Boring but important disclaimers:

If you are not getting this from the GitHub repository or the associated Canvas page (e.g. CourseHero, Chegg etc.), you are probably getting the substandard version of these slides Don't pay money for those, because you can get the most updated version for free at

https://github.com/julianmak/academic-notes

The repository principally contains the compiled products rather than the source for size reasons.

- Associated Python code (as Jupyter notebooks mostly) will be held on the same repository. The source data however might be big, so I am going to be naughty and possibly just refer you to where you might get the data if that is the case (e.g. JRA-55 data). I know I should make properly reproducible binders etc., but I didn't...
- ▶ I do not claim the compiled products and/or code are completely mistake free (e.g. I know I don't write Pythonic code). Use the material however you like, but use it at your own risk.
- As said on the repository, I have tried to honestly use content that is self made, open source or explicitly open for fair use, and citations should be there. If however you are the copyright holder and you want the material taken down, please flag up the issue accordingly and I will happily try and swap out the relevant material.

OCES 2003 : Descriptive Physical Oceanography

(a.k.a. physical oceanography by drawing pictures)

Lecture 10: Mechanical forcing 4 (diffusion and friction)



Outline

- energy + momentum
 - \rightarrow recap of concept
 - \rightarrow sources and sinks
- concept of diffusion and friction
 - \rightarrow inherently small-scale
 - → largely driven by dynamics!
 - → Reynolds/Ekman number (cf. Rossby number)
 - → Ekman boundary layers revisited

Key terms: friction, diffusion, viscosity/diffusivity, boundary layers, Reynolds/Ekman number



Recap: equations of motion

Denoting u = (u, v) and $u_3 = (u, v, w)$, to <u>numerous</u> approximations (!!!) (see OCES 3203) ocean dynamics is governed by

$$\rho_0 \left(\frac{\partial u}{\partial t} + u \cdot \nabla u + 2\Omega \times u \right) = -\nabla p + F_u + \frac{D_u}{D_u}$$
 (1)

$$\frac{\partial p}{\partial z} = -\rho g \tag{2}$$

$$\nabla \cdot \boldsymbol{u}_3 = 0 \tag{3}$$

$$\left(\frac{\partial T}{\partial t} + \mathbf{u}_3 \cdot \nabla T\right) = F_T + \mathbf{D}_T \tag{4}$$

$$\left(\frac{\partial S}{\partial t} + u_3 \cdot \nabla S\right) = F_S + \frac{D_S}{} \tag{5}$$

$$\rho = \rho(T, S, p) \tag{6}$$

Respectively, (1) momentum equation, (2) hydrostatic balance, (3) incompressibility, (4) temperature equation, (5) salinity equation, and (6) equation of state (EOS)



Recap: wind forcing



Figure: Schematic of Ekman layer (boundary denoted by orange).

- wind a source of momentum for the ocean
- but influence has vertical limit
- direct influence only over the Ekman (boundary) layer
- difference in wind/current speed ⇒ transfer of momentum ocean (usually into ocean and hence source; why?)
 - \rightarrow molecular diffusive rate \Rightarrow very slow! (see next Lec.)
 - \rightarrow instabilities \Rightarrow much faster (because on dynamical time-scales; see Lec. 17)

Wind forcing the chief source of momentum into ocean, but where is the sink? (vertical transfer in Lec 13)



Recap: wind forcing

- Ekman transport (mass flux) perpendicular to wind vector
- Ekman suction/pumping (i.e. up/downwelling) related to wind stress curl

$$w_e \sim \frac{1}{f} e_z \cdot (\nabla \times \boldsymbol{\tau})$$

- \rightarrow related to fluid di/convergence $\nabla \cdot \boldsymbol{u}$
- ightarrow analogous relations with vertical component of fluid vorticity $\omega = e_7 \cdot (\nabla \times u)$

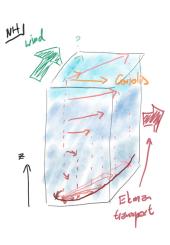


Figure: Schematic of Ekman spiral.

What sets Ekman layer depth?

Recall that (Lec. 4)

- linear) momentum p = mv (units = ???)
 - \rightarrow here it will be $p = \rho u$
- ▶ Newton's law refers to momentum <u>not</u> velocity

Recall that (Lec. 4)

- ► (linear) momentum p = mv (units = ???)

 → here it will be $p = \rho u$
- ▶ Newton's law refers to momentum not velocity
- ▶ kinetic energy is $K = (\rho/2)|u|^2 = (\rho/2)u \cdot u$ (units = ???)
 - \rightarrow note this is a scalar (why?)
 - \rightarrow might sometimes see it without a factor of ρ (units = ???)
- potential energy is (!) $P = \rho gz$ (but need a reference)

Recall that (Lec. 4)

- ► (linear) momentum p = mv (units = ???)

 → here it will be $p = \rho u$
- ▶ Newton's law refers to momentum <u>not</u> velocity
- ▶ kinetic energy is $K = (\rho/2)|u|^2 = (\rho/2)u \cdot u$ (units = ???)
 - \rightarrow note this is a scalar (why?)
 - \rightarrow might sometimes see it without a factor of ρ (units = ???)
- potential energy is (!) $P = \rho gz$ (but need a reference)
- wind puts momentum + energy into ocean
- ▶ thermodynamic forcing puts energy in ($P \nearrow but P \leftrightarrow K$)

Recall that (Lec. 4)

- ► (linear) momentum p = mv (units = ???)

 → here it will be $p = \rho u$
- ▶ Newton's law refers to momentum not velocity
- ▶ kinetic energy is $K = (\rho/2)|u|^2 = (\rho/2)u \cdot u$ (units = ???)
 - \rightarrow note this is a scalar (why?)
 - \rightarrow might sometimes see it without a factor of ρ (units = ???)
- potential energy is (!) $P = \rho gz$ (but need a reference)
- wind puts momentum + energy into ocean
- ▶ thermodynamic forcing puts energy in ($P \nearrow but P \leftrightarrow K$)

dynamics ↔ mass/momentum/energy transfer + conservation





Figure: Making a mess of coffee + milk...

Beyond advection (more later), momentum, energy and other tracers (e.g. temperature, salinity, oxygen) can transfer by diffusion

- if no stirring then, milk will gradually spread
 - \rightarrow but very **slowly**!



Figure: Making a mess of coffee + milk...

Beyond advection (more later), momentum, energy and other tracers (e.g. temperature, salinity, oxygen) can transfer by diffusion

- if no stirring then, milk will gradually spread
 - \rightarrow but very **slowly**!

- tendency to spread = diffusion
 - → diffusion aims to erase **gradients** of "stuff"
 - \rightarrow a **redistribution**, i.e. total amount of "stuff' **conserved**

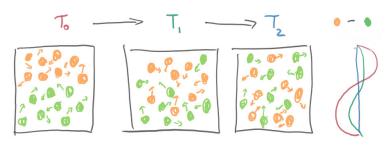


Figure: Schematic of microscopic motion leading to macroscopic diffusion. Here it is actual particle types, but can also imagine particles carrying "stuff", bumping into each other and transferring "stuff", and eventually the distribution of "stuff" evens out (but the total is conserved).

- ▶ at small-scale (!), "random" motion of particles (Brownian motion)
- collision ⇒ transfers ⇒ spreading
 → physics (!) ⇒ gradient erasing transfers more likely than gradient increasing ones (e.g. entropy)
- molecular diffusion, inherently small-scale



- lacktriangle diffusive rate \sim magnitude of **gradients** and **diffusivity** κ
 - \rightarrow thermal diffusivity κ_T , salinity diffusivity κ_S etc.

- ▶ diffusive rate \sim magnitude of **gradients** and diffusivity κ \rightarrow thermal diffusivity κ_T , salinity diffusivity κ_S etc.
- to a good approximation (!!!) momentum "spreading" can be treated as a diffusion

```
but called viscosity, denoted \nu or \mu = \rho \nu (I'm going to use \nu, the kinematic viscosity)
```

 \rightarrow think viscosity as how **sticky** something is

kinematic viscosity)

- ▶ diffusive rate \sim magnitude of **gradients** and diffusivity κ \rightarrow thermal diffusivity κ_T , salinity diffusivity κ_S etc.
- to a good approximation (!!!) momentum "spreading" can be treated as a diffusion but called viscosity, denoted ν or $\mu = \rho \nu$ (I'm going to use ν , the
 - \rightarrow think viscosity as how **sticky** something is
- ► **molecular** diffusivities/viscosities (units of m² s⁻¹) are substance dependent and **known**

material	κ_T	κ_S	ν
seawater	10^{-7}	10^{-9}	10^{-6}
air	10^{-5}		10^{-5}
honey	10^{-6}		10^{-2}
lava	10^{-7} (!)		depends, 10 ⁰ ?
steel	10^{1}	<u> </u>	big (!?)

Table: Table of kinematic molecular diffusivity/viscosity values at some control conditions. All numerical entries have units of ${\rm m^2~s^{-1}}$.

material	κ_T	κ_S	ν
seawater	10^{-7}	10^{-9}	10^{-6}
air	10^{-5}		10^{-5}
honey	10^{-6}		10^{-2}
lava	10^{-7} (!)		depends, 10 ⁰ ?
steel	10^{1}	_	big (!?)

Table: Table of kinematic molecular diffusivity/viscosity values at some control conditions. All numerical entries have units of $\mathrm{m}^2\,\mathrm{s}^{-1}$.

 \triangleright κ has units of m² s⁻¹, define a diffusion time as

$$t_m = L^2/\kappa$$

- \rightarrow time t_m to travel distance L is transport is only by molecular diffusion with diffusivity κ
- \rightarrow for sea water, if L=100 m (surface to thermocline say), temperature diffusion time is about 3000 yrs! (homework)

clearly something else is responsible for tracers moving beyond diffusion (generally speaking, not just in the ocean)

clearly something else is responsible for tracers moving beyond diffusion (generally speaking, not just in the ocean)



Figure: Making a mess of coffee + milk...

clearly something else is responsible for tracers moving beyond diffusion (generally speaking, not just in the ocean)



Figure: Making a mess of coffee + milk...

▶ stirring leads to larger-than-molecular transport, spreading and erasing gradients (!!! eventually anyway...) \rightarrow quantify as an effective diffusivity $\kappa_{\ell} \gg \kappa_{\text{molecular}}$

But we don't know κ_e , because it is dynamics dependent!



► resulting dynamic time-scales $t_d = \kappa_e \ll t_m$ \rightarrow e.g., for $\kappa_{e,z} = O(10^{-2} \text{ m}^2 \text{ s}^{-1})$, then for example above t_d is about **10 days** ($\ll t_m = 3000 \text{ yrs!}$)

- ► resulting dynamic time-scales $t_d = \kappa_e \ll t_m$ \rightarrow e.g., for $\kappa_{e,z} = O(10^{-2} \text{ m}^2 \text{ s}^{-1})$, then for example above t_d is about **10 days** ($\ll t_m = 3000 \text{ yrs!}$)
- the ocean has anisotropic effective diffusivities

$$ightarrow \kappa_{e,z} = 10^{-4} - 10^{-5} \; \mathrm{m^2 \; s^{-1}}$$

$$ightarrow \kappa_{e,h} = 10^{-1} - 10^3 \ \mathrm{m^2 \ s^{-1}}$$

- → ocean stratification inhibits vertical motion/dynamics (well really cross-isopycnal motion; see Lec. 5, 6, 17)
- → both can be larger in boundary layers (see later)

- ► resulting dynamic time-scales $t_d = \kappa_e \ll t_m$ \rightarrow e.g., for $\kappa_{e,z} = O(10^{-2} \text{ m}^2 \text{ s}^{-1})$, then for example above t_d is about **10 days** ($\ll t_m = 3000 \text{ yrs!}$)
- the ocean has anisotropic effective diffusivities

$$ightarrow \kappa_{e,z} = 10^{-4} - 10^{-5} \; \mathrm{m^2 \; s^{-1}}$$

$$ightarrow \kappa_{e,h} = 10^{-1} - 10^3 \ \mathrm{m^2 \ s^{-1}}$$

- ightarrow ocean stratification inhibits vertical motion/dynamics (well really cross-isopycnal motion; see Lec. 5, 6, 17)
- → both can be larger in boundary layers (see later)

Understanding dynamics fundamental in getting ocean transport/circulation right! (see Lec. 15-20)



Friction

Friction = resistance to **relative** motion

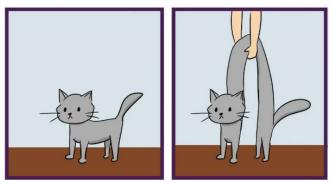
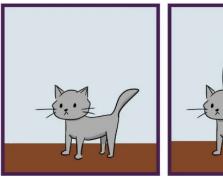


Figure: Cat physics. Picture from Meowingtons.

Friction

Friction = resistance to **relative** motion



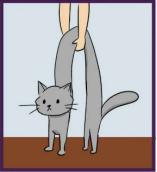


Figure: Cat physics. Picture from Meowingtons.

- ightharpoonup resistance to **relative** motion \Rightarrow needs **motion**
 - ightarrow function of velocity u (κ can sometimes be function of u too...)
 - ightarrow e.g. linear drag -ru (moving a pig along the floor)
 - ightarrow e.g. quadratic drag $-C_d|u|u$ (air resistance after a certain point)



Friction and diffusion

- internal stresses of object resisting relative motion
 - → contributes to fluid viscosity (cats are viscous then?)
- viscosity redistributes momentum (but total momentum conserved), and friction removes velocity differences

Friction and diffusion

- internal stresses of object resisting relative motion
 - → contributes to fluid viscosity (cats are viscous then?)
- viscosity redistributes momentum (but total momentum conserved), and friction removes velocity differences



Figure: Momentum diffusion and friction. Friction arises because there is relative motion, so it might look like there is momentum loss if we are looking only at the ocean as it's own system. But really what is happening is that the ocean is transferring momentum into the land, and from the ocean + earth system point of view there is no momentum lost

Diffusion and friction



Figure: Same as above but for atmosphere and land. Atmosphere acts as a **source** of momentum for ocean most of the time (equivalently, the ocean acts as a **sink** of momentum for the atmosphere as well as a **drag** for the atmospheric winds).

Similar picture but between fluid of different density too
 → more details about interfacial + topographic form stress later (see Lec 13)

Diffusion and friction

Some key points:

- diffusion reduces gradients + variance of "stuff", but conserves "stuff"
 - ightarrow e.g. viscosity reduces K and ∇p (bit sloppy here with notation), but conserves total p
- effective diffusion = molecular + dynamics, dynamics usually the most important contribution across the length-scales (see Reynolds number later)
- ► friction removes velocity differences (and therefore *K*)
 - → but momentum of the "complete" system should be conserved

Diffusion and friction

Some key points:

- diffusion reduces gradients + variance of "stuff", but conserves "stuff"
 - ightharpoonup e.g. viscosity reduces K and ∇p (bit sloppy here with notation), but conserves total p
- effective diffusion = molecular + dynamics, dynamics usually the most important contribution across the length-scales (see Reynolds number later)
- ► friction removes velocity differences (and therefore *K*)
 - → but momentum of the "complete" system should be conserved

Fundamentally dependent on dynamics!



Boundary layers



Figure: Schematic of Ekman layer (boundary denoted by orange).

- momentum input by atmosphere via frictional stress
 - \rightarrow details actually complicated and current research topic

(e.g. search for bulk formulas, mixed layer dynamics etc.)

- Q. what sets the boundary layer extent?
- Q. where/when is diffusion (and thus friction) important?



Non-dimensional numbers

Recall Ro = (U/L)/f is a measure of relative importance of advection and rotation

- ▶ Ro ≪ 1 means rotation important
 - \rightarrow could use it as "how small does *L* need to be for rotation to not matter" (cf. diffusion t_m above)

Non-dimensional numbers

Recall Ro = (U/L)/f is a measure of relative importance of advection and rotation

- ▶ Ro \ll 1 means rotation important \rightarrow could use it as "how small does L need to be for rotation to not matter" (cf. diffusion t_m above)
- ▶ for general diffusion the Péclet number is defined as

$$Pe = \frac{UL}{\kappa} = \frac{\text{advective transport}}{\text{diffusive transport}}$$

Reynolds + Ekman number

when dealing with viscosity it is the Reynolds number

$$Re = \frac{UL}{\nu}$$

ightarrow if Re \ll 1 (e.g. L and/or U small, ν large) then **viscous** effects are important

for rotating fluid systems the Ekman number

$$Ek = \frac{\nu}{fL^2} = \frac{(U/L)/f}{UL/\nu} = \frac{Ro}{Re}$$

is usually the important one

 \rightarrow might see other conventions and analogues (e.g. plasma physics?)



Boundary layers

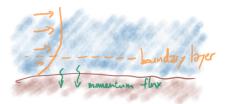


Figure: Schematic of a Prandtl boundary layer.

b boundary layer is where diffusive effects important, define it as e.g. Re \approx 1, i.e.

$$Re = \frac{UL}{\nu} \qquad \Rightarrow \qquad L_{bl} \sim \frac{\nu}{U}$$

- \rightarrow complication since *U* depends on *L*
- \rightarrow one standard solution has $L_{bl} \sim 1/\sqrt{\text{Re}}$, i.e. boundary layer width **decreases** with increasing Reynolds number



Boundary layers

▶ similarly, the Ekman layer depth could be when Ek \approx 1, i.e.

$$\mathrm{Ek} = \frac{\nu}{fL^2} \qquad \Rightarrow \qquad L_{ek} \sim \sqrt{\frac{\nu}{f}}$$

 \rightarrow since (!) $f \sim 10^{-4} \, \mathrm{s}^{-1}$, if we take $\nu = \nu_{e,z} = 10^{-2} \, \mathrm{m}^2 \, \mathrm{s}^{-1}$, we have

$$L_{ek} \sim \sqrt{\frac{10^{-2}}{10^{-4}}} = \sqrt{10^2} = 10 \text{ m}$$

- if we use the molecular ν the Ekman layer depth as well as Prandtl boundary layer width will be tiny (homework)
 - \rightarrow Re \gg 1 for geophysical flows
 - ightarrow Ek \ll 1 because ν and f small

Summary

- diffusion is spreading (erasing gradients)
 - ightarrow redistributing things (variance decreasing but conserved total)
 - \rightarrow depends on magnitude of gradients and diffusivity κ
 - → molecular vs. effective (latter dynamics dependent)
- friction resists relative motion
 - \rightarrow motion dependent
 - \rightarrow sink/source of momentum/energy of ocean (to land/atmosphere) (see Lec. 13)
- ▶ inter-system transfers through boundary layers
 - \rightarrow where Re and/or Ek = O(1)
 - \rightarrow these are usually **thin** (important consequences!)

another reason why dynamics is important!

