## Partial Differential Equations in Finance

Your Name

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Introduction to PDEs in Financial

**Applications** 

#### **Overview of Topics**

- Charting theory and nature of boundary and initial conditions
- Explicit solutions, including original Black-Scholes formula
- Special problems arising when there are free boundaries

#### **Key Questions in Financial PDEs**

- Physical interpretation of the equations?
- Mathematical properties of the solution?
- Techniques for obtaining explicit solutions?

#### PDEs in finance:

- Fundamental equations (e.g., Black-Scholes)
- Linear vs. nonlinear problems

#### Considerations for PDEs in Finance

- 1. Does the equation make sense as a well-posed problem?
  - Appropriate boundary or initial/final conditions?
  - Nature of the mathematical problem?
  - Smooth or discontinuous solutions?
- 2. Can we develop analytical tools to solve the equation?
- 3. How should we solve the equation numerically if necessary?

### First Order Linear PDEs

#### First Order Linear PDE

Consider the equation:

$$\alpha(s,t)\frac{\partial u}{\partial s} + \beta(s,t)\frac{\partial u}{\partial t} = \gamma(s,t)u(t)$$

**Linearity Property** 

If  $u_1, u_2$  are solutions, then  $c_1u_1 + c_2u_2$  is also a solution for constants  $c_1, c_2$ .

#### Proof.

Substitute  $u = c_1 u_1 + c_2 u_2$  into the PDE and use linearity of derivatives.

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#### **Constant Coefficient Case**

When  $\alpha, \beta$  are constants and  $\gamma \equiv 0$ :

$$\alpha_0 \frac{\partial u}{\partial s} + \beta_0 \frac{\partial u}{\partial t} = 0 \quad (4.2)$$

Define the vector  $\vec{v} = (\alpha_0, \beta_0)$  with  $\beta_0 \alpha_0 \neq 0$ .

The directional derivative:

$$\nabla_{\vec{v}}u = \langle \alpha_0, \beta_0 \rangle \cdot \langle \frac{\partial u}{\partial s}, \frac{\partial u}{\partial t} \rangle$$

Thus,  $(4.2) \Leftrightarrow \nabla_{\vec{v}} u \equiv 0$ 

#### **Characteristic Coordinates**

Introduce new coordinates:

$$\xi = \beta_0 S + \alpha_0 t, \quad \zeta = \beta_0 S - \alpha_0 t$$

The Jacobian determinant:

$$\begin{vmatrix} \beta_0 & \alpha_0 \\ \beta_0 & -\alpha_0 \end{vmatrix} = -2\alpha_0\beta_0 \neq 0$$

Let's solve the system for S, t:

$$S = \frac{1}{2\beta_0} (\xi + \zeta) \tag{1}$$

$$S = \frac{1}{2\beta_0} (\xi + \zeta)$$

$$t = \frac{1}{2\alpha_0} (\xi - \zeta)$$
(1)

Then

$$\frac{\partial S}{\partial \xi} = \frac{1}{2\beta_0} \tag{3}$$

$$\frac{\partial t}{\partial \xi} = \frac{1}{2\alpha_0} \tag{4}$$

## Computing $\frac{\partial u}{\partial \xi}$ :

$$\frac{\partial u}{\partial \xi} = \frac{\partial u}{\partial S} \frac{\partial S}{\partial \xi} + \frac{\partial u}{\partial t} \frac{\partial t}{\partial \xi}$$
 (5)

$$=\frac{1}{2\beta_0}\frac{\partial u}{\partial S} + \frac{1}{2\alpha_0}\frac{\partial u}{\partial t} \tag{6}$$

$$=\frac{1}{2\beta_0\alpha_0}\left(\alpha_0\frac{\partial u}{\partial S}+\beta_0\frac{\partial u}{\partial t}\right) \tag{7}$$

$$=0 (8)$$

Therefore, u depends only on  $\zeta$ :  $u = F(\zeta)$ 

## **Characteristics and Information Propagation**

- $\xi = \beta_0 s \alpha_0 t$  is called a **characteristic**
- Characteristics represent directions in which information propagates
- Analogous to constants of integration in ODEs

#### **Important Note**

Without boundary conditions, the solution remains arbitrary (like the constant  $\mathcal{C}$  in ODEs)

### Second-Order Linear PDEs

#### General Second-Order Linear PDE

Consider the general form in two variables:

$$a(x,t)u_{xx} + 2b(x,t)u_{xt} + c(x,t)u_{tt} + d(x,t)u_x + e(x,t)u_t + f(x,t)u = g(x,t)$$
(9)

where  $a,b,c,d,e,f,g:\mathbb{R}\times\mathbb{R}\to\mathbb{R}$  are coefficient functions.

#### **Principal Part**

The highest-order terms form the principal part:

$$L_p(u) = au_{xx} + 2bu_{xt} + cu_{tt}$$

#### Classification of PDEs

Define the discriminant:

$$\Delta(L) = b^2 - ac : (x, t) \mapsto b(x, t)^2 - a(x, t)c(x, t)$$

The PDE is classified as:

- **Hyperbolic** if  $\Delta(L) > 0$  (e.g., wave equation)
- Parabolic if  $\Delta(L) = 0$  (e.g., heat equation)
- Elliptic if  $\Delta(L) < 0$  (e.g., Laplace equation)

#### **Coordinate Transformations**

#### **Definition**

A transformation  $(\xi, \eta) = (\xi(x, t), \eta(x, t))$  is called a change of coordinates if the Jacobian is non-zero:

$$J := \xi_x \eta_t - \xi_t \eta_x \neq 0$$

#### Theorem

The type of a linear second-order PDE in two variables is invariant under a change of coordinates.

#### **Proof of Invariance**

Under transformation  $(\xi, \eta)$ , the principal part becomes:

$$Au_{\xi\xi} + 2Bu_{\xi\eta} + Cu_{\eta\eta}$$

where:

$$A = a\xi_x^2 + 2b\xi_x\xi_t + c\xi_t^2$$

$$B = a\xi_x\eta_x + b(\xi_x\eta_t + \xi_t\eta_x) + c\xi_t\eta_t$$

$$C = a\eta_x^2 + 2b\eta_x\eta_t + c\eta_t^2$$

The new discriminant satisfies:

$$\Delta_{new} = AC - B^2 = J^2(b^2 - ac) = J^2 \Delta_{original}$$

#### **Canonical Forms**

- **Hyperbolic**:  $u_{\xi\xi} u_{\eta\eta} = \text{lower order terms (wave equation)}$
- **Parabolic**:  $u_{\xi\xi} = \text{lower order terms (heat equation)}$
- **Elliptic**:  $u_{\xi\xi} + u_{\eta\eta} = \text{lower order terms (Laplace equation)}$

#### **Examples**

- Wave equation:  $u_{tt} = c^2(u_{xx} + u_{yy})$
- Heat equation:  $u_t = k(u_{xx} + u_{yy})$
- Laplace equation:  $u_{xx} + u_{yy} = 0$

#### **Simplification Through Coordinates**

A proper linear change of coordinates can transform:

- Hyperbolic equations to  $u_{\xi\eta}=0$
- Parabolic equations to  $u_{\xi\xi}=0$
- Elliptic equations to  $u_{\xi\xi} + u_{\eta\eta} = 0$

#### **Important Note**

The classification determines:

- Nature of characteristics
- Appropriate boundary conditions
- Solution techniques

## Heat Equation Analysis

#### **Heat Equation**

The fundamental heat equation:

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}$$

Physical interpretation:

- u(x, t) represents temperature distribution along a thin insulated rod
- Models heat diffusion over time

#### **Mathematical Properties**

#### Key characteristics:

- Linear equation (superposition principle applies)
- Second order in space (x)
- First order in time (t)
- ullet Parabolic type (discriminant  $\Delta=0$ )

#### **Analytic Solutions**

Solutions are analytic functions of x:

- For any  $t^* > 0$  and  $a < x_0 < b$
- $u(x, t^*)$  can be expressed as a convergent power series in  $(x x_0)$

#### **Solution Behavior**

#### Important features:

- Smoothing property: Solutions become instantly smooth for t > 0
- Infinite propagation speed: Initial disturbances spread infinitely fast
- Maximum principle: Extreme values occur only at initial/boundary conditions

# Financial Interpretation Analogous to:

- · Diffusion of volatility in option pricing
- Smoothing of price shocks over time

**Initial Value Problems with** 

**Boundary Conditions** 

#### **Dirichlet Boundary Value Problem**

Heat equation on finite interval:

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}, \quad -L < x < L, \ t > 0$$

With conditions:

- Initial condition:  $u(x,0) = u_0(x)$
- Boundary conditions:
  - $u(-L, t) = g_1(t)$  (left end temperature)
  - $u(L, t) = g_2(t)$  (right end temperature)

#### **Physical Interpretation**

Models heat diffusion in a finite rod with controlled end temperatures

#### **Neumann Boundary Value Problem**

Same heat equation with flux conditions:

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}, \quad -L < x < L, \ t > 0$$

With conditions:

- Initial condition:  $u(x,0) = u_0(x)$
- Flux boundary conditions:
  - $-\frac{\partial u}{\partial x}(-L,t) = h_1(t)$  (left end flux)
  - $\frac{\partial u}{\partial x}(L,t) = h_2(t)$  (right end flux)

#### **Physical Interpretation**

Models heat diffusion with controlled heat flow at boundaries

#### Infinite Domain Problem

Taking  $L \to \infty$ :

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2}, \quad x \in \mathbb{R}, \ t > 0$$

With initial condition:

$$u(x,0)=u_0(x)$$

#### **Growth Conditions**

For well-posedness:

- $u_0(x)$  piecewise continuous (finite jumps allowed)
- $\lim_{|x|\to\infty} \frac{u_0(x)}{e^{ax^2}} = 0 \quad \forall a > 0$
- $\lim_{|x|\to\infty} \frac{u(x,t)}{e^{ax^2}} = 0 \quad \forall a, t > 0$

#### Semi-Infinite Domain with Mixed Conditions

#### Example problem:

$$u_t = u_{xx}, \quad 0 < x < \infty, \ t > 0$$

#### With conditions:

- Initial condition:  $u(x,0) = u_0(x)$  (sufficiently smooth)
- Boundary condition: u(0, t) = g(t)
- Growth conditions at infinity

#### **Well-Posedness**

The problem is well-posed when:

- Initial/boundary conditions are compatible
- Growth conditions are satisfied
- Solution exists, is unique, and depends continuously on data

# Backward Heat Equation

#### **Time Reversal Transformation**

Consider the time transformation:

$$t = T - \tau$$
, T constant

The time derivative transforms as:

$$\frac{\partial u}{\partial \tau} = \frac{\partial u}{\partial t} \frac{\partial t}{\partial \tau} = -\frac{\partial u}{\partial t}$$

Applying to the heat equation:

$$\frac{\partial u}{\partial \tau} = -\frac{\partial^2 u}{\partial x^2}$$
 (Backward Heat Equation)

#### **Important Property**

The backward heat equation is generally **ill-posed** for most initial conditions

#### **Fundamental Solution**

The fundamental solution to the (forward) heat equation:

$$u_g(x,t) = \begin{cases} \frac{1}{2\sqrt{\pi t}} e^{-x^2/(4t)} & t > 0\\ \delta_0(x) & t = 0 \end{cases}$$

Consider the variant:

$$u(x,t) = \frac{1}{2\sqrt{\pi(\tau-t)}}e^{-\frac{x^2}{4(\tau-t)}}$$

- Solves the backward heat equation
- Initial condition:  $u_0(x) = \frac{1}{2\sqrt{\pi\tau}}e^{-x^2/(4\tau)}$
- Blows up at  $t = \tau$  (singularity)

#### **Final Value Problem**

The backward heat equation with final condition:

$$u_T(x) = u_g(x, \tau)$$

#### **Well-Posed Case**

This specific final value problem is well-posed when:

- ullet Properly transformed to a forward problem via  $t\mapsto T-t$
- The final condition decays sufficiently fast
- Solution exists in a limited time domain  $t < \tau$

#### **Example**

In finance, appears in:

- Option pricing with terminal payoff conditions
- Implicit finite difference methods

#### **Mathematical Implications**

Key observations about the backward heat equation:

- III-posedness: Small changes in final conditions can lead to large solution changes
- Regularization needed for practical applications
- Time reversal requires special treatment

#### Numerical Challenges Requires:

- Special discretization schemes
- Stability analysis
- Often needs additional constraints