

Electron detachment from alkali-metal negative ions by electron collisions

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Abstract. The Bethe–Born and Born–Ochkur approximations have been used to calculate the cross sections for electron detachment from the alkali metal negative ions Li^- , Na^- , K^- , Rb^- and Cs^- .

1. Introduction

Interest in reactions involving alkali metal negative ions has arisen from work on alkali metal vapours and gaseous plasmas. There is little information on the collisional detachment from alkali metal negative ions by electron impact, and the purpose of the present paper is to provide estimates of the cross sections. We use a simplified version of the Born–Ochkur approximation. This approach treats the incident and scattered electrons as plane waves but allows for distortion for the detached electron. This method has proved to be reliable in the case of the negative hydrogen ion (John and Williams 1973), their results were in good agreement with more elaborate calculations and experiment.

A discussion of the method is given in § 2 and the results in § 3.

2. A simplified version of the Born–Ochkur approximation

The cross section for the collisional ionization of an alkali metal negative ion is given by the alternative formulae

$$\left. \begin{aligned} Q_1(k_1^2) &= \frac{4}{\pi k_1^2} \int_0^{\epsilon_{\max}} \int_{k_1-k_3}^{k_1+k_3} K^{-3} |I(K, k_2^2)|^2 dK dk_2^2 \\ Q_2(k_1^2) &= \frac{8}{\pi k_1^2} \int_0^{\frac{1}{2}\epsilon_{\max}} \int_{k_1-k_3}^{k_1+k_3} K^{-3} |I(K, k_2^2)|^2 dK dk_2^2 \end{aligned} \right\} \quad (\text{in units of } \pi a_0^2) \quad (1)$$

where k_1 , k_2 and k_3 are the momenta of the incident, detached and scattered electrons respectively, $K = k_3 - k_1$, γ^2 is the attachment energy of the negative ion and $\epsilon_{\max} = k_1^2 - \gamma^2$. Q_1 and Q_2 give identical results for exact calculations (see Peterkop 1961). However, differences arise in the Born approximation for which $I(K, k^2)$ is given by

$$I(K, k^2) = \int e^{iKz} \phi_{\gamma^2}^* \phi_{k^2} d\tau \quad (2)$$

ϕ_{γ^2} , ϕ_{k^2} being the bound state and continuous state wavefunctions of the negative ion.

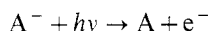
In the Bethe–Born approximation (2) is replaced by

$$|I(K, k^2)| = K \int z \phi_{\gamma^2} \phi_{k^2} d\tau \quad (3)$$

so that

$$\int_{k_1-k_3}^{k_1+k_3} K^{-3} |I(K, k^2)|^2 dK = \frac{a(k^2)}{4\pi\alpha a_0^2(\gamma^2 + k^2)} \ln \left| \frac{k_1 + k_3}{k_1 - k_3} \right| \quad (4)$$

$a(k^2)$ being the photoionization cross section for the reaction



α the fine structure constant and a_0 the Bohr radius. The photo-detachment cross section of alkali metal negative ions has been calculated (John 1972, John and Williams 1972) by the Bethe–Longmire formula

$$a(k^2) = \frac{32\pi\alpha}{3} a_0^2 \gamma N^2 \left[\frac{k}{\gamma^2 + k^2} \right]^3 \left(\cos \eta_1^+ + \frac{\gamma(\gamma^2 + 3k^2)}{2k^3} \sin \eta_1^+ \right)^2. \quad (5)$$

These cross sections are in satisfactory agreement with more elaborate calculations of Norcross and Moores (1972) for Li^- and Na^- and the Li^- measurements of Feldman (1972 private communication). The essential feature of the Bethe–Longmire method is that the wavefunction for the valence electron undergoing the transition is approximated by

$$P_{\gamma^2}(r) = \left[\frac{\gamma}{2\pi} \right]^{1/2} N \frac{e^{-\gamma r}}{r} \quad (\text{bound state})$$

$$P_{k^2}(r) = \left[\frac{3}{4\pi} \right]^{1/2} \frac{\cos \theta}{r} k^{-1/2} \left\{ \frac{\sin(kr + \eta_1^+)}{kr} - \cos(kr + \eta_1^+) \right\} \quad (\text{free state}) \quad (6)$$

N being a normalizing constant.

If we use the wavefunctions (6) to calculate $I(K, k^2)$ we get

$$|I(K, k^2)| = \left(\frac{6\gamma}{k} \right)^{1/2} N K^{-1} \int_0^\infty \frac{e^{-\gamma r}}{r} \left(\frac{\sin Kr}{Kr} - \cos Kr \right) \left(\frac{\sin(kr + \eta_1^+)}{kr} - \cos(kr + \eta_1^+) \right) dr$$

$$= \left(\frac{6\gamma}{k} \right)^{1/2} N K^{-1} (P(K, k^2) \cos \eta_1^+ + Q(K, k^2) \sin \eta_1^+). \quad (7)$$

Details of the derivation and explicit forms of P and Q are given in an appendix. The Bethe–Born version can be obtained by expanding P and Q in a series in K , neglecting terms of higher order than K^2 :

$$P(K, k^2) \simeq \frac{2K^2}{3} \frac{k^2}{(\gamma^2 + k^2)^2}$$

$$Q(K, k^2) \simeq \frac{\gamma K^2}{3k(\gamma^2 + k^2)^2} (\gamma^2 + 3k^2) \quad (8)$$

consistent with (4) and (5).

Exchange effects can be included in the Born approximation by the Ochkur correction (Ochkur 1964), the relevant Born–Ochkur expression for $|I(K, k^2)|^2$ is

$$|I(K, k^2)|^2 = 6\gamma N^2 K^{-2} k^{-1} \{P(K, k^2) \cos \eta_1^+ + Q(K, k^2) \sin \eta_1^+\}^2 \left(1 - \frac{K^2}{k_1^2} + \frac{K^4}{k_1^4}\right). \quad (9)$$

3. Discussion of results

We use (9) and (5) to calculate the Born–Ochkur and Bethe–Born versions of the detachment cross section, we find that $Q_2 \simeq 2Q_1$ at nearly all energies. From work on the detachment of H^- , we found the Born–Ochkur Q_2 gave the most accurate values and we expect the same to be true for the alkali metal negative ions. In the Bethe–Born approximation Q_1 was the most reliable form and is in agreement with the Born–Ochkur Q_2 . We adopted the same data as was used in the photodetachment calculations, that is, the electron affinities of Weiss (1961, 1966) and the phases of Burke and Taylor (1969) for e–Li scattering, Karule and Peterkop (1965) for e–Na, e–K and e–Cs and Balling (1969) for e–Rb scattering. N^2 was calculated from the oscillator strength sum rule

$$\int_0^\infty a(k^2) dk^2 = 8\pi^2 \alpha a_0^2. \quad (10)$$

Our results are given in table 1, the cross sections of the negative ions in the various approximations, apart from Rb^- , show a similar pattern; for Rb^- the Bethe–Born and Born–Ochkur Q_2 differ by about 9 per cent at high energies.

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Appendix

The integrals arising in the Born approximation are of the form

$$J = \frac{-1}{2kK} \int_0^\infty e^{-\gamma r} (\sin Kr - Kr \cos Kr) \{\sin(kr + \eta) - kr \cos(kr + \eta)\} dr \quad (A.1)$$

integrating by parts we can reduce J to integrals of the standard types:

$$\begin{aligned} & \int_0^\infty r^{-1} e^{-\gamma r} \sin Kr \sin(kr + \eta) dr \\ & \int_0^\infty r^n e^{-\gamma r} \begin{pmatrix} \sin Kr \\ \cos Kr \end{pmatrix} \begin{pmatrix} \sin(kr + \eta) \\ \cos(kr + \eta) \end{pmatrix} dr \quad (n = 0, 1) \end{aligned}$$

which can be evaluated by elementary methods, giving

$$J = P(K, k^2) \cos \eta + Q(K, k^2) \sin \eta \quad (A.2)$$

Table 1. The cross section for detachment of the alkali metal negatives by electron collisions (in units πa_0^2)

Lithium			Sodium			Potassium			Caesium			Rubidium		
γ^2 (Ryd 13.6 eV) 0.0452			0.0396			0.0344			0.0287			0.0309		
N^2			1.839			2.550			1.176			0.879		
k_1^2 (eV)	Born-Ochkur Q_1	Bethe-Born Q_1	Born-Ochkur Q_2	Bethe-Born Q_1	Born-Ochkur Q_2	Born-Ochkur Q_1	Bethe-Born Q_1	Born-Ochkur Q_2	Born-Ochkur Q_1	Bethe-Born Q_1	Born-Ochkur Q_2	Born-Ochkur Q_1	Bethe-Born Q_1	Born-Ochkur Q_2
3.3 198	353	479	—	—	—	265	—	—	573	—	—	3.7	—	—
6.1 154	305	359	—	—	—	—	—	—	—	—	—	3.8	805	1583
11.5 106	213	237	518	—	—	234	451	—	498	972	—	5.6	—	—
27.8 59.3	118	124	—	—	—	—	—	—	—	—	—	7.2	504	998
55 35.1	70.1	72.7	—	—	—	210	—	—	438	—	—	9.25	—	—
109.4 20.3	40.5	41.2	—	—	—	—	—	—	—	—	—	10.6	372	—
136 16.9	33.8	34.7	—	—	—	—	—	—	—	—	—	—	—	—
190.4 12.8	25.5	26.2	—	—	—	187	—	—	384	—	—	14	299	596
299.2 8.7	17.4	17.8	—	—	—	166	—	—	337	—	—	20.8	—	—
408 6.7	13.4	13.7	—	—	—	135	271	—	272	544	—	27.2	174	—
489.6 5.7	11.4	11.7	—	—	—	125	230	—	230	459	—	27.6	—	—
571.2 5.0	10.0	10.2	—	—	—	81	162	—	159	318	—	54.4	97.4	195
652.8 4.5	9.0	9.1	—	—	—	68.2	136	—	133	267	—	81.6	69.0	138
734.4 4.0	8.0	8.2	—	—	—	39.3	78.6	—	75.8	152	—	136	44.5	88.9
816 3.7	7.4	7.5	—	—	—	28.3	56.5	—	54.1	108	—	108.8	53.9	108
897.6 3.4	6.8	6.9	—	—	—	22.3	44.5	—	42.4	84.9	—	190.4	33.2	66.4
979.2 3.2	6.3	6.4	—	—	—	18.5	36.9	—	35.6	70.3	—	299.2	22.4	44.8
						13.9	27.8	—	26.4	52.7	—	408	17.0	34.0
						9.5	18.9	—	17.8	35.7	—	489.6	14.5	29.0
						7.3	14.5	—	13.7	27.3	—	571.2	12.7	25.3
						6.2	12.4	—	11.6	23.3	—	652.8	11.2	22.4
						5.4	10.8	—	10.2	20.3	—	734.4	10.1	20.2
						4.8	9.6	—	9.0	18.1	—	816	9.2	18.4
						4.4	8.7	—	8.2	16.3	—	897.6	8.5	17.0
						3.9	7.9	—	7.4	14.9	—	979.2	7.8	15.6
						3.7	7.3	—	6.8	13.7	—			
						3.4	6.8	—	6.3	12.6	—			

where

$$\begin{aligned} P(K, k^2) &= \left(\frac{\gamma^2 + k^2 + K^2}{8kK} \right) \ln \left(\frac{\gamma^2 + (k+K)^2}{\gamma^2 + (k-K)^2} \right) - \frac{1}{2} \\ Q(K, k^2) &= \left(\frac{\gamma^2 + k^2 + K^2}{4kK} \right) \left\{ \tan^{-1} \left(\frac{k+K}{\gamma} \right) - \tan^{-1} \left(\frac{k-K}{\gamma} \right) \right\} - \frac{\gamma}{2k}. \end{aligned} \quad (\text{A.2})$$

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