

CHAPTER 4

THE CONTINUOUS-TIME FOURIER
TRANSFORM

0 Introduction

Representation of Aperiodic Signals: The Continuous-Time Fourier Transform

The Fourier Transform for Periodic Signals

Properties of The Continuous-Time Fourier Transform

The Convolution Property

The Multiplication Property

Systems Characterized By Linear Constant Coefficient Differential Equations

7 Frequency-Selective Filters

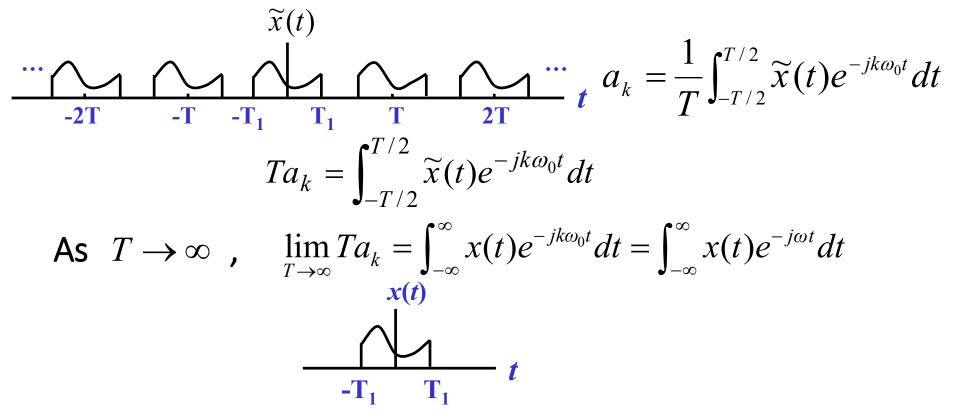
Transmission Without Distortion

9 Sampling

4.0 Introduction

- Representation of continuous-time aperiodic signals as linear combination of complex exponentials — Inverse Fourier Transform
- Frequency spectrum of aperiodic signals Fourier Transform
- > Applications of continuous-time Fourier transform

4.1.1 Fourier Transform and Inverse Fourier Transform



Use $X(j\omega)$ to denote this integral, then we have:

Frequency
$$X(j\omega) = \int_{-\infty}^{\infty} x(t)e^{-j\omega t}dt$$
 Fourier transform of $X(t)$

Since
$$X(j\omega) = \lim_{T \to \infty} Ta_k = \lim_{\omega_0 \to 0} 2\pi \frac{a_k}{\omega_0}$$
,

 $X(j\omega)$ is actually spectrum-density function(频谱密度函数).

$$\tilde{x}(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} = \sum_{k=-\infty}^{\infty} \left[\frac{1}{T} \int_{-T/2}^{T/2} \tilde{x}(t) e^{-jk\omega_0 t} dt \right] e^{jk\omega_0 t}$$

$$= \frac{\omega_0}{2\pi} \sum_{k=-\infty}^{\infty} \left[\int_{-T/2}^{T/2} \tilde{x}(t) e^{-jk\omega_0 t} dt \right] e^{jk\omega_0 t}$$

As
$$T o \infty$$
 ,

As
$$T \to \infty$$
, $x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\omega) e^{j\omega t} d\omega$ — Inverse Fourier transform

analysis equation: $X(j\omega) = \int_{-\infty}^{\infty} x(t)e^{-j\omega t}dt$

synthesis equation:
$$x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\omega) e^{j\omega t} d\omega$$

transform

Fourier Transform Pair

Comparing the synthesis equations in:

CTFS:
$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t}$$
 $CTFT$: $x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\omega) e^{j\omega t} d\omega$

$$a_k \Leftrightarrow \frac{X(j\omega)d\omega}{2\pi}$$

This fact means: Aperiodic signals can also be decomposed as linear combination of infinite numbers of complex exponentials, which occur at a continuum of frequencies, but have amplitudes infinitesimally small — approaching zero!

 $|X(j\omega)|$ indicate the relative amplitudes of all components, and angle $\angle X(j\omega)$ indicate the phases of all components.

An useful relationship: $a_k = \frac{1}{T}X(j\omega)\big|_{\omega=k\omega_0}$ Where, $\tilde{x}(t) \stackrel{FS}{\longleftrightarrow} a_k$, $x(t) \stackrel{FT}{\longleftrightarrow} X(j\omega)$.

This relationship is also valid for aperiodic signals with unlimited duration!

- Convergence of Fourier Transforms
- Dirichlet conditions:
- 1. x(t) is absolutely integrable; that is $\int_{-\infty}^{\infty} |x(t)| dt < \infty$
- 2. x(t) have a finite number of maxima and minima within any finite interval.
- 3. x(t) have a finite number of discontinuities within any finite interval.

Furthermore, each of these discontinuities must be finite.

If impulse functions are permitted in the transform, some signals which are not absolutely integrable over an infinite interval, can also be considered to have Fourier transforms. This will be convenient in the discussion of Fourier methods.

4.1.2 Examples

Example 4.1

Consider the signal $x(t) = e^{-\alpha t}u(t)$ $\alpha > 0$.

$$X(j\omega) = \int_{0}^{\infty} e^{-\alpha t} e^{-j\omega t} dt = -\frac{1}{\alpha + j\omega} e^{-(\alpha + j\omega)t} \Big|_{0}^{\infty} = \frac{1}{\alpha + j\omega}, \quad \alpha > 0$$

$$|X(j\omega)| = \frac{1}{\sqrt{\alpha^2 + \omega^2}}, \quad \angle X(j\omega) = -\tan^{-1}\left(\frac{\omega}{\alpha}\right)$$

$$\angle X(j\omega)|_{1/\alpha}$$

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$$-\frac{1}{\alpha}$$

$$= \frac{1}{\alpha + j\omega}, \quad \alpha > 0$$

$$\angle X(j\omega)|_{1/\alpha}$$

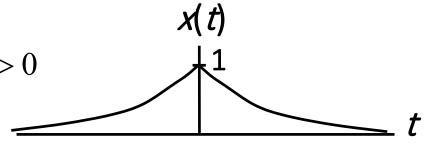
$$= \frac{\pi/2}{-\pi/4}$$

$$= \frac{\pi}{-\pi/4}$$

$$= \frac{\pi}{-\pi/4}$$

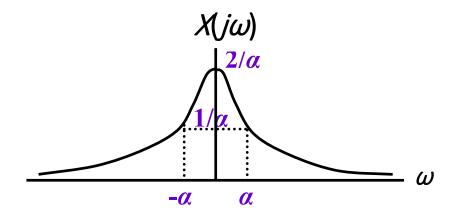
Example 4.2

Consider the signal $x(t) = e^{-\alpha|t|}, \alpha > 0$



Sol: The Fourier transform of the signal is

$$X(j\omega) = \int_{-\infty}^{\infty} e^{-\alpha|t|} e^{-j\omega t} dt = \int_{-\infty}^{0} e^{\alpha t} e^{-j\omega t} dt + \int_{0}^{\infty} e^{-\alpha t} e^{-j\omega t} dt$$
$$= \frac{1}{\alpha - j\omega} + \frac{1}{\alpha + j\omega} = \frac{2\alpha}{\alpha^2 + \omega^2}$$



Example 4.3

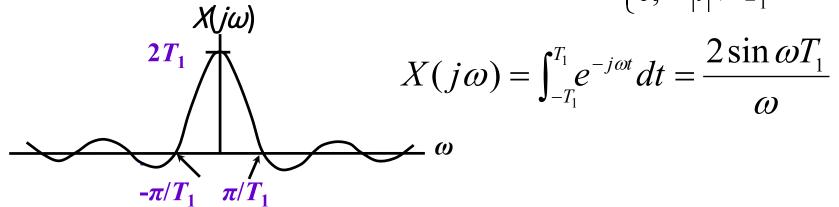
Determine the Fourier transform of the unit impulse $x(t) = \delta(t)$

$$X(j\omega) = \int_{-\infty}^{\infty} \delta(t)e^{-j\omega t}dt = 1$$

That is, the unit impulse has a Fourier transform consisting of equal contribution at all frequencies. This spectrum is referred to as white-spectrum.

Example 4.4

Consider the rectangular pulse signal $x(t) = \begin{cases} 1, & |t| < T_1 \\ 0, & |t| > T_1 \end{cases}$



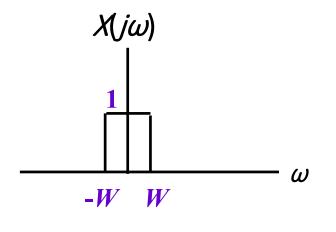
Example 4.5

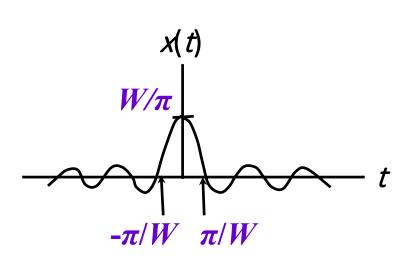
Consider the signal x(t) whose Fourier transform is

$$X(j\omega) = \begin{cases} 1, & |\omega| < W \\ 0, & |\omega| > W \end{cases}$$

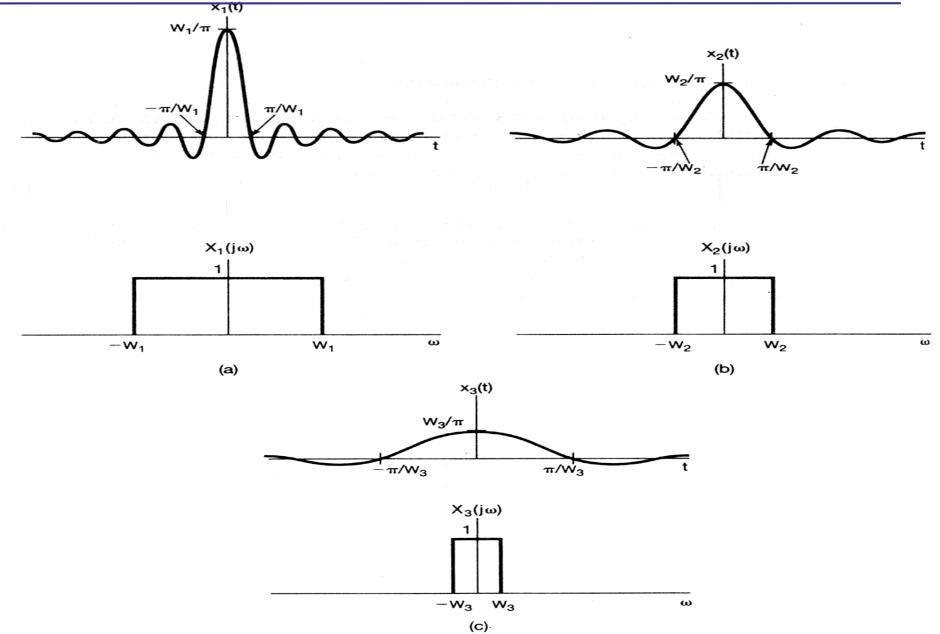
Sol: Using the synthesis equation, we can determine

$$x(t) = \frac{1}{2\pi} \int_{-W}^{W} e^{j\omega t} d\omega = \frac{\sin Wt}{\pi t}$$









$$\mathscr{F}\{e^{j\omega_0t}\}=?$$

From the analysis equation,

$$X(j\omega) = \int_{-\infty}^{\infty} e^{j\omega_0 t} e^{-j\omega t} dt = \frac{1}{j(\omega_0 - \omega)} e^{j(\omega_0 - \omega)t} \Big|_{-\infty}^{\infty} \text{Does not converge!}$$

We have obtained
$$\delta(t) \leftrightarrow 1$$
, so $\delta(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} 1 \cdot e^{j\omega t} d\omega$

$$2\pi\delta(t) = \int_{-\infty}^{\infty} 1 \cdot e^{j\omega t} d\omega$$

$$2\pi\delta(-\omega) = \int_{-\infty}^{\infty} e^{-ja\omega} da \implies 1 \xleftarrow{FT} 2\pi\delta(-\omega) = 2\pi\delta(\omega)$$

$$2\pi\delta(\omega_0 - \omega) = \int_{-\infty}^{\infty} e^{-jt(\omega - \omega_0)} dt = \int_{-\infty}^{\infty} e^{j\omega_0 t} e^{-j\omega t} dt \implies \mathcal{F}\{e^{j\omega_0 t}\} = 2\pi\delta(\omega - \omega_0)$$

The Fourier transform of the complex exponential signal $e^{j\omega_0 t}$ is an impulse located at $\omega = \omega_0$ with its area 2π .

For an arbitrary periodic signal x(t), if it can be represented by Fourier series as

 $x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t}$

We can obtain its Fourier transform as

$$X(j\omega) = \sum_{k=-\infty}^{\infty} a_k 2\pi \delta(\omega - k\omega_0) = \sum_{k=-\infty}^{\infty} 2\pi a_k \delta(\omega - k\omega_0)$$

The Fourier transform of a periodic signal with Fourier series coefficients $\{a_k\}$ can be interpreted as a train of impulses occurring at the harmonically related frequencies and for which the area of the impulse at the kth harmonic frequency $k\omega_0$ is 2π times the kth Fourier series coefficient a_k .

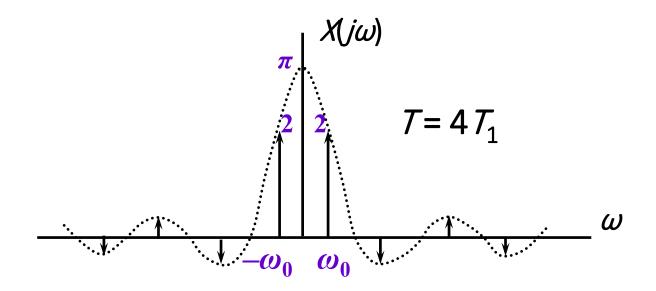
Example 4.6

Consider the periodic square wave with FS coefficients are

$$a_k = \frac{\sin k\omega_0 T_1}{k\pi}$$

Sol: From the formula its Fourier transform is

$$X(j\omega) = \sum_{k=-\infty}^{\infty} 2\pi \cdot \frac{\sin k\omega_0 T_1}{k\pi} \delta(\omega - k\omega_0) = \sum_{k=-\infty}^{\infty} \frac{2\sin k\omega_0 T_1}{k} \delta(\omega - k\omega_0)$$



Example 4.7

Find the Fourier Transforms of $x_1(t) = \sin \omega_0 t$ and $x_2(t) = \cos \omega_0 t$.

Sol: The Fourier series coefficients for $x_1(t)$ are

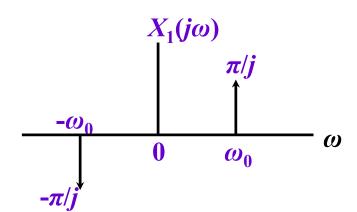
$$a_1 = \frac{1}{2i}, \qquad a_{-1} = -\frac{1}{2i}, \qquad a_k = 0, \quad k \neq \pm 1$$

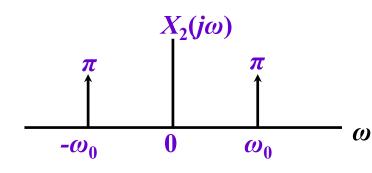
The Fourier series coefficients for $x_2(t)$ are

$$a_1 = a_{-1} = \frac{1}{2}, \qquad a_k = 0, \quad k \neq \pm 1$$

$$\sin \omega_0 t \longleftrightarrow j\pi [\delta(\omega + \omega_0) - \delta(\omega - \omega_0)]$$

$$\cos \omega_0 t \longleftrightarrow \pi \left[\delta(\omega + \omega_0) + \delta(\omega - \omega_0) \right]$$

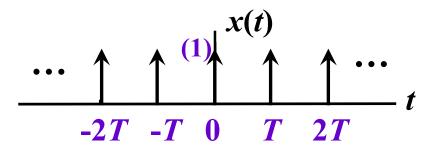




Example 4.8

Consider the impulse train

$$x(t) = \sum_{k=-\infty}^{\infty} \delta(t - kT)$$

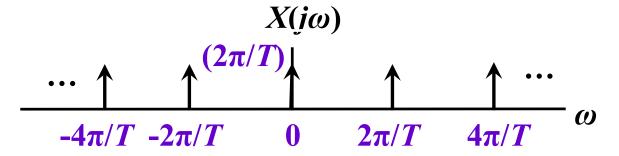


Sol: The Fourier series coefficients for this signal are

$$a_k = \frac{1}{T} \int_{-T/2}^{T/2} \delta(t) e^{-jk\omega_0 t} dt = \frac{1}{T}$$

Thus, its Fourier transform is

$$X(j\omega) = \frac{2\pi}{T} \sum_{k=-\infty}^{\infty} \delta(\omega - \frac{2k\pi}{T}) = \omega_0 \sum_{k=-\infty}^{\infty} \delta(\omega - k\omega_0)$$



4.3.1 Linearity

If
$$x(t) \stackrel{FT}{\longleftrightarrow} X(j\omega)$$
 and $y(t) \stackrel{FT}{\longleftrightarrow} Y(j\omega)$
then $ax(t) + by(t) \stackrel{FT}{\longleftrightarrow} aX(j\omega) + bY(j\omega)$

4.3.2 Time Shifting

If
$$x(t) \stackrel{FT}{\longleftrightarrow} X(j\omega)$$
, $4\{e^{-j\omega t_0}X(j\omega)\} = 4\{X(j\omega)\} - \omega t_0$
then $x(t-t_0) \stackrel{FT}{\longleftrightarrow} e^{-j\omega t_0}X(j\omega)$

A signal which is shifted in time does not have the magnitude of its Fourier transform altered. The effect of a time shift on a signal is to introduce into its transform a phase shift, namely, $-\omega t_0$.

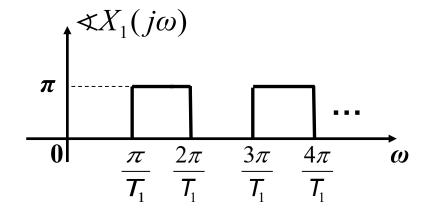
Example 4.9

Determine the Fourier transform of x(t) shown in the figure.

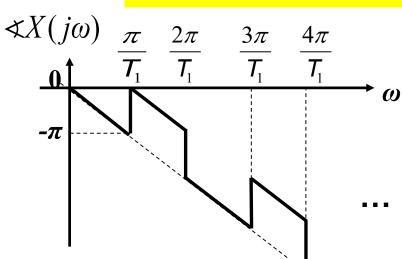
$$x_{1}(t) = u(t + T_{1}) - u(t - T_{1}) \leftrightarrow X_{1}(j\omega) = \frac{2\sin\omega T_{1}}{\omega}$$
and $x(t) = Ax_{1}(t - T_{1})$

From the time shifting property, we have

$$X(j\omega) = AX_1(j\omega)e^{-j\omega T_1} = \frac{2A\sin\omega T_1}{\omega}e^{-j\omega T_1}$$







4.3.3 Conjugation and Conjugate Symmetry

If
$$x(t) \stackrel{FT}{\longleftrightarrow} X(j\omega)$$
, then $x^*(t) \stackrel{FT}{\longleftrightarrow} X^*(-j\omega)$

> If x(t) is real, then $X(-j\omega) = X^*(j\omega)$

- ightharpoonup If x(t) is both real and even, so is $X(j\omega)$, i.e., $X(j\omega) = \text{Re}\{X(j\omega)\}$
- If x(t) is real and odd, $X(j\omega)$ is purely imaginary and odd, i.e., $X(j\omega)=j\cdot Im\{X(j\omega)\}$

$$x(t) = x_e(t) + x_o(t)$$

$$x_e(t) \stackrel{FT}{\longleftrightarrow} \text{Re}\{X(j\omega)\} \qquad x_o(t) \stackrel{FT}{\longleftrightarrow} j \text{Im}\{X(j\omega)\}$$

Example 4.10

Consider again the Fourier transform evaluation of $x(t) = e^{-\alpha|t|}$

Sol:
$$e^{-\alpha|t|} = e^{-\alpha t}u(t) + e^{\alpha t}u(-t) = 2\left[\frac{e^{-\alpha t}u(t) + e^{\alpha t}u(-t)}{2}\right] = 2Ev\left\{e^{-\alpha t}u(t)\right\}$$
and $e^{-\alpha t}u(t) \longleftrightarrow \frac{FT}{\alpha + i\omega}$

From the symmetry properties of the Fourier transform, we have

$$X(j\omega) = 2 \operatorname{Re} \left\{ \frac{1}{\alpha + j\omega} \right\} = \frac{2\alpha}{\alpha^2 + \omega^2}$$

4.3.4 Differentiation and Integration

If
$$x(t) \stackrel{FT}{\longleftrightarrow} X(j\omega)$$
, then $\frac{dx(t)}{dt} \stackrel{FT}{\longleftrightarrow} j\omega X(j\omega)$

$$\frac{d^n x(t)}{dt^n} \stackrel{FT}{\longleftrightarrow} (j\omega)^n X(j\omega)$$

Consequence: Strengthening the high-frequencies in the signal.

$$\int_{-\infty}^{t} x(\tau)d\tau \xleftarrow{FT} \frac{1}{j\omega} X(j\omega) + \pi X(0)\delta(\omega)$$

Consequence: Strengthening the low-frequencies in the signal.

Example 4.11

Determine the Fourier transform of the unit step x(t) = u(t).

$$Sol: u(t) = \frac{1}{2} + \frac{1}{2} \operatorname{sgn}(t) \qquad \operatorname{sgn}(t) = \lim_{\alpha \to 0} \left[e^{-\alpha t} u(t) - e^{\alpha t} u(-t) \right] (\alpha > 0)$$

$$\operatorname{Esgn}(t) = \lim_{\alpha \to 0} \left(\frac{1}{\alpha + j\omega} - \frac{1}{\alpha - j\omega} \right) = \frac{2}{j\omega} \qquad X(j\omega) = \frac{1}{j\omega} + \pi \delta(\omega)$$

Example 4.12

Calculate the Fourier transform $X(j\omega)$ for the signal x(t) displayed in the figure: x(t)

Sol: Let
$$g(t) = \frac{dx(t)}{dt} = \frac{1}{-1} t + \frac{1}{-1} t$$

$$X_1(j\omega) = \left(\frac{2\sin\omega}{\omega}\right), \quad X_2(j\omega) = -e^{j\omega} - e^{-j\omega}$$

$$X(j\omega) = \frac{2\sin\omega}{j\omega^2} - \frac{2\cos\omega}{j\omega}$$

4.3.5 Time and Frequency Scaling

If
$$x(t) \stackrel{FT}{\longleftrightarrow} X(j\omega)$$
, then $x(at) \stackrel{FT}{\longleftrightarrow} \frac{1}{|a|} X(\frac{j\omega}{a})$, a is real

A linear scaling in time by a factor of a corresponds to a linear scaling in frequency by a factor of 1/a, and vice versa.

Specially, when
$$a = -1$$
, we have $x(-t) \stackrel{FT}{\longleftrightarrow} X(-j\omega)$

4.3.6 Duality

If
$$x(t) \stackrel{FT}{\longleftrightarrow} X(j\omega)$$
, then $X(jt) \stackrel{FT}{\longleftrightarrow} 2\pi x(-\omega)$

Example 4.13

Using duality to find $G(j\omega)$ of the signal $g(t) = \frac{2}{1+t^2}$.

Sol: From pair
$$x(t) = e^{-|t|} \longleftrightarrow X(j\omega) = \frac{2}{1 + \omega^2}$$

By duality property, we can write
$$g(t) = \frac{2}{1+t^2} \longleftrightarrow G(j\omega) = 2\pi e^{-|-\omega|} = 2\pi e^{-|\omega|}$$

Duality property shows that for any Fourier transform pair there is a *dual pair* with the time and frequency variables interchanged.

Differentiation in Frequency-domain:

$$-jtx(t) \longleftrightarrow \frac{dX(j\omega)}{d\omega}$$

Integration in Frequency-domain:

$$-\frac{1}{it}x(t) + \pi x(0)\delta(t) \longleftrightarrow_{-\infty}^{FT} \int_{-\infty}^{\infty} X(\eta)d\eta$$

Frequency Shifting:

$$e^{j\omega_0 t} x(t) \stackrel{FT}{\longleftrightarrow} X(j(\omega - \omega_0))$$

4.3.7 Parseval's Relation

$$\int_{-\infty}^{\infty} |x(t)|^2 dt = \frac{1}{2\pi} \int_{-\infty}^{\infty} |X(j\omega)|^2 d\omega$$
energy-density spectrum
(能量密度谱)

Parseval's relation says that this total energy may be determined either by computing energy per unit time ($|x(t)|^2$) and integrating over all time or by computing the energy per unit frequency ($|X(j\omega)|^2/2\pi$) and integrating over all frequencies.

Example 4.14

For each of the Fourier transforms shown in figures (a) and (b), evaluate the following time-domain expressions:

evaluate the following time-domain expressions:
$$E = \int_{-\infty}^{\infty} \left| x(t) \right|^2 dt, \ D = \frac{d}{dt} x(t) \big|_{t=0}$$

$$X(j\omega) \qquad \qquad X(j\omega) \qquad X(j\omega) \qquad \qquad X(j\omega)$$

$$D = g(0) = \frac{1}{2\pi} \int_{-\infty}^{\infty} G(j\omega) d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} j\omega X(j\omega) d\omega$$

$$D_a = 0 \qquad D_b = -\frac{\sqrt{\pi}}{2\pi}$$

Consider the convolution integral: $y(t) = x(t) * h(t) = \int_{-\infty}^{\infty} x(\tau)h(t-\tau)d\tau$ The Fourier transform of y(t) is:

$$Y(j\omega) = \mathcal{F}\{y(t)\} = \int_{-\infty}^{\infty} \left[\int_{-\infty}^{\infty} x(\tau)h(t-\tau)d\tau \right] e^{-j\omega t} dt$$

Interchanging the order of integration and noting that $x(\tau)$ does not depend on t, we have $Y(j\omega) = \int_{-\infty}^{\infty} x(\tau) \left[\int_{-\infty}^{\infty} h(t-\tau)e^{-j\omega t} dt \right] d\tau$

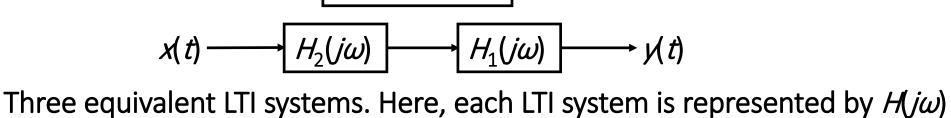
$$Y(j\omega) = \int_{-\infty}^{\infty} x(\tau)e^{-j\omega\tau}H(j\omega)d\tau = H(j\omega)\int_{-\infty}^{\infty} x(\tau)e^{-j\omega\tau}d\tau = H(j\omega)X(j\omega)$$
$$y(t) = h(t) * x(t) \xleftarrow{FT} Y(j\omega) = H(j\omega)X(j\omega)$$

The Fourier transform maps the convolution of two signals into the product of their Fourier transforms. $H(j\omega)$, the frequency response, is the Fourier transform of the impulse response h(t). It captures the change in complex amplitude of the Fourier transform of the input at each frequency ω .

 $H(j\omega)$.

 \triangleright The frequency response $H(j\omega)$ also can characterize an LTI system, just as its inverse transform, the unit impulse response h(t). $x(t) \longrightarrow H_1(j\omega) \longmapsto H_2(j\omega) \longmapsto y(t)$

 $x(t) \longrightarrow H_1(j\omega) H_2(j\omega) \longrightarrow y(t)$



- \rightarrow $H(j\omega)$ cannot be defined for every LTI system.
- >Since essentially all physical or practical signals satisfy the last two conditions in Dirichlet conditions, the condition of absolutely integrable becomes the determining factor which can guarantee the existence of the Fourier transform $H(j\omega)$ of h(t). That is, only a stable LTI system has a frequency response

Example 4.15

Consider an integrator — that is, an LTI system specified by the equation $y(t) = \int_{-\infty}^{t} x(\tau) d\tau$

Sol: Since
$$y(t) = x(t) * u(t) = \int_{-\infty}^{\infty} x(\tau)u(t-\tau)d\tau = \int_{-\infty}^{t} x(\tau)d\tau$$

So the impulse response for this system is the unit step u(t).

The frequency response of the system is $H(j\omega) = \frac{1}{j\omega} + \pi\delta(\omega)$

Using the convolution property, we have

$$Y(j\omega) = H(j\omega)X(j\omega)$$

$$= \frac{1}{j\omega}X(j\omega) + \pi X(j\omega)\delta(\omega)$$

$$= \frac{1}{j\omega}X(j\omega) + \pi X(0)\delta(\omega)$$

Example 4.16

Find the response of an LTI system with $h(t) = e^{-at}u(t)$, a > 0 to the input signal $x(t) = e^{-bt}u(t)$, b > 0.

Sol:
$$X(j\omega) = \frac{1}{b+j\omega}$$
, $H(j\omega) = \frac{1}{a+j\omega} \Rightarrow Y(j\omega) = \frac{1}{(a+j\omega)(b+j\omega)}$

Expand $Y(j\omega)$ in a partial-fraction expansion (部分分式展开法).

When
$$b \neq a$$
, let $Y(j\omega) = \frac{A}{(a+j\omega)} + \frac{B}{(b+j\omega)} = \frac{1}{b-a} \left[\frac{1}{a+j\omega} - \frac{1}{b+j\omega} \right]$

$$A = (a+j\omega)Y(j\omega)\Big|_{j\omega=-a} = \frac{1}{b+j\omega}\Big|_{j\omega=-a} = \frac{1}{b-a}$$

$$B = (b+j\omega)Y(j\omega)\Big|_{j\omega=-b} = \frac{1}{a+j\omega}\Big|_{j\omega=-b} = \frac{-1}{b-a}$$

$$y(t) = \frac{1}{b-a} \left[e^{-at}u(t) - e^{-bt}u(t) \right] \quad b \neq a$$

When
$$b = a$$
, $Y(j\omega) = \frac{1}{(a+j\omega)^2} = j\frac{d}{d\omega} \left| \frac{1}{(a+j\omega)} \right|$

From the differentiation in the frequency-domain property,

$$e^{-at}u(t) \longleftrightarrow \frac{1}{a+j\omega}$$

$$te^{-at}u(t) \longleftrightarrow \frac{FT}{d\omega} \left[\frac{1}{a+j\omega}\right] = \frac{1}{(a+j\omega)^2}$$

Consequently,

$$y(t) = te^{-at}u(t)$$
 $b = a$

Example 4.17

Determine the response of an ideal low-pass filter to an input signal x(t) that has the form of a *sinc* function, $x(t) = \frac{\sin \omega_i t}{\pi t}$

Sol: The impulse response of the ideal low-pass filter is of a similar form: $h(t) = \frac{\sin \omega_c t}{2}$

$$X(j\omega) = \begin{cases} 1 & |\omega| \le \omega_i \\ 0 & elsewhere \end{cases}, \qquad H(j\omega) = \begin{cases} 1 & |\omega| \le \omega_c \\ 0 & elsewhere \end{cases}$$

Therefore, $Y(j\omega) = \begin{cases} 1 & |\omega| \le \omega_0 \\ 0 & elsewhere \end{cases}$ where $\omega_0 = \min(\omega_i, \omega_c)$.

Finally, the inverse Fourier transform of $Y(j\omega)$ is given by

$$y(t) = \begin{cases} \frac{\sin \omega_c t}{\pi t} = h(t) & \text{if } \omega_c \le \omega_i \\ \frac{\sin \omega_i t}{\pi t} = x(t) & \text{if } \omega_i \le \omega_c \end{cases}$$

4.5 The Multiplication Property

$$r(t) = s(t)p(t) \stackrel{FT}{\longleftrightarrow} R(j\omega) = \frac{1}{2\pi} [S(j\omega) * P(j\omega)]$$

amplitude modulation property (幅度调制定理)

Example 4.18

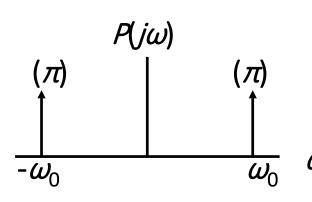
Let s(t) be a signal whose spectrum $S(j\omega)$ is depicted in the following Figure. Also, consider the signal $p(t) = \cos \omega_0 t$,

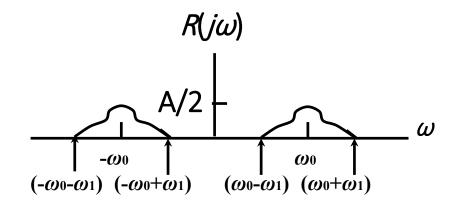
$$S(j\omega)$$
 A
 ω

Sol: Since
$$P(j\omega) = \pi\delta(\omega - \omega_0) + \pi\delta(\omega + \omega_0)$$

From the multiplication property:

$$R(j\omega) = \frac{1}{2\pi}S(j\omega) * P(j\omega) = \frac{1}{2}S(j(\omega - \omega_0)) + \frac{1}{2}S(j(\omega + \omega_0))$$

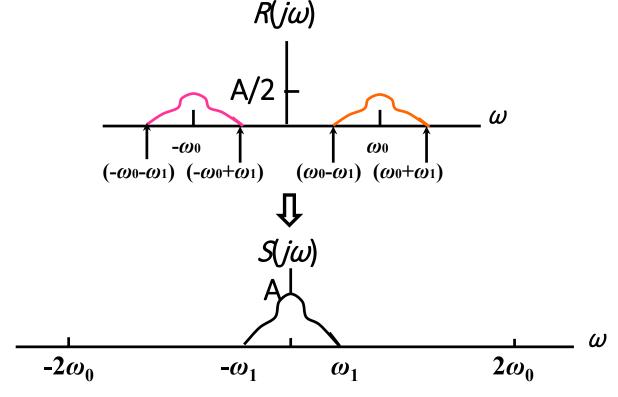




4.5 The Multiplication Property

Example 4.19

Let us consider r(t) as obtained in Example 4.18, and let g(t) = r(t)p(t), we will show how to recover the modulated signal s(t).



For a *stable* LTI system which is described by a linear constant-coefficient differential equation of the form :

$$\sum_{k=0}^{N} a_k \frac{d^k y(t)}{dt^k} = \sum_{k=0}^{M} b_k \frac{d^k x(t)}{dt^k}$$

$$\mathcal{F}\left\{\sum_{k=0}^{N} a_k \frac{d^k y(t)}{dt^k}\right\} = \mathcal{F}\left\{\sum_{k=0}^{M} b_k \frac{d^k x(t)}{dt^k}\right\} \sum_{k=0}^{N} a_k \mathcal{F}\left\{\frac{d^k y(t)}{dt^k}\right\} = \sum_{k=0}^{M} b_k \mathcal{F}\left\{\frac{d^k x(t)}{dt^k}\right\}$$

$$\sum_{k=0}^{N} a_k (j\omega)^k Y(j\omega) = \sum_{k=0}^{M} b_k (j\omega)^k X(j\omega) \quad \text{(differentiation)}$$
Equivalently,
$$Y(j\omega) \left[\sum_{k=0}^{N} a_k (j\omega)^k\right] = X(j\omega) \left[\sum_{k=0}^{M} b_k (j\omega)^k\right]$$

$$H(j\omega) = \frac{Y(j\omega)}{X(j\omega)} \sum_{k=0}^{M} b_k (j\omega)^k \quad \text{(Convolution)}$$

 \nearrow $H(j\omega)$ is a ratio of polynomials in $(j\omega)$, so it is a rational function.

Example 4.20

Consider a stable LTI system that is characterized by the differential equation $\frac{d^2y(t)}{dt^2} + 4\frac{dy(t)}{dt} + 3y(t) = \frac{dx(t)}{dt} + 2x(t)$, determine its impulse response.

Sol: The frequency response is

$$H(j\omega) = \frac{(j\omega) + 2}{(j\omega)^2 + 4(j\omega) + 3}$$

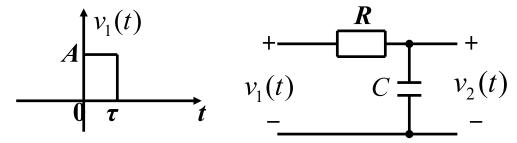
To determine the corresponding impulse response, we use the method of partial-fraction expansion:

Thus, the impulse response is

$$H(j\omega) = \frac{\frac{1}{2}}{j\omega + 1} + \frac{\frac{1}{2}}{j\omega + 3}$$
$$h(t) = \frac{1}{2}e^{-t}u(t) + \frac{1}{2}e^{-3t}u(t)$$

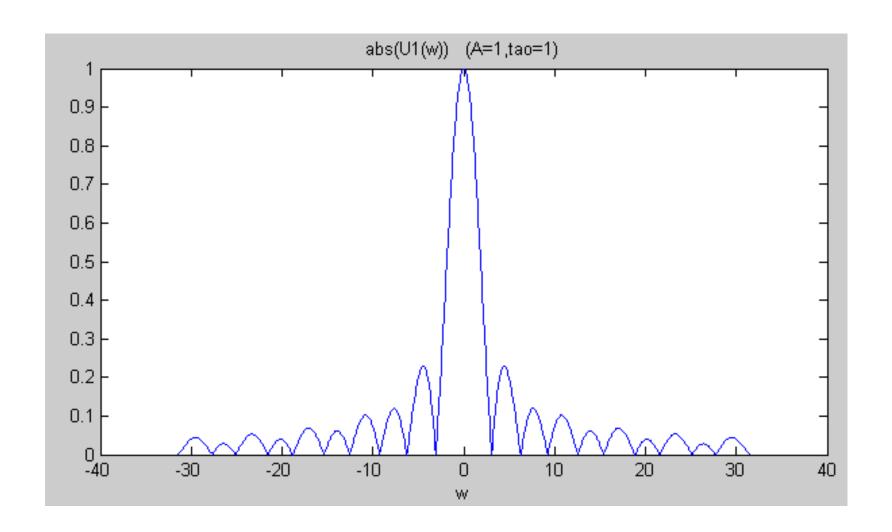
Example 4.21

A RC circuit is as follows, $v_1(t)$ is the input, determine $v_2(t)$.



Sol: 1) Compute $V_1(j\omega)$

$$\begin{split} v_1(t) &= A[u(t) - u(t - \tau)] \\ V_1(j\omega) &= A[\pi\delta(\omega) + \frac{1}{j\omega} - \pi\delta(\omega)e^{-j\omega\tau} - \frac{e^{-j\omega\tau}}{j\omega}] \\ &= \frac{A}{j\omega}(1 - e^{-j\omega\tau}) = \frac{2A}{\omega}\sin\frac{\omega\tau}{2}e^{-j\frac{\omega\tau}{2}} = A\tau Sa(\frac{\omega\tau}{2})e^{-j\frac{\omega\tau}{2}} \end{split}$$



2) Compute $H(j\omega)$

$$H(j\omega) = \frac{V_2(j\omega)}{V_1(j\omega)} = \frac{\frac{1}{j\omega C}}{R + \frac{1}{j\omega C}} = \frac{1}{1 + j\omega RC} = \frac{\frac{1}{RC}}{\frac{1}{RC} + j\omega}$$

Let
$$\frac{1}{RC} = \alpha = \frac{1}{\tau_0}$$

$$H(j\omega) = \frac{\alpha}{\alpha + j\omega} = \frac{\alpha}{\sqrt{\alpha^2 + \omega^2}} e^{-j\arctan\frac{\omega}{\alpha}}$$

3) Compute $V_2(j\omega)$

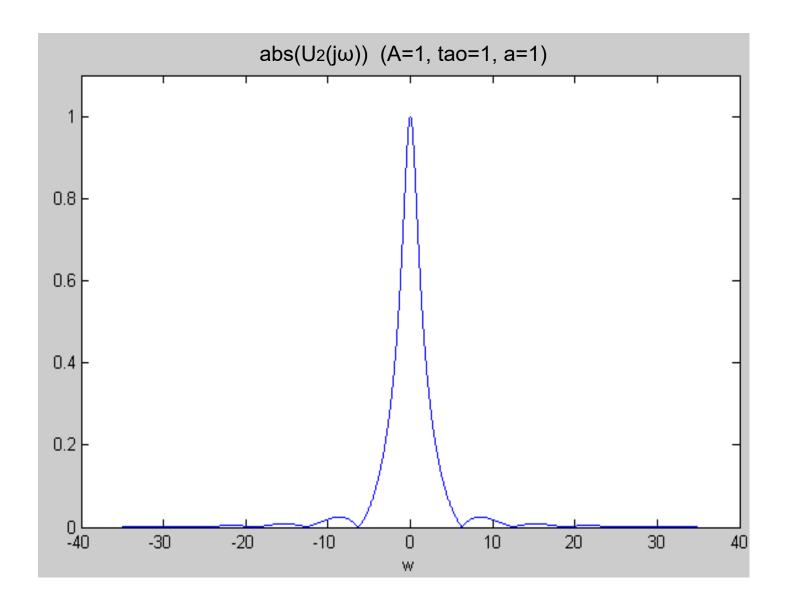
$$V_{2}(j\omega) = V_{1}(j\omega)H(j\omega)$$

$$= A \tau Sa(\frac{\omega\tau}{2})e^{-j\frac{\omega\tau}{2}} \cdot \frac{\alpha}{\sqrt{\alpha^{2} + \omega^{2}}}e^{-j\arctan\frac{\omega}{\alpha}}$$

$$= \frac{\alpha A \tau}{\sqrt{\alpha^{2} + \omega^{2}}}Sa(\frac{\omega\tau}{2})e^{-j(\frac{\omega\tau}{2} + \arctan\frac{\omega}{\alpha})}$$

$$|V_{2}(j\omega)| = \frac{\alpha A \tau}{\sqrt{\alpha^{2} + \omega^{2}}}|Sa(\frac{\omega\tau}{2})| = \frac{2\alpha A\left|\sin(\frac{\omega\tau}{2})\right|}{\omega\sqrt{\alpha^{2} + \omega^{2}}}$$

$$|V_{2}(j\omega)| = \begin{cases} -(\frac{\omega\tau}{2} + \arctan\frac{\omega}{\alpha}), \sin(\frac{\omega\tau}{2}) > 0\\ \pm \pi - (\frac{\omega\tau}{2} + \arctan\frac{\omega}{\alpha}), \sin(\frac{\omega\tau}{2}) < 0 \end{cases}$$



Equations

4) Compute
$$v_2(t) = \mathcal{F}^{-1}[V_2(j\omega)]$$

$$V_1(j\omega) = \frac{A}{j\omega} (1 - e^{-j\omega\tau})$$

Compute
$$v_2(t) = \mathcal{F}^{-1}[V_2(j\omega)]$$

$$V_1(j\omega) = \frac{A}{j\omega}(1 - e^{-j\omega\tau})$$

$$V_2(j\omega) = A(1 - e^{-j\omega\tau}) \frac{\alpha}{j\omega(\alpha + j\omega)} = A(1 - e^{-j\omega\tau})(\frac{1}{j\omega} - \frac{1}{j\omega + \alpha})$$

$$= \frac{A(1 - e^{-j\omega\tau})}{j\omega} - \frac{A}{\alpha + j\omega} + \frac{A}{\alpha + j\omega}e^{-j\omega\tau}$$

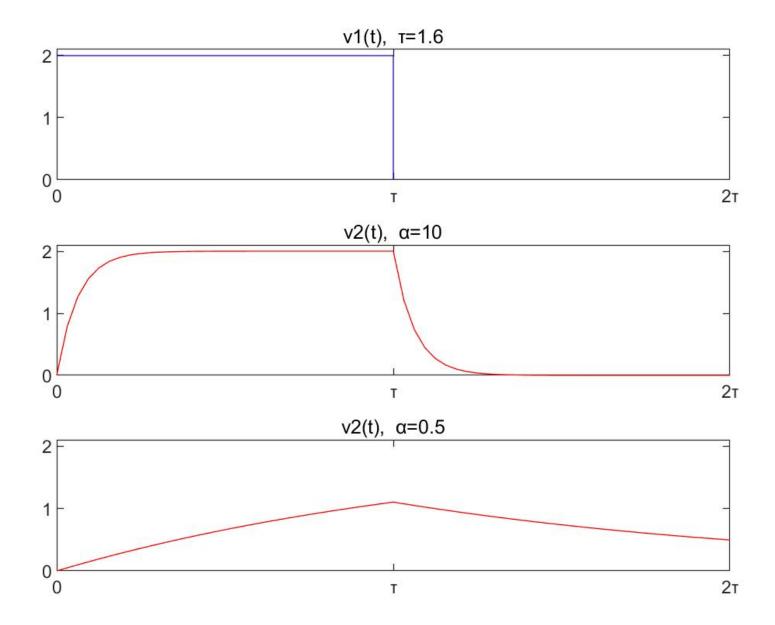
$$\therefore v_2(t) = A[u(t) - u(t - \tau)] - Ae^{-\alpha t}u(t) + Ae^{-\alpha(t - \tau)}u(t - \tau)$$

$$\therefore v_2(t) = A[u(t) - u(t - \tau)] - Ae^{-\alpha t}u(t) + Ae^{-\alpha(t - \tau)}u(t - \tau)$$
$$= A(1 - e^{-\alpha t})u(t) - A[1 - e^{-\alpha(t - \tau)}]u(t - \tau)$$

Rising and falling characteristics in time domain:

 $v_1(t)$: $t = 0, \tau$ Change very fast (jump)—Abundant high frequencies $v_2(t)$: $t = 0, \tau$ Change slowly, need a period of time to rise or fall

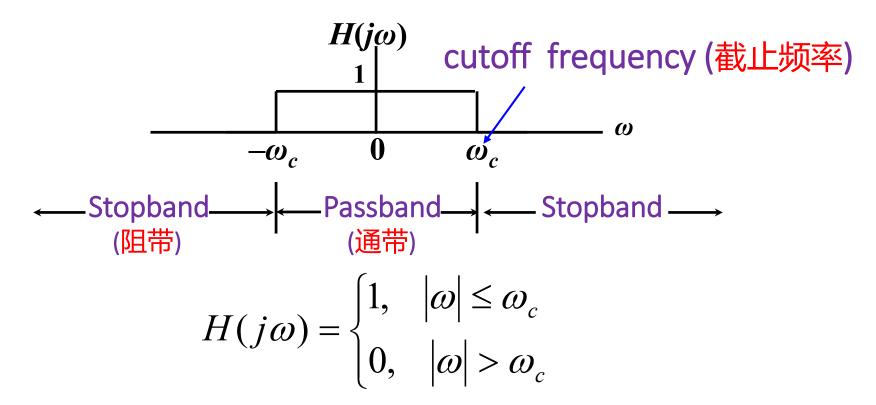
High frequencies are attenuated



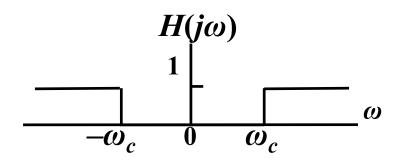
4.7.1 Introduction to Ideal Frequency-Selective Filters

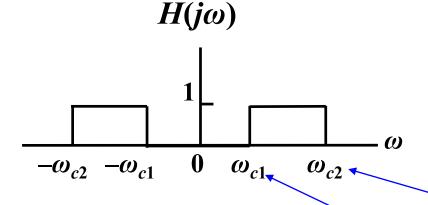
- Filtering: a process in which the relative amplitudes of the frequency components in a signal are changed or some frequency components are eliminated entirely.
- Frequency-selective filters (频选滤波器): systems that are designed to pass some frequencies essentially undistorted and significantly attenuate or eliminate others.
- Types of frequency-selective filters
 - ✓ low-pass filter (低通滤波器)
 - ✓ high-pass filter (高通滤波器)
 - ✓ band-pass filter (带通滤波器)
 - ✓ band-stop filter (带阻滤波器)

The frequency responses of a zero-phase idea/ low-pass filter, a zero-phase idea/ high-pass filter, a zero-phase idea/ band-pass filter and a zero-phase idea/ band-stop filter are illustrated in the following figures, respectively:

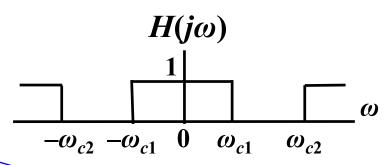


$$H(j\omega) = \begin{cases} 1, & |\omega| \ge \omega_c \\ 0, & |\omega| < \omega_c \end{cases}$$





$$H(j\omega) = \begin{cases} 1, & \omega_{c_2} \ge |\omega| \ge \omega_{c_1} \\ 0, & |\omega| < \omega_{c_1} \text{ or } |\omega| > \omega_{c_2} \end{cases}$$



upper cutoff frequency (上截止频率) lower cutoff frequency

(下截止频率)

The impulse response of the ideal low-pass filter is:

$$h(t) = \frac{1}{2\pi} \int_{-\omega_c}^{\omega_c} e^{j\omega t} d\omega = \frac{\sin \omega_c t}{\pi t}$$

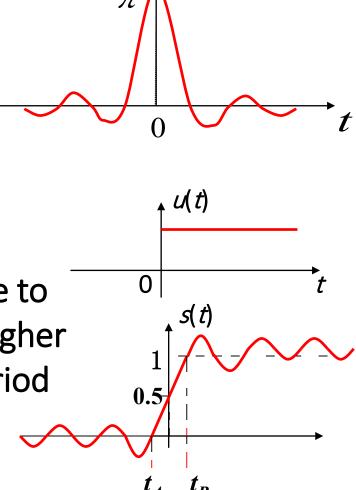
The step response is:

$$s(t) = \frac{1}{2} + \frac{1}{\pi} Si(\omega_c t)$$

Here,
$$Si(\omega_c t) = \int_0^{\omega_c t} \frac{\sin x}{x} dx$$

Different from the input signal — u(t), output signal s(t) needs a period of time to rise from 0 to 1, because frequencies higher than ω_c in u(t) are rejected. And the period

$$t_r = t_B - t_A = \frac{3.84}{\omega_c}$$



The frequency response of the ideal high-pass filter can be represented in terms of the frequency response of the lowpass filter as:

$$H_h(j\omega) = \begin{cases} 1, & |\omega| \ge \omega_c \\ 0, & |\omega| < \omega_c \end{cases} = 1 - \begin{cases} 1, & |\omega| \le \omega_c \\ 0, & |\omega| > \omega_c \end{cases} = 1 - H_l(j\omega)$$

Thus, the impulse response of the ideal high-pass filter is:

$$h(t) = \delta(t) - \frac{\sin \omega_c t}{\pi t}$$

Clearly h(t) is not causal, so the ideal high-pass filter cannot be implemented.

4.7.2 Realizable Systems and Binding Characteristic of $H(j\omega)$

- \triangleright A physically realizable system must have its impulse response h(t) satisfy $h(t) = h(t) \cdot u(t)$ (sufficient and necessary condition) Time Domain condition
- Imposing causality on the system restricts the $H(j\omega)$ in significant ways. The *Paley-Wiener Criterion* says if $|H(j\omega)|^2$ is integrable, i.e., $\int_{-\infty}^{\infty} |H(j\omega)|^2 d\omega < \infty$, it can be proved that the necessary condition for a realizable $|H(j\omega)|$ (causal system) is:

$$\int_{-\infty}^{\infty} \frac{\left| \ln |H(j\omega)| \right|}{1 + \omega^2} d\omega < \infty \tag{I}$$

 \checkmark To satisfy above condition, H(jω) cannot be zero in any band of frequencies (although it can be zero at a finite number of frequencies).

✓ The decay rate of $|H(j\omega)|$ cannot be greater than that of exponentials.

Consider the causality of system with $|H(j\omega)| = e^{-|\omega|}$,

$$\lim_{B \to \infty} \int_{-B}^{B} \frac{\left| \ln |H(j\omega)| \right|}{1 + \omega^{2}} d\omega = \lim_{B \to \infty} \int_{-B}^{B} \frac{\left| \ln e^{-|\omega|} \right|}{1 + \omega^{2}} d\omega = \lim_{B \to \infty} \int_{-B}^{B} \frac{|\omega|}{1 + \omega^{2}} d\omega$$

$$= 2 \times \lim_{B \to \infty} \int_{0}^{B} \frac{\omega}{1 + \omega^{2}} d\omega = \lim_{B \to \infty} \ln(1 + \omega^{2}) \Big|_{0}^{B} \to \infty$$

- This shows that exponential magnitude response $|H(j\omega)| = e^{-|\omega|}$ does not satisfy the equation (I). Consequently, a system with a magnitude response function which decays faster than exponential must be non-causal so that cannot be implemented.
 - It can be proved that a magnitude response function $|H(j\omega)|$ that is composed of rational polynomials can satisfy equation (I).

Because of the causality restriction $(h(t)=h(t)\cdot u(t))$ there is some kind of mutually binding character between the real and imaginary parts or the magnitude and phase of $H(j\omega)$. Specifically, if let $H(j\omega)=H_R(j\omega)+jH_I(j\omega)$

then

$$H_{R}(j\omega) = \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{H_{I}(j\eta)}{\omega - \eta} d\eta \quad (II)$$

$$H_{I}(j\omega) = -\frac{1}{\pi} \int_{-\infty}^{\infty} \frac{H_{R}(j\eta)}{\omega - \eta} d\eta \quad (III)$$

Equations (II) and (III) are referred to as *Hilbert Transform pair*.

Conclusion: The real and imaginary parts of the transform of a real, causal impulse response h(t) can be determined from one another using the *Hilbert Transform*. (Problems 4.47 and 4.48) So do the phase and logarithm of the magnitude of $H(j\omega)$.

4.7.3 A Real Simple RCLow-pass Filter with rational $H(j\omega)$

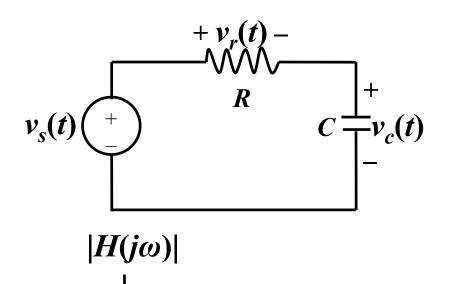
Input: source voltage $v_s(t)$;
Output: capacitor voltage $v_c(t)$

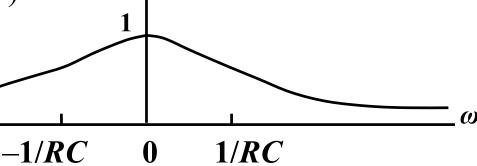
$$RC\frac{dv_c(t)}{dt} + v_c(t) = v_s(t)$$

$$RCj\omega V_c(j\omega) + V_c(j\omega) = V_s(j\omega)$$

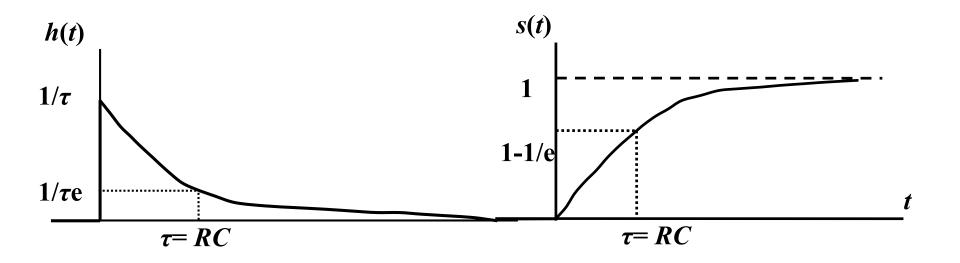
$$H_{lp}(j\omega) = \frac{V_c(j\omega)}{V_s(j\omega)}$$

$$=\frac{1}{1+RCj\omega}=\frac{\frac{1}{RC}}{\frac{1}{RC}+j\omega}$$





$$h(t) = \frac{1}{RC}e^{-t/RC}u(t)$$
 $s(t) = [1 - e^{-t/RC}]u(t)$



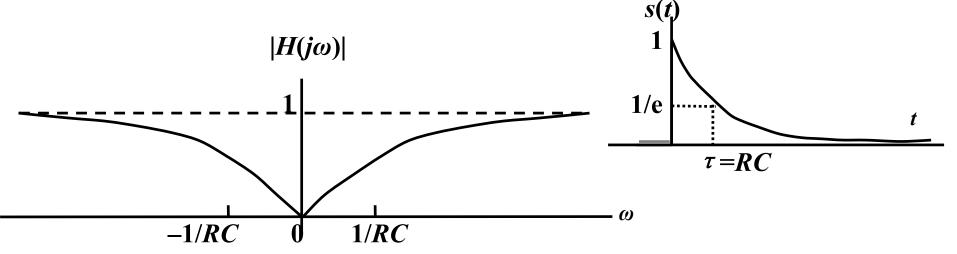
4.7.4 A Real Simple RC High-pass Filter with rational $H(j\omega)$

Input: source voltage $v_s(t)$; Output: resistor voltage $v_s(t)$

$$H_{hp}(j\omega) = \frac{j\omega RC}{1+j\omega RC} = 1 - \frac{\frac{dv_r(t)}{dt} + v_r(t) = RC\frac{dv_s(t)}{dt}}{\frac{1}{RC} + j\omega}$$

$$S(t) = e^{-t/RC}u(t)$$

Non-ideal high-pass filter



4.8 Transmission without Distortion

$$\begin{array}{c} \chi(t) \\ \chi(j\omega) \end{array} \longrightarrow \begin{array}{c} \chi(t) \\ \chi(j\omega) \end{array}$$

If the waveform of y(t) is different from that of x(t), distortion occurs.

- Linear systems can only introduce *linear distortion*, in which there isn't new frequencies generated.
- We can see from $Y(j\omega) = |H(j\omega)||X(j\omega)|e^{j(argH(j\omega)+argX(j\omega))}$ that linear distortions include magnitude distortion and phase distortion.

In the case of $y(t)=Kx(t-t_0)$, x(t) is transmitted without distortion.

Applying Fourier transform we have $Y(j\omega) = KX(j\omega)e^{-j\omega t_0}$

$$H(j\omega) = Ke^{-j\omega t_0}$$
 or
$$\begin{cases} |H(j\omega)| = K \\ \not\sim H(j\omega) = -\omega t_0 \end{cases}$$

4.8 Transmission without Distortion

Suppose
$$x(t) = E_1 \sin \omega_0 t + E_2 \sin 2\omega_0 t$$

Then

$$y(t) = KE_1 \sin(\omega_0 t + \varphi_1) + KE_2 \sin(2\omega_0 t + \varphi_2)$$

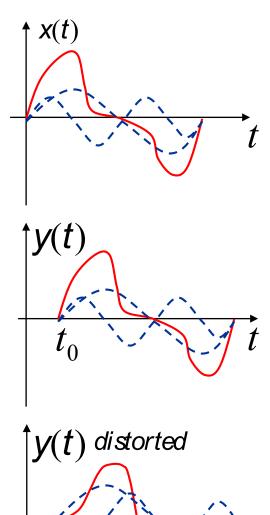
$$= KE_1 \sin \omega_0 (t + \frac{\varphi_1}{\omega_0}) + KE_2 \sin 2\omega_0 (t + \frac{\varphi_2}{2\omega_0})$$

To guarantee undistorted,

$$\frac{\varphi_1}{\omega_0} = \frac{\varphi_2}{2\omega_0} = const = -t_0$$

Consider non-periodic inputs case, we can write without loss of generality that

$$\varphi(\omega) = -\omega t_0 = \angle H(j\omega)$$



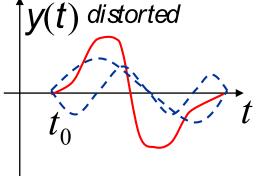
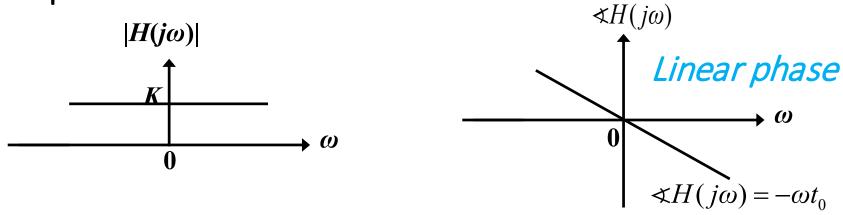


Illustration of phase distortion

4.8 Transmission without Distortion

> Conclusions: a linear system that can transmit signals applying to it as input without distortion must have a constant magnitude response and a phase response directly proportional to frequencies.



Linear phase shifts lead to very simple and easily understood change in a signal.

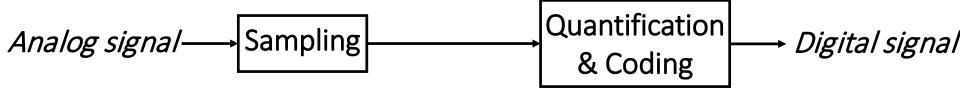
Group Delay:
$$\tau(\omega) = -\frac{d\{ \langle H(j\omega) \}}{d\omega}$$

More about this refer to Section 6.2.

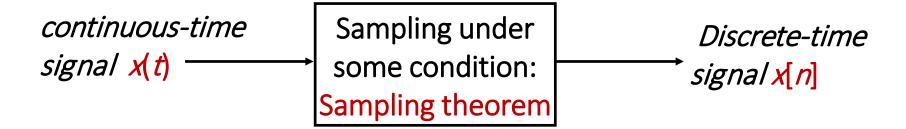
$$h(t) = K\delta(t - t_0)$$

Systems with this impulse response is ideal!

4.9.1 Introduction

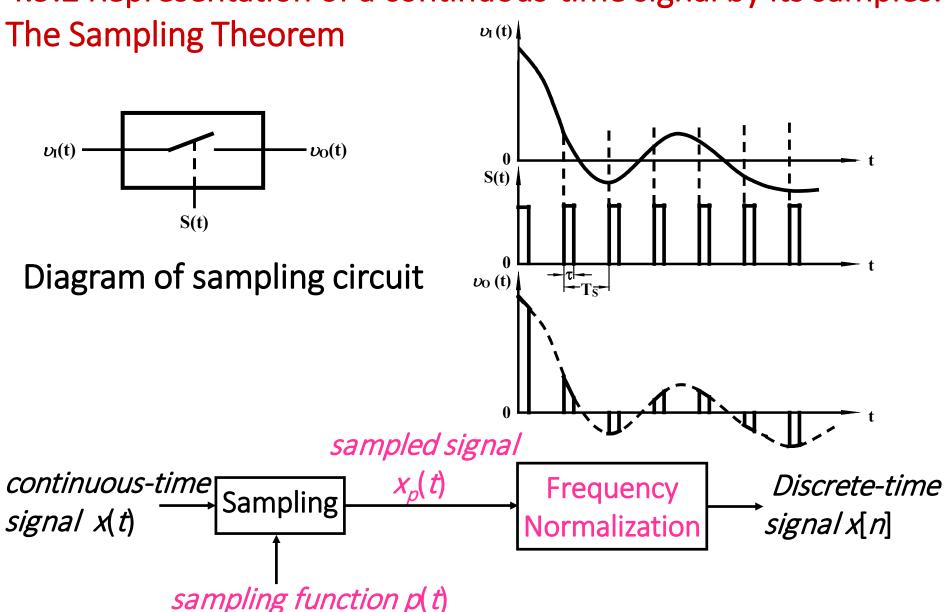


Analog to digital signals conversion process



Sampling theorem is a bridge between continuous-time signals and discrete-time signals.

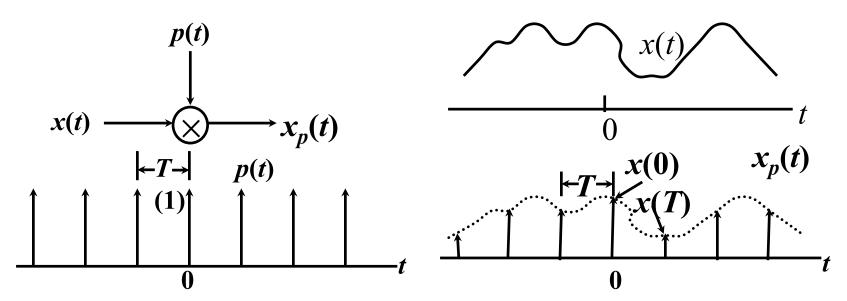
4.9.2 Representation of a continuous-time signal by its samples:



To make sure that x[n] can keep most information in x(t), or x(t) can be recovered from x[n], what are the conditions a sampling procedure must satisfy? Let's consider the problems:

- 1. What is the Fourier transform of $x_p(t)$?
- 2. How is it related to that of x(t)?

4.9.2.1 Impulse-train sampling (冲激串采样)

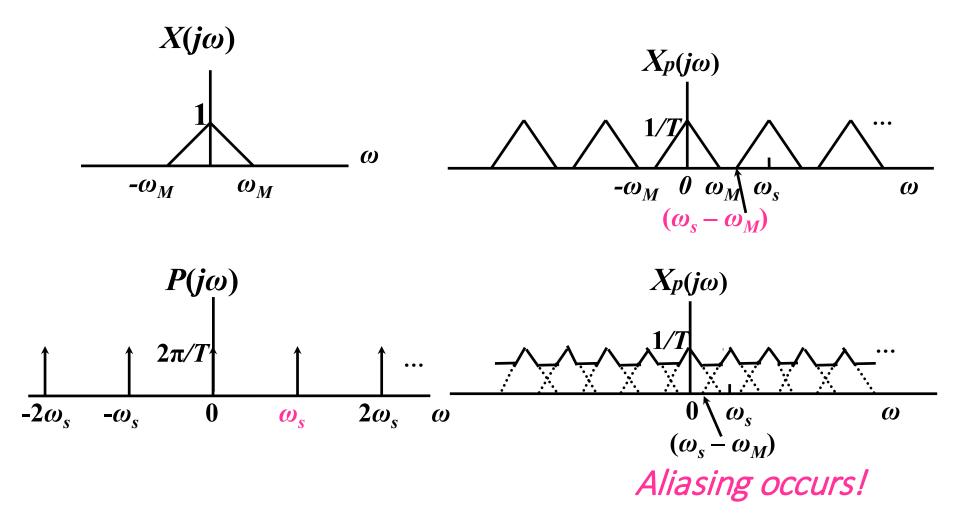


Mechanism of Impulse-train sampling

$$x_p(t) = x(t) \cdot \sum_{n = -\infty}^{\infty} \delta(t - nT) \quad X_p(j\omega) = \frac{1}{2\pi} \left[X(j\omega) * P(j\omega) \right] \quad P(j\omega) = \frac{2\pi}{T} \sum_{k = -\infty}^{\infty} \delta(\omega - k\omega_s)$$

$$X_{p}(j\omega) = \frac{1}{T} \sum_{k=-\infty}^{\infty} X(j(\omega - k\omega_{s}))$$

Answer for Q1&2: $X_p(j\omega)$ is a periodic function of ω consisting of a superposition of shifted replicas of $X(j\omega)$, scaled by $\frac{1}{r}$.



Effect in the frequency domain of sampling in the time domain

Sampling Theorem (采样定理)

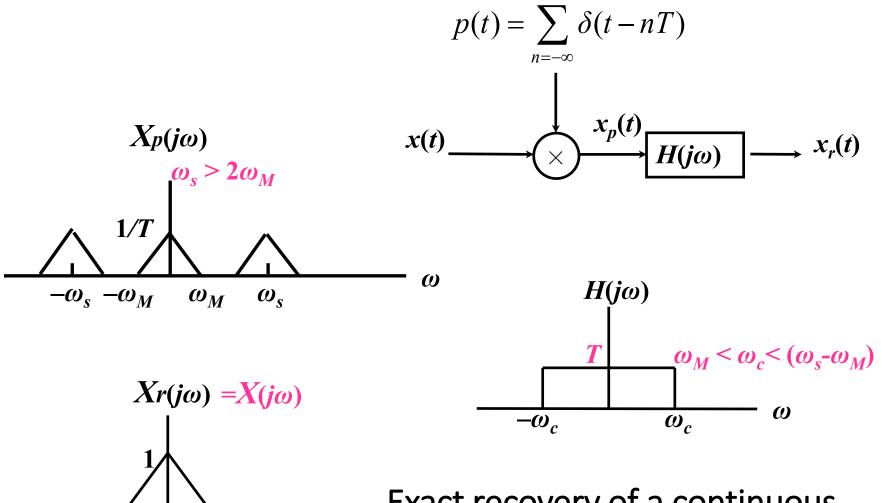
Let x(t) be a band-limited signal with $X(j\omega)=0$ for $|\omega|>\omega_M$. Then x(t) is uniquely determined by its samples x(n,T), n=0, ± 1 , ± 2 , ..., if $\omega_s > 2\omega_M$

where

Given these samples, we can reconstruct x(t) by generating a periodic impulse train in which successive impulses have amplitudes that are successive sample values. This impulse train is then processed through an ideal low-pass filter with gain T and cutoff frequency greater than ω_M and less than $\omega_s - \omega_M$. The resulting output signal will exactly equal x(t).

 $-\omega_M$

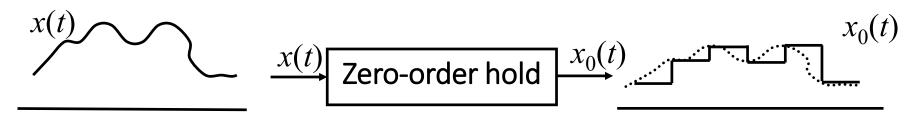
 ω_{M}



W

Exact recovery of a continuoustime signal from its samples using an ideal low-pass filter

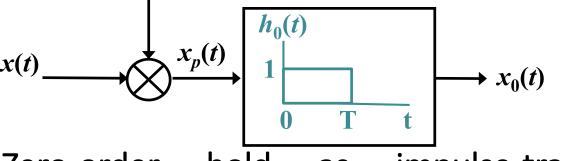
4.9.2.2 Sampling with a Zero-Order Hold



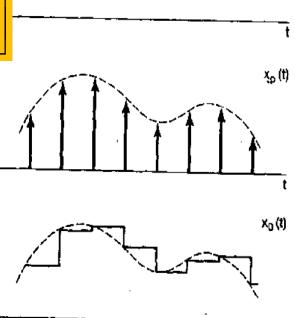
Sampling utilizing a zero-order hold

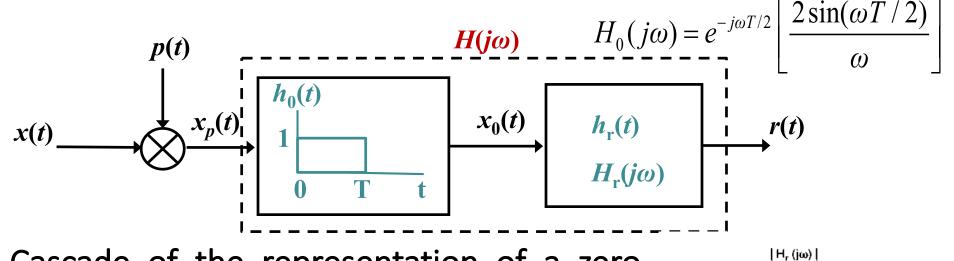
 \rightarrow How to generate $x_0(t)$?

$$p(t) = \sum_{n=-\infty}^{\infty} \delta(t - nT) \quad H_0(j\omega) = e^{-j\omega T/2} \left[\frac{2\sin(\omega T/2)}{\omega} \right]$$



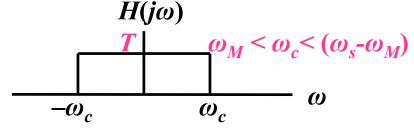
Zero-order hold as impulse-train sampling followed by an LTI system with a rectangular impulse response.





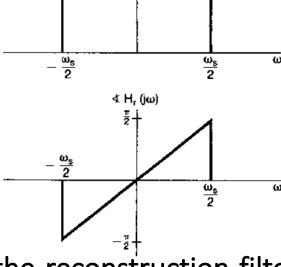
Cascade of the representation of a zeroorder hold with a reconstruction filter.

Here, $H(j\omega)$ is the ideal low-pass filter



$$H_r(j\omega) = \frac{e^{j\omega T/2}H(j\omega)}{\frac{2\sin(\omega T/2)}{\omega}}$$

Magnitude and phase for the reconstruction filter for a zero-order hold with the cut-off frequency of $H(j\omega)$ equal to $\omega / 2$.



- 4.9.3 Reconstruction of A Signal From Its Samples Using Interpolation
- > Interpolation (插值): a procedure in which the fitting (拟合) of a continuous-time signal to a set of sample values. It is commonly used for reconstruction of a function from samples.
- > Ways to interpolate: zero-order hold interpolation \ linear interpolation band-limited interpolation ...

$$p(t) = \sum_{n=-\infty}^{\infty} \delta(t - nT)$$

$$x(t)$$

$$x_p(t)$$

$$x_p(t)$$

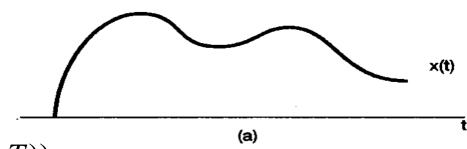
$$x_p(t)$$

$$x_p(t)$$

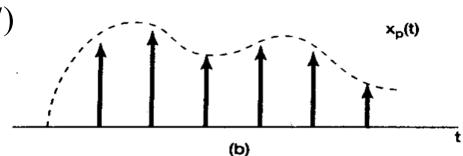
$$x_r(t) = x_p(t) * h(t) = \sum_{n=-\infty}^{\infty} x(nT)h(t-nT)$$

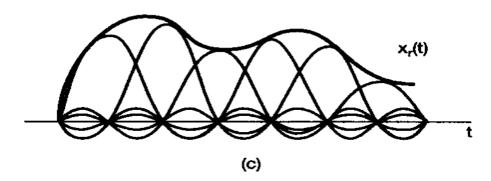
$$h(t) = \frac{\omega_c T \sin(\omega_c t)}{\pi \omega_c t}$$

$$x_{r}(t) = \sum_{n=-\infty}^{\infty} x(nT) \frac{\omega_{c}T}{\pi} \frac{\sin(\omega_{c}(t-nT))}{\omega_{c}(t-nT)}$$
Interpolation formula



$$x_r(t) = \sum_{n=-\infty}^{\infty} x(nT) \frac{\omega_c T}{\pi} \frac{\sin(\omega_c (t - nT))}{\omega_c (t - nT)}$$





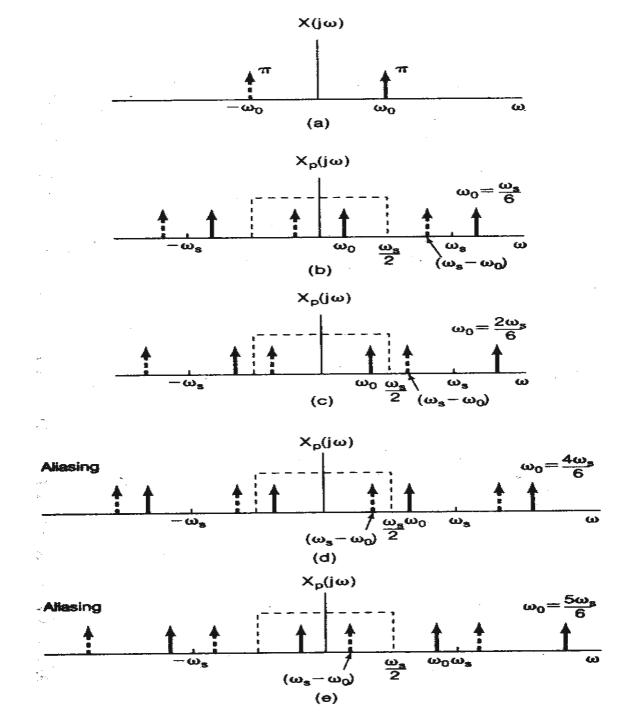
Ideal band-limited interpolation using the *sinc* function with $\omega_c = \frac{\omega_s}{2}$

4.9.4 The Effect of Under-Sampling (欠采样): Aliasing(混叠)

$$x(t) = \cos \omega_0 t,$$

 $x(t) = \cos 1000 t$ is sampled with

$$\omega_{\rm s}$$
 = 1500 rad / s and $\omega_{\rm s}$ = 1200 rad / s respectively.



4.10 SUMMARY

- The Fourier transform for both non-periodic and periodic continuous-time signals;
- The properties of the Fourier transform (relationships between characteristics of a continuous-time signal in time and frequency domains);
- Fourier analysis (Frequency domain analysis) for continuoustime LTI systems including both characteristics of systems and responses to some input signals;
- Frequency response and the way to obtain it;
- Continuous-time signals' sampling and their reconstruction.

Homework