590979
0.75	0.0986343	0.0986764	0.04262419
0.70	0.0817787	0.0816850	-0.1147088
0.65	0.0647388	0.0644987	-0.3723174
0.60	0.0474760	0.0472117	-0.5598186
0.55	0.0344048	0.0340867	-0.9334432
0.50	0.0194978	0.0192216	-1.4371830
0.45	0.0074592	0.0071675	-4.0704532
0.40	-0.0094205	-0.0096449	2.3259011
0.35	-0.0276257	-0.0278360	0.7554943
0.30	-0.0509145	-0.0510775	0.3190640
0.25	-0.0814199	-0.0815223	0.1257077
0.20	-0.1101176	-0.1101736	0.0508288
0.15	-0.1395627	-0.1394019	-0.1153499
0.10	-0.1394197	-0.1390995	-0.2301949
0.05	-0.0870163	-0.0867933	-0.2568744
0.00	0.0000000	0.0000000	0.0000000

Table 28- Comparison of Non-dimensional values of Vertical Velocity along the Horizontal Centerline between First Order and Second Order Time Accurate Approximations ($Ra = 10^5$, $\Delta \bar{t} = 1.0E-5$ and 21×21)

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Vertical Velocity, \bar{v} (1st Order)	Non-dimensional Vertical Velocity, \bar{v} (2nd Order)	% Deviation in \bar{v}
1.00	0.0000000	0.0000000	0.0000000
0.95	-0.2440669	-0.2440447	-0.0090966
0.90	-0.2001967	-0.2002889	0.0460335
0.85	-0.1041439	-0.1043035	0.1530149
0.80	-0.0356895	-0.0357540	0.1803991
0.75	-0.0030238	-0.0031484	3.9582403
0.70	0.0087026	0.0086785	-0.2781932
0.65	0.0107321	0.0106349	-0.9147241
0.60	0.0096055	0.0095869	-0.1936064
0.55	0.0071374	0.0070514	-1.2197429
0.50	0.0043102	0.0043028	-0.1719107
0.45	0.0015751	0.0015100	-4.3091294
0.40	-0.0003151	-0.0003056	-3.1057804
0.35	0.0007948	0.0007269	-9.3348089
0.30	0.0079977	0.0079743	-0.2939556
0.25	0.0277972	0.0276678	-0.4676911
0.20	0.0664985	0.0664488	-0.0748997
0.15	0.1278577	0.1278468	-0.0085258
0.10	0.1946601	0.1949693	0.1585890
0.05	0.2022876	0.2027110	0.2088687
0.00	0.0000000	0.0000000	0.0000000

Table 29- Comparison of Non-dimensional values of Temperature along the Horizontal Centerline between First Order and Second Order Time Accurate Approximations ($Ra = 10^5$, $\Delta \bar{t} = 1.0E-5$ and 21×21)

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Temperature, \bar{T} (1st Order)	Non-dimensional Temperature, \bar{T} (2nd Order)	% Deviation in \bar{T}
1.00	0.8919274	0.8919274	0.0000000
0.95	0.9541786	0.9541405	-0.0039931
0.90	0.9990383	0.9988405	-0.0198029
0.85	1.0216900	1.0214050	-0.0279027
0.80	1.0293790	1.0290520	-0.0317768
0.75	1.0303840	1.0300350	-0.0338823
0.70	1.0295550	1.0291910	-0.0353675
0.65	1.0287620	1.0283880	-0.0363675
0.60	1.0283270	1.0279430	-0.0373561
0.55	1.0281210	1.0277350	-0.0375583
0.50	1.0278720	1.0274820	-0.0379568
0.45	1.0273150	1.0269280	-0.0376852
0.40	1.0262080	1.0258250	-0.0373358
0.35	1.0243960	1.0240220	-0.0365226
0.30	1.0219900	1.0216240	-0.0358253
0.25	1.0196770	1.0193210	-0.0349252
0.20	1.0193410	1.0189950	-0.0339550
0.15	1.0245320	1.0242100	-0.0314388
0.10	1.0402610	1.0399920	-0.0258655
0.05	1.0699080	1.0697490	-0.0148633
0.00	1.1080730	1.1080730	0.0000000

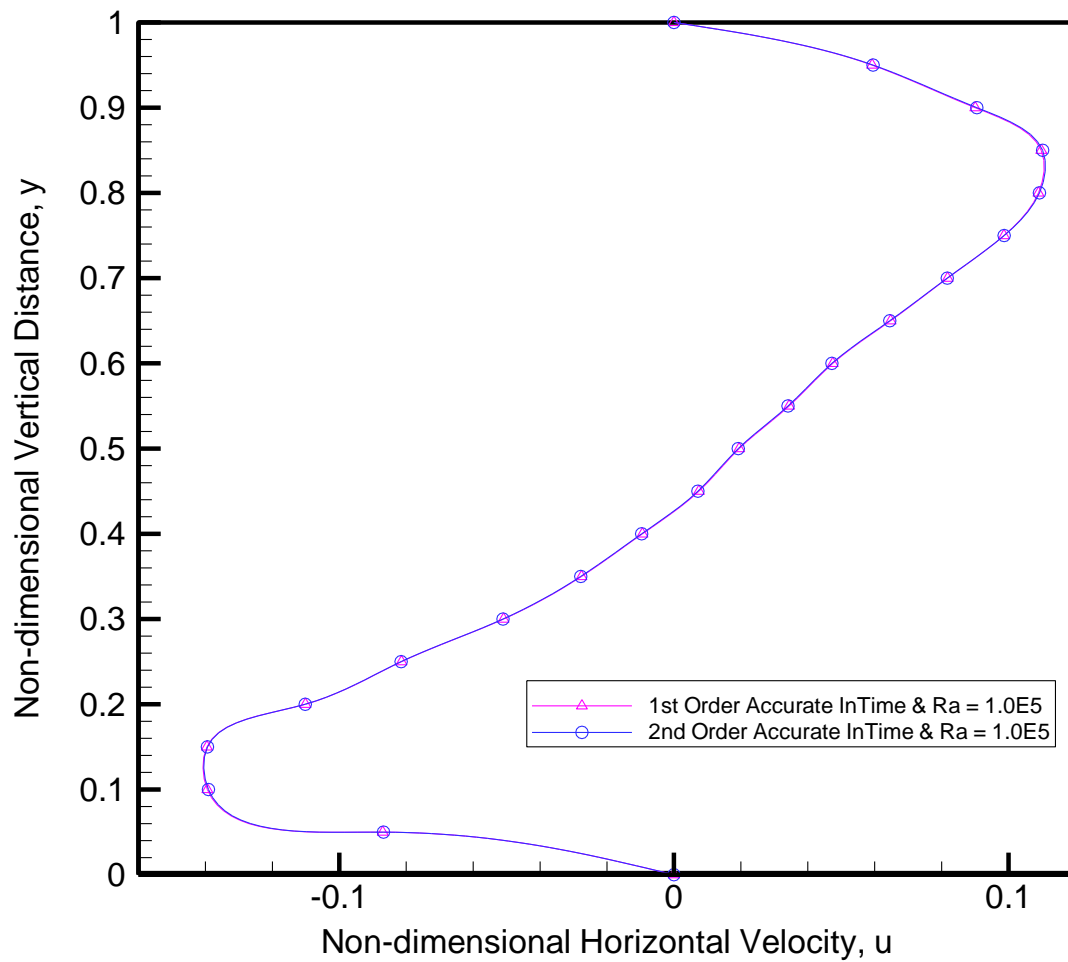


Figure 52- Comparison of Steady State Horizontal Velocity Distribution along the Vertical centerline of the Enclosure between First Order and Second Order Time Accurate Approximation for $Ra= 10^5$

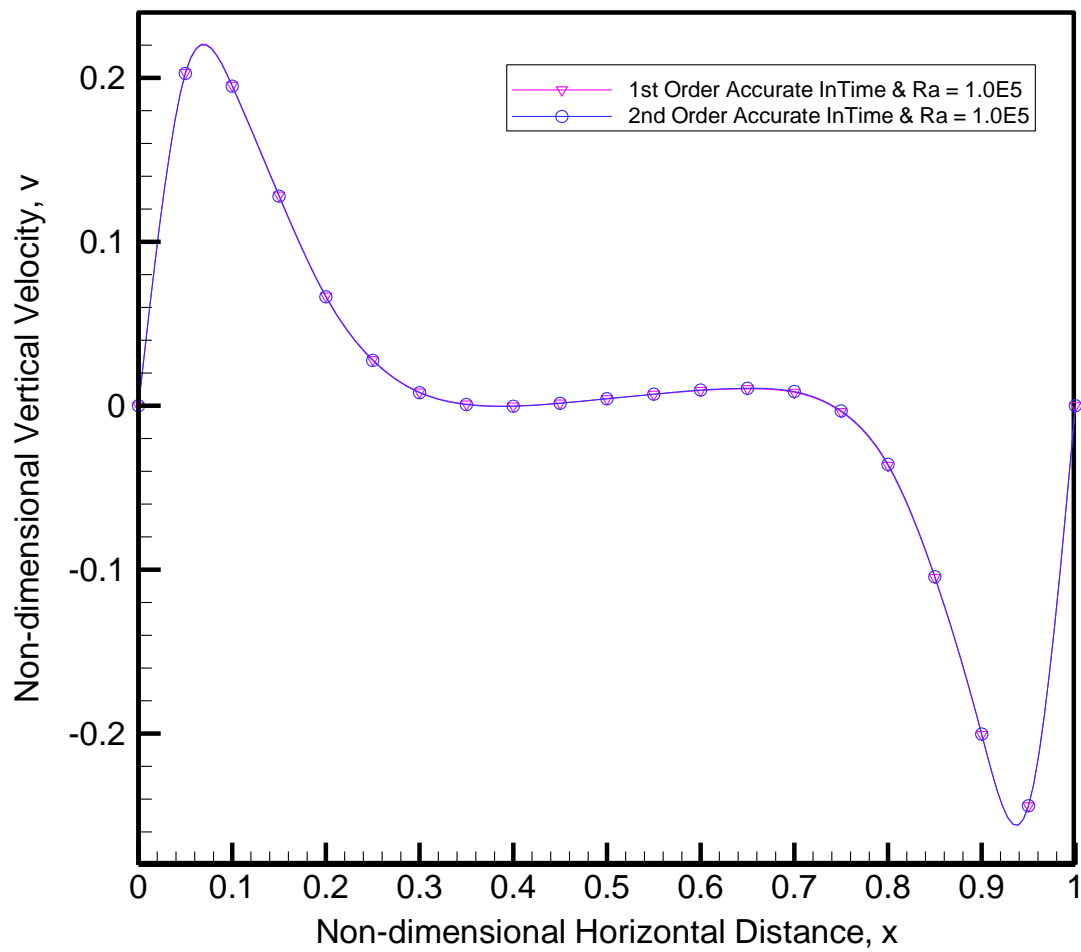


Figure 53- Comparison of Steady State Vertical Velocity Distribution along the Horizontal Centerline of the Enclosure between First Order and Second Order Time Accurate Approximation for $Ra= 10^5$

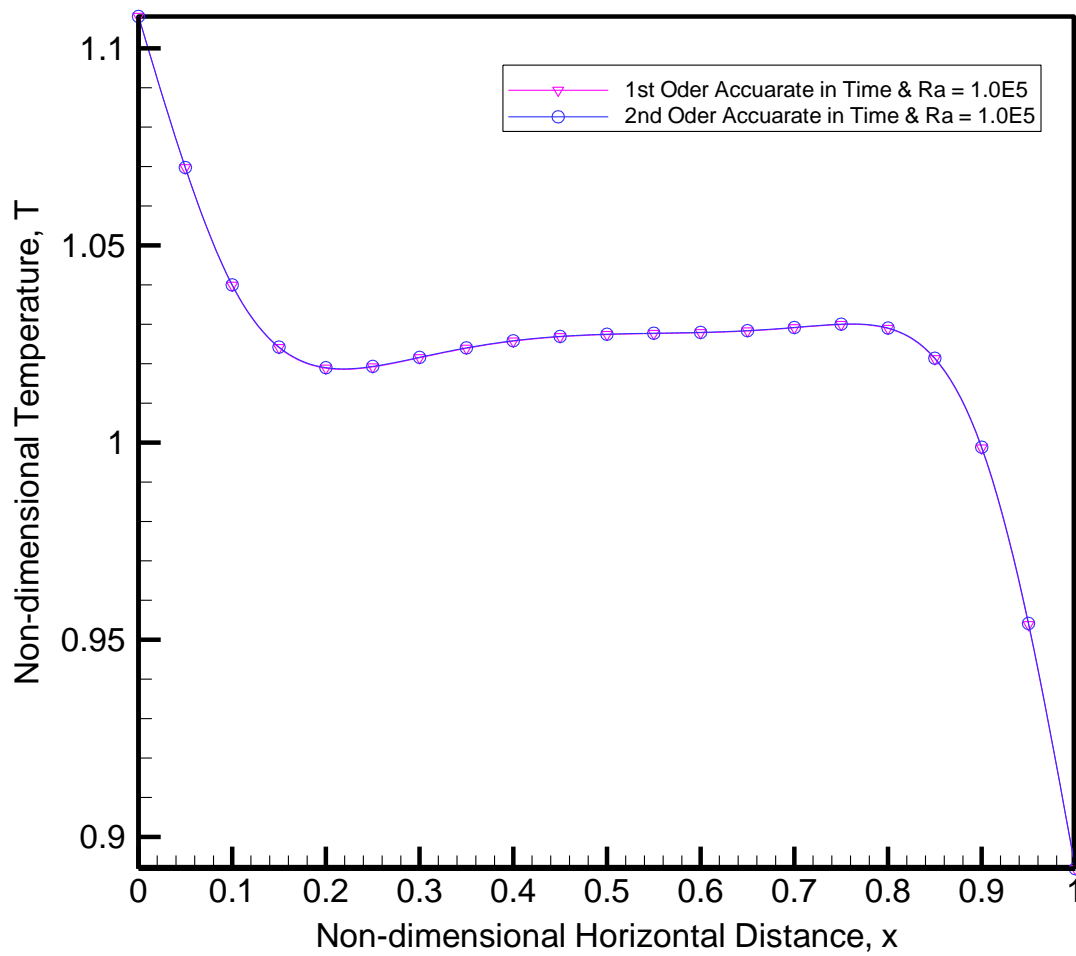


Figure 54- Comparison of Steady State Temperature Distribution along the Horizontal Centerline of the Enclosure between First Order and Second Order Time Accurate Approximation for $Ra = 10^5$

Table 30- Comparison of Non-dimensional values of Horizontal Velocity along the Vertical Centerline between First Order and Second Order Time Accurate Approximations ($Ra = 10^4$, $\Delta \bar{t} = 1.0E-6$ and 21×21)

Non-dimensional Vertical Distance, \bar{y}	Non-dimensional Horizontal Velocity, \bar{u} (1st Oder)	Non-dimensional Horizontal Velocity, \bar{u} (2nd Order)	% Deviation in \bar{u}
1.00	0.0000000	0.0000000	0.0000000
0.95	0.0718289	0.0716946	-0.1873362
0.90	0.1109938	0.1108348	-0.1434567
0.85	0.1421789	0.1418300	-0.2459987
0.80	0.1526182	0.1521436	-0.3119421
0.75	0.1559837	0.1553692	-0.3955095
0.70	0.1456733	0.1449529	-0.4969890
0.65	0.1320233	0.1312004	-0.6272084
0.60	0.1075737	0.1067081	-0.8111849
0.55	0.0826481	0.0816952	-1.1664944
0.50	0.0481424	0.0472270	-1.9384242
0.45	0.0139338	0.0129922	-7.2470014
0.40	-0.0285451	-0.0292848	2.5260817
0.35	-0.0714584	-0.0719673	0.7071402
0.30	-0.1186585	-0.1186277	-0.0259635
0.25	-0.1640687	-0.1633397	-0.4463091
0.20	-0.2022790	-0.2006958	-0.7888555
0.15	-0.2229911	-0.2206229	-1.0734153
0.10	-0.2130995	-0.2103346	-1.3145245
0.05	-0.1350920	-0.1333185	-1.3302729
0.00	0.0000000	0.0000000	0.0000000

Table 31- Comparison of Non-dimensional values of Vertical Velocity along the Horizontal Centerline between First Order and Second Order Time Accurate Approximations ($Ra = 10^4$, $\Delta \bar{t} = 1.0E-6$ and 21×21)

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Vertical Velocity, \bar{v} (1 st Order)	Non-dimensional Vertical Velocity, \bar{v} (2nd Order)	% Deviation in \bar{v}
1.00	0.0000000	0.0000000	0.0000000
0.95	-0.1774396	-0.1748586	-1.4760497
0.90	-0.2428447	-0.2397430	-1.2937602
0.85	-0.2299650	-0.2278353	-0.9347541
0.80	-0.1823231	-0.1809458	-0.7611671
0.75	-0.1288377	-0.1284063	-0.3359648
0.70	-0.0824661	-0.0822997	-0.2021511
0.65	-0.0464804	-0.0467008	0.4720043
0.60	-0.0195359	-0.0196341	0.5004028
0.55	0.0016007	0.0013556	-18.078290
0.50	0.0195719	0.0195468	-0.1283070
0.45	0.0374285	0.0372685	-0.4293166
0.40	0.0566598	0.0566889	0.0512798
0.35	0.0796021	0.0793940	-0.2621734
0.30	0.1061552	0.1060370	-0.1114705
0.25	0.1366630	0.1361237	-0.3961837
0.20	0.1663514	0.1657783	-0.3457026
0.15	0.1886132	0.1874899	-0.5991256
0.10	0.1873610	0.1863276	-0.5546145
0.05	0.1414387	0.1402807	-0.8254877
0.00	0.0000000	0.0000000	0.0000000

Table 32- Comparison of Non-dimensional values of Temperature along the Horizontal Centerline between First Order and Second Order Time Accurate Approximations ($Ra = 10^4$, $\Delta \bar{t} = 1.0E-6$ and 21×21)

Non-dimensional Horizontal Distance, \bar{x}	Non-dimensional Temperature, \bar{T} (1st Order)	Non-dimensional Temperature, \bar{T} (2nd Order)	% Deviation in \bar{T}
1.00	0.9891927	0.9891927	0.0000000
0.95	0.9925951	0.9928062	0.0212629
0.90	0.9959832	0.9961624	0.0179890
0.85	0.9989185	0.9990303	0.0111908
0.80	1.0011520	1.0012220	0.0069914
0.75	1.0026770	1.0027260	0.0048866
0.70	1.0036120	1.0036520	0.0039854
0.65	1.0041200	1.0041590	0.0038838
0.60	1.0043510	1.0043920	0.0040820
0.55	1.0044200	1.0044650	0.0044799
0.50	1.0044090	1.0044590	0.0049778
0.45	1.0043770	1.0044310	0.0053761
0.40	1.0043710	1.0044290	0.0057744
0.35	1.0044400	1.0045010	0.0060726
0.30	1.0046400	1.0047040	0.0063700
0.25	1.0050360	1.0051000	0.0063675
0.20	1.0056900	1.0057500	0.0059656
0.15	1.0066410	1.0066930	0.0051654
0.10	1.0078790	1.0079160	0.0036709
0.05	1.0093140	1.0093330	0.0018824
0.00	1.0108070	1.0108070	0.0000000

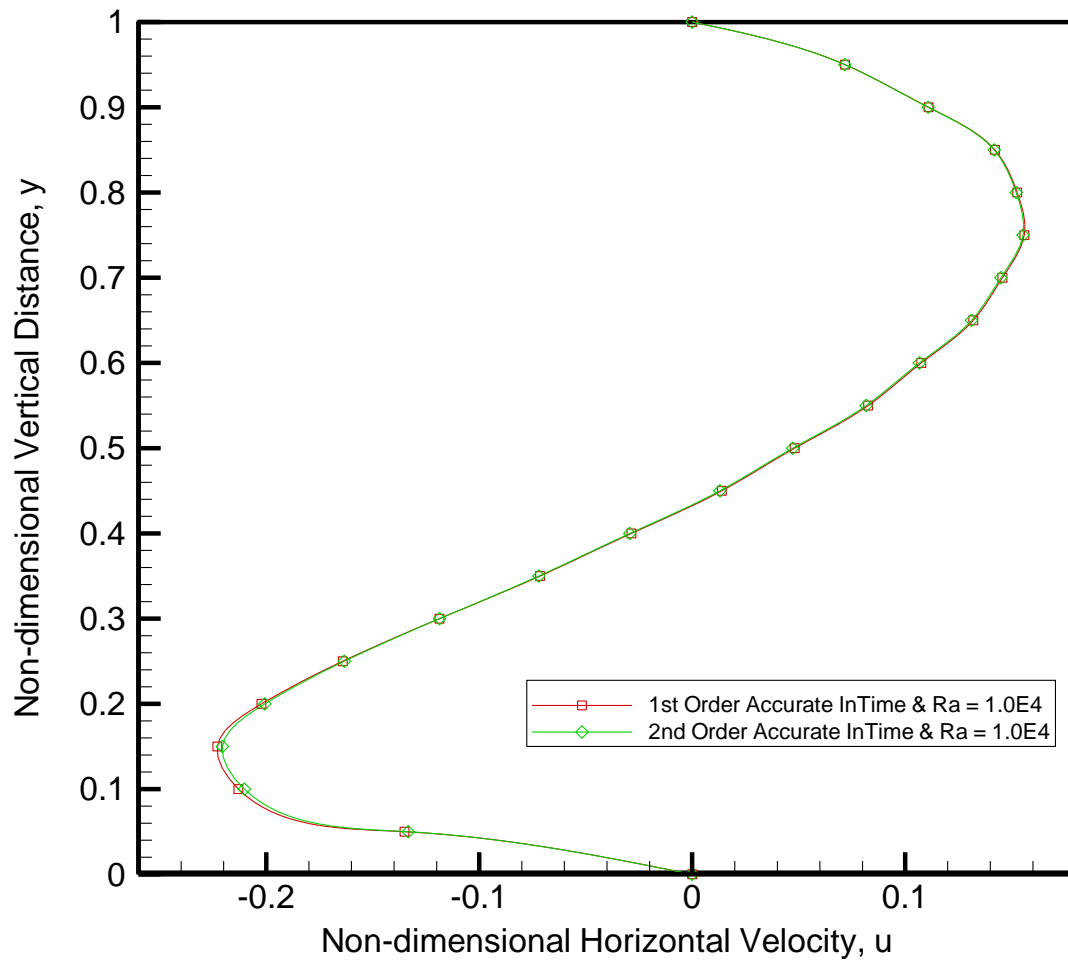


Figure 55- Comparison of Steady State Horizontal Velocity Distribution along the Vertical Centerline of the Enclosure between First Order and Second Order Time Accurate Approximation for $Ra = 10^4$

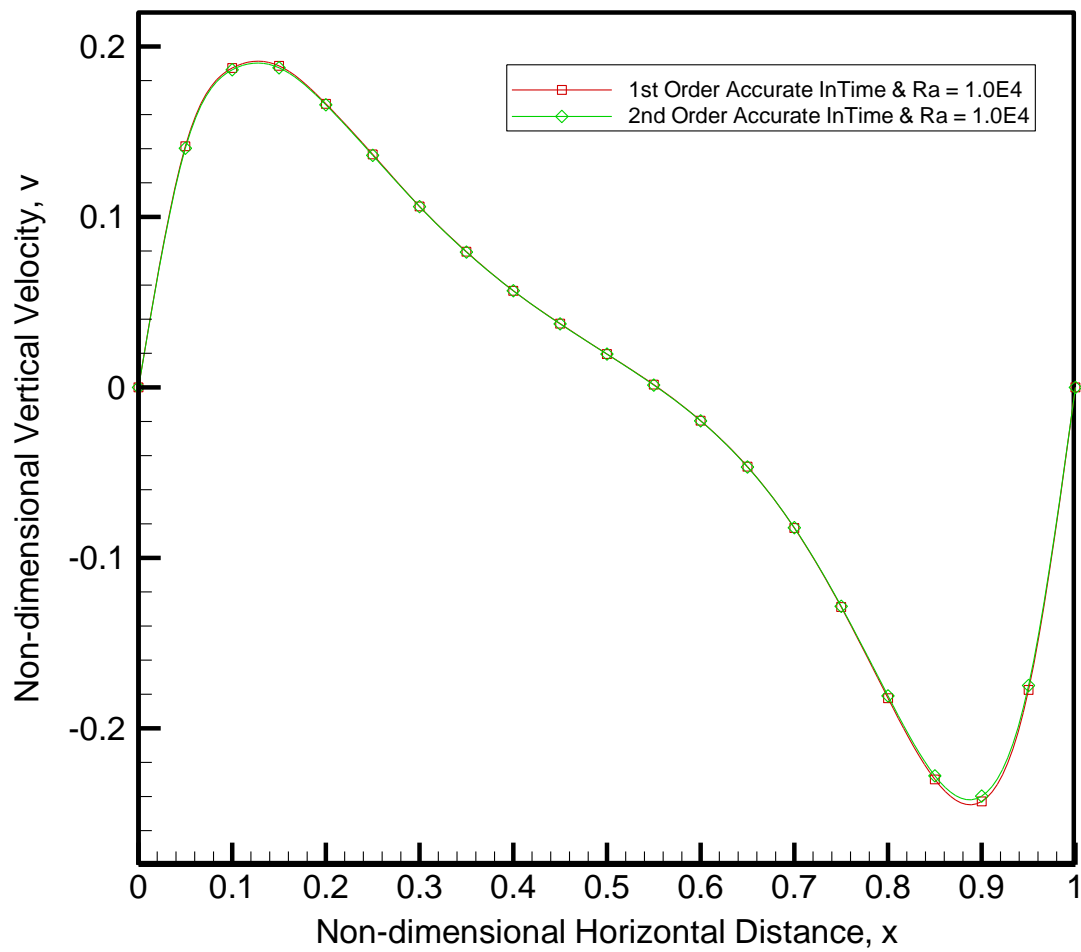


Figure 56- Comparison of Steady State Vertical Velocity Distribution along the Horizontal Centerline of the Enclosure between First Order and Second Order Time Accurate Approximation for $Ra= 10^4$

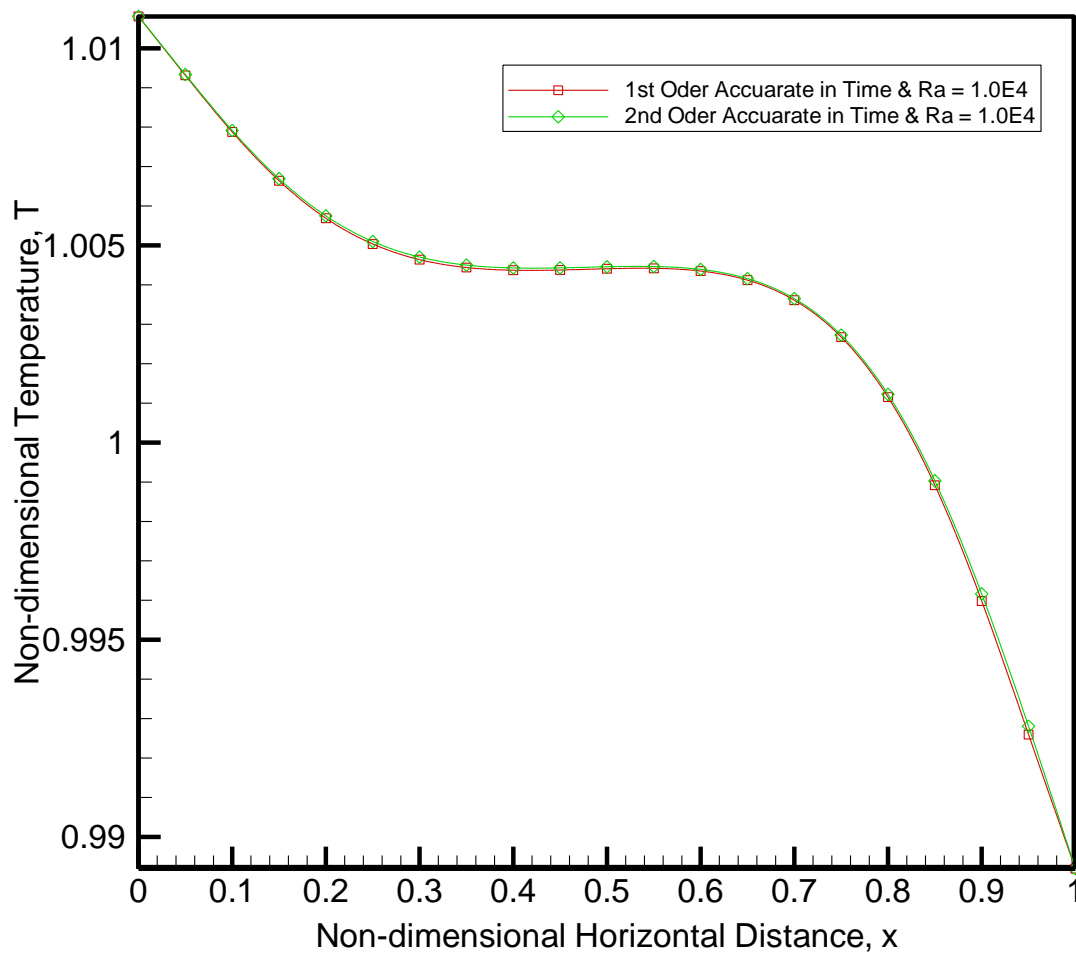


Figure 57- Comparison of Steady State Temperature Distribution along the Horizontal Centerline of the Enclosure between First Order and Second Order Time Accurate Approximation for $Ra = 10^4$

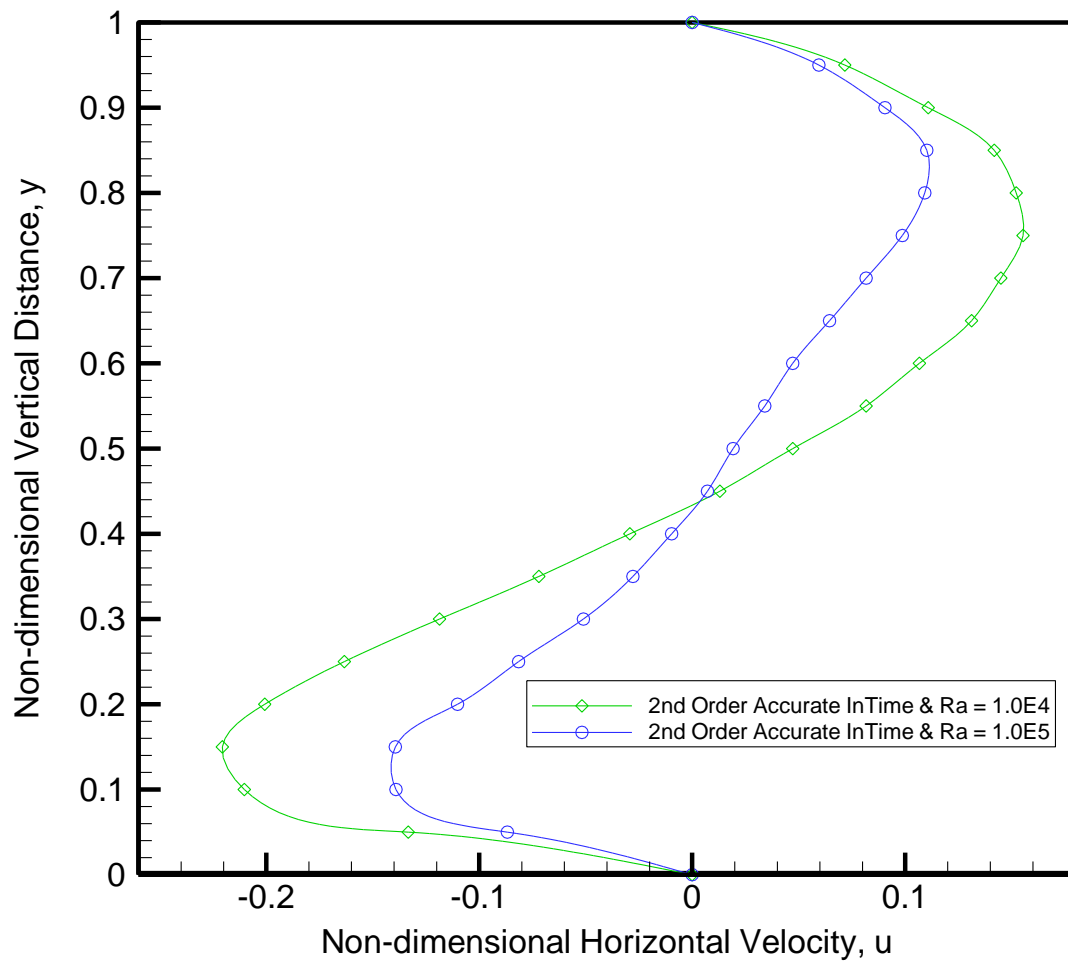


Figure 58- Comparison of Steady State Horizontal Velocity Distribution along the Vertical Centerline of the Enclosure between $Ra= 10^4$ and $Ra= 10^5$ in Second Order Time Accurate Approximation

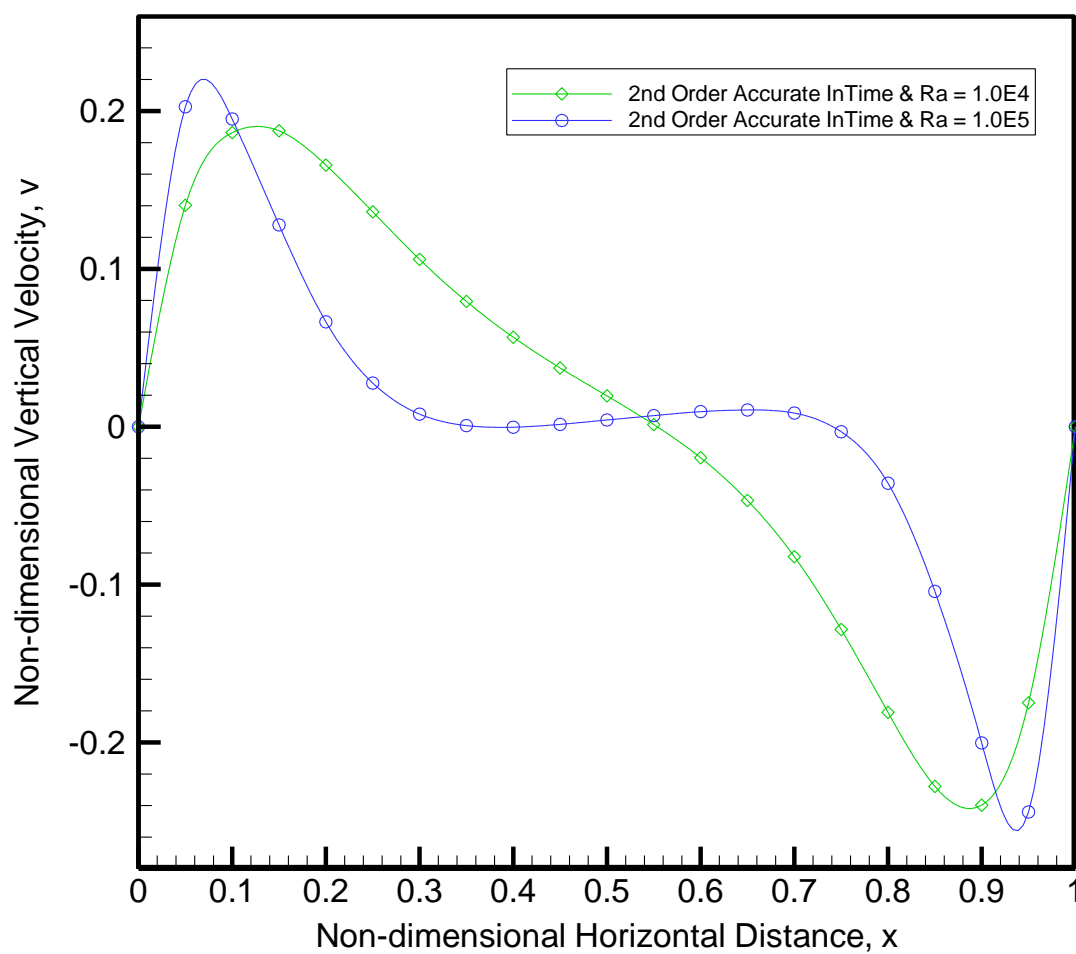


Figure 59- Comparison of Steady State Vertical Velocity Distribution along the Horizontal Centerline of the Enclosure between $Ra= 10^4$ and $Ra= 10^5$ in Second Order Time Accurate Approximation

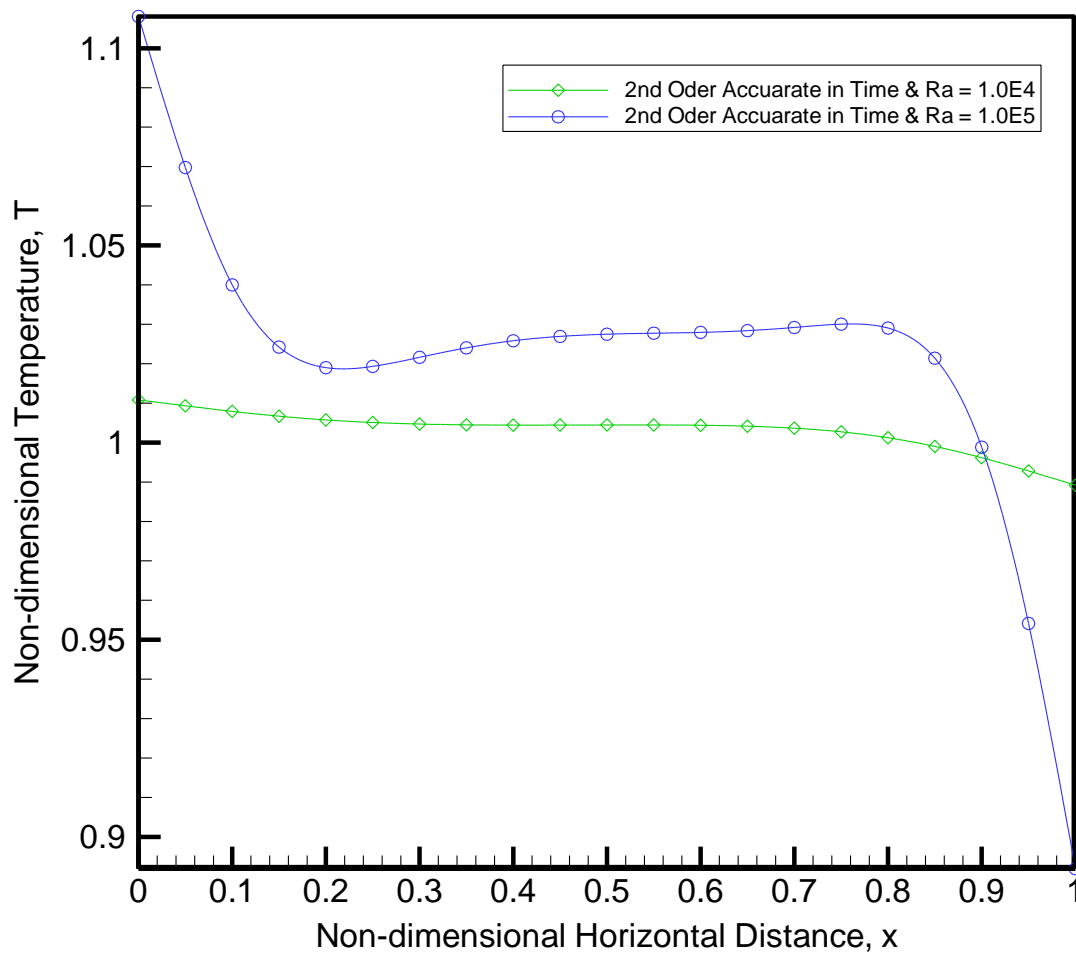


Figure 60- Comparison of Steady State Temperature Distribution along the Horizontal Centerline of the Enclosure between $Ra = 10^4$ and $Ra = 10^5$ in Second Order Time Accurate Approximation

Chapter 11

The Study of Unsteady Natural Convection in Square Enclosure

The unsteady thermal and hydrodynamic behavior of the working fluid was studied by imposing a sudden wall temperature change in a square enclosure. The unsteady development of the velocity vectors and temperatures for $Ra = 10^5$ and grid size of 41×41 for different non-dimensional time levels (1 s, 5 s, 10 s, 15 s, 20 s, and 50 s) using second order accurate in time model are presented in this chapter. The density of the air near the hot wall is lower than that of the cold wall. As a result, the hot air rises near the left hot wall and falls near the cold right wall ensuring in a clockwise circular motion. These results are presented in Figures 62 through 67 in terms of velocity vectors. The flow is not unicellular as evident that the formation of two secondary vertices. The temperature contours are presented in Figures 68 through 73. These plots show the development of the expected temperature stratification at vertical centerline of the square enclosure. The isotherms are parallel to the side walls of the enclosure depicting that conduction is the major mode of heat transfer before imposing a sudden wall temperature change. They compare favorable to ones given by Kublbeck [36]. A histogram of non-dimensional temperature at location $\bar{x} = 0.05$ and $\bar{y} = 0.5$ is given in Figure 61.

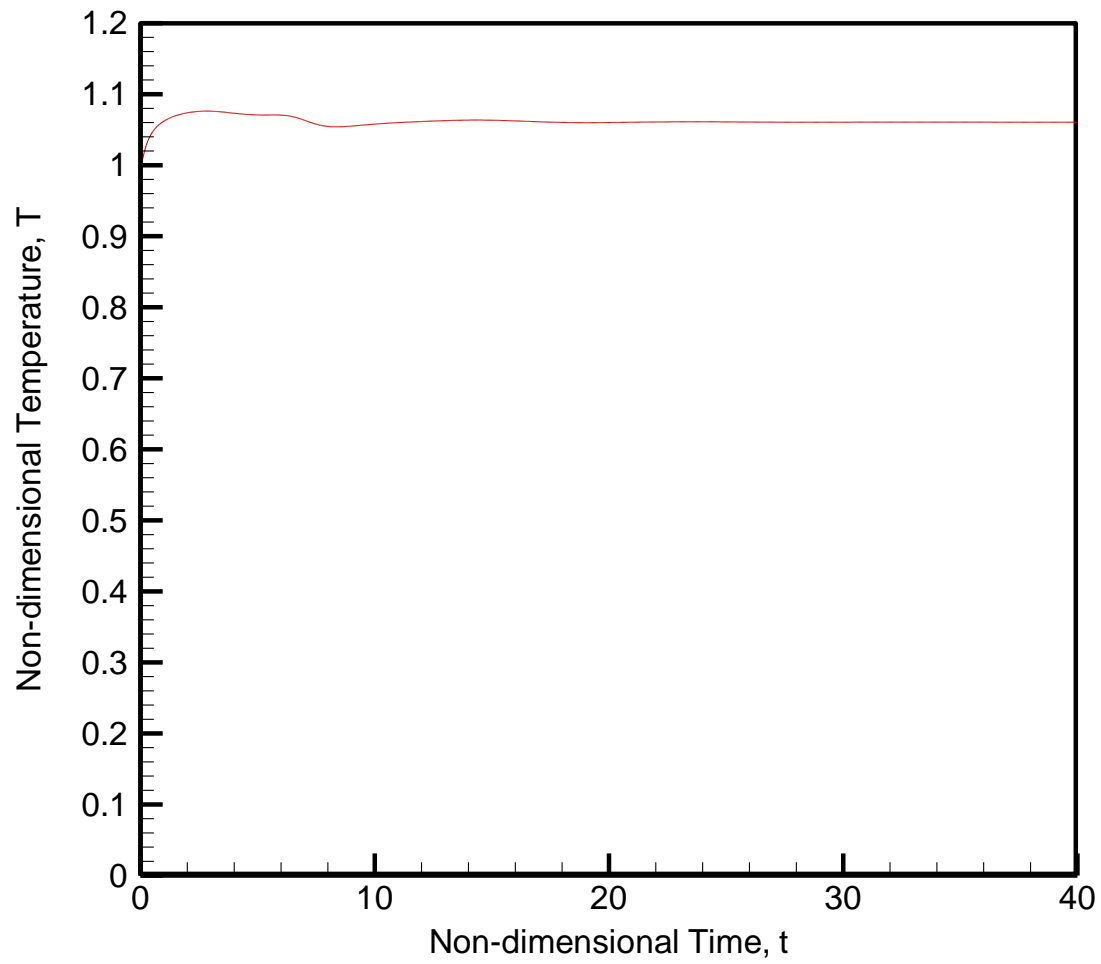


Figure 61- Non-dimensional Temperature Histogram for $Ra = 10^5$ at $\bar{x} = 0.05$ and $\bar{y} = 0.5$
(41×41 , $\Delta \bar{t} = 5.0E-6$) in Second Order Time Accurate Approximation

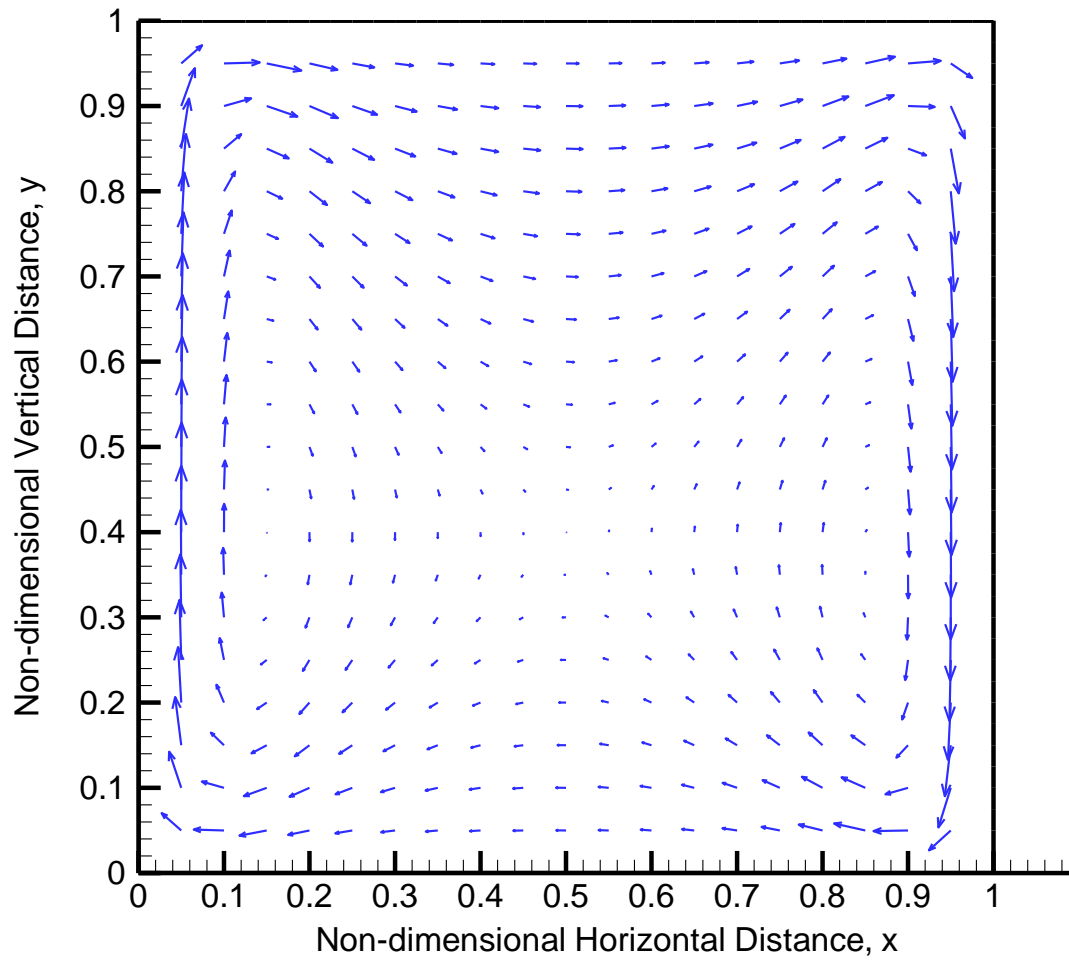


Figure 62- Unsteady State Non-dimensional Velocity Vectors for $Ra=10^5$ at Non-dimensional Time, $\bar{t} = 1$ (41×41 , $\Delta \bar{t} = 5.0E-6$) in Second Order Time Accurate Approximation

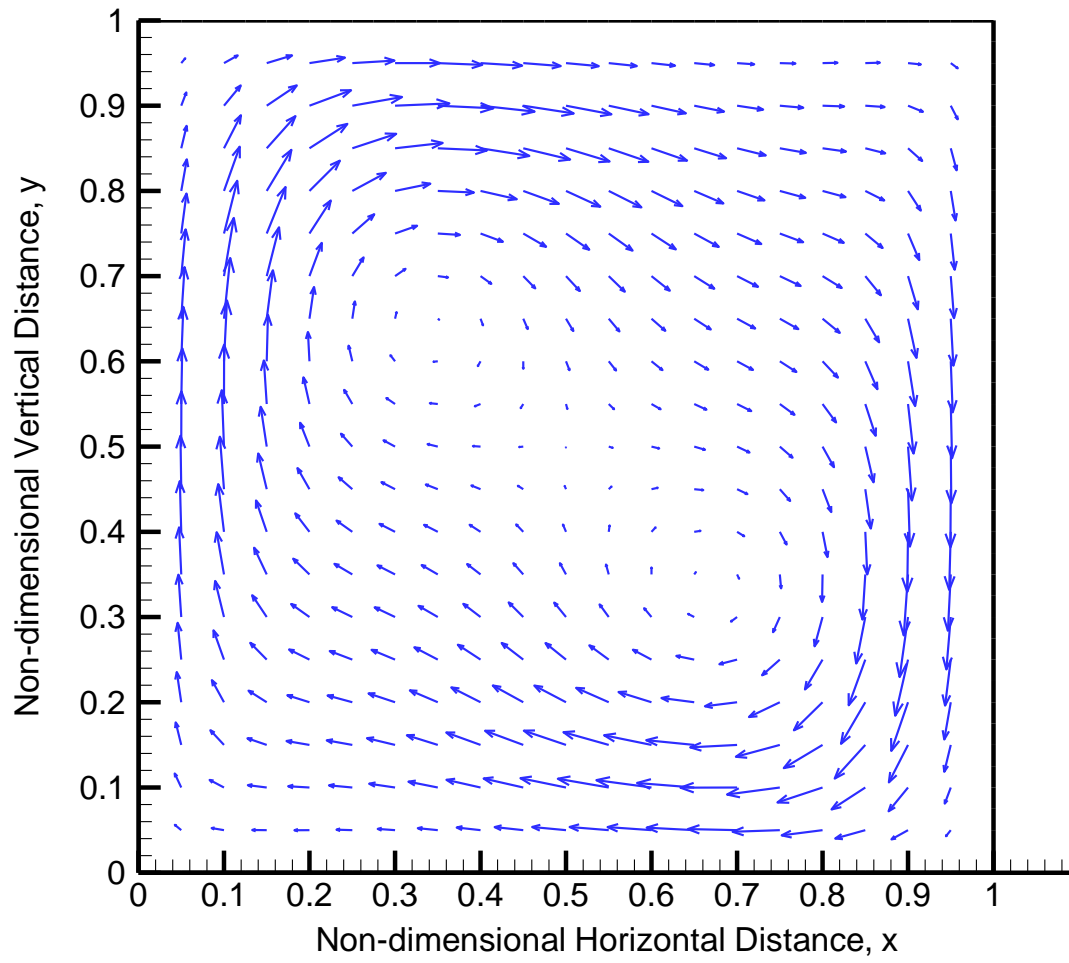


Figure 63- Unsteady State Non-dimensional Velocity Vectors for $Ra=10^5$ at Non-dimensional Time, $\bar{t} = 5$ (41×41 , $\Delta\bar{t} = 5.0E-6$) in Second Order Time Accurate Approximation

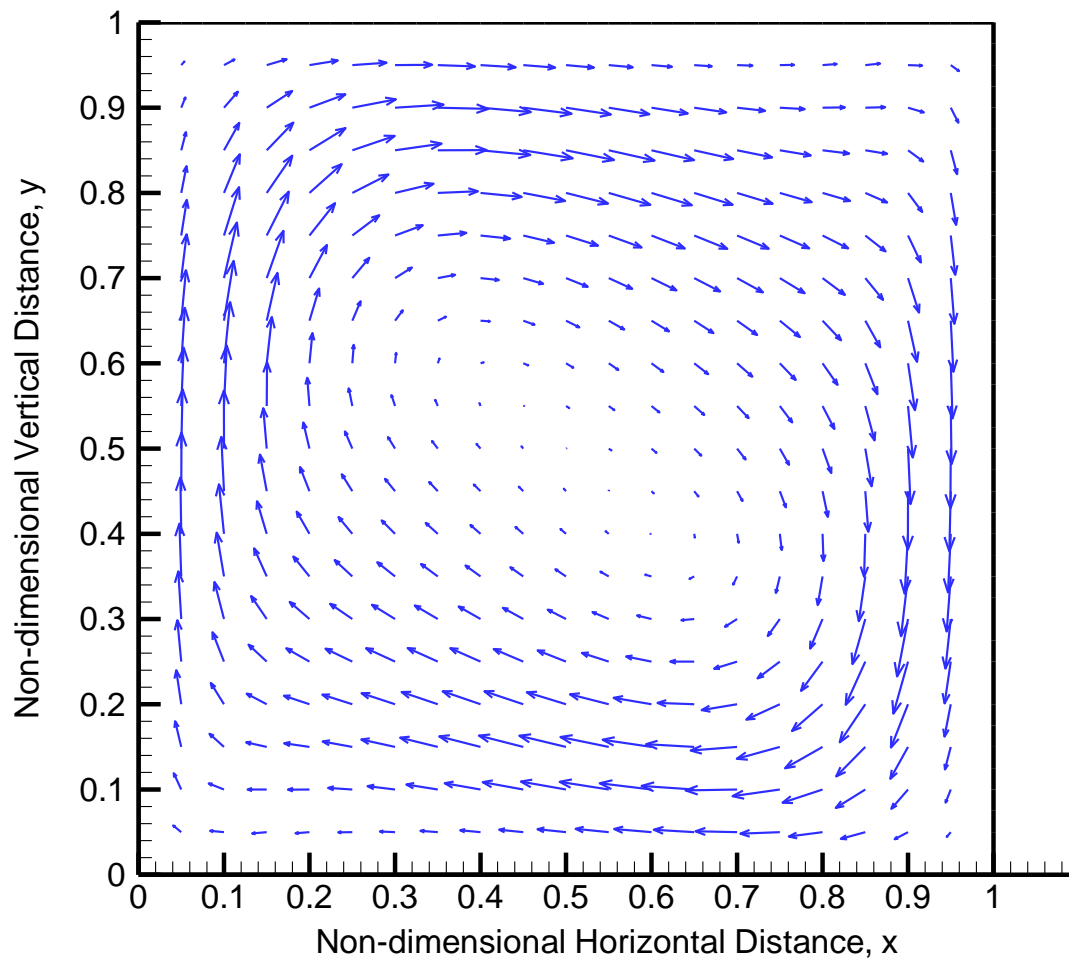


Figure 64- Unsteady State Non-dimensional Velocity Vectors for $Ra=10^5$ at Non-dimensional Time, $\bar{t} = 10$ (41×41 , $\Delta\bar{t} = 5.0E-6$) in Second Order Time Accurate Approximation

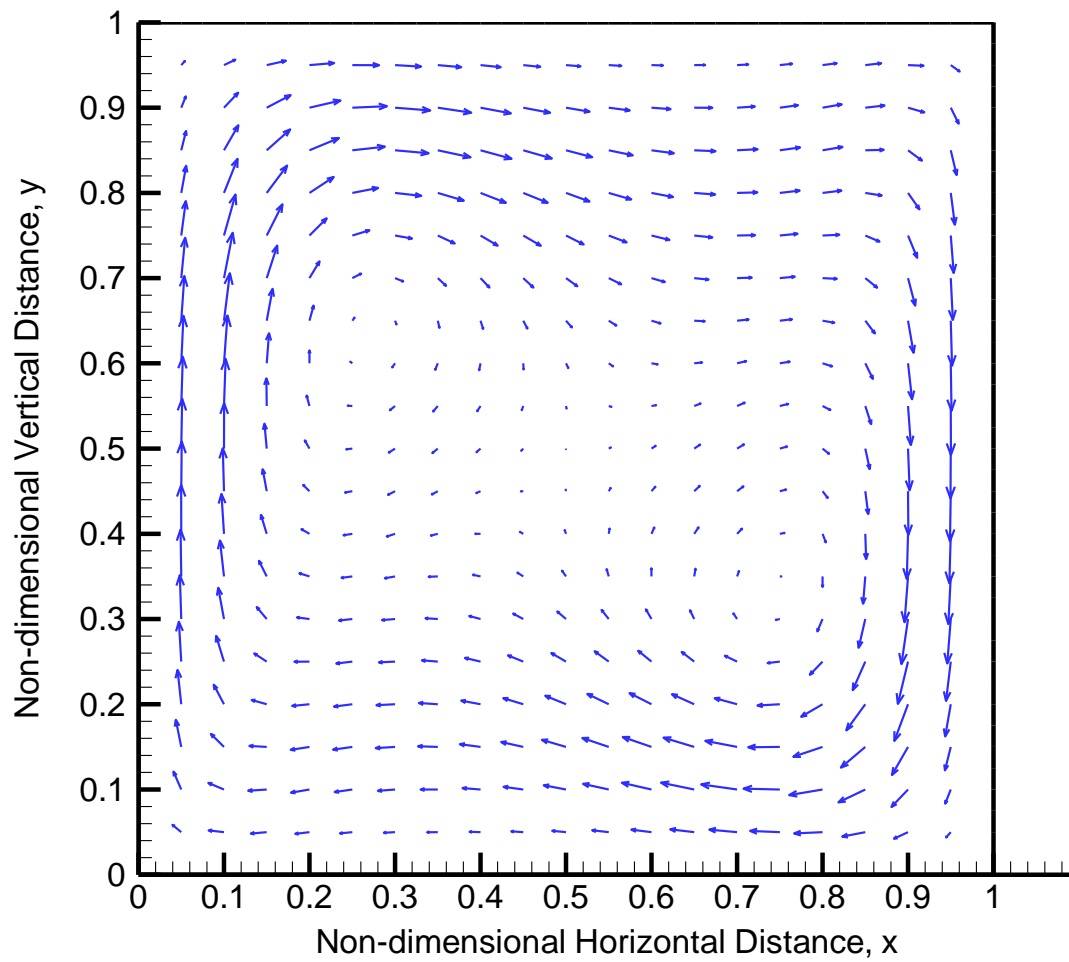


Figure 65- Unsteady State Non-dimensional Velocity Vectors for $Ra=10^5$ at Non-dimensional Time, $\bar{t} = 15$ (41×41 , $\Delta\bar{t} = 5.0E-6$) in Second Order Time Accurate Approximation

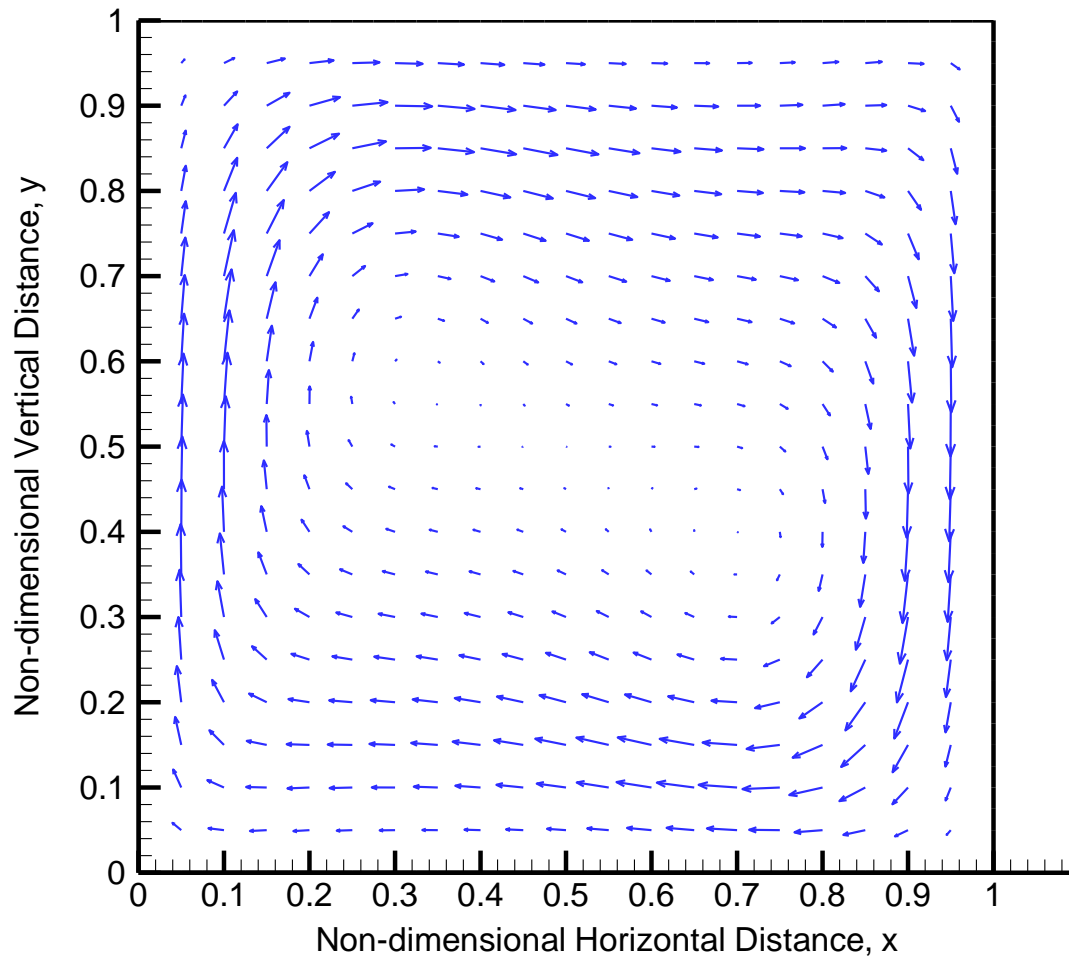


Figure 66- Unsteady State Non-dimensional Velocity Vectors for $Ra=10^5$ at Non-dimensional Time, $\bar{t} = 20$ (41×41 , $\Delta\bar{t} = 5.0E-6$) in Second Order Time Accurate Approximation

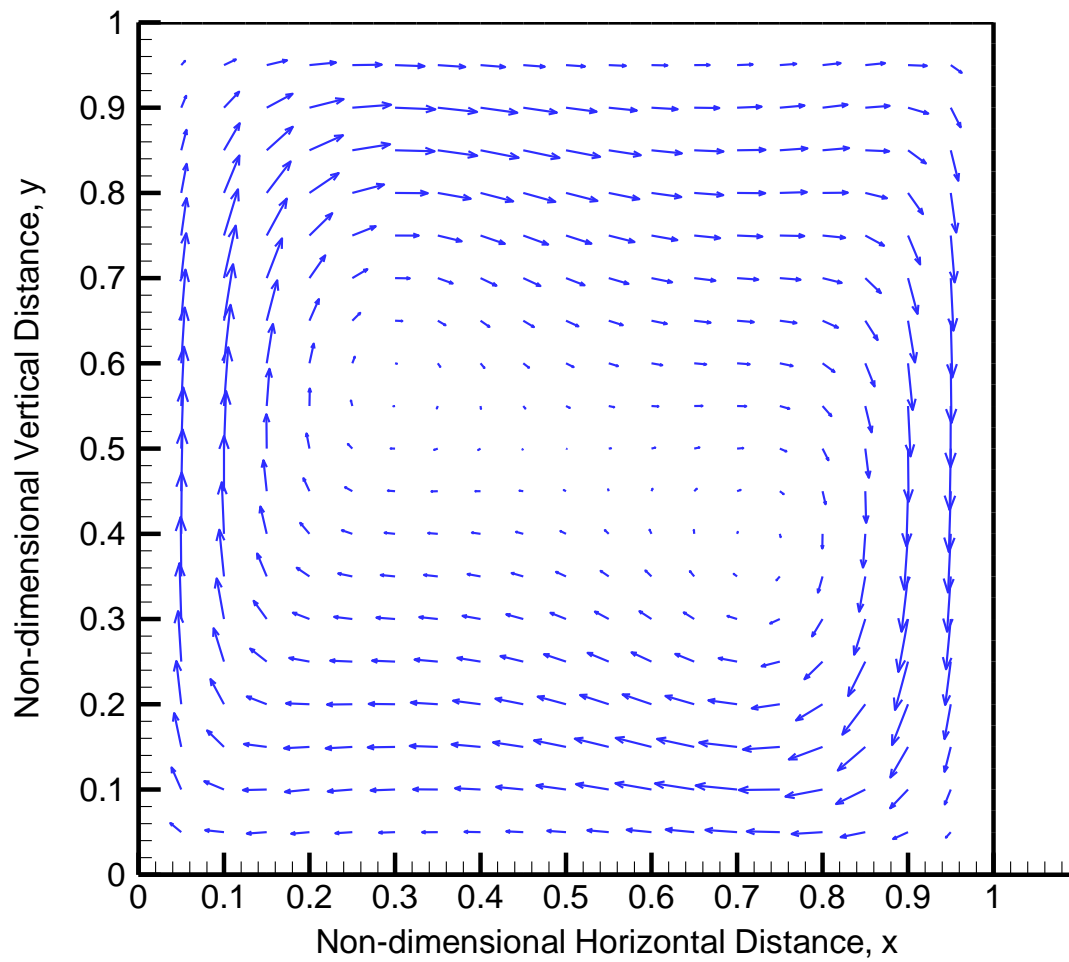


Figure 67- Unsteady State Non-dimensional Velocity Vectors for $Ra=10^5$ at Non-dimensional Time, $\bar{t} = 50$ (41×41 , $\Delta\bar{t} = 5.0E-6$) in Second Order Time Accurate Approximation

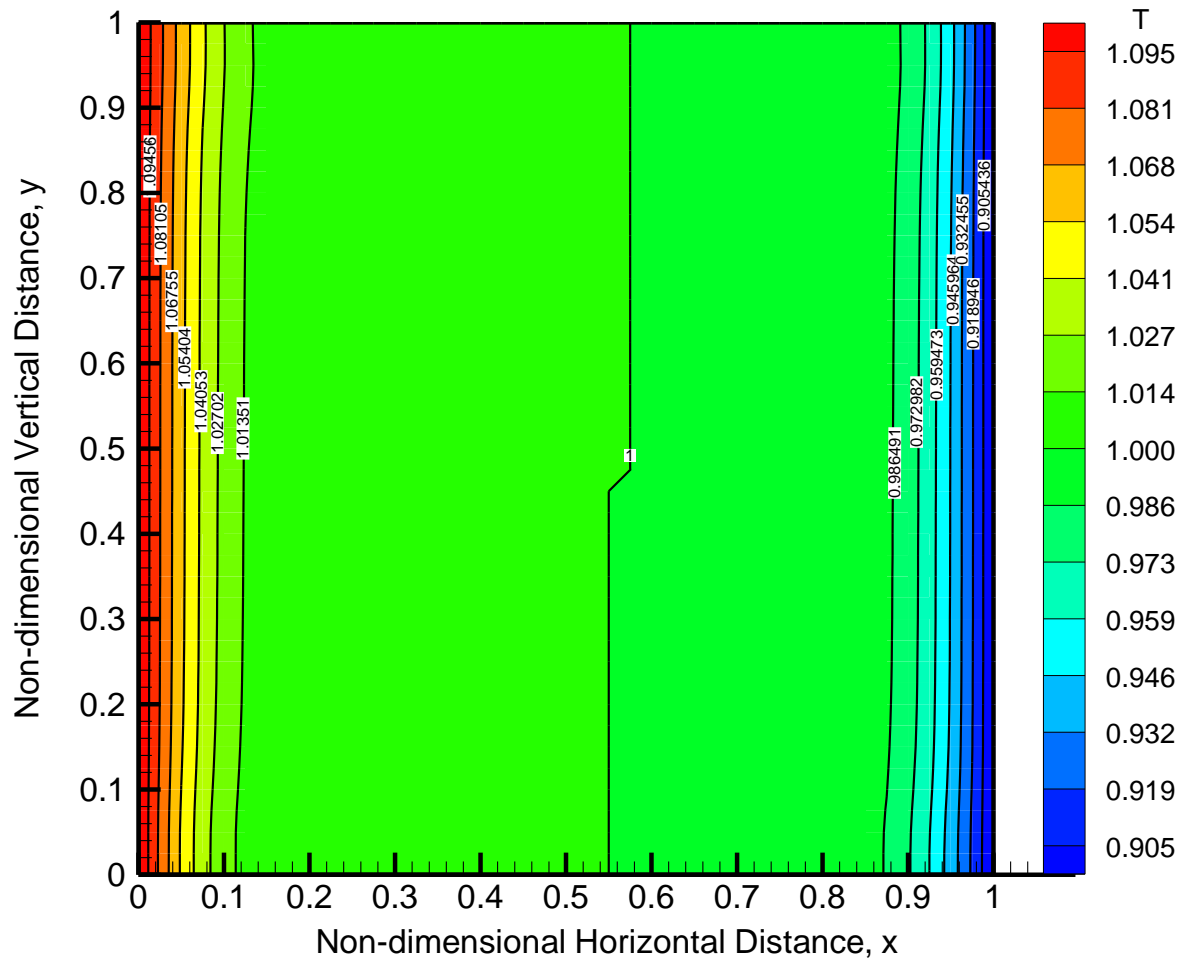


Figure 68- Unsteady State Non-dimensional Temperature Contours for $Ra=10^5$ at Non-dimensional Time, $\bar{t} = 1$ (41×41 , $\Delta\bar{t} = 5.0E-6$) in Second Order Time Accurate Approximation

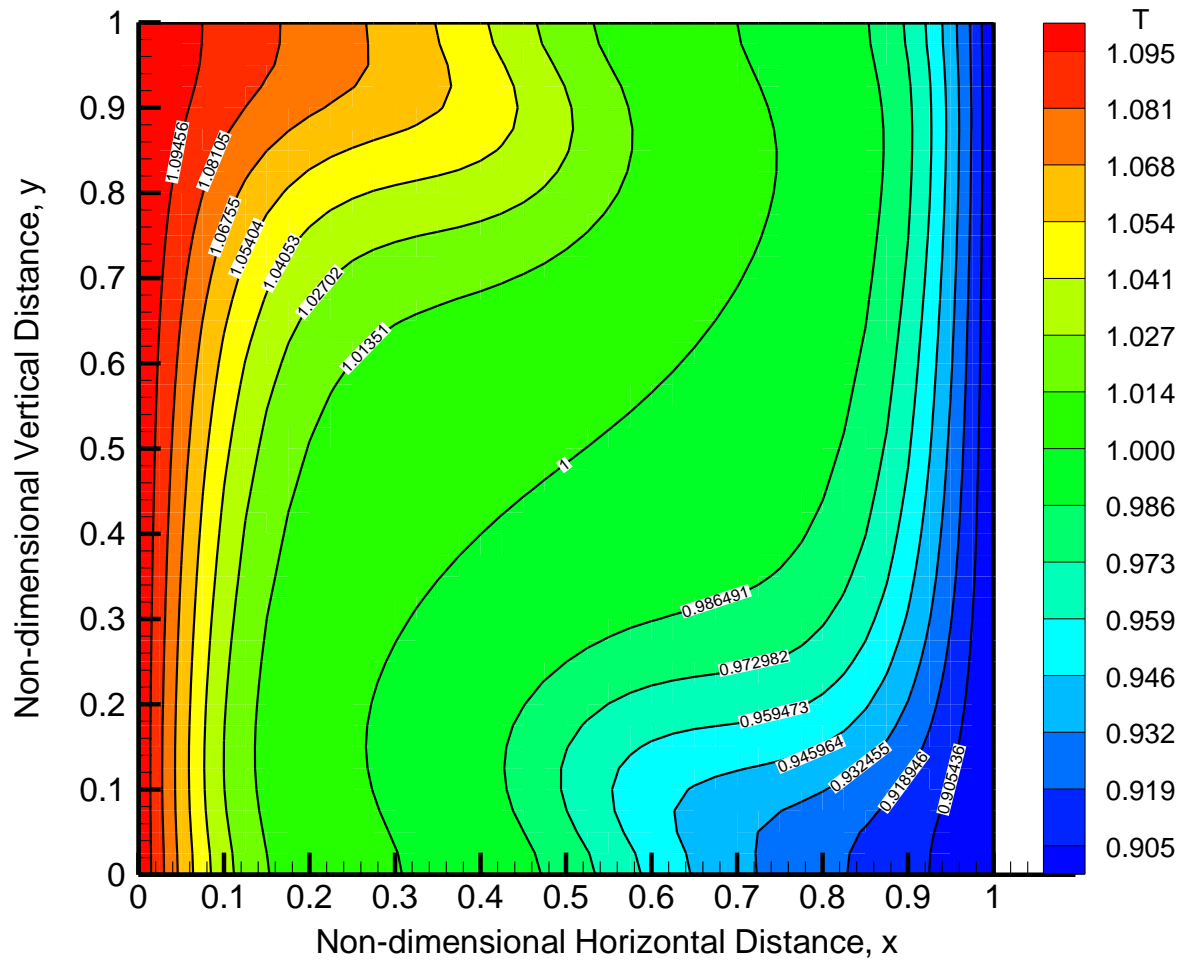


Figure 69- Unsteady State Non-dimensional Temperature Contours for $Ra=10^5$ at Non-dimensional Time, $\bar{t} = 5$ (41×41 , $\Delta \bar{t} = 5.0E-6$) in Second Order Time Accurate Approximation

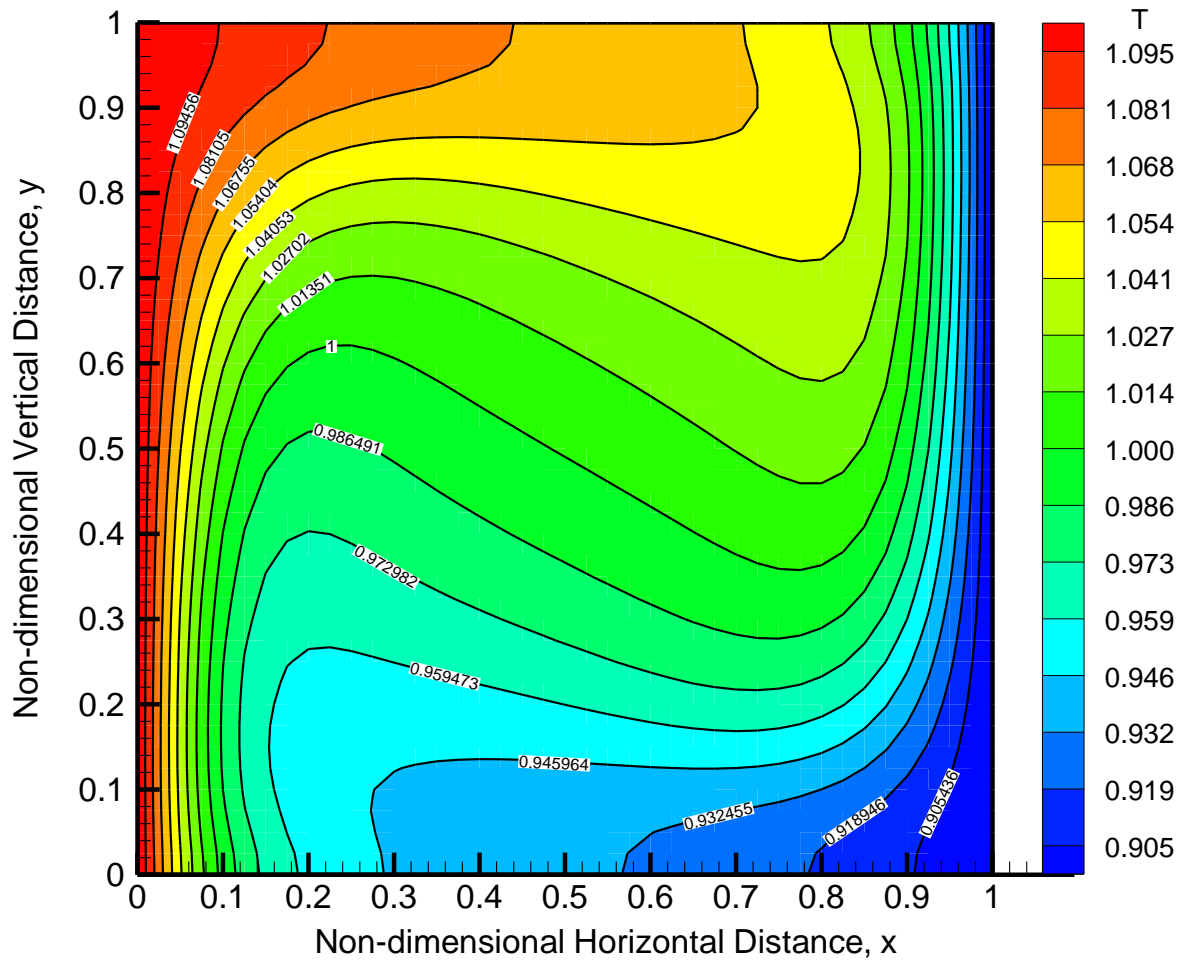


Figure 70- Unsteady State Non-dimensional Temperature Contours for $Ra=10^5$ at Non-dimensional Time, $\bar{t} = 10$ (41×41 , $\Delta \bar{t} = 5.0E-6$) in Second Order Time Accurate Approximation

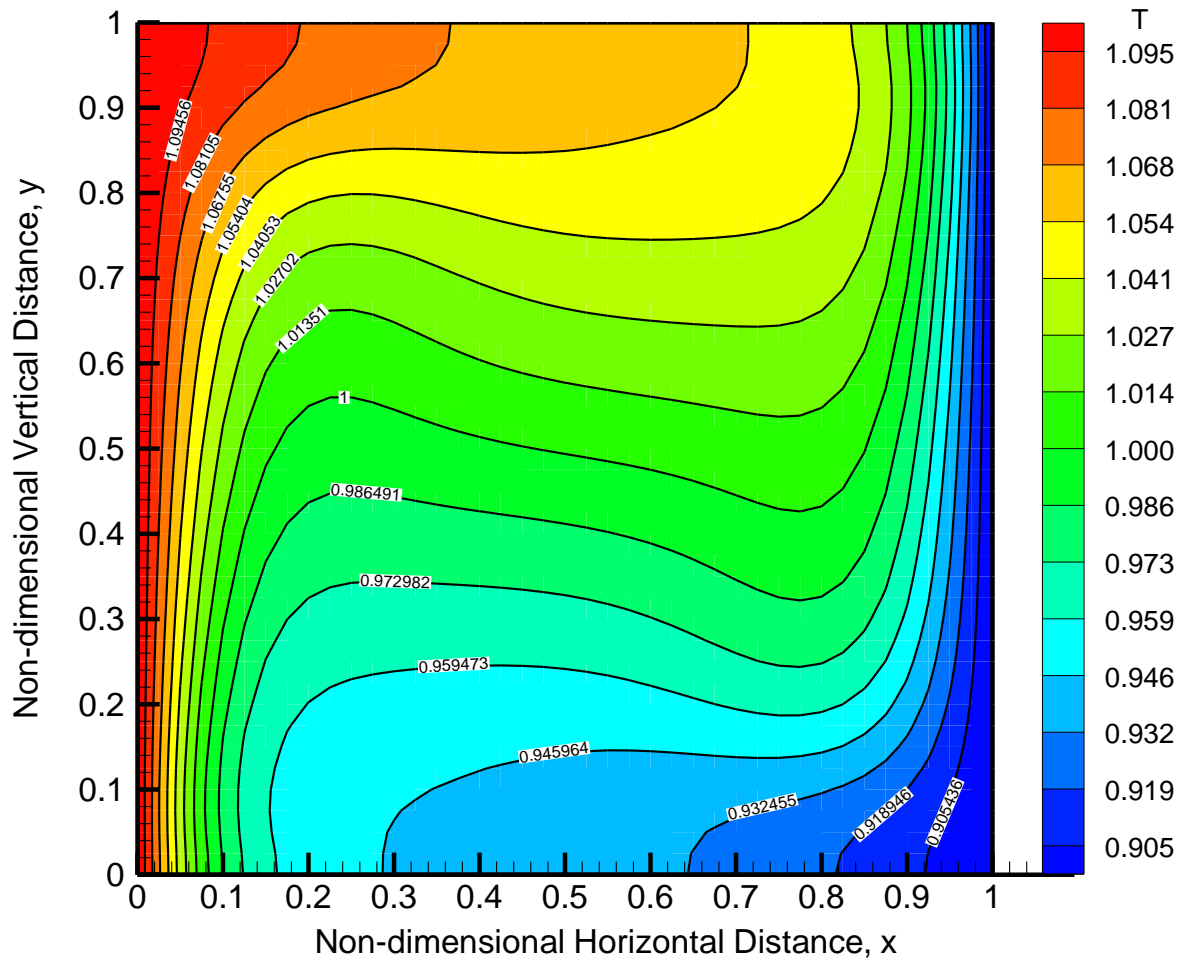


Figure 71- Unsteady State Non-dimensional Temperature Contours for $Ra=10^5$ at Non-dimensional Time, $\bar{t} = 15$ (41×41 , $\Delta\bar{t} = 5.0E-6$) in Second Order Time Accurate Approximation

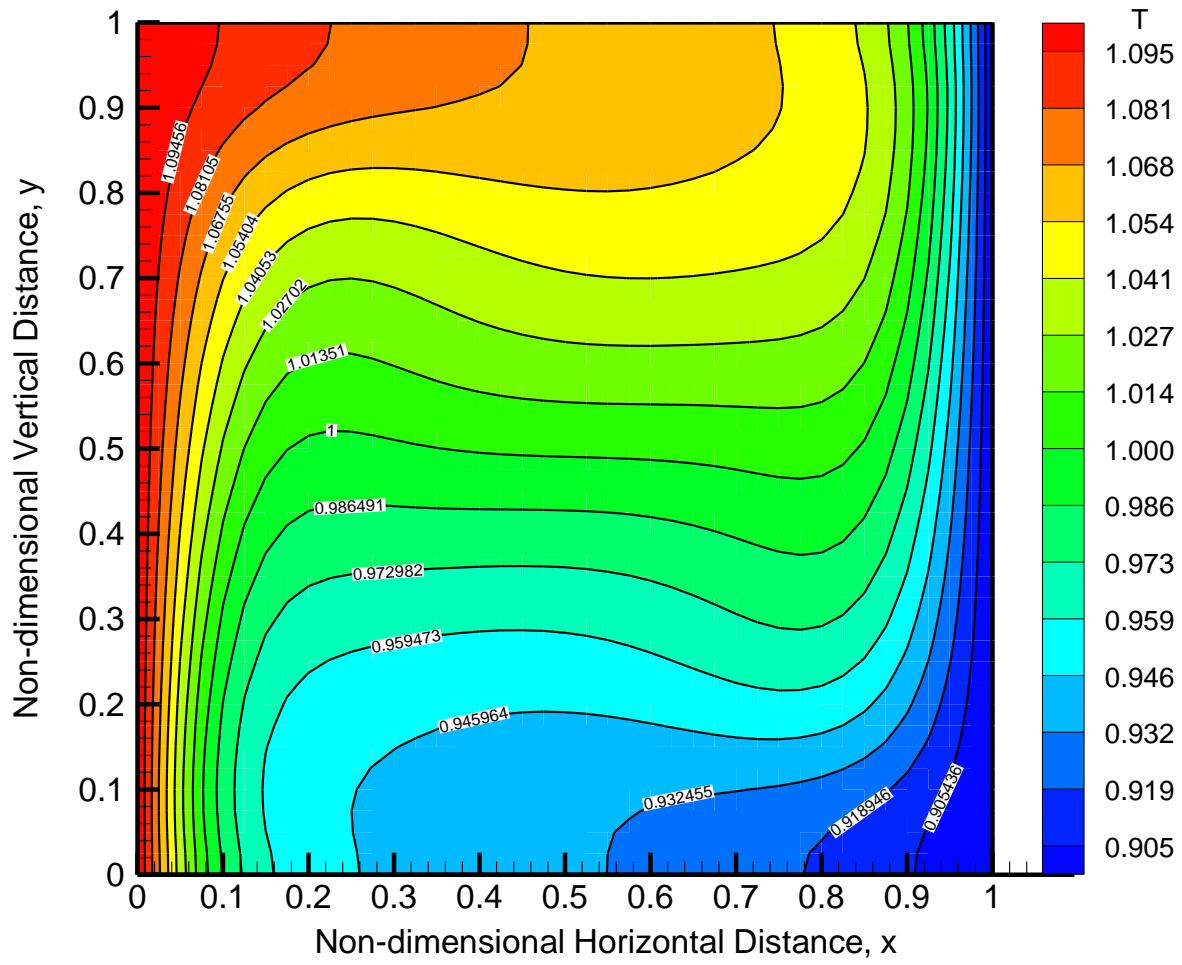


Figure 72- Unsteady State Non-dimensional Temperature Contours for $Ra=10^5$ at Non-dimensional Time, $\bar{t} = 20$ (41×41 , $\Delta\bar{t} = 5.0E-6$) in Second Order Time Accurate Approximation

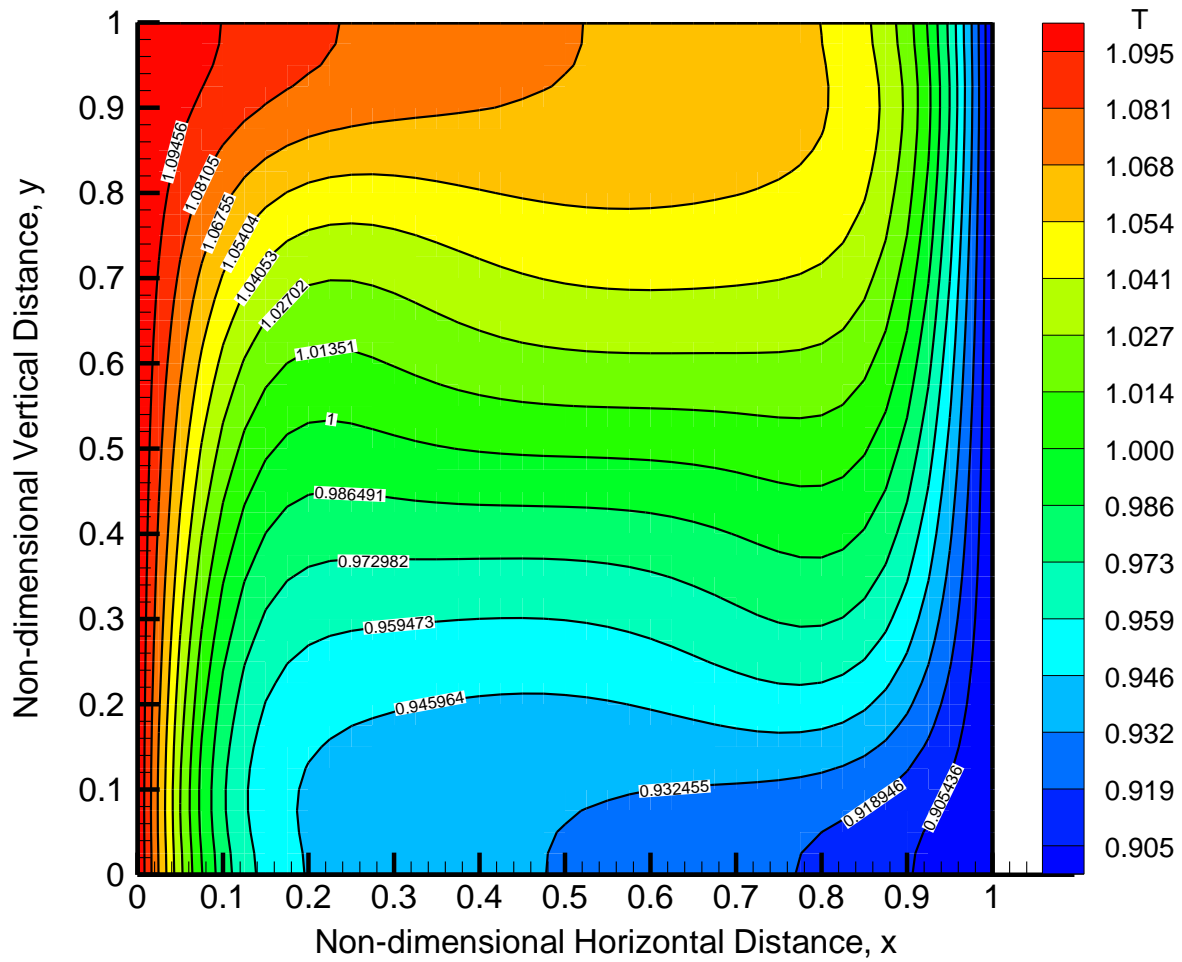


Figure 73- Unsteady State Non-dimensional Temperature Contours for $Ra=10^5$ at Non-dimensional Time, $\bar{t} = 50$ (41×41 , $\Delta\bar{t} = 5.0E-6$) in Second Order Time Accurate Approximation

Chapter 12

Calculations of Wall Heat Flux and Nusselt Number

The calculations of heat flux at the walls of the enclosure (the high temperature wall in this case) and Nusselt number are presented in this chapter.

The rate of heat transfer is observed to be proportional to the temperature difference and is expressed by Newton's law of cooling as

$$\dot{q}_{conv} = h(T_w - T_\infty) \quad (9.1)$$

where h is the heat transfer coefficient in $\text{W/m}^2\text{-K}$, T_w is temperature at hot wall side of the enclosure, and T_∞ is the reference temperature.

The heat transfer from the solid surface to the fluid adjacent to the surface is by pure conduction can be expressed as

$$\dot{q}_{cond} = -k \left(\frac{dT}{dx} \right)_{x=0} \quad (9.2)$$

where T is the temperature distribution in the fluid, $\left(\frac{dT}{dx} \right)_{x=0}$ is the temperature gradient at the surface of hot wall, and k is thermal conductivity of the fluid in W/m-K

Conduction heat flux from the solid surface to the fluid is equal to the convection heat flux to the fluid adjacent to surface, at $x = 0$. Therefore, equating equations 9.1 and 9.2 and solving for the heat transfer coefficient yields,

$$h = \frac{-k \left(\frac{dT}{dx} \right)_{x=0}}{T_w - T_\infty} \quad (9.3)$$

Non-dimensional heat transfer coefficient (or) Nusselt number is defined as,

$$Nu = \frac{hH}{k} \quad (9.4)$$

The wall heat fluxes and Nu numbers for $\Delta T = 64.8$ deg K and $\Delta T = 356.6$ deg K at $Ra = 10^5$ were calculated from the temperature profiles (see Figures 74 through 77) near the high temperature wall as predicted by both models, the constant property and the variable property. The results of these calculations are presented in Table 33.

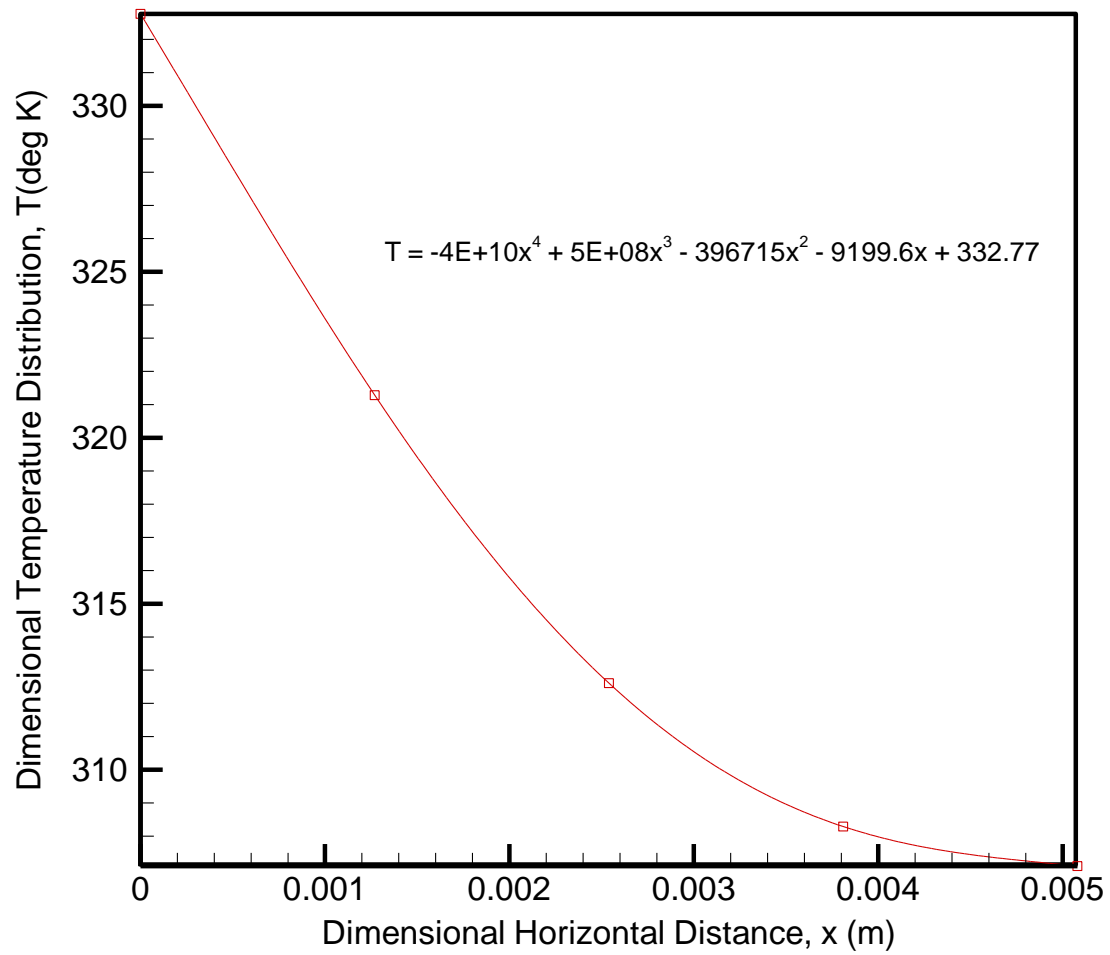


Figure 74- Dimensional Temperature Distribution along the Horizontal Distance for Constant Properties ($\Delta T = 64.8$ deg K)

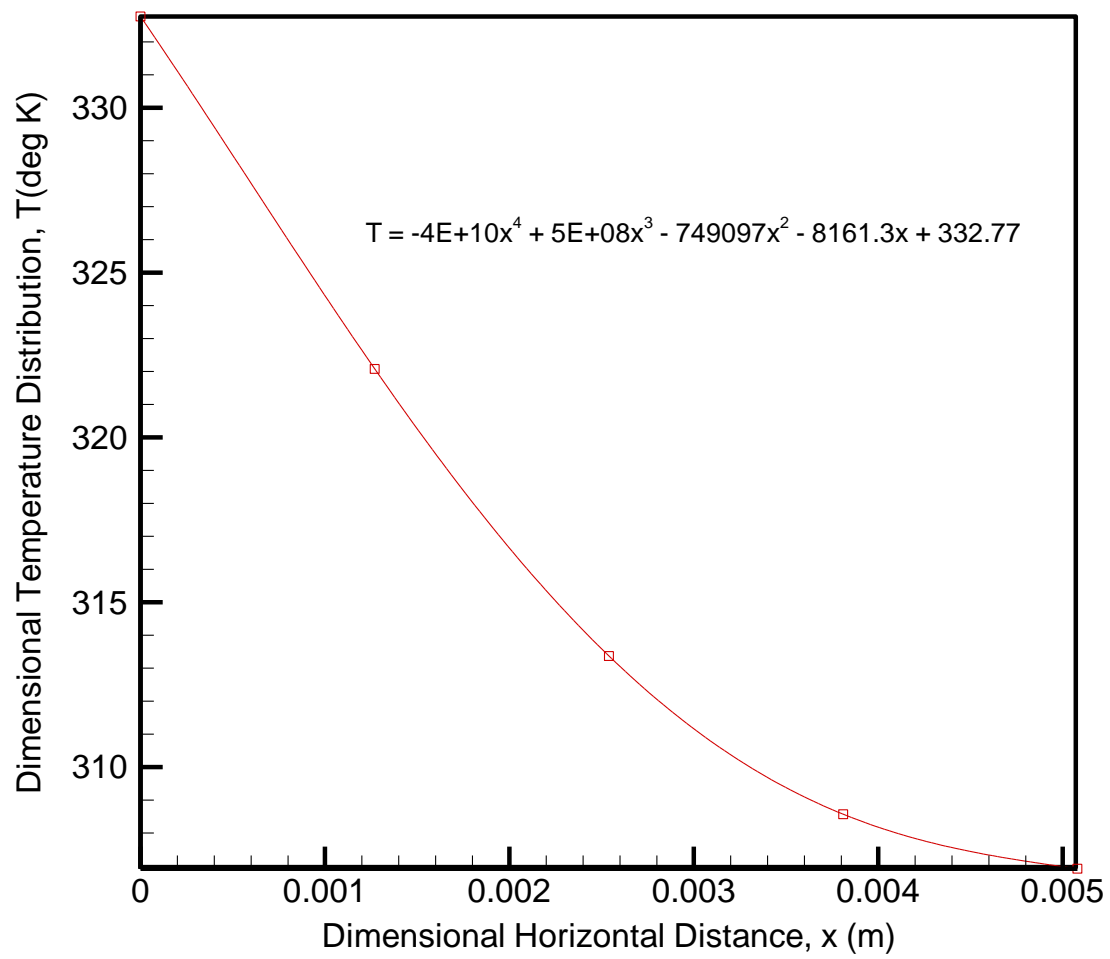


Figure 75- Dimensional Temperature Distribution along the Horizontal Distance for Variable Properties ($\Delta T = 64.8$ deg K)

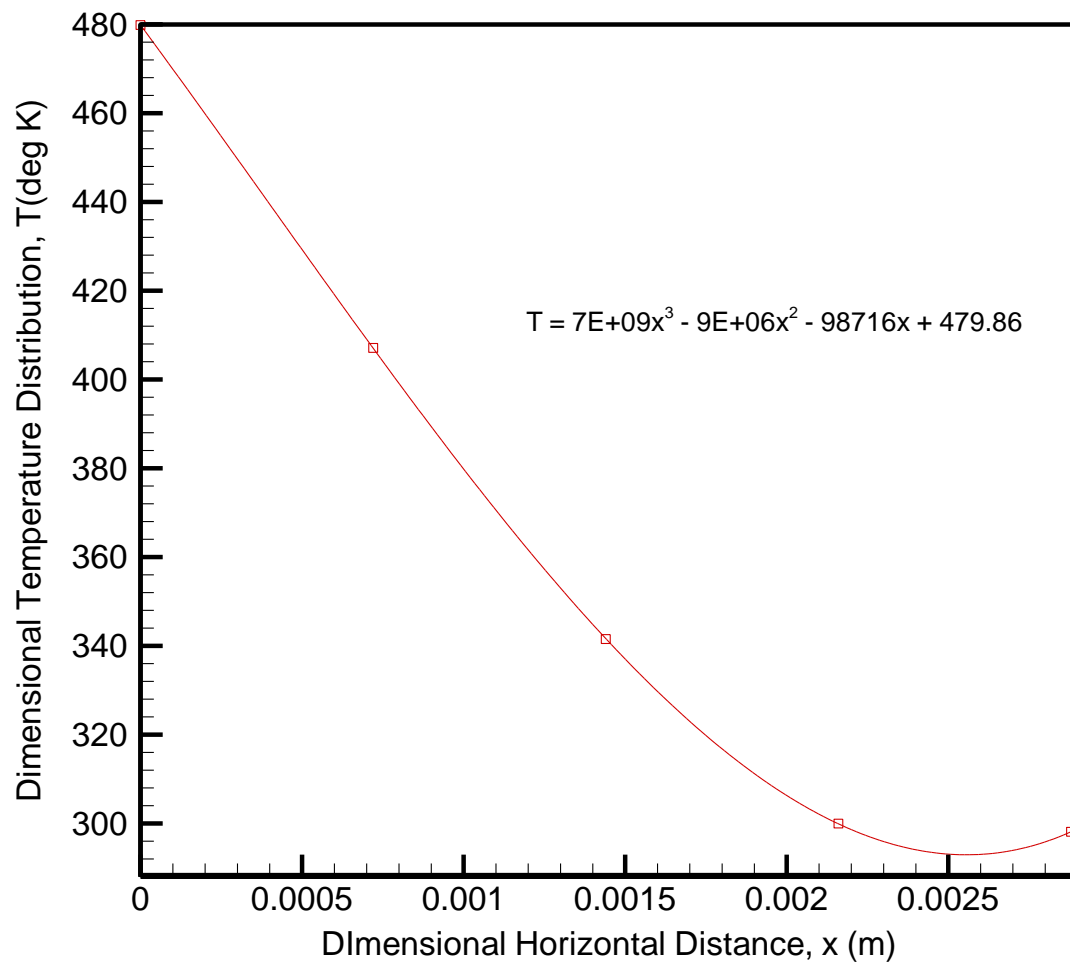


Figure 76- Dimensional Temperature Distribution along the Horizontal Distance for Constant Properties ($\Delta T = 356.6$ deg K)

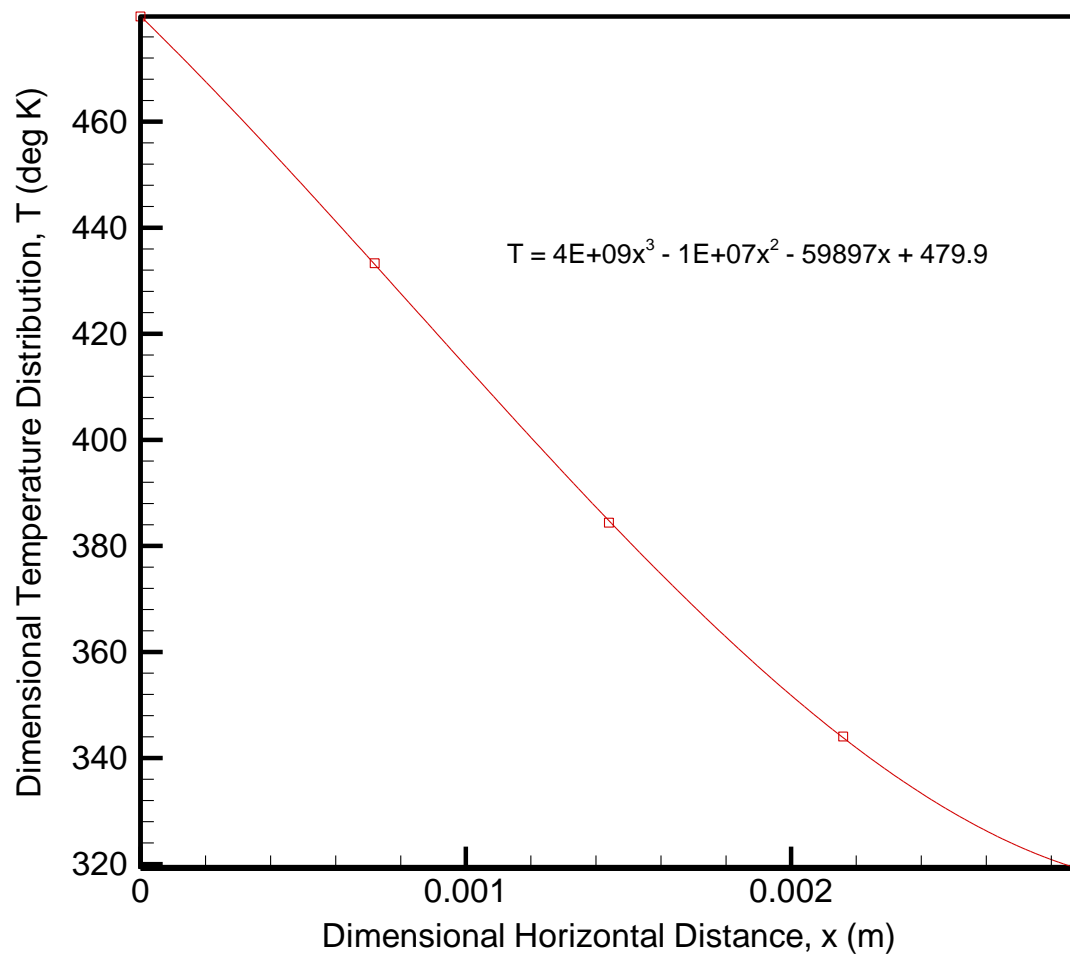


Figure 77- Dimensional Temperature Distribution along the Horizontal Distance for Variable Properties ($\Delta T = 356.6$ deg K)

Table 33 - Comparison of the Results of the Present Study (Variable properties) and Constant Properties for Different Wall Temperature differences at $Ra=10^5$

$\Delta T (^{\circ}K)$	Variable Properties $\dot{q}_w (W/m^2)$	Constant Properties $\dot{q}_w (W/m^2)$	% Deviation in \dot{q}_w
64.8	234.23	264.03	- 12.80
356.6	2180.00	3593.30	- 64.82

$\Delta T (^{\circ}K)$	Variable Properties Nu	Constant Properties Nu	% Deviation in Nu
64.8	6.40	7.21	- 12.65
356.6	4.83	7.97	- 65.02

As shown in the above table, the variable property model predicts the wall heat flux and the Nu number values lower than the constant property model. The wall heat flux predicted by the variable property model deviates by - 12.8 % from the values predicted by the variable property model for $\Delta T = 64.8^{\circ}K$. However, this difference is more for higher wall temperature differences. Similarly, the computed value of Nusselt number of the variable property model is 12.6% less than the corresponding value predicted by the constant property model, and this difference is high for higher wall temperature differences.

Chapter 13

Conclusions

A numerical study of unsteady natural convection inside a square enclosure has been carried out to determine the effect of variable properties, i.e., thermodynamic and transport properties, on circulation patterns, velocity profiles in vertical and horizontal directions, and temperature characteristics for a fluid with Pr number 0.72 and a Ra number ranging from 10^3 to 10^6 . The following conclusions are drawn from this study.

1. A newly updated Coupled Modified Strongly Implicit Procedure (CMSIP) was successfully employed to solve the governing equations of unsteady laminar natural convection inside an enclosure using a first order and a second order accurate in time finite difference approximations.
2. The results of the numerical study carried out using the mathematical model and solution procedure proposed in this thesis fair well when compared to the results of the benchmark case given in the literature.
3. There is a notable difference between the results of the variable property and the constant property models. The variable property model predicts the wall heat fluxes and the Nu number values lower than the constant property model.
4. The effect of variable properties on circulation patterns and velocity profiles is more pronounced for high Ra numbers when the temperature difference between the walls of the enclosure is high.
5. There is negligible difference between the results of the first order and second order accurate in time models.

6. Transient development of circulation patterns and temperature isotherms were qualitatively comparable to those found in literature.

Chapter 14

Recommendations

The following recommendations are made to improve the present model in simulating natural convection inside an enclosure with differentially heated side walls.

1. A finer mesh should be utilized to obtain more accurate results, especially at high Rayleigh numbers.
2. Due to the boundary layer nature of the flow for high Rayleigh numbers, a stretched grid should be incorporated near the walls of the enclosure to improve the accuracy of the predicted results.
3. Effects of different boundary conditions such as constant and variable heat flux should be considered.
4. A turbulence model $k-\epsilon$ should be incorporated to study natural convection in the turbulent region, that is for $Ra > 10^6$.
5. Density should be considered as a primitive variable instead of pressure so that the model can be applied to any fluid (liquid or gas).

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Appendices

Appendix I

Vector Form of Governing Differential Equations

The continuity equation is given by:

$$\frac{D\rho}{Dt} + \rho(\vec{\nabla} \cdot \vec{V}) = 0 \quad (\text{I.1})$$

The momentum equations are given by:

$$\rho \frac{D\vec{V}}{Dt} = \rho \vec{g} - \vec{\nabla} p + \frac{\partial}{\partial x_j} \left[\mu \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) - \frac{2}{3} \delta_{i,j} \frac{\partial u_k}{\partial x_k} \right] \quad (\text{I.2})$$

where

$$\delta_{i,j} = \begin{cases} 1 & \text{if } i=j \\ 0 & \text{if } i \neq j \end{cases}$$

The energy equation is given by:

$$\rho \frac{\partial E_t}{\partial t} + \vec{\nabla} \cdot E_t \vec{V} = \frac{\partial Q}{\partial t} - \vec{\nabla} \cdot \vec{q} + \rho \vec{f} \cdot \vec{V} + \vec{\nabla} \cdot (\vec{\Pi}_{i,j} \cdot \vec{V}) \quad (\text{I.3})$$

where total energy, E_t ,

$$E_t = \rho \left(e + \frac{V^2}{2} + PE + \dots \right)$$

and the stress tensor,

$$\vec{\Pi}_{i,j} = -p \delta_{i,j} + \mu \left[\left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) - \frac{2}{3} \delta_{i,j} \frac{\partial u_k}{\partial x_k} \right] \quad \text{where } i, j, k = 1, 2, 3$$

The heat transfer by conduction in Eq. I.3 is

$$\vec{q} = k \vec{\nabla} T$$

Equation of state for ideal gas is given by:

$$p = \rho RT \quad (I.4)$$

Appendix II

Dimensional Form of Governing Differential Equations

(A. Conservative Form)

The continuity equation is given by:

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x}(\rho u) + \frac{\partial}{\partial y}(\rho v) = 0 \quad (\text{IIA.1})$$

The momentum equation in the x-direction is given by:

$$\begin{aligned} \frac{\partial}{\partial t}(\rho u) + \frac{\partial}{\partial x}(\rho u^2) + \frac{\partial}{\partial y}(\rho uv) + \frac{\partial p}{\partial x} - \frac{\partial}{\partial x} \left[\frac{2}{3} \mu \left(2 \frac{\partial u}{\partial x} - \frac{\partial v}{\partial y} \right) \right] \\ - \frac{\partial}{\partial y} \left[\mu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \right] + \rho g \sin \theta = 0 \end{aligned} \quad (\text{IIA.2})$$

The momentum equation in the y-direction is given by:

$$\begin{aligned} \frac{\partial}{\partial t}(\rho v) + \frac{\partial}{\partial x}(\rho uv) + \frac{\partial}{\partial y}(\rho v^2) + \frac{\partial p}{\partial y} \\ - \frac{\partial}{\partial x} \left[\mu \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \right] - \frac{\partial}{\partial y} \left[\frac{2}{3} \mu \left(2 \frac{\partial v}{\partial y} - \frac{\partial u}{\partial x} \right) \right] + \rho g \cos \theta = 0 \end{aligned} \quad (\text{IIA.3})$$

The energy equation is given by:

$$\begin{aligned} \frac{\partial}{\partial t}(\rho c_p T) + \frac{\partial}{\partial x}(\rho c_p T u) + \frac{\partial}{\partial y}(\rho c_p T v) \\ - \frac{\partial}{\partial x} \left(k \frac{\partial T}{\partial x} \right) - \frac{\partial}{\partial y} \left(k \frac{\partial T}{\partial y} \right) = 0 \end{aligned} \quad (\text{IIA.4})$$

Appendix II

Dimensional Form of Governing Differential Equations

(B. Using in terms of Primitive Variables u, v, p, T)

The continuity equation is given by:

$$\frac{\partial}{\partial t}\left(\frac{p}{T}\right) + \frac{\partial}{\partial x}\left(\frac{pu}{T}\right) + \frac{\partial}{\partial y}\left(\frac{pv}{T}\right) = 0 \quad (\text{IIB.1})$$

The momentum equation in the x-direction is given by:

$$\begin{aligned} \frac{\partial}{\partial t}\left(\frac{pu}{T}\right) + \frac{\partial}{\partial x}\left(\frac{puu}{T}\right) + \frac{\partial}{\partial y}\left(\frac{puv}{T}\right) + \frac{\partial(Rp)}{\partial x} - \frac{\partial}{\partial x}\left[\frac{2}{3}R\mu\left(2\frac{\partial u}{\partial x} - \frac{\partial v}{\partial y}\right)\right] \\ - \frac{\partial}{\partial y}\left[R\mu\left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}\right)\right] + \left(\frac{p}{T}\right)g \sin \theta = 0 \end{aligned} \quad (\text{IIB.2})$$

The momentum equation in the y-direction is given by:

$$\begin{aligned} \frac{\partial}{\partial t}\left(\frac{pv}{T}\right) + \frac{\partial}{\partial x}\left(\frac{puv}{T}\right) + \frac{\partial}{\partial y}\left(\frac{pvv}{T}\right) + \frac{\partial(Rp)}{\partial y} \\ - \frac{\partial}{\partial x}\left[R\mu\left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}\right)\right] - \frac{\partial}{\partial y}\left[\frac{2}{3}R\mu\left(2\frac{\partial v}{\partial y} - \frac{\partial u}{\partial x}\right)\right] + \left(\frac{p}{T}\right)g \cos \theta = 0 \end{aligned} \quad (\text{IIB.3})$$

The energy equation is given by:

$$\frac{\partial}{\partial t}(c_p p) + \frac{\partial}{\partial x}(c_p pu) + \frac{\partial}{\partial y}(c_p pv) - \frac{\partial}{\partial x}\left(Rk \frac{\partial T}{\partial x}\right) - \frac{\partial}{\partial y}\left(Rk \frac{\partial T}{\partial y}\right) = 0 \quad (\text{IIB.4})$$

Appendix III

Sub Models of Thermodynamic and Transport Properties

1. Variable Specific Heat at Constant Pressure ($C_p(T)$)

Constant-volume and constant-pressure specific heats are assumed to be functions of temperature. That is,

$$C_v = f(T) \text{ and } C_p = f(T)$$

because all gases approach ideal-gas behavior as the pressure approaches zero.

The specific heat at constant pressure function of temperature for an ideal gas, C_{p0} can be given as (refer to Sonntag, Engineering Thermodynamics)

$$C_{p0} = C_0 + C_1\phi + C_2\phi^2 + C_3\phi^3 \quad \text{KJ/Kg K} \quad (\text{III.1})$$

where

$$\phi = \frac{T}{1000}$$

For Air, values of constants in the above equation are as follows

$$C_0 = 1.05, C_1 = -0.365, C_2 = 0.85, C_3 = -0.39$$

This approximate form can be valid from 250K to 1200 K

2. Variable Viscosity ($\mu(T)$)

A widely used approximation resulted from a kinetic theory by Sutherland (1893) using an idealized intermolecular-force potential. The formula is (refer to Frank M. White, Viscous Flow)

$$\frac{\mu}{\mu_0} = \left(\frac{T}{T_0} \right)^{3/2} \frac{T_0 + S}{T + S} \quad (\text{III.2})$$

where

μ is the Viscosity, N.s/m²

μ_0 is the Reference value, N.s/m²

T is the Absolute Temperature, K

T_0 is the Reference Value, K and

S is the Sutherland Constant, which is characteristic of the gas, K

For Air, Values of constants in the above equation are as follows

$$\mu_0 = 1.716 \times 10^{-5} \text{ N.s/m}^2$$

$$T_0 = 273 \text{ K}$$

$$S = 111 \text{ K}$$

Appendix IV

Non-Dimensional Form of Governing Differential Equations

Continuity Equation

$$\frac{\partial}{\partial \bar{t}} \left(\frac{\bar{p}}{\bar{T}} \right) + \frac{\partial}{\partial \bar{x}} \left(\frac{\bar{p}\bar{u}}{\bar{T}} \right) + \frac{\partial}{\partial \bar{y}} \left(\frac{\bar{p}\bar{v}}{\bar{T}} \right) = 0 \quad (\text{IV.1})$$

Momentum equation in x-direction

$$\begin{aligned} \frac{\partial}{\partial \bar{t}} \left(\frac{\bar{p}\bar{u}}{\bar{T}} \right) + \frac{\partial}{\partial \bar{x}} \left(\frac{\bar{p}\bar{u}\bar{u}}{\bar{T}} \right) + \frac{\partial}{\partial \bar{y}} \left(\frac{\bar{p}\bar{u}\bar{v}}{\bar{T}} \right) + \frac{\partial}{\partial \bar{x}} (\bar{R}\bar{p}) \\ + \frac{1}{Gr^{1/2}} \frac{\partial}{\partial \bar{x}} \left[\frac{2}{3} \bar{R}\bar{\mu} \left(2 \frac{\partial \bar{u}}{\partial \bar{x}} - \frac{\partial \bar{v}}{\partial \bar{y}} \right) \right] - \frac{1}{Gr^{1/2}} \frac{\partial}{\partial \bar{y}} \left[\bar{R}\bar{\mu} \left(\frac{\partial \bar{u}}{\partial \bar{y}} + \frac{\partial \bar{v}}{\partial \bar{x}} \right) \right] + \frac{1}{Fr} \left(\frac{\bar{p}}{\bar{T}} \right) \sin \theta = 0 \end{aligned} \quad (\text{IV.2})$$

Momentum equation in y-direction

$$\begin{aligned} \frac{\partial}{\partial \bar{t}} \left(\frac{\bar{p}\bar{v}}{\bar{T}} \right) + \frac{\partial}{\partial \bar{x}} \left(\frac{\bar{p}\bar{u}\bar{v}}{\bar{T}} \right) + \frac{\partial}{\partial \bar{y}} \left(\frac{\bar{p}\bar{v}\bar{v}}{\bar{T}} \right) + \frac{\partial}{\partial \bar{y}} (\bar{R}\bar{p}) \\ - \frac{1}{Gr^{1/2}} \frac{\partial}{\partial \bar{x}} \left[\bar{R}\bar{\mu} \left(\frac{\partial \bar{v}}{\partial \bar{x}} + \frac{\partial \bar{u}}{\partial \bar{y}} \right) \right] - \frac{1}{Gr^{1/2}} \frac{\partial}{\partial \bar{y}} \left[\frac{2}{3} \bar{R}\bar{\mu} \left(2 \frac{\partial \bar{v}}{\partial \bar{y}} - \frac{\partial \bar{u}}{\partial \bar{x}} \right) \right] + \frac{1}{Fr} \left(\frac{\bar{p}}{\bar{T}} \right) \cos \theta = 0 \end{aligned} \quad (\text{IV.3})$$

Energy equation

$$\begin{aligned} \frac{\partial}{\partial \bar{t}} (\bar{c}_p \bar{p}) + \frac{\partial}{\partial \bar{x}} (\bar{c}_p \bar{p} \bar{u}) + \frac{\partial}{\partial \bar{y}} (\bar{c}_p \bar{p} \bar{v}) - \frac{1}{\text{Pr} Gr^{1/2}} \frac{\partial}{\partial \bar{x}} \left(\bar{R} \bar{c}_p \bar{\mu} \frac{\partial \bar{T}}{\partial \bar{x}} \right) \\ - \frac{1}{\text{Pr} Gr^{1/2}} \frac{\partial}{\partial \bar{y}} \left(\bar{R} \bar{c}_p \bar{\mu} \frac{\partial \bar{T}}{\partial \bar{y}} \right) = 0 \end{aligned} \quad (\text{IV.4})$$

$$U =$$

The elements of these matrices are calculated using the following formulas:

$$a_{i,j} = A_{i,j}^6 \quad (\text{V.1})$$

$$b_{i,j} = (A_{i,j}^5 - a_{i,j} f_{i-1,j-1} - \alpha A_{i,j}^4 f_{i+1,j-1}) (1 - \alpha f_{i,j-1} f_{i+1,j-1})^{-1} \quad (\text{V.2})$$

$$c_{i,j} = A_{i,j}^4 - b_{i,j} f_{i,j-1} \quad (\text{V.3})$$

$$d_{i,j} = (A_{i,j}^7 - a_{i,j} h_{i-1,j-1} - b_{i,j} g_{i,j-1} - 2\alpha a_{i,j} g_{i-1,j}) (1 + \alpha g_{i-1,j})^{-1} \quad (\text{V.4})$$

$$e_{i,j} = (A_{i,j}^9 - a_{i,j} h_{i-1,j-1} - b_{i,j} g_{i,j-1} - 2\alpha a_{i,j} g_{i-1,j} - d_{i,j} f_{i-1,j}) + \alpha(2\varphi_{i,j}^1 + \varphi_{i,j}^2 + \varphi_{i,j}^3 + 2\varphi_{i,j}^4) \quad (\text{V.5})$$

$$f_{i,j} = (e_{i,j})^{-1} (A_{i,j}^3 - b_{i,j} s_{i,j-1} - c_{i,j} h_{i+1,j-1} - 2\alpha(\varphi_{i,j}^1 + \varphi_{i,j}^3)) \quad (\text{V.6})$$

$$g_{i,j} = (e_{i,j})^{-1} (A_{i,j}^8 - d_{i,j} h_{i-1,j}) \quad (\text{V.7})$$

$$g_{i,j} = (e_{i,j})^{-1} (A_{i,j}^3 - b_{i,j} s_{i,j-1} - c_{i,j} h_{i+1,j-1} - 2\alpha(\varphi_{i,j}^1 + \varphi_{i,j}^3)) \quad (\text{V.8})$$

$$h_{i,j} = (e_{i,j})^{-1} (A_{i,j}^1 - d_{i,j} s_{i-1,j} - \alpha \varphi_{i,j}^4) \quad (\text{V.9})$$

$$s_{i,j} = (e_{i,j})^{-1} (A_{i,j}^2) \quad (\text{V.10})$$

where

$$\varphi_{i,j}^1 = c_{i,j} f_{i+1,j-1} \quad (\text{V.11})$$

$$\varphi_{i,j}^2 = a_{i,j} g_{i-1,j-1} \quad (\text{V.12})$$

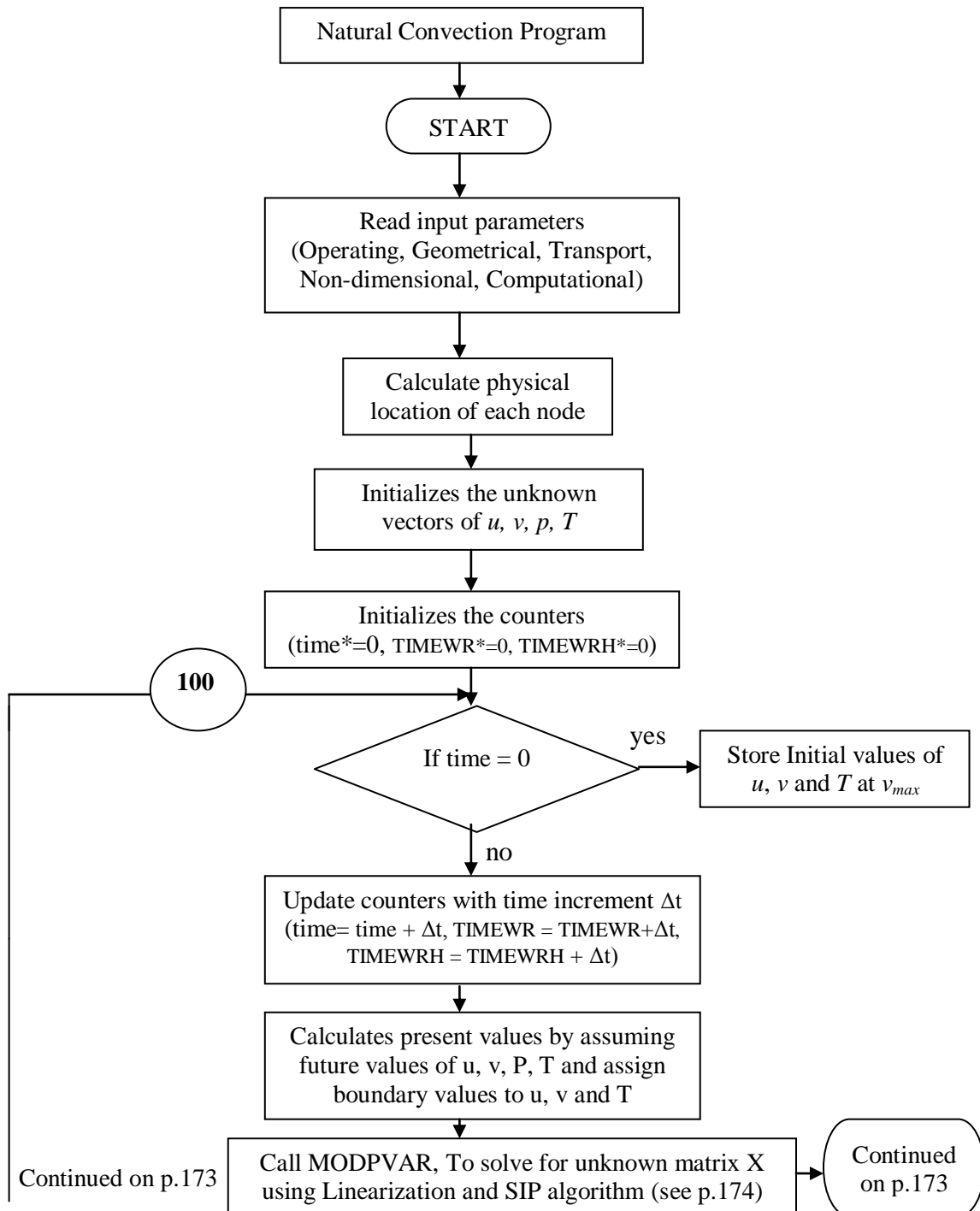
$$\varphi_{i,j}^3 = c_{i,j} s_{i+1,j-1} \quad (\text{V.13})$$

$$\varphi_{i,j}^4 = d_{i,j} g_{i-1,j} \quad (\text{V.14})$$

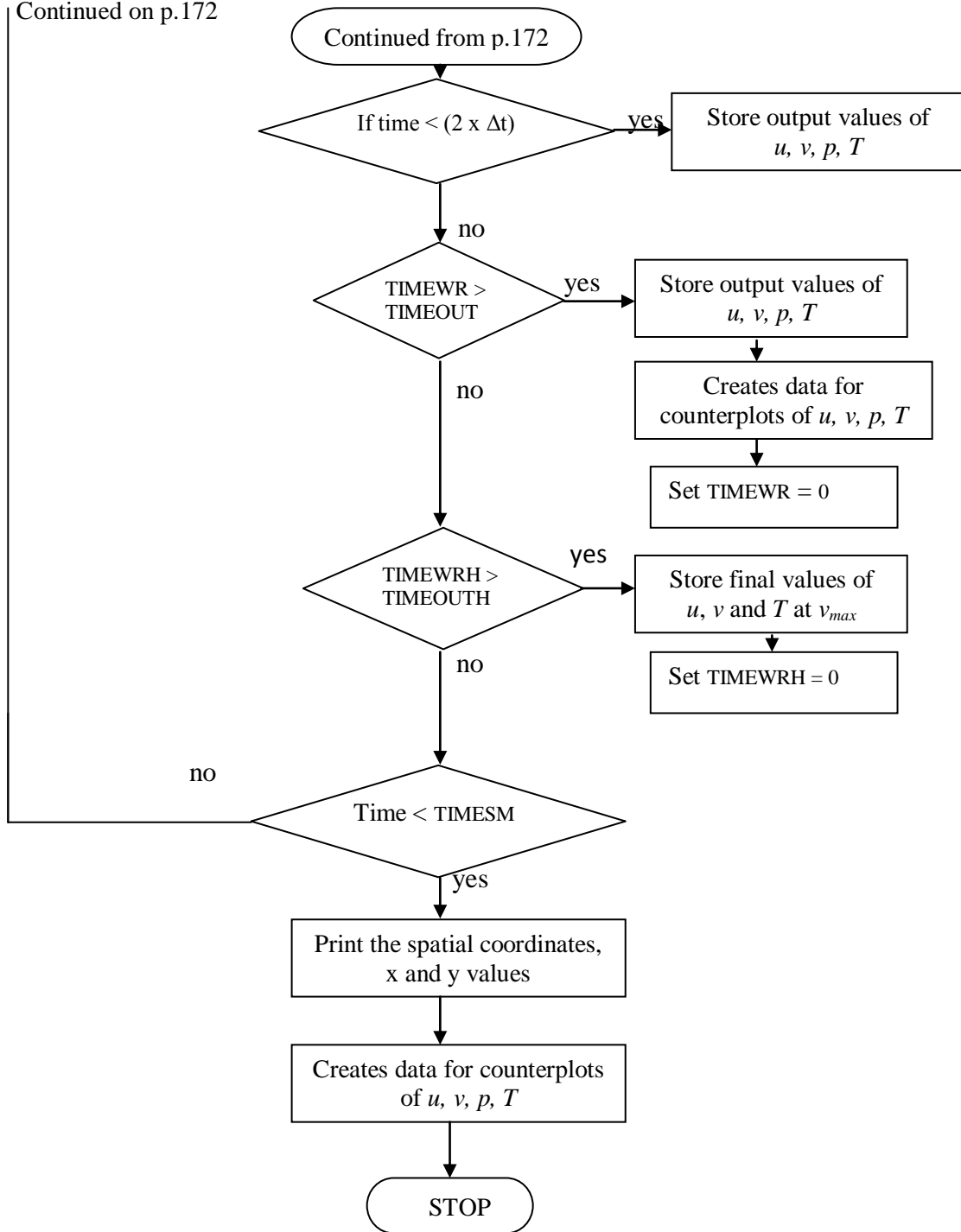
Appendix VI

Flow Chart for the Computer Program

(A. First Order Accurate in Time Model)

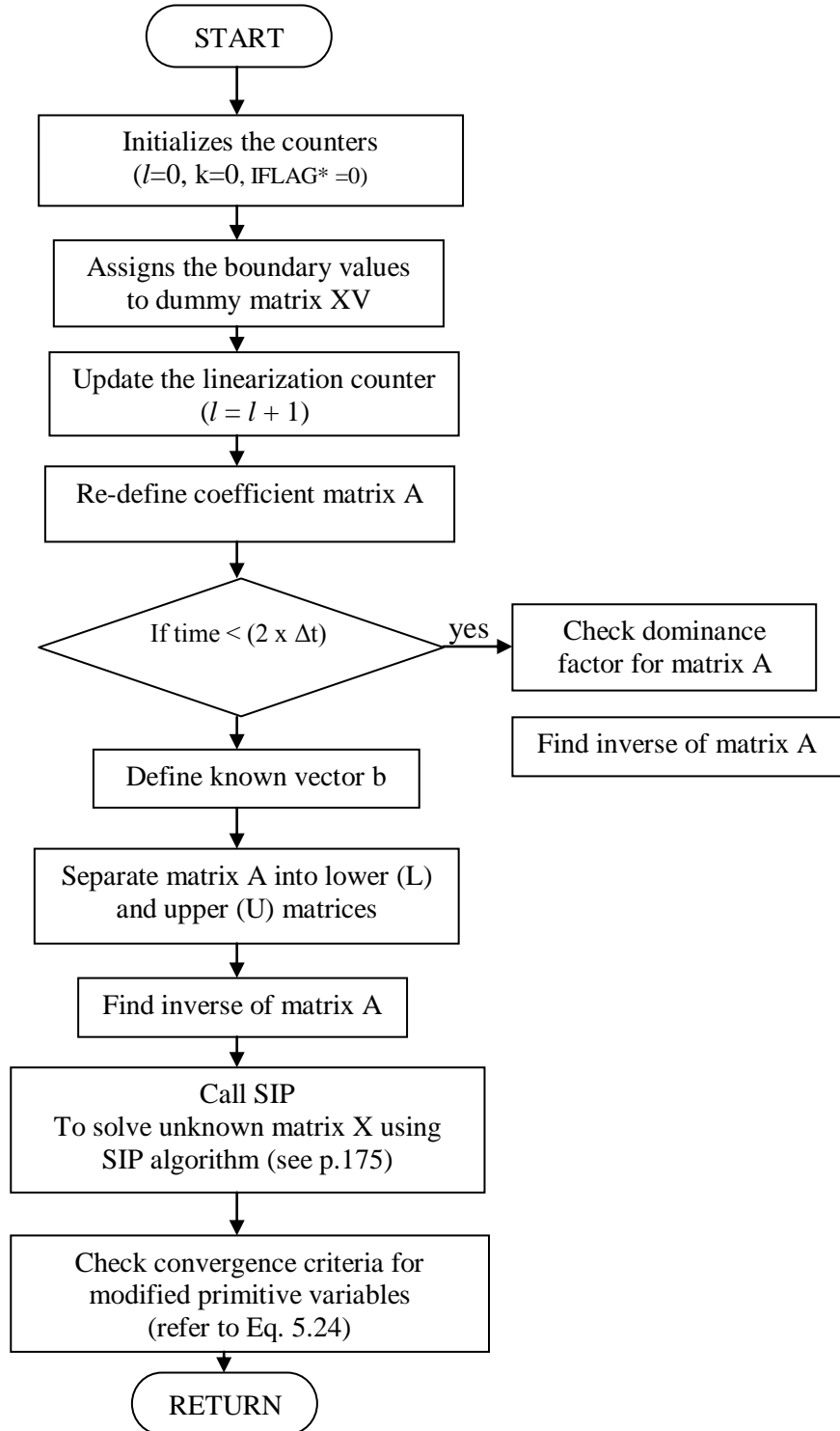


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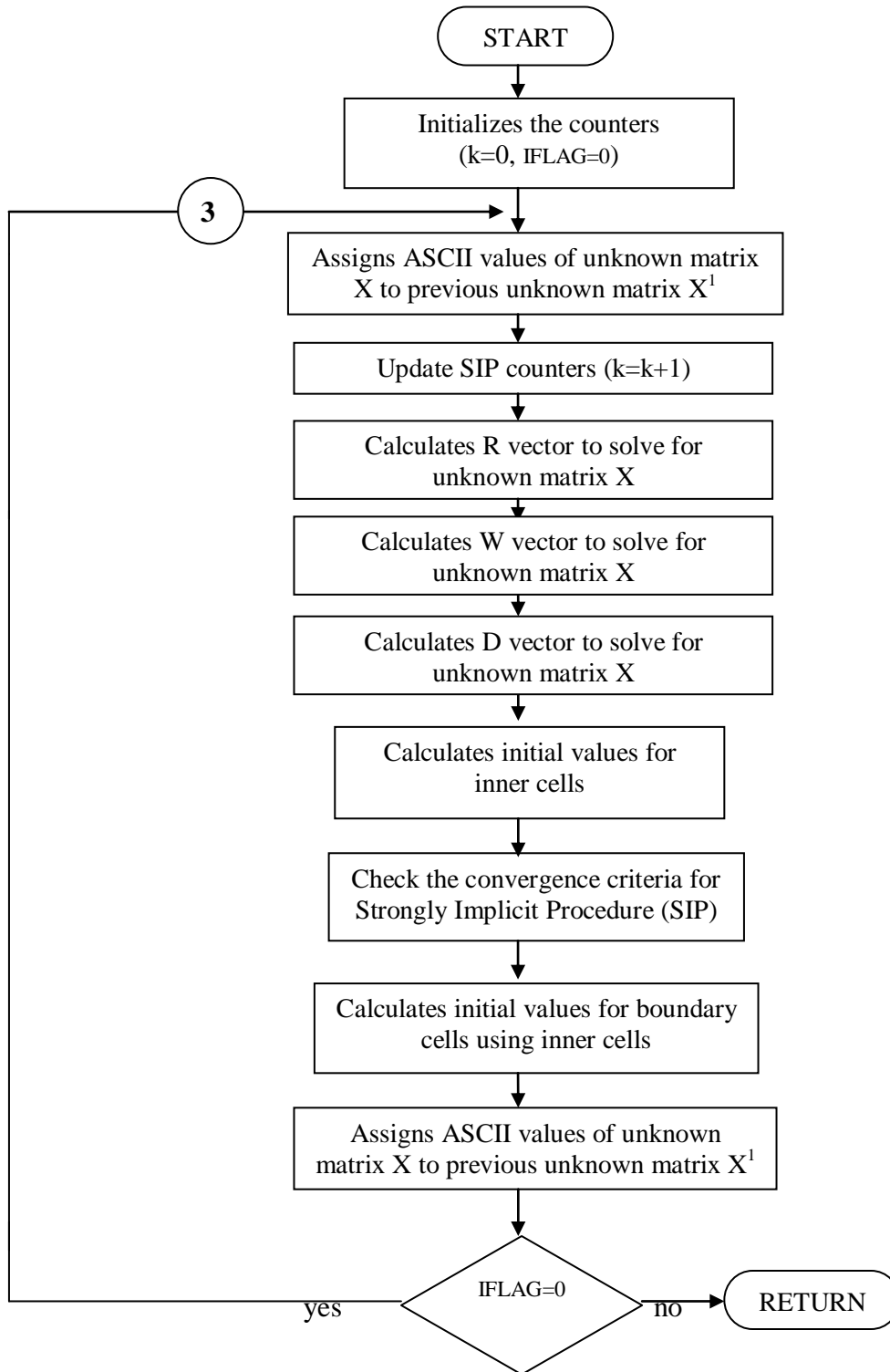
MODPVAR

(Subroutine for Modification of Primitive Variables using Linearization Iteration)



SIP

(Subroutine to Solve Unknown Primitive Variables using Strongly Implicit Procedure)



* List of Program Variables used in the Flow Chart

Time – Non-dimensional time

TIMESH - Total time of simulation

TIMEOUT - Time interval for prints

TIMEOUTH - Time interval for histograms

IFLAG - Conditional counter

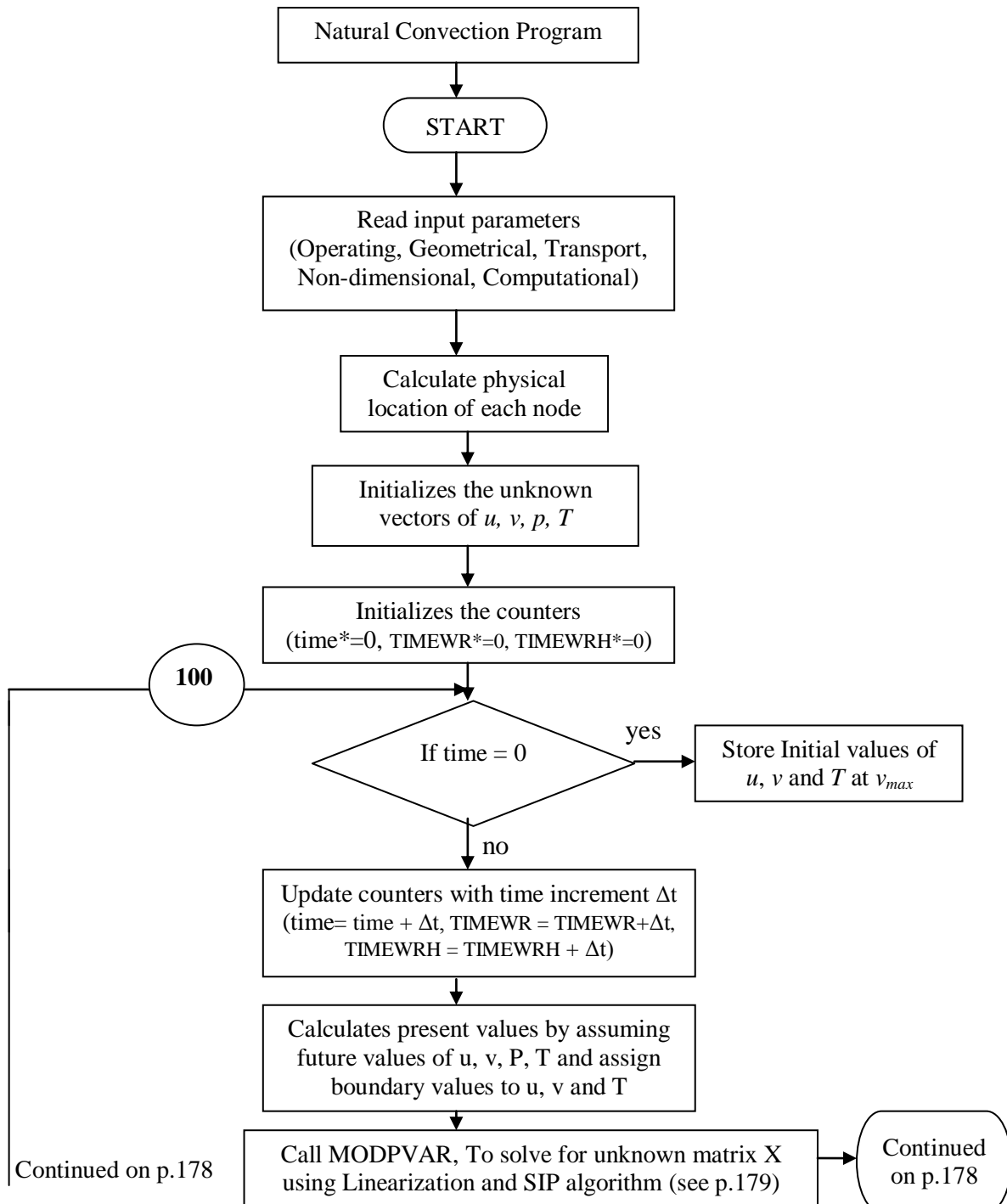
TIMEWR - Time to write ASCII values of output primitive variables

TIMEWRH - Time to write ASCII values of output primitive variables as histogram data

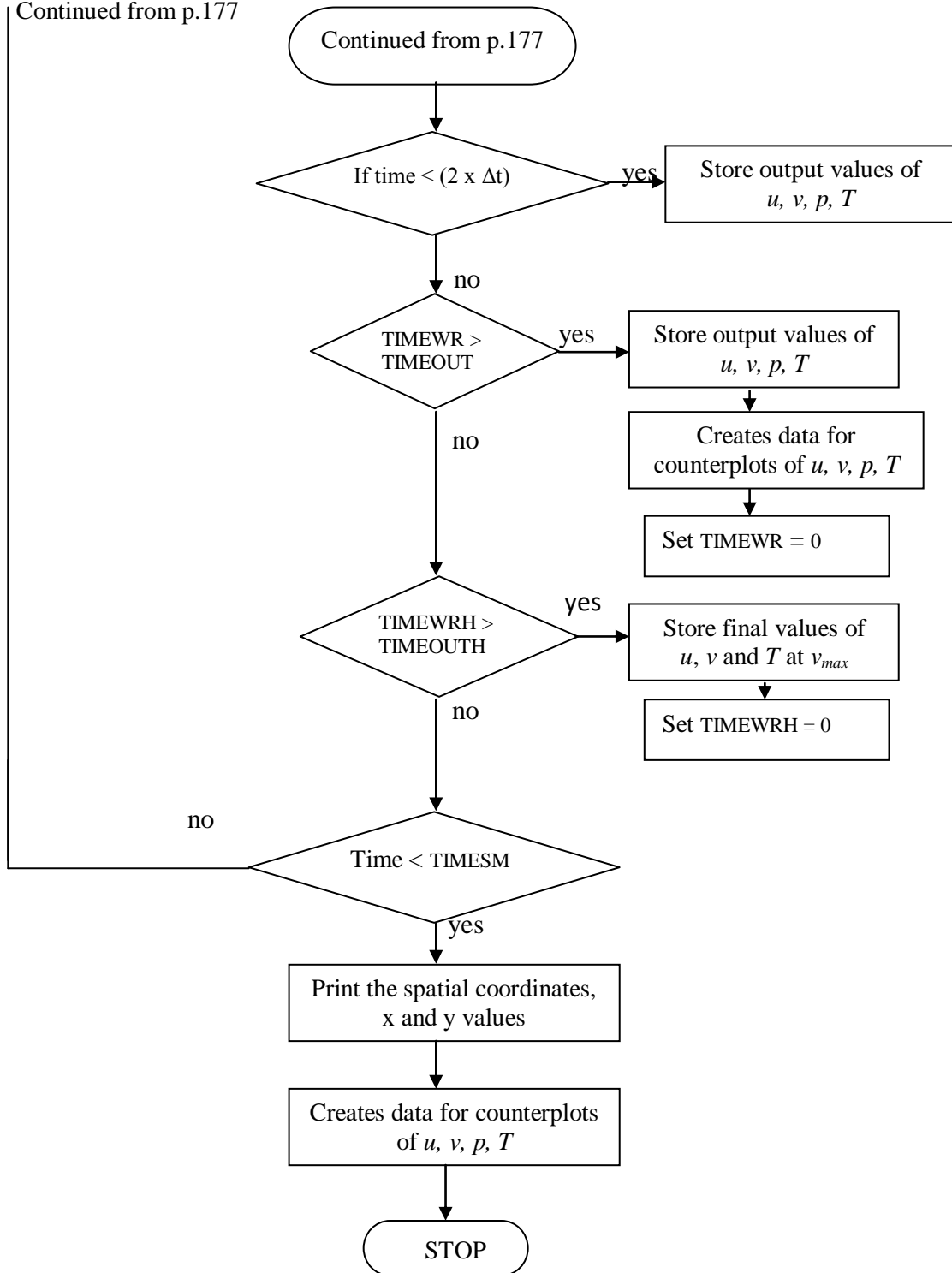
Appendix VI

Flow Chart for the Computer Program

(B. Second Order Accurate in Time Model)

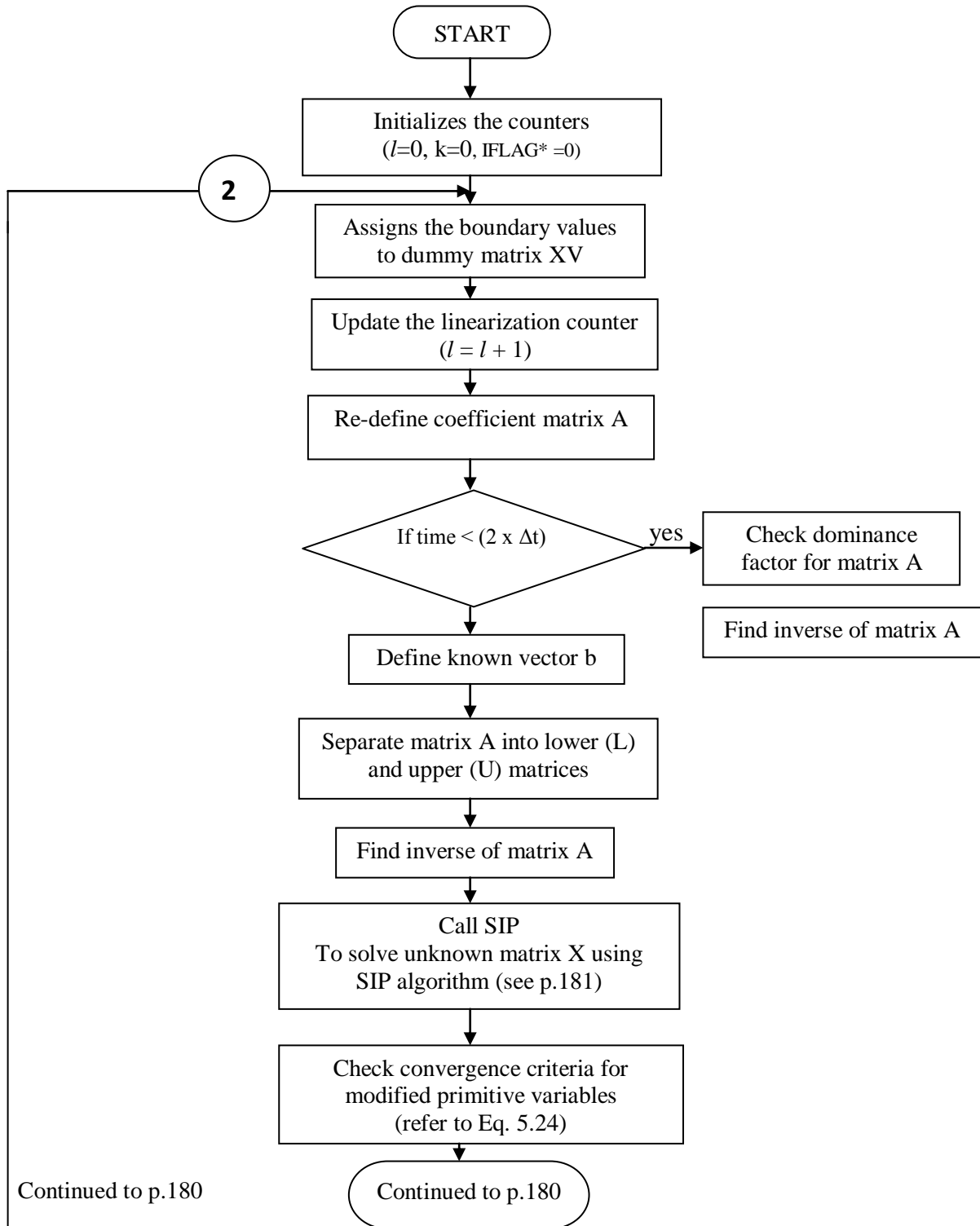


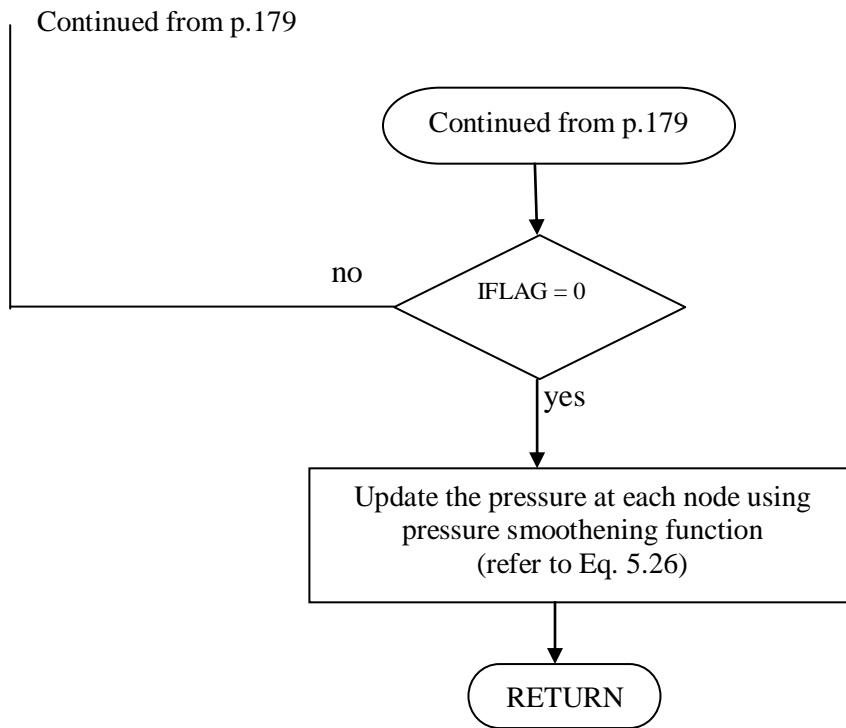
Continued from p.177



MODPVAR

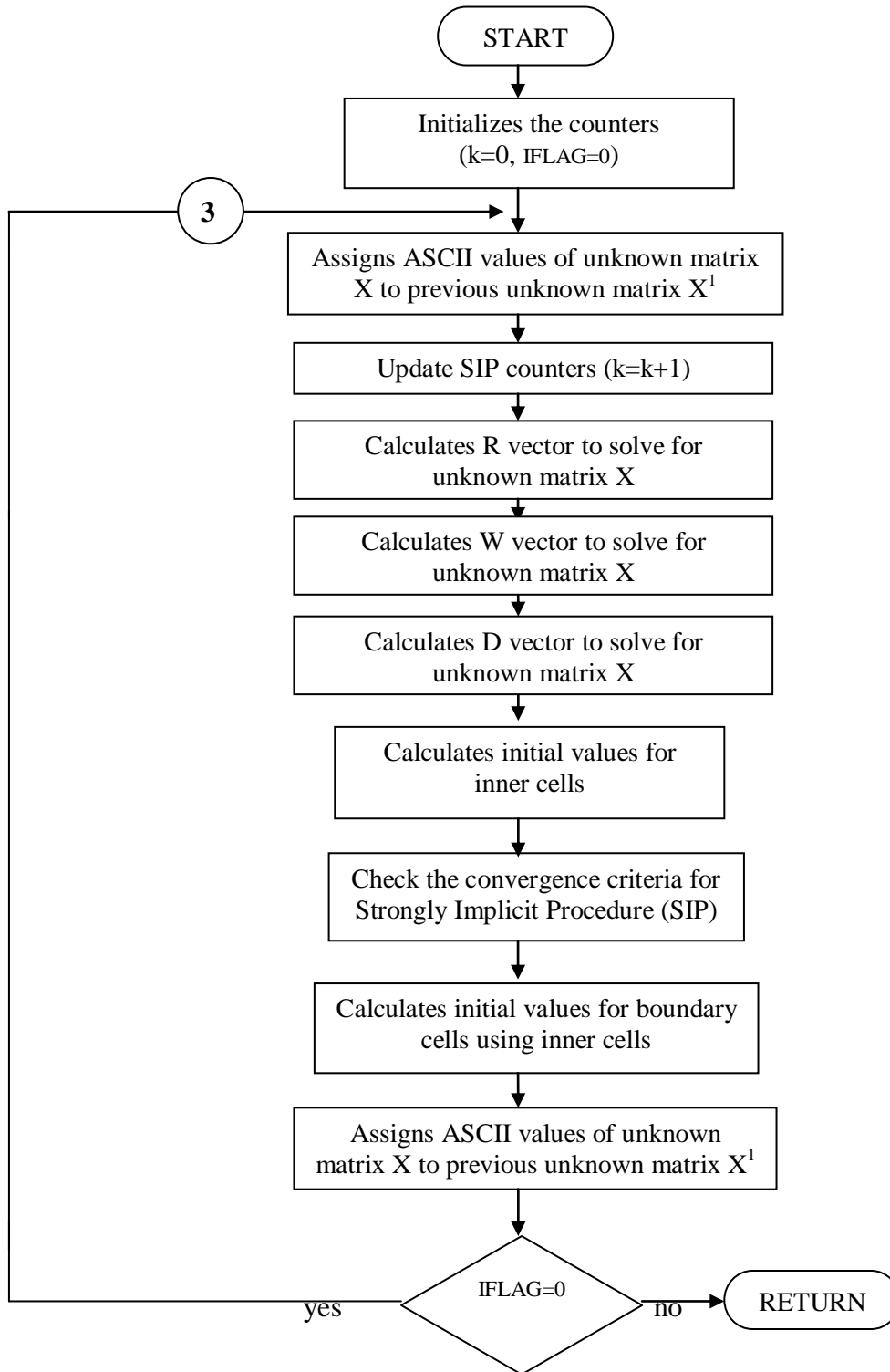
(Subroutine for Modification of Primitive Variables using Linearization Iteration)





SIP

(Subroutine to Solve Unknown Primitive Variables using Strongly Implicit Procedure)



* List of Program Variables used in the Flow Chart

Time – Non-dimensional time

TIMESH - Total time of simulation

TIMEOUT - Time interval for prints

TIMEOUTH - Time interval for histograms

IFLAG - Conditional counter

TIMEWR - Time to write ASCII values of output primitive variables

TIMEWRH - Time to write ASCII values of output primitive variables as histogram data

Appendix VII

Program Parameters – Common to all Runs

List of program Parameters used in the Computer Program that are common to all runs:

$$AR = 1$$

$$g = 9.81 \text{ m/s}^2$$

$$\alpha = 2.21 \times 10^5 \text{ m/s}^2$$

$$R = 287.0 \text{ J/Kg-K}$$

$$c_{p0} = 1006.5 \text{ J/Kg-K}$$

$$\mu_0 = 1.716 \times 10^{-5} \text{ J/Kg-K}$$

$$Pr = 0.72$$

$$\gamma = 1.4$$

$$\alpha = 0.2$$

$$\varepsilon = 1.0 \times 10^{-7}$$

$$\varepsilon_p = 1.0 \times 10^{-7}$$

Appendix VIII

Program Parameters - Specific to Each Run

List of parameters that are specific to the Base Run:

Mesh size: 21 x 21

$$\Delta \bar{t} = 1 \times 10^{-5}$$

$$Ra = 1 \times 10^5$$

$$Gr = 1.413265 \times 10^5$$

$$p_{ref} = 1 \times 10^5 \text{ N/m}^2$$

$$T_{ref} = 300 \text{ K}$$

$$u_{ref} = 0.232073 \text{ m/s (calculated)}$$

$$L_{ref} = 0.0254 \text{ m}$$

$$\Delta T = 64.84 \text{ K}$$

$$T_h = 332.42 \text{ K (calculated)}$$

$$T_c = 267.58 \text{ K (calculated)}$$

$$Fr = 0.464914$$

Vita

Manohar Chidurala was born on August 31, 1984 in Mutharam, Andhra Pradesh, India. He received his bachelor's degree in mechanical engineering from the Jawaharlal Nehru Technological University, Hyderabad, India in June, 2006. He was admitted to the University of New Orleans as a graduate student in mechanical engineering. During the course of his study, he worked as a graduate teaching assistant and had a chance to assistant in fluid mechanics and thermal sciences Laboratories. His research interests include fluid mechanics, heat transfer and computational fluid dynamics. In August 2009, he was granted a Master of Science degree in mechanical engineering from the University of New Orleans.