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First unbiased determination of the parton distributions of lead nuclei

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Abstract A nice abstract

We present

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1 Introduction

In a series of papers [1–11], the NNPDF collaboration has introduced a methodology aimed at reducing as much as possible this procedural uncertainty.

Now we want to apply it to the case of nuclear PDFs

2 Dataset

In this section we summarize the data that has been used, the kinematic cuts, and the treatment of experimental uncertainties

The data in our analysis consist of $l + A$ DIS measurements. These nuclear data comprise the NMC [12–15], SLAC 139 [16] and EMC [17] results for ratios of the DIS structure function $F_2^A(x, Q^2)$ for several heavy nuclei to those of deuterium, lithium or carbon, see Tab. 1.

The kinematical variables in DIS are the Bjorken- x and the virtuality of the photon Q^2 . As the compression of the nuclear data to lead structure function is done using EPS09 NLO nPDFs, the kinematical cut $Q^2 \geq 1.69$ is imposed [18].

The experimental uncertainty, $d\sigma_i^{exp}$, is obtained adding in quadrature the systematic and statistical errors for each point. The theoretical uncertainty, $d\sigma_i^{th}$, is the one due to the conversion factor. We require that only points where $d\sigma_i^{th} \leq 1.5 d\sigma_i^{exp}$ are included in the fit.

3 Compression of nuclear data to lead structure functions

Here we explain how all the nuclear ratio data for different values of A is compressed to the lead structure function.

We write the relevant conversion formula, discuss error propagation, and show representative results for the various experiments.

$$F_2^{Pb}(x, Q^2) = \frac{F_2^{Pb}(x, Q^2)}{F_2^A(x, Q^2)} R_{F_2^{D,C,Li}}^A(x, Q^2) F_2^{D,C,Li}(x, Q^2), \quad (1)$$

where $R_{F_2^{D,C,Li}}^A(x, Q^2)$ are the experimental data for ratios of the DIS structure function $F_2^A(x, Q^2)$ for various heavy nuclei to those for deuterium, lithium or carbon,

$$R_{F_2^{D,C,Li}}^A(x, Q^2) = \frac{F_2^A(x, Q^2)}{F_2^{D,C,Li}(x, Q^2)}, \quad (2)$$

see Tab. 1.

For $F_2^{Pb}(x, Q^2)$, $F_2^A(x, Q^2)$, $F_2^{Li}(x, Q^2)$ and $F_2^C(x, Q^2)$ we use EPS09 NLO nPDFs [18] and, in order to obtain the DIS structure functions for deuterium, $F_2^D(x, Q^2)$, we neglect any nuclear effect, assume isospin symmetry ($u^p = d^n$ and $d^p = u^n$) and use the free proton NLO PDFs of MSTW [19].

Therefore, the idea is to transform the DIS nuclear data, $R_{F_2^{D,C,Li}}^A(x, Q^2)$, into lead structure functions $F_2^{Pb}(x, Q^2)$ using equation (1). We define, the conversion factor of this transformation as,

$$\mathcal{C}(x, Q^2) = \frac{F_2^{Pb}(x, Q^2)}{F_2^A(x, Q^2)} F_2^{D,C,Li}(x, Q^2) \quad (3)$$

and, consequently, $F_2^{Pb}(x, Q^2) = \mathcal{C}(x, Q^2) R_{F_2^{D,C,Li}}^A(x, Q^2)$.

We use the Hessian method [20] to calculate the uncertainty of the conversion factor. In the Hessian approach, the corresponding uncertainty for a PDF-dependent quantity $\mathcal{O} = \mathcal{O}[f]$ can be computed as,

$$(\Delta\mathcal{O})^2 = \frac{1}{4} \sum_k (\mathcal{O}[S_k^+] - \mathcal{O}[S_k^-])^2, \quad (4)$$

Experiment	Nuclei	Data points	ref.
SLAC E-139	He(4)/D	9	[16]
NMC 95, re.	He/D	8	[12]
NMC 95	Li(6)/D	10	[13]
NMC 95, Q^2 dependence	Li/D	144	[13]
SLAC E-139	Be(9)/D	9	[16]
NMC 96	Be/C	12	[14]
CERN EMC	C(12)/D	9	[17]
SLAC E-139	C/D	3	[16]
NMC 95, NMC 95, re.	C/D	10	[12, 13]
NMC 95, Q^2 dependence	C/D	159	[13]
NMC 95, re.	C/Li	6	[12]
SLAC E-139	Al(27)/D	13	[16]
NMC 96	Al/C	15	[14]
SLAC E-139	Ca(40)/D	5	[16]
NMC 95, re.	Ca/D	15	[12]
NMC 95, re.	Ca/Li	7	[12]
NMC 96	Ca/C	15	[14]
SLAC E-139	Fe(56)/D	23	[16]
NMC 96	Fe/C	15	[14]
CERN EMC	Cu(64)/D	19	[17]
SLAC E-139	Ag(108)/D	7	[16]
CERN EMC	Sn(117)/C	8	[17]
NMC 96	Sn/C	10	[14]
NMC 96, Q^2 dependence	Sn/C	139	[15]
SLAC E-139	Au(197)/D	17	[16]
NMC 96	Pb/C	15	[14]
Total		702	

Table 1: Data sets included in the analysis. The mass numbers are indicated in parentheses. The number of data points refers to those falling within our cuts: $Q^2 \geq 1.69$ and $d\sigma_i^{th} \leq 1.5 d\sigma_i^{exp}$.

where $\mathcal{O}[S_k^+]$ and $\mathcal{O}[S_k^-]$ denote the values of the quantity \mathcal{O} , computed by the nPDF error sets S_k^+ and S_k^- and S_0 is the central PDF set.

For the lithium and carbon cases, the uncertainty of the conversion factor is given by

$$(\Delta\mathcal{C})^2 = \frac{1}{4} \sum_k \left(\frac{F_2^{Pb}[S_k^+]}{F_2^A[S_k^+]} F_2^{C,Li}[S_k^+] - \frac{F_2^{Pb}[S_k^-]}{F_2^A[S_k^-]} F_2^{C,Li}[S_k^-] \right)^2. \quad (5)$$

Neglecting the uncertainty due to the deuterium PDF, the uncertainty of the conversion factor is in this case,

$$(\Delta\mathcal{C})^2 = \frac{1}{4} \sum_k \left(\frac{F_2^{Pb}[S_k^+]}{F_2^A[S_k^+]} F_2^D[S_0] - \frac{F_2^{Pb}[S_k^-]}{F_2^A[S_k^-]} F_2^D[S_0] \right)^2, \quad (6)$$

where $F_2^D[S_0]$ is the value of F_2 computed by the central NLO MSTW PDF set [19] and $F_2^{Pb,C,Li,A}[S_k^{\pm,+}]$ are computed by the NLO EPS09 sets [18].

Therefore, the theoretical uncertainty of the lead structure functions is

$$d\sigma^{th} = \Delta\mathcal{C} R_{F_2^{D,C,Li}}^A. \quad (7)$$

4 Fitting methodology

Here we discuss the fitting methodology, the sum rules, the flavor decomposition, positivity, the parameterization of nuclear PDFs as ratios etc

5 Results

Here we show, well, the results of the fit

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