

## MEAN-FIELD THEORY

The mean-field Nambu-Gor'kov Hamiltonian is

$$H = \frac{1}{2} \iint d^3\mathbf{r} d^3\mathbf{r}' \widehat{\Psi}^\dagger(\mathbf{r}') \begin{pmatrix} \xi(\mathbf{r})\delta(\mathbf{r}-\mathbf{r}')\hat{1} & \widehat{\Delta}(\mathbf{r}, \mathbf{r}') \\ \widehat{\Delta}^\dagger(\mathbf{r}, \mathbf{r}') & -\xi(\mathbf{r})\delta(\mathbf{r}-\mathbf{r}')\hat{1} \end{pmatrix} \widehat{\Psi}(\mathbf{r}) \quad (1)$$

$$\equiv \frac{1}{2} \iint d^3\mathbf{r} d^3\mathbf{r}' \widehat{\Psi}^\dagger(\mathbf{r}') \widehat{\tau}_3 \left( \xi(\mathbf{r})\delta(\mathbf{r}-\mathbf{r}') + \widehat{\Delta}(\mathbf{r}, \mathbf{r}') \right) \widehat{\Psi}(\mathbf{r}). \quad (2)$$

where

$$\widehat{\Delta}(\mathbf{r}, \mathbf{r}') = \begin{pmatrix} 0 & \widehat{\Delta}(\mathbf{r}, \mathbf{r}') \\ -\widehat{\Delta}^\dagger(\mathbf{r}, \mathbf{r}') & 0 \end{pmatrix} \quad (3)$$

When the system is homogeneous, the two-point gap function only depends on the relative coordinate, i.e.,  $\widehat{\Delta}(\mathbf{r}, \mathbf{r}') = \widehat{\Delta}(\mathbf{r} - \mathbf{r}')$ . In this case, we can transform to momentum space, and the Hamiltonian becomes

$$H = \frac{1}{2} \sum_{\mathbf{p}} \widehat{\Psi}_{\mathbf{p}}^\dagger \widehat{\tau}_3 (\xi_{\mathbf{p}} + \widehat{\Delta}(\mathbf{p})) \widehat{\Psi}_{\mathbf{p}}. \quad (4)$$

where  $\widehat{\Psi}_{\mathbf{p}}^\dagger = (c_{\mathbf{p}\uparrow}^\dagger, c_{\mathbf{p}\downarrow}^\dagger, c_{-\mathbf{p}\uparrow}, c_{-\mathbf{p}\downarrow})$ .

## FINITE TEMPERATURE

The imaginary time Green's function is

$$\widehat{G}(x, x') = - \left\langle T_\tau \widehat{\Psi}(x) \widehat{\Psi}^\dagger(x') \right\rangle. \quad (5)$$

where  $x = (\mathbf{r}, \tau)$ . The unperturbed Matsubara Green's function is

$$\widehat{G}_0(\mathbf{p}, i\epsilon_n) = (i\epsilon_n - \widehat{\tau}_3 \xi_{\mathbf{p}})^{-1} = -\frac{i\epsilon_n + \widehat{\tau}_3 \xi_{\mathbf{p}}}{\epsilon_n^2 + \xi_{\mathbf{p}}^2}. \quad (6)$$

The Gorkov Green's function satisfies

$$\widehat{G}^{-1}(\mathbf{p}, i\epsilon_n) = \widehat{G}_0^{-1}(\mathbf{p}, i\epsilon_n) - \widehat{\tau}_3 \widehat{\Delta}(\mathbf{p}) = i\epsilon_n - \widehat{\tau}_3 \xi_{\mathbf{p}} - \widehat{\tau}_3 \widehat{\Delta}(\mathbf{p}). \quad (7)$$

which gives

$$\widehat{G}(\mathbf{p}, i\epsilon_n) = -\frac{i\epsilon_n + \widehat{\tau}_3 \xi_{\mathbf{p}} + \widehat{\tau}_3 \widehat{\Delta}(\mathbf{p})}{\epsilon_n^2 + \xi_{\mathbf{p}}^2 + |\Delta|^2} \quad (8)$$

The quasiclassical Green's function is

$$\widehat{g}(\mathbf{p}, i\epsilon_n) = \int d\xi_p \widehat{G}(\mathbf{p}, i\epsilon_n) \widehat{\tau}_3 = -\pi \frac{i\epsilon_n \widehat{\tau}_3 - \widehat{\Delta}(\mathbf{p})}{\sqrt{|\Delta|^2 + \epsilon_n^2}} = \frac{\pi}{\sqrt{|\Delta|^2 + \epsilon_n^2}} \begin{pmatrix} -i\epsilon_n & \widehat{\Delta} \\ -\widehat{\Delta}^\dagger & i\epsilon_n \end{pmatrix} \equiv \begin{pmatrix} \widehat{g} & \widehat{f} \\ -\widehat{f}^\dagger & -\widehat{g} \end{pmatrix} \quad (9)$$

and it satisfies the normalization condition

$$\widehat{g}(\mathbf{p}, i\epsilon_n)^2 = -\pi^2 \widehat{1}. \quad (10)$$

The bulk quasiclassical Green's function satisfies

$$[i\epsilon_n \widehat{\tau}_3 - \widehat{\Delta}, \widehat{g}] = 0 \quad (11)$$

But when near the boundary or impurities, the quasiclassical Green's function  $\widehat{g}(\mathbf{r}, \mathbf{p}; \epsilon)$  and the gap function  $\widehat{\Delta}(\mathbf{r}, \mathbf{p})$  are inhomogeneous, and are described by the Eilenberger equation and gap equation.

$$[\epsilon \widehat{\tau}_3 - \widehat{\Delta}(\mathbf{r}, \mathbf{p}), \widehat{g}(\mathbf{r}, \mathbf{p}; \epsilon)] + i\hbar \mathbf{v}_p \cdot \nabla \widehat{g}(\mathbf{r}, \mathbf{p}; \epsilon) = 0 \quad (12)$$

and

$$\widehat{\Delta}(\mathbf{r}, \mathbf{p}) = \left\langle v(\mathbf{p}, \mathbf{p}') T \sum_{\epsilon_n=-\omega_c}^{\omega_c} \widehat{f}(\mathbf{r}, \mathbf{p}', \epsilon_n) \right\rangle_{\mathbf{p}'} \quad (13)$$

### **${}^3\text{HE-A}$**

For  ${}^3\text{He-A}$ , the spin structure is  $\hat{\mathbf{d}} \cdot (i\vec{\sigma}\hat{\sigma}_y) = \hat{\sigma}_x$ , and we have two components of the gap function,

$$\widehat{\Delta}(\mathbf{r}, \mathbf{p}) = \hat{\sigma}_x (\Delta_1(\mathbf{r}) p_x + \Delta_2(\mathbf{r}) p_y) \quad (14)$$

In bulk we have  $\widehat{\Delta}(\mathbf{r}, \mathbf{p}) = \hat{\sigma}_x \Delta(p_x + ip_y)$ .

### **S-WAVE PAIRING**

For s-wave pairing, the spin structure is  $i\hat{\sigma}_y$ , and we have

$$\widehat{\Delta}(\mathbf{r}, \mathbf{p}) = i\hat{\sigma}_y \Delta(\mathbf{r}) \quad (15)$$

In bulk we have  $\widehat{\Delta}(\mathbf{r}, \mathbf{p}) = i\hat{\sigma}_y \Delta$ .

## EDGE GAP PROFILE

We assume translational invariance along the edge, i.e. the  $y$  direction. We also assume the gap to be real(imaginary) along the  $p_x(p_y)$  direction.

$$\hat{\Delta}(x, \mathbf{p}) = \hat{\sigma}_x(\Delta_1(x)p_x + i\Delta_2(x)p_y) \quad (16)$$

Here  $\Delta_1(x)$  and  $\Delta_2(x)$  are real functions, and we have

$$\hat{\Delta}(\mathbf{r}, \mathbf{p}) = \begin{pmatrix} 0 & \hat{\sigma}_x(\Delta_1(x)p_x + i\Delta_2(x)p_y) \\ -\hat{\sigma}_x(\Delta_1(x)p_x - i\Delta_2(x)p_y) & 0 \end{pmatrix} \quad (17)$$

$$= i\hat{\sigma}_x(\Delta_2 p_y \hat{\tau}_1 + \Delta_1 p_x \hat{\tau}_2) \quad (18)$$

We can also write the anomalous Green's function as

$$\hat{f}(\mathbf{r}, \mathbf{p}, \epsilon_n) = \hat{\sigma}_x \left( f_1(\mathbf{r}, \mathbf{p}, \epsilon_n) + i f_2(\mathbf{r}, \mathbf{p}, \epsilon_n) \right) \quad (19)$$

and we have

$$\hat{g}(\mathbf{r}, \mathbf{p}, \epsilon_n) = \hat{g}\hat{\tau}_3 + i\hat{\sigma}_x(f_2\hat{\tau}_1 + f_1\hat{\tau}_2) \quad (20)$$

## NUMERICAL SOLUTION

The (1, 1) and (1, 2) elements of the Eilenberger equation are

$$i\epsilon_n \hat{g} + \hat{\Delta} \hat{f}^\dagger + i\hbar \mathbf{v}_p \cdot \nabla \hat{g} = 0 \quad (21)$$

$$i\epsilon_n \hat{f} + \hat{\Delta} \hat{g} + i\hbar \mathbf{v}_p \cdot \nabla \hat{f} = 0 \quad (22)$$

Note that  $\hat{\Delta}$  and  $\hat{f}$  has the same spin structure, so we can substitute in  $\hat{g}$  and  $\hat{f}$ , divide by  $T_c$ , and we get

$$i(2n+1)\pi t g(\mathbf{r}, \hat{\mathbf{p}}, n) + \delta(\mathbf{r}, \hat{\mathbf{p}}) f^*(\mathbf{r}, \hat{\mathbf{p}}, n) + i(\hat{\mathbf{p}}_x \partial_x + \hat{\mathbf{p}}_y \partial_y) g(\mathbf{r}, \hat{\mathbf{p}}, n) = 0 \quad (23)$$

$$i(2n+1)\pi t f(\mathbf{r}, \hat{\mathbf{p}}, n) + \delta(\mathbf{r}, \hat{\mathbf{p}}) g(\mathbf{r}, \hat{\mathbf{p}}, n) + i(\hat{\mathbf{p}}_x \partial_x + \hat{\mathbf{p}}_y \partial_y) f(\mathbf{r}, \hat{\mathbf{p}}, n) = 0 \quad (24)$$

where  $t = T/T_c$ ,  $\delta = \Delta/T_c$ ,  $\xi = \hbar v_f/T_c$  and  $\vec{\partial} = \xi \nabla$ .

For the gap equation, we ...

$$\delta = \dots \quad (25)$$

We first assume an initial guess for the gap function  $\delta(\mathbf{r}, \hat{\mathbf{p}})$ . For each trajectory  $\hat{\mathbf{p}}$  along which the quasiparticle propagates, we can solve the Eilenberger equation to get the Green's function  $g(\mathbf{r}, \hat{\mathbf{p}}, n)$  and the anomalous Green's function  $f(\mathbf{r}, \hat{\mathbf{p}}, n)$ . Then we iterate to update the gap function  $\delta(\mathbf{r}, \hat{\mathbf{p}})$  by solving the gap equation. Repeat until convergence.