

## 20. Approximating functions globally: Periodic functions and Fourier analysis

## Summary of the previous class

- Numerical methods for ODEs
- Taylor methods: calculate the Taylor series of the exact solution
- Adaptivity: Choosing the step size to control the error

## Goals for today

- Approximating functions globally
- Fourier series for periodic functions
- Rate of decay of Fourier coefficients
- Trigonometric interpolation

## Manipulating functions

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  - Derivative:  $f'$
  - Definite integral:  $\int_a^b f$
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  - Derivative:  $f'$
  - Definite integral:  $\int_a^b f$
  - Find roots of  $f$  in  $[a, b]$
- Idea: Manipulate an **approximation** of  $f$  instead

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- Taylor series:
  - A polynomial that reproduces  $f$  locally near a point
  - Works well only near that point
- Lagrange interpolation:
  - Expensive to calculate and manipulate
  - Need to be careful about location of interpolation points

# Collaboration I

## Representing functions globally

- 1 What other way could we represent a function on an interval?
- 2 If  $f$  is a  $2\pi$ -periodic function, i.e.  $f(t + 2\pi) = f(t) \quad \forall t$ , how can we represent it?
- 3 What do we need to calculate in order to do so?

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- We will do “linear algebra with functions” – i.e. **functional analysis**

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- $T$  is the **period** of function



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- Alternatively

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} [a_n \cos(nx) + b_n \sin(nx)]$$

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- Can *any* periodic function can be expressed like this?
- We need to specify which **space of functions** we allow

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- Infinite-dimensional generalization of Euclidean space
- It can be shown that the  $\phi_n$  form a **complete basis** of  $L^2$ :
- Eigenfunctions of Laplacian operator  $\frac{\partial^2}{\partial x^2}$
- Sturm–Liouville theory
- Spectral theorem for compact operators

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- Provided the basis vectors are **orthogonal**
- We need to be able to talk about orthogonality of *functions*

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- For  $z = x + iy \in \mathbb{C}$ , the **complex conjugate** is  $\bar{z} := x - iy$
- Properties:  $(f, g)$  is linear in  $g$ , and  $(f, f) \geq 0$
- Two functions are **orthogonal** if  $(f, g) = 0$

## Orthogonality

- Let's look at inner products of our basis functions:

$$\begin{aligned}
 (\phi_n, \phi_m) &= \int_0^{2\pi} e^{-i n x} e^{i m x} dx \\
 &= \int_0^{2\pi} e^{i(m-n)x} dx \\
 &= \frac{1}{i(m-n)} \left[ e^{i(m-n)x} \right]_{x=0}^{2\pi}
 \end{aligned}$$

- So  $(\phi_n, \phi_m) = 0$  for  $m \neq n$  and  $= 2\pi$  for  $m = n$
- Hence the basis functions  $(\phi_n)$  are orthogonal

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- Thus

$$\hat{f}_m = \frac{1}{2\pi} \int_0^{2\pi} e^{-imx} f(x) dx$$

## Representation on a computer

- The natural idea to represent a periodic function on a computer is then to truncate to a *finite* sum

$$f \simeq f_N := \sum_{n=-N}^N \hat{f}(n) \phi_n$$



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- It depends on *how fast* the  $\hat{f}(n)$  decay as  $n \rightarrow \infty$

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- Integrate by parts to get

$$\widehat{f'}_n = in \widehat{f}_n$$

## Rate of decay of Fourier coefficients

- Can show: If  $f$  is integrable, e.g. continuous, on an interval, then

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- Hence  $\hat{f}_n = o(n^{-k})$

## Rate of decay II

- If  $f$  is **smooth** ( $C^\infty$ ) then decays faster than any polynomial
- In fact, if  $f$  is analytic in a suitable region then

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- This is called **spectral convergence**

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- We get an approximation of our general periodic function to  $\sim$  machine epsilon!
- **Implementations:**
  - The **Chebfun** package in Matlab
  - The **ApproxFun.jl** package in Julia

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# Trigonometric interpolation

- Calculating Fourier integrals numerically is not too nice
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- How else could we look for a Fourier series that approximates  $f$ ?
- Return to an idea from earlier in the course: **interpolation**
- Let's try to **interpolate**  $f$  using trigonometric functions

## Trigonometric interpolation II

- Recall the setting of interpolation:

*Given  $x_j$  and  $f_j = f(x_j)$  for  $j = 0, \dots, N$ , find  $g$  (in some class) such that  $g(x_j) = f_j$*

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- In the context of periodic functions we take  $x_N = x_0$
- We will look for trigonometric functions

$$g = \sum_{k=0}^{N-1} g_k \phi_k$$

- For technical reasons restrict to odd  $N$

# Interpolation II

- Condition for interpolation:

$$g(x_j) = f_j$$

- So interpolate  $N$  points  $(x_j, f_j)$  with a sum of basis functions with  $N$  unknown coefficients  $g_n$ :

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- Which points should we interpolate in?
- For polynomials, equally-spaced points were **bad**

# Interpolation III

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- Take  $N$  **equally-spaced** points  $x_j := jh$  with  $h := \frac{2\pi}{N}$
- The interpolation condition becomes

$$\sum_{k=0}^{N-1} g_k \exp(ikjh) = f_j$$

## Collaboration II

- 1 What kind of equation do we have?
- 2 How can we solve it?
- 3 What number of operations do we expect to need to do so?

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- When is this easy to solve?

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- Hence

$$\mathbf{g} = \frac{1}{N} M^* \mathbf{f}$$

# Discrete Fourier transform (DFT)

- The **Discrete Fourier transform**  $\mathcal{F}$  maps *from* function values *to* Fourier coefficients:

$$g_k = \frac{1}{N} \sum_j e^{-ijk \frac{2\pi}{N}} f_j$$

- The **Inverse Discrete Fourier Transform** maps *from* Fourier coefficients *to* function values:

$$f_j = \sum_k e^{ijk \frac{2\pi}{N}} g_k$$



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- It looks like we need to do a matrix–vector multiplication to calculate the DFT –  $\mathcal{O}(n^2)$  operations
- One of the great algorithmic discoveries of the 19th/20th century: **Fast Fourier Transform** (FFT)
- This uses *structure* in  $M$  to calculate in  $\mathcal{O}(n \log n)$  time
- Has myriad applications throughout applied mathematics, signal processing, physics, ...

# Summary

- Periodic functions may be expanded in Fourier series
- Infinite linear combination of orthogonal basis functions
- The Fourier coefficients  $\hat{f}_n$  decay at a rate determined by the smoothness of  $f$
- Trigonometric interpolation corresponds to a discrete Fourier transform