

Reproduction and analysis of: *R*-matrix type parametrization of the Jost function for extracting the resonance parameters from scattering data[†]

Juliana Prada Suarez

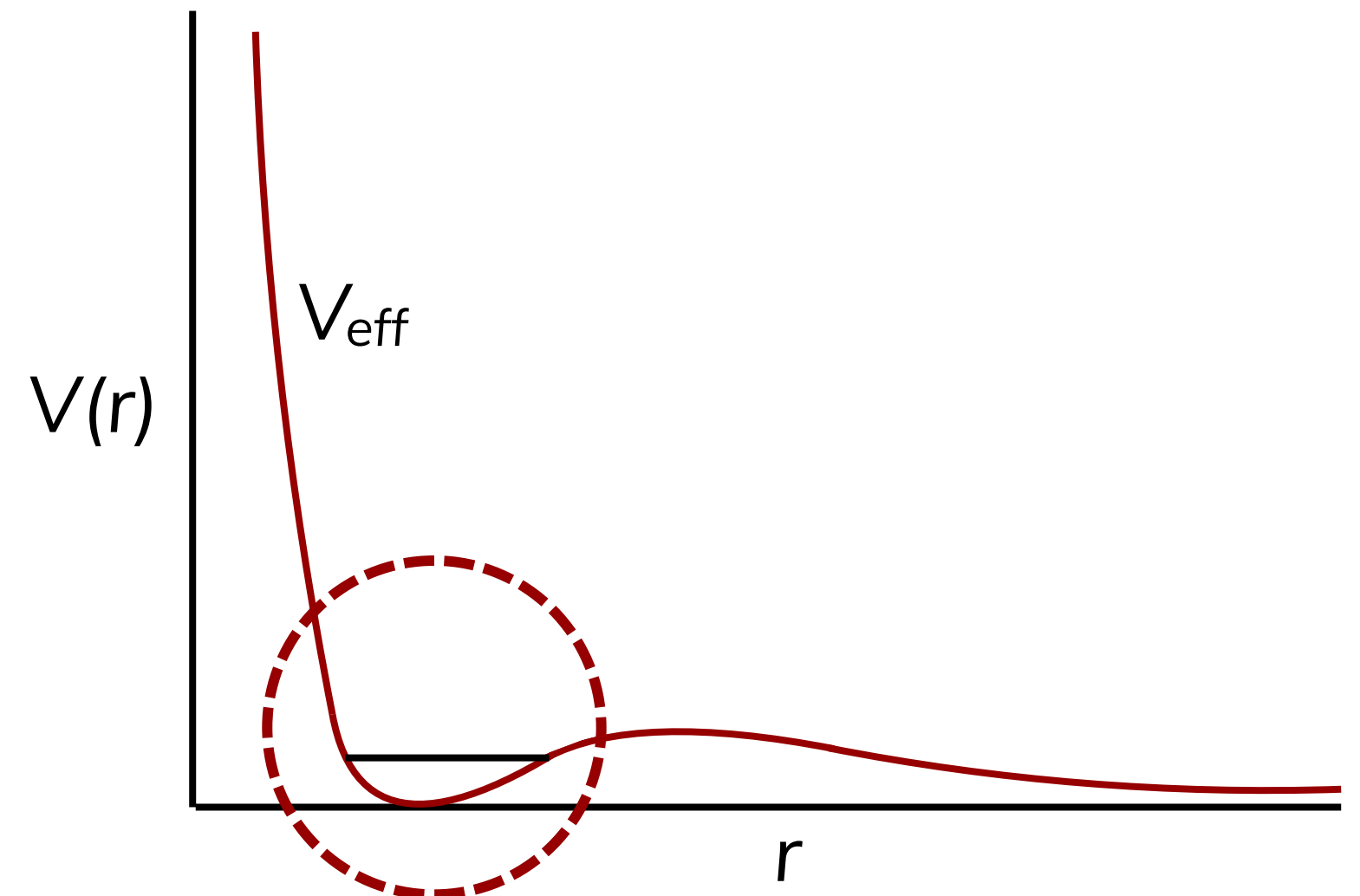
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Facultad de Ciencias

Departamento de Física

Introducción a la Teoría de Reacciones Nucleares



OUTLINE

- 1 Author's motivation and problem
- 2 Some approaches to the problem
- 3 Author's proposed approach
- 4 Theoretical background
- 5 Understanding the method
- 6 Results
- 7 Conclusions

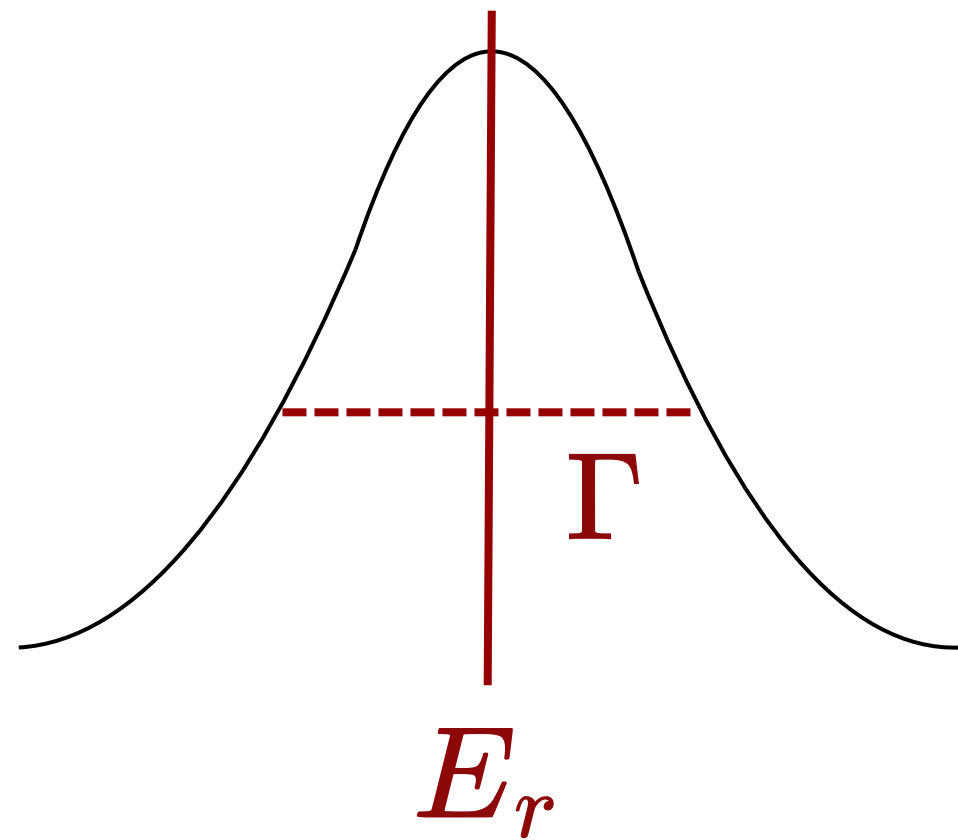
Reference list

- † Vaandrager, P., Lekala, M.L., & Rakityansky, S.A. (2025). R-matrix type parametrization of the Jost function for extracting the resonance parameters from scattering data. European Physical Journal A, 61, 79.

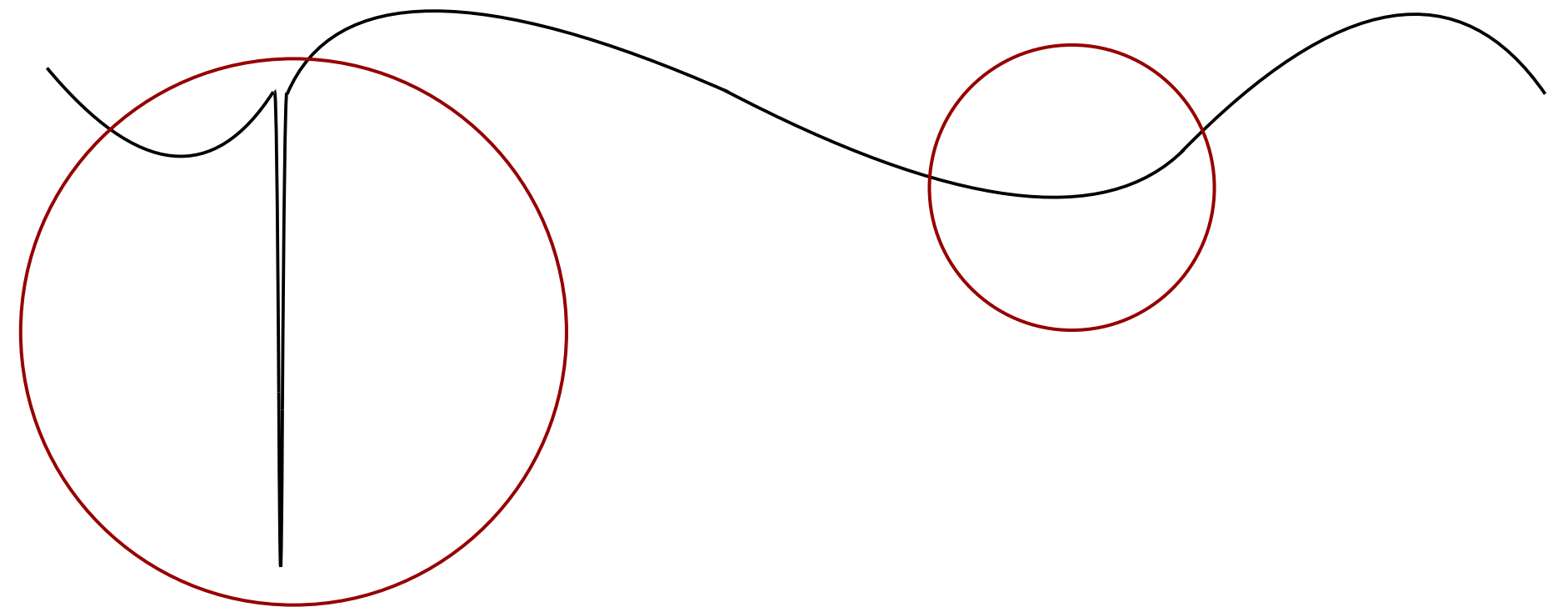
Author's motivation and problem[†]

Resonances: irregularities in the scattering cross-section

Bell-shaped resonances



Non bell-shaped resonances



Author's motivation and problem[†]

Almost all methods are based on:

Each resonance corresponds to a pole of the S-matrix at complex energy

$$E = E_r - \frac{i}{2}\Gamma$$

There is a special interest in:

Parametrizing the S-matrix to fit the data

One detail to consider:

The functional form of the parametrization must be adequate

Some approaches to the problem[†]

Jost functions method

a Jost functions

$$f_{\ell}^{(\text{in})}(E) \quad f_{\ell}^{(\text{out})}(E)$$

b S-matrix

$$S(E) = f_{\ell}^{(\text{out})}(E) \left[f_{\ell}^{(\text{in})}(E) \right]^{-1}$$

Why does this method work?

Almost all the factors in this functions are given explicitly and in an exact way

The remaining unknown factors are always single-valued and analytic functions of E

Why isn't it enough?

The number of fitting parameters is on the order of 50

Some approaches to the problem[†]

R-matrix method

a Phenomenological R-matrix

$$R_{\ell}(E, B_R) = \sum_{n=1}^N \frac{[\gamma_{n\ell}(B_R)]^2}{E_{n\ell} - E}$$

Why does this method work?

Few fitting parameters

Why isn't it enough?

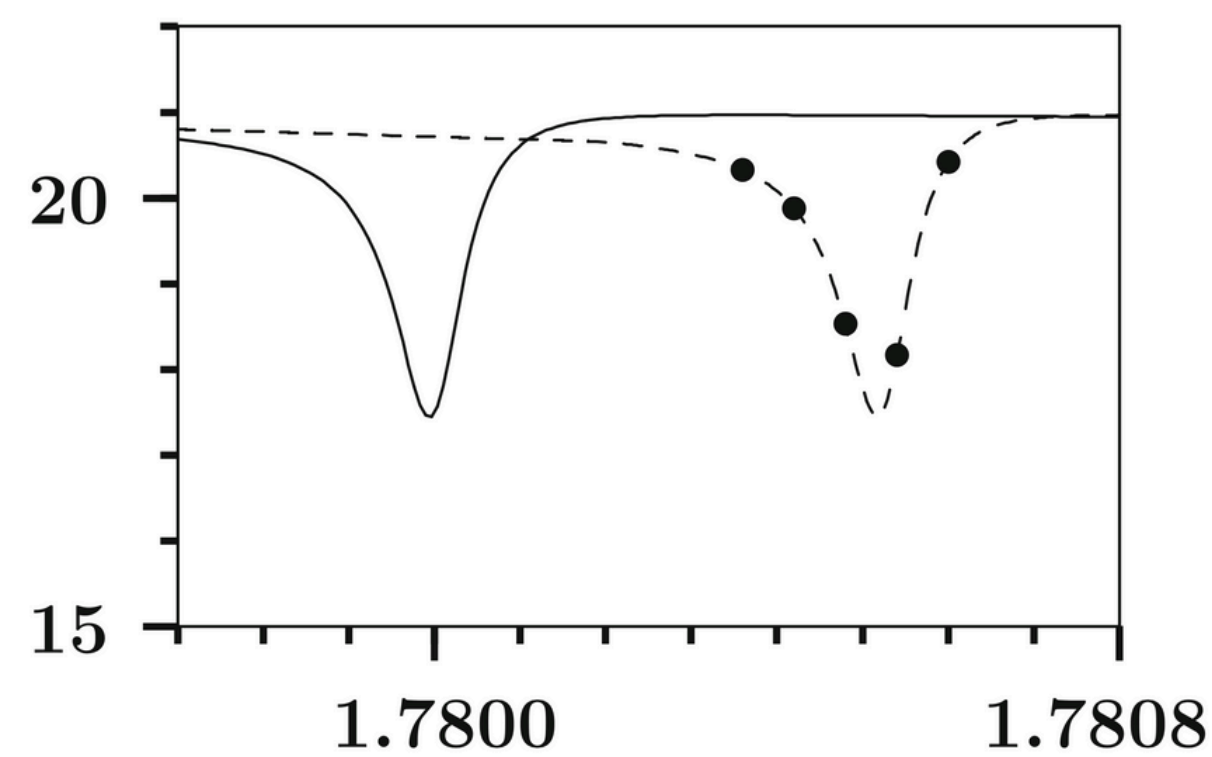
Arbitrary choice of the channel radius

Problematic functional form of the S-matrix

Some approaches to the problem[†]

Why isn't it enough?

Arbitrary choice of the channel radius



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ON THE PRACTICAL APPLICATION

Int

R-matrix Methods with an application to $^{12}\text{C}(\alpha, \gamma)^{16}\text{O}$

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Eff

2024 Fall Meeting of the APS Division of Nuclear Physics
Sunday–Thursday, October 6–10, 2024; Boston, Massachusetts

Session F17: Nuclear Theory III
2:00 PM–3:48 PM, Tuesday, October 8, 2024
Hilton Boston Park Plaza Room: Beacon Hill, 4th Floor

Chair: Kyle Beyer, Michigan State University

Abstract: F17.00006 : Investigations of channel radius sensitivity in R-matrix analysis*
3:00 PM–3:12 PM

(Received 9 August 2013; published 31 January 2014)

Author's proposed approach[†]

Combine both methods!

Preliminary considerations

- ✓ Single-channel problem
- ✓ Artificial experimental data
- ✓ Noro-Taylor potential with a Coulomb tail

Reference list

- † Vaandrager, P., Lekala, M.L., & Rakityansky, S.A. (2025). R-matrix type parametrization of the Jost function for extracting the resonance parameters from scattering data. European Physical Journal A, 61, 79.
- * S.A. Rakityansky, Jost Functions in Quantum Mechanics: A Unified Approach to Scattering, Bound, and Resonant State Problems (Springer, Berlin, 2022)
- ‡ Canto, L.F., & Hussein, M.S. (2013). Scattering Theory of Molecules, Atoms and Nuclei. WORLD SCIENTIFIC. ISBN 978-9814329835

Theoretical background^{†‡*}

Jost functions

Regular solution of the radial Schrödinger equation in the asymptotic limit

$$\Phi_\ell(k, r) \rightarrow \left[H_\ell^{(+)}(\eta, kr) f_\ell^{(\text{in})}(\eta, k) + H_\ell^{(-)}(\eta, kr) f_\ell^{(\text{out})}(\eta, k) \right]$$

Formal definition

Amplitudes of the incoming and outgoing Haenkel-Coulomb functions in the asymptotic behaviour of this solution

Theoretical background^{†‡*}

How does the Jost function look?

$$f_{\ell}^{(\text{in/out})}(E) = e^{\mp i\delta_{\ell}^c} k^{\ell} \left\{ \frac{k}{D_{\ell}(\eta, k)} A_{\ell}(E) - [M(k) \pm i] D_{\ell}(\eta, k) B_{\ell}(E) \right\}$$

$$\delta_{\ell}^c(\eta) = \frac{1}{2i} \ln \left(\frac{\Gamma(\ell + 1 + i\eta)}{\Gamma(\ell + 1 - i\eta)} \right) \quad D_{\ell}(\eta, k) = C_{\ell}(\eta) k^{\ell+1} \quad M(k) = \frac{2\eta h(\eta)}{C_0^2(\eta)}$$

$$C_{\ell}(\eta) = \frac{e^{-\pi\eta/2}}{\Gamma(\ell + 1)} \exp \left\{ \frac{1}{2} [\ln \Gamma(\ell + 1 + i\eta) + \ln \Gamma(\ell + 1 - i\eta)] \right\}$$

$$h(\eta) = \frac{1}{2} [\psi(1 + i\eta) + \psi(1 - i\eta)] - \ln \eta \quad \psi(z) = \frac{\Gamma'(z)}{\Gamma(z)}$$

Theoretical background^{†‡*}

How does the Jost function look?

$$f_{\ell}^{(\text{in/out})}(E) = e^{\mp i\delta_{\ell}^c} k^{\ell} \left\{ \frac{k}{D_{\ell}(\eta, k)} A_{\ell}(E) - [M(k) \pm i] D_{\ell}(\eta, k) B_{\ell}(E) \right\}$$

$$A_{\ell}(E) \approx \sum_{n=0}^N a_n(\ell, E_0) (E - E_0)^n \qquad B_{\ell}(E) \approx \sum_{n=0}^N b_n(\ell, E_0) (E - E_0)^n$$

Fitting parameters!

Theoretical background^{†‡*}

This gives us a clue about the proposed modification

$$~~A_\ell(E) \approx \sum_{n=0}^N a_n(\ell, E_0) (E - E_0)^n~~$$

$$~~B_\ell(E) \approx \sum_{n=0}^N b_n(\ell, E_0) (E - E_0)^n~~$$

This Taylor series needs to be replaced!

The R-matrix is the solution

Theoretical background^{†‡*}

Relation between the R-matrix and the Jost functions

$$f_{\ell}^{(\text{in/out})}(E) = \frac{\pm i e^{\pm i \delta_{\ell}} k^{\ell}}{Q_{\ell}(E, B_R)} \left\{ H_{\ell}^{(\pm)}(\eta, ka) - \left[a \frac{d}{dr} H_{\ell}^{(\pm)}(\eta, kr) \Big|_{r=a} - B_R H_{\ell}^{(\pm)}(\eta, ka) \right] R_{\ell}(E, B_R) \right\}$$

$$Q_{\ell}(E, \mathbf{B}_R) = \sum_{n=1}^N \frac{\lambda_{n\ell}(\mathbf{B}_R) \gamma_{n\ell}(\mathbf{B}_R)}{E_{n\ell} - E}$$

Let's remember our objective

$$f_{\ell}^{(\text{in/out})}(E) = e^{\mp i \delta_{\ell}^c} k^{\ell} \left\{ \frac{k}{D_{\ell}(\eta, k)} A_{\ell}(E) - [M(k) \pm i] D_{\ell}(\eta, k) B_{\ell}(E) \right\}$$

Theoretical background^{†‡*}

Relation between the R-matrix and the Jost functions

- a** Comparison of the terms in both expressions
- b** Rewriting Q and the R-matrix as sums of polynomials
- c** Replace those expressions in the coefficients

$$A_\ell(E) = \sum_{n=0}^N \alpha_{n\ell} \mathcal{P}_n(E)$$

$$B_\ell(E) = \sum_{n=0}^N \beta_{n\ell} \mathcal{P}_n(E)$$

$$\mathcal{P}_n(E) = \frac{\mathcal{P}_0(E)}{E_{n\ell} - E}$$

$$\mathcal{P}_0(E) = \prod_{n=1}^N (E_{n\ell} - E)$$

New fitting parameters!

Understanding the method[†]

In a typical scenario

$$\sigma_{\text{total}}(E_i) \pm \Delta_i$$
$$i = 1, 2, \dots, N^{(\text{data})}$$

The steps of the method are

 Parametrize the functions

$$A_\ell(E) = \sum_{n=0}^N \alpha_{n\ell} \mathcal{P}_n(E) \qquad B_\ell(E) = \sum_{n=0}^N \beta_{n\ell} \mathcal{P}_n(E)$$

Understanding the method[†]

b Substituting them in the S-matrix

$$S_\ell(E) = e^{2i\delta_\ell^c} \frac{kA_\ell(E) - [M(k) - i]D_\ell^2(\eta, k)B_\ell(E)}{kA_\ell(E) - [M(k) + i]D_\ell^2(\eta, k)B_\ell(E)}$$

c Obtain the total cross-section by summing over several partial cross-sections

$$\sigma_{\text{total}}(E) = \sum_{\ell=0}^{\ell_{\text{max}}} \sigma_\ell(E) \qquad \sigma_\ell(E) = \frac{\pi}{k^2} (2\ell + 1) |S_\ell(E) - 1|^2$$

Understanding the method[†]

Little recap

What do we want?

Fitting the total cross-section experimental data to extract resonance parameters

How are we going to do it?

By fitting the parameters in the S-matrix

$$A_\ell(E) = \sum_{n=0}^N \alpha_{n\ell} \mathcal{P}_n(E) \quad B_\ell(E) = \sum_{n=0}^N \beta_{n\ell} \mathcal{P}_n(E) \quad \mathcal{P}_n(E) = \frac{\mathcal{P}_0(E)}{E_{n\ell} - E}$$

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- ‡ Canto, L.F., & Hussein, M.S. (2013). Scattering Theory of Molecules, Atoms and Nuclei. WORLD SCIENTIFIC. ISBN 978-9814329835
- # S.A. Sofianos, S.A. Rakityansky, Exact method for locating potential resonances and Regge trajectories. J. Phys. A: Math. Gen. 30, 3725–3737 (1997)

Understanding the method[†]

Let's talk about the implementation

Artificial experimental data

- Exact total cross-section using a given potential

$$V(r) = \frac{\hbar^2}{2\mu} \left(7.5r^2 e^{-r} + \frac{2k\eta}{r} \right)$$

Understanding the method[†]#

Let's talk about the implementation

Artificial experimental data

Partial cross-section can be calculated using any of the many known methods for solving a two-body scattering problem

Direct calculation of the Jost functions

Complex coordinate method to solve the Schrödinger equation

Advantajes

No need of approximations, expansions or variational procedures

Understanding the method^{†#}

Direct calculation of the Jost functions

$$r \in [0, 1]$$

$$r \in [1, 30]$$

a Complex continuation of the radial coordinate

b We look for a solution near $r=0$ in the form

$$\Phi_\ell(k, r) = F_\ell(\eta, kr)A_\ell(\eta, k, x, \theta) + G_\ell(\eta, kr)B_\ell(\eta, k, x, \theta)$$

c We look for a solution for large r in the form

$$\Phi_\ell(k, r) = \frac{1}{2} \left[H_\ell^{(+)}(\eta, kr) \mathcal{F}_\ell^{(+)}(k, x, \theta) + H_\ell^{(-)}(\eta, kr) \mathcal{F}_\ell^{(-)}(k, x, \theta) \right]$$

Understanding the method^{†#}

Direct calculation of the Jost functions

d We need to solve this first-order coupled differential equations for each solution

$$\partial_x \mathcal{F}_\ell^{(+)}(k, x, \theta) = -\frac{e^{i\theta}}{2ik} H_\ell^{(+)}(\eta, kr) V(r) \left[H_\ell^{(+)}(\eta, kr) \mathcal{F}_\ell^{(+)}(k, x, \theta) + H_\ell^{(-)}(\eta, kr) \mathcal{F}_\ell^{(-)}(k, x, \theta) \right]$$

$$\partial_x \mathcal{F}_\ell^{(-)}(k, x, \theta) = \frac{e^{i\theta}}{2ik} H_\ell^{(-)}(\eta, kr) V(r) \left[H_\ell^{(+)}(\eta, kr) \mathcal{F}_\ell^{(+)}(k, x, \theta) + H_\ell^{(-)}(\eta, kr) \mathcal{F}_\ell^{(-)}(k, x, \theta) \right]$$

$$\partial_x A_\ell(\eta, k, x, \theta) = \frac{e^{i\theta}}{k} G_\ell(\eta, kr) V(r) [F_\ell(\eta, kr) A_\ell(\eta, k, x, \theta) + G_\ell(\eta, kr) B_\ell(\eta, k, x, \theta)]$$

$$\partial_x B_\ell(\eta, k, x, \theta) = -\frac{e^{i\theta}}{k} F_\ell(\eta, kr) V(r) \left[A_\ell(\eta, k, 0, \theta) = 1, \quad B_\ell(\eta, k, 0, \theta) = 0 \right]$$

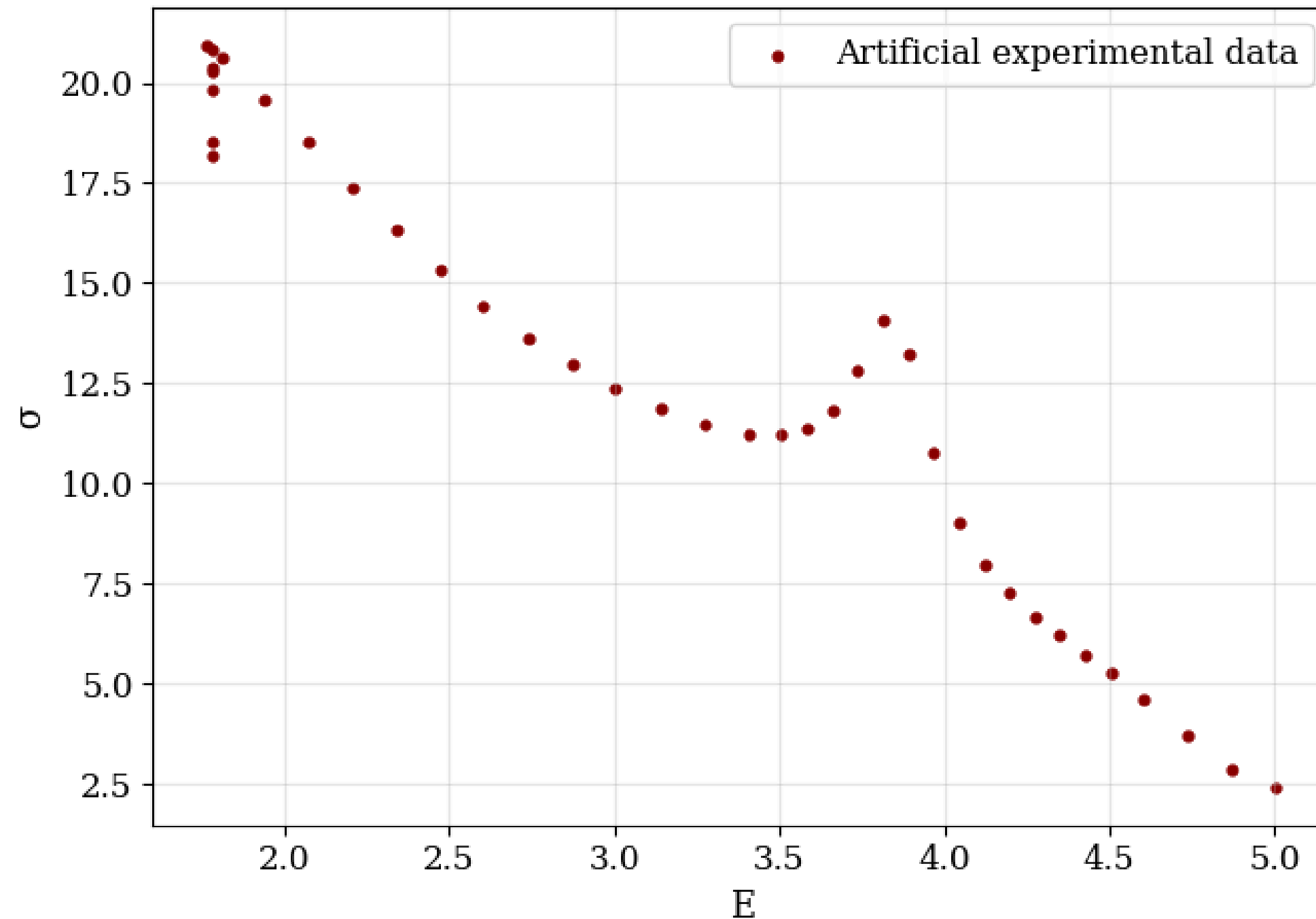
Understanding the method

Comments on the reproduction of the exact curve

- a** Coupling of two differential equation systems in Python
- b** Problems with Coulomb wave functions in cmath library
- c** Number of points to get significant results

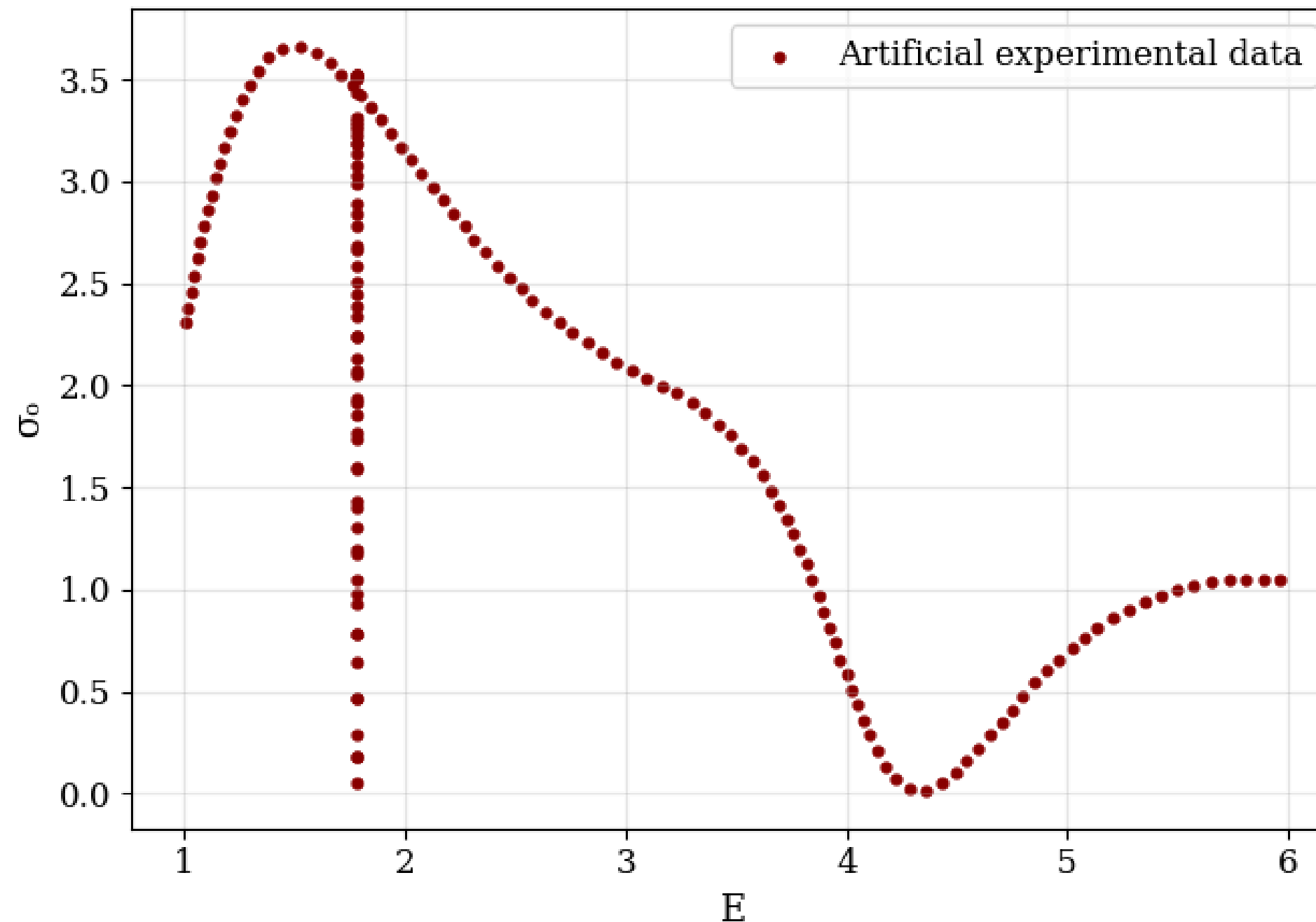
Understanding the method

Exact total cross-section curve



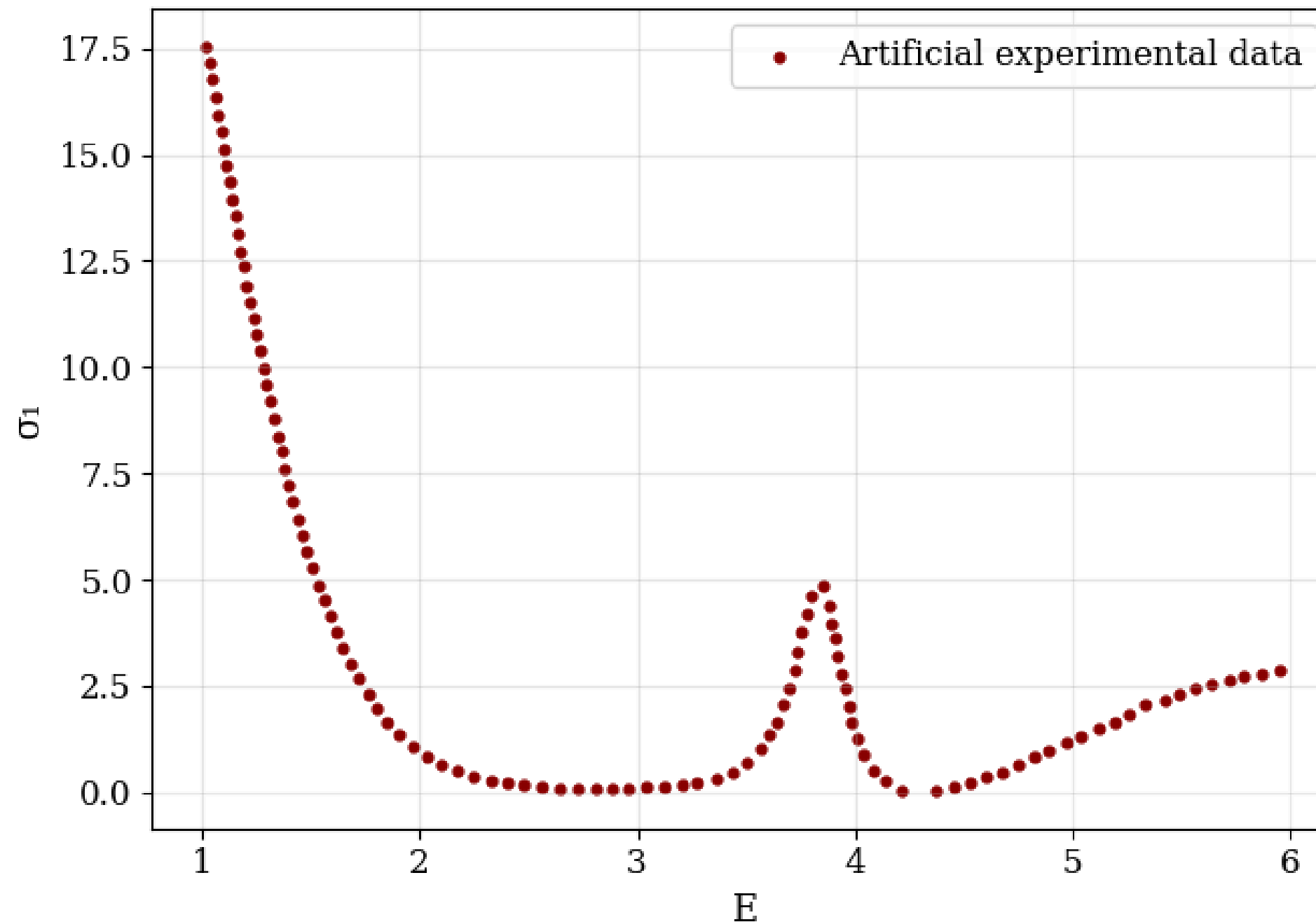
Understanding the method

Exact partial cross-section curves



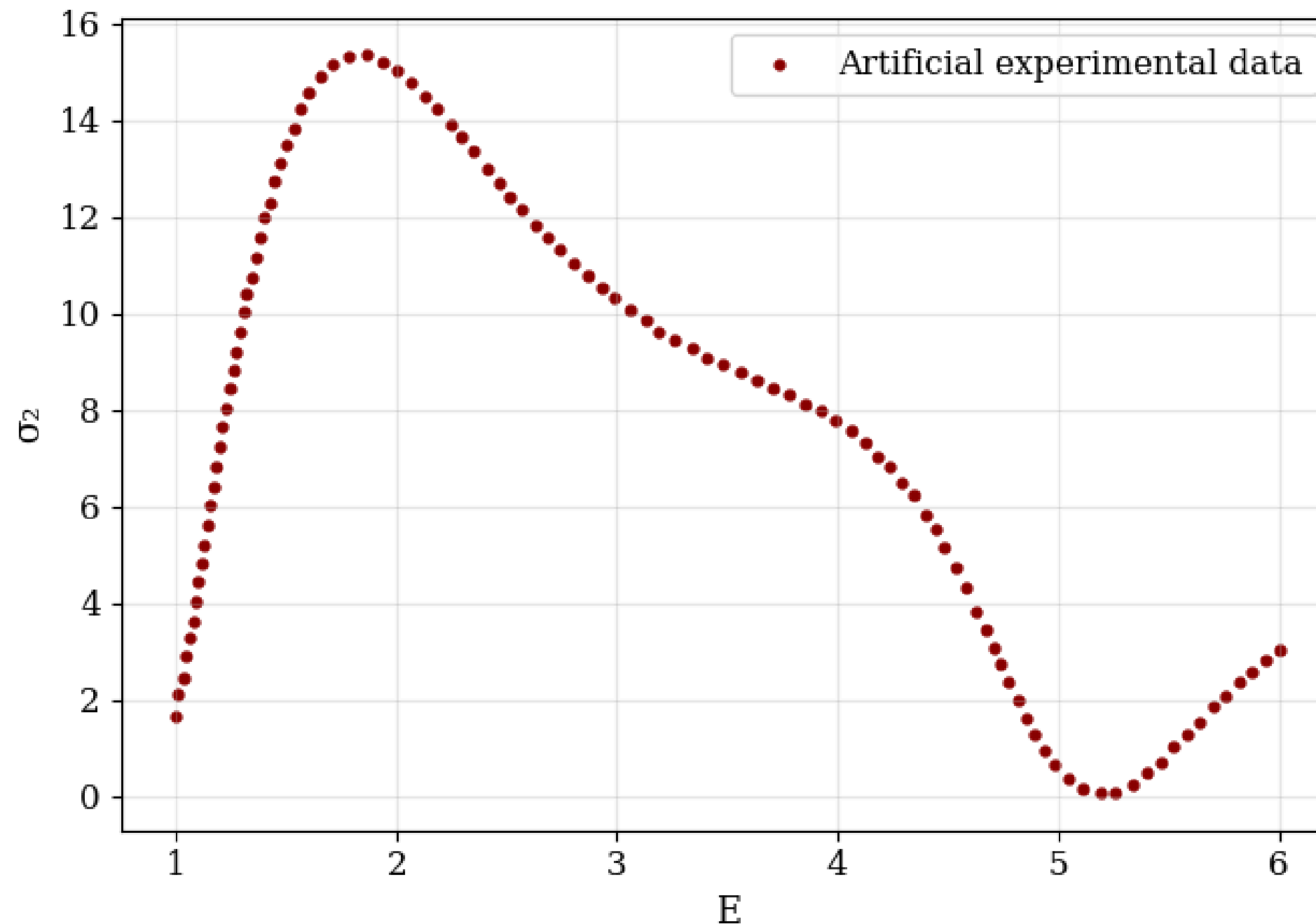
Understanding the method

Exact partial cross-section curves



Understanding the method

Exact partial cross-section curves



Understanding the method[†]#

Fitting procedure

- a** 40 energy data points between 1.7-5 to fit
- b** Minimization routine MINUIT
 - i** Find optimal values by minimizing

$$\chi^2 = \sum_{i=1}^{N(\text{data})} \left[\frac{\sigma_{\text{total}}(E_i) - \sigma_{\text{fit}}(E_i)}{\Delta_i} \right]^2$$

- c** Find the zeros of the Jost functions
- d** Extract resonance parameters

Understanding the method

Comments on the reproduction of the fitting curve

- a** We tested all the functions required for the fitting
- b** The Coulomb barrier factor was wrong

$$C_\ell(\eta) = \frac{e^{-\pi\eta/2}}{\Gamma(\ell+1)} \exp \left\{ \frac{1}{2} [\ln \Gamma(\ell+1+i\eta) + \ln \Gamma(\ell+1-i\eta)] \right\} \frac{1}{(2\ell+1)!!}$$

$$C_\ell(\eta) = \frac{2^\ell e^{-\pi\eta/2}}{\Gamma(2\ell+2)} |\Gamma(\ell+1 \pm i\eta)|$$

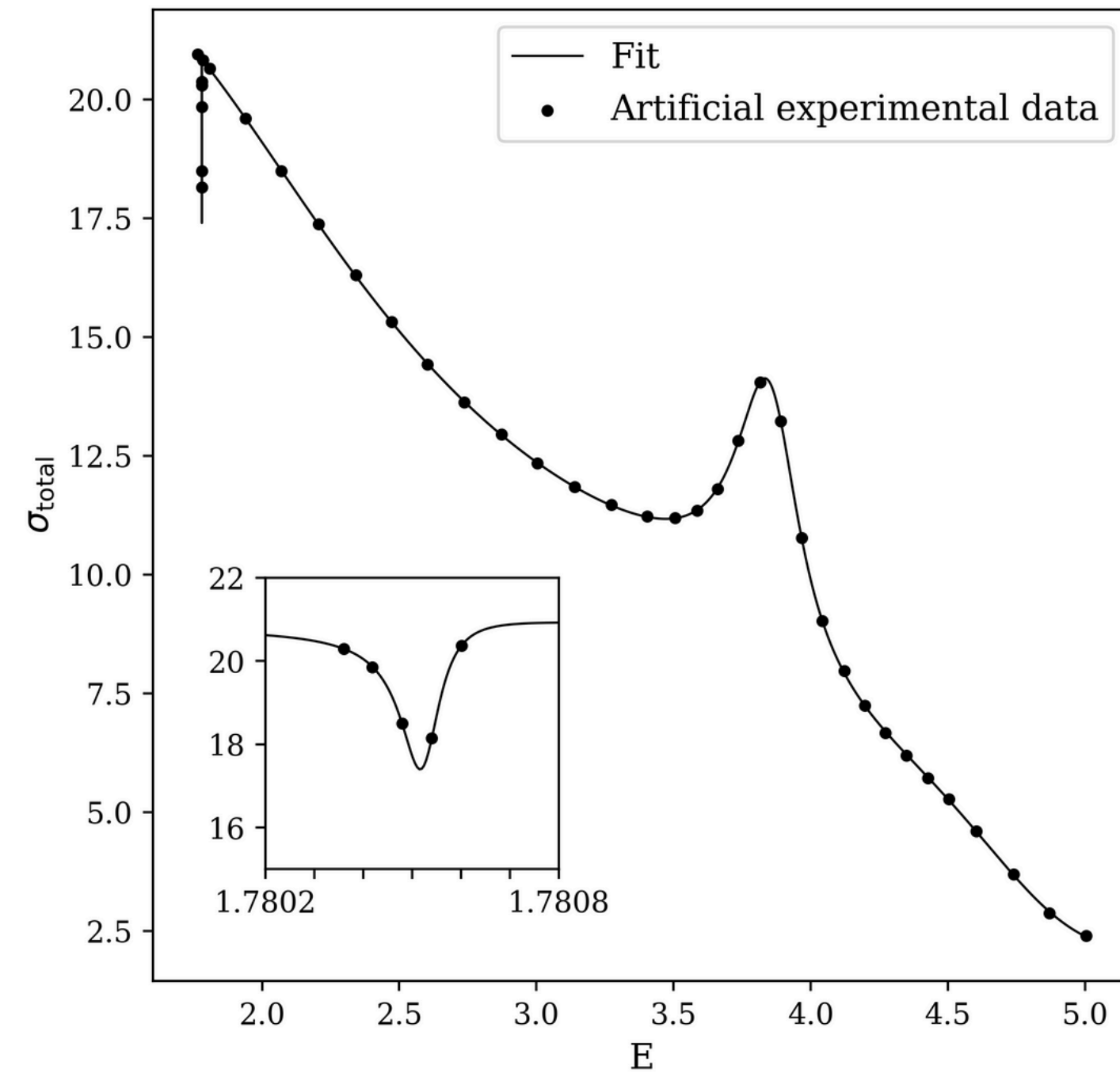
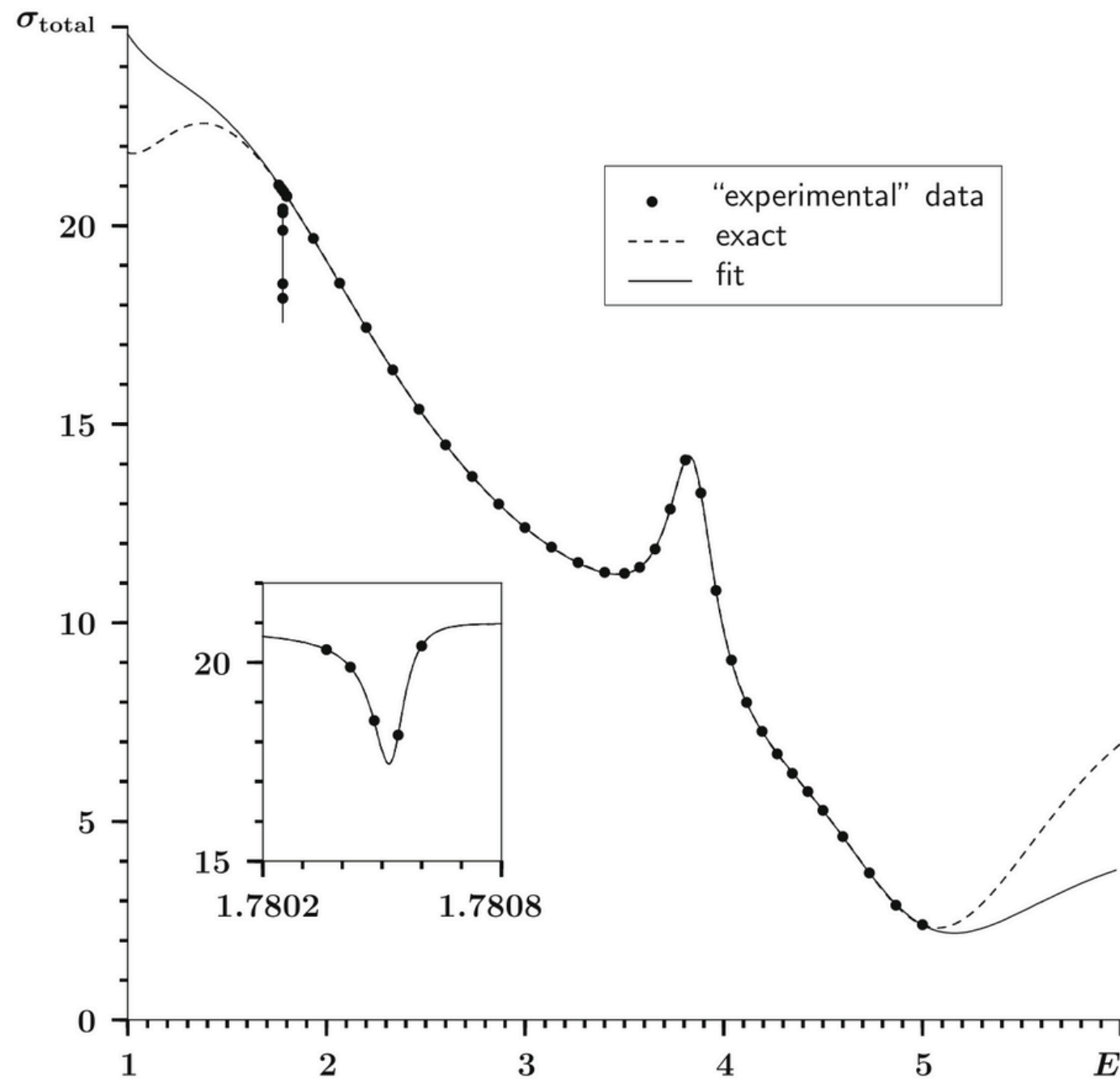
- c** η was incorrectly defined

$$\eta = -\frac{1}{2k}$$

$$\eta = -\frac{1}{k}$$

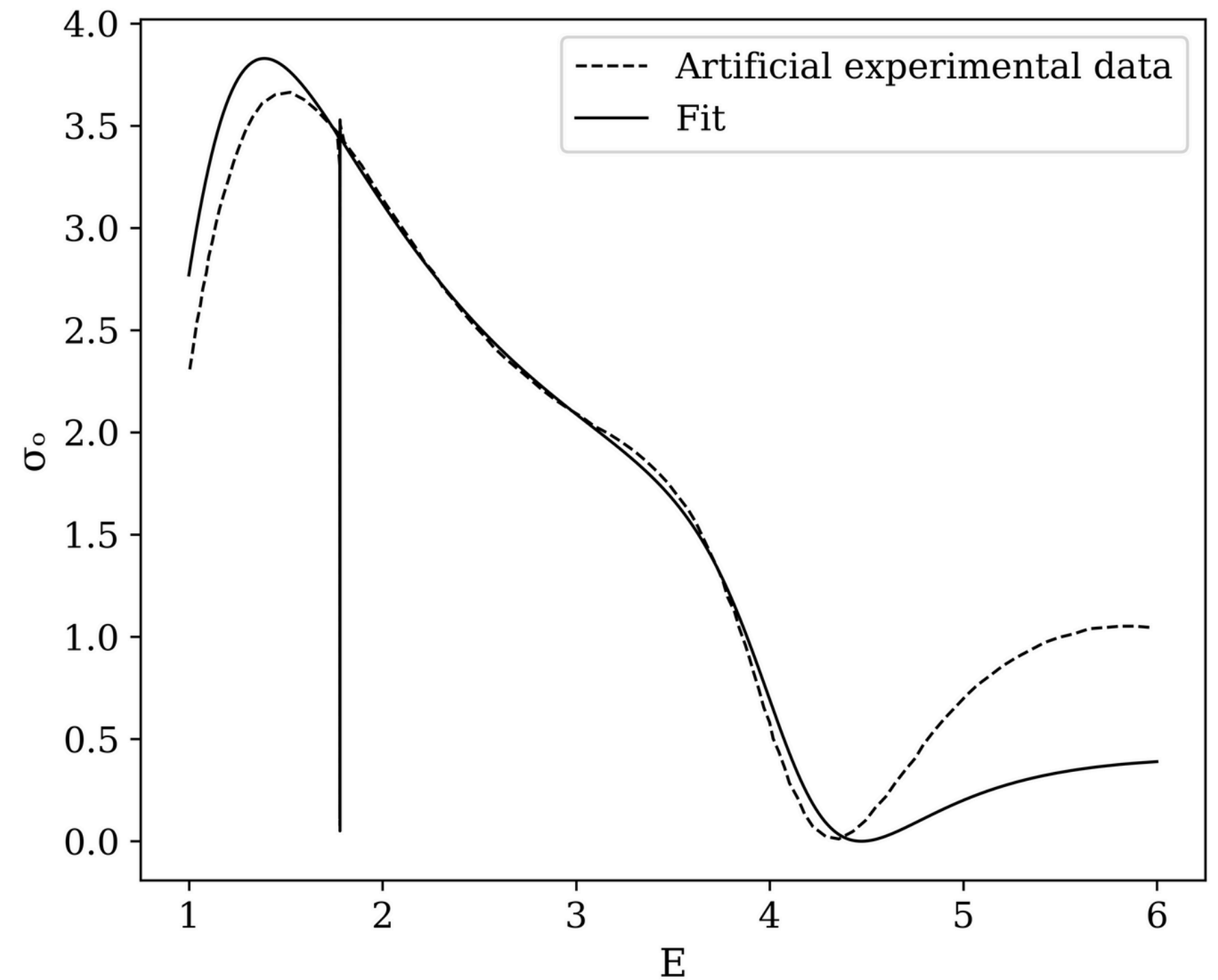
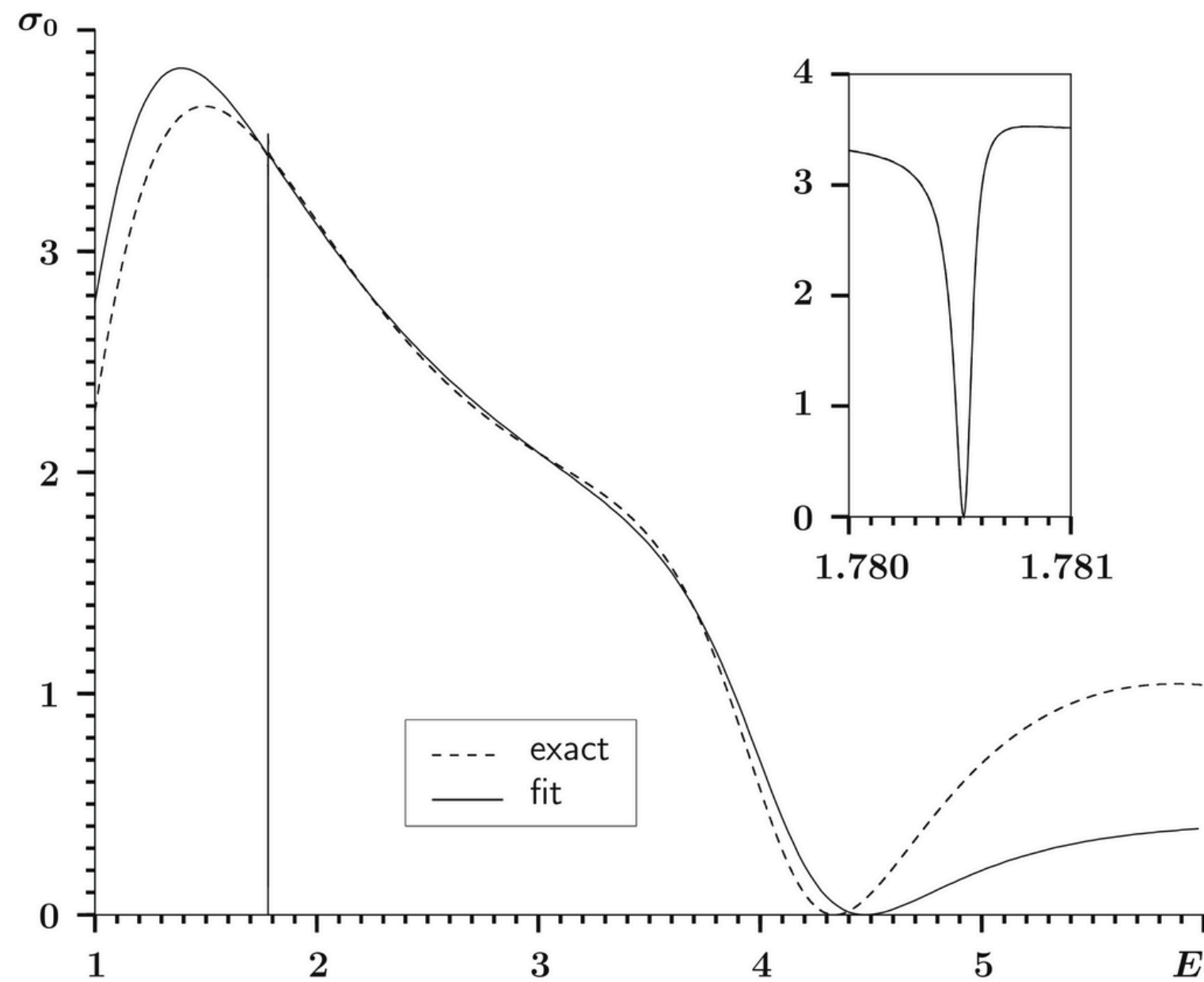
Results

Total cross-section



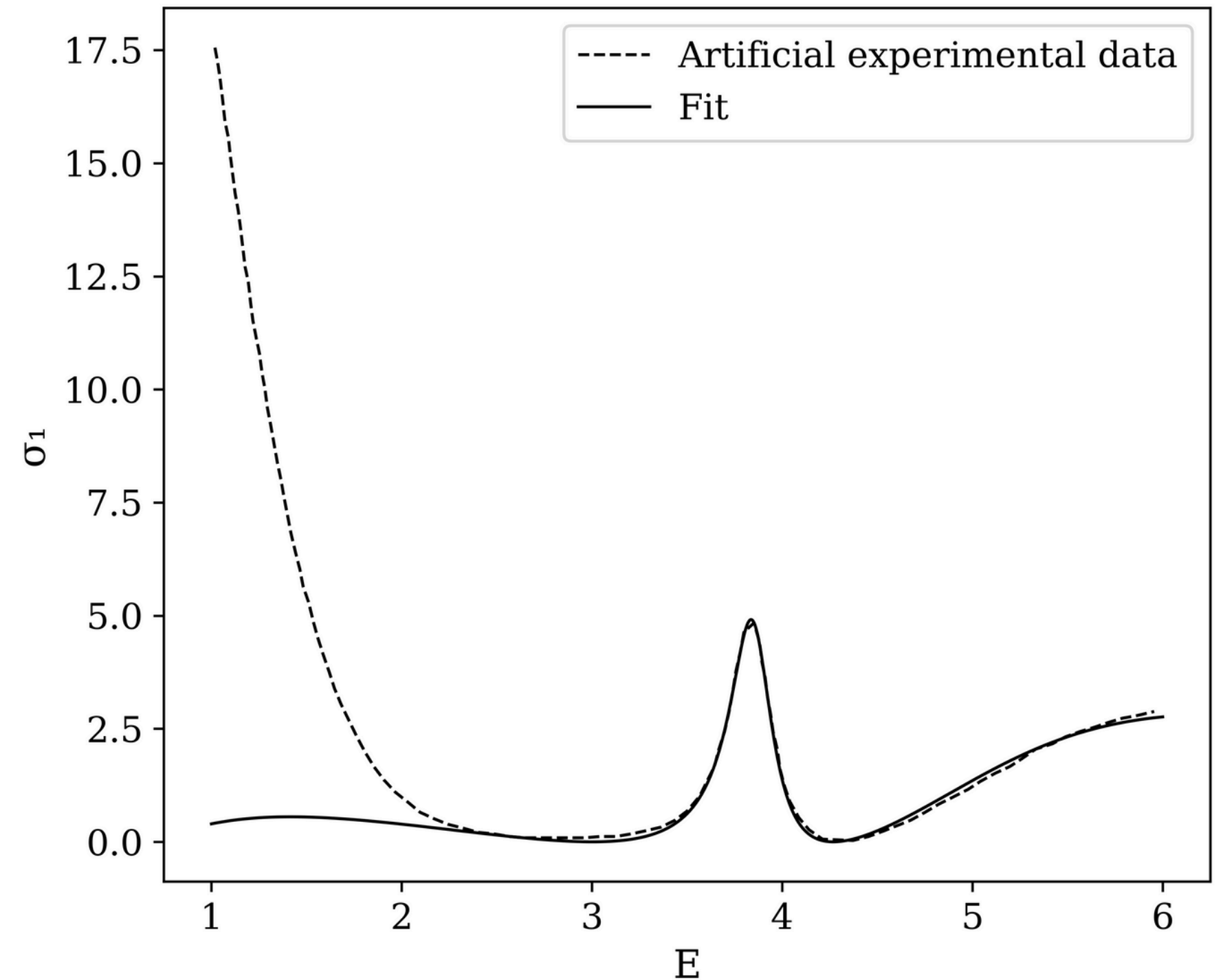
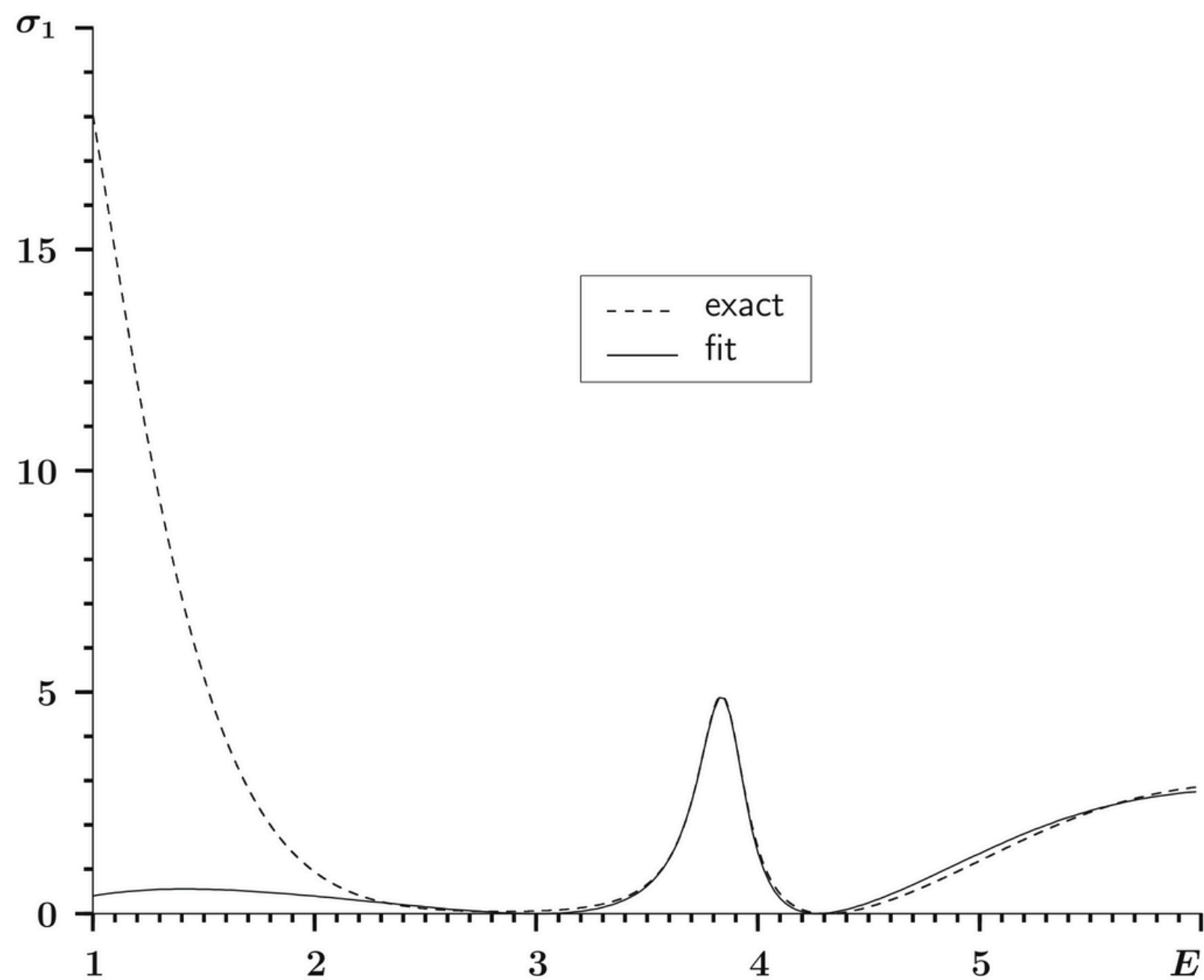
Results

Partial cross-sections



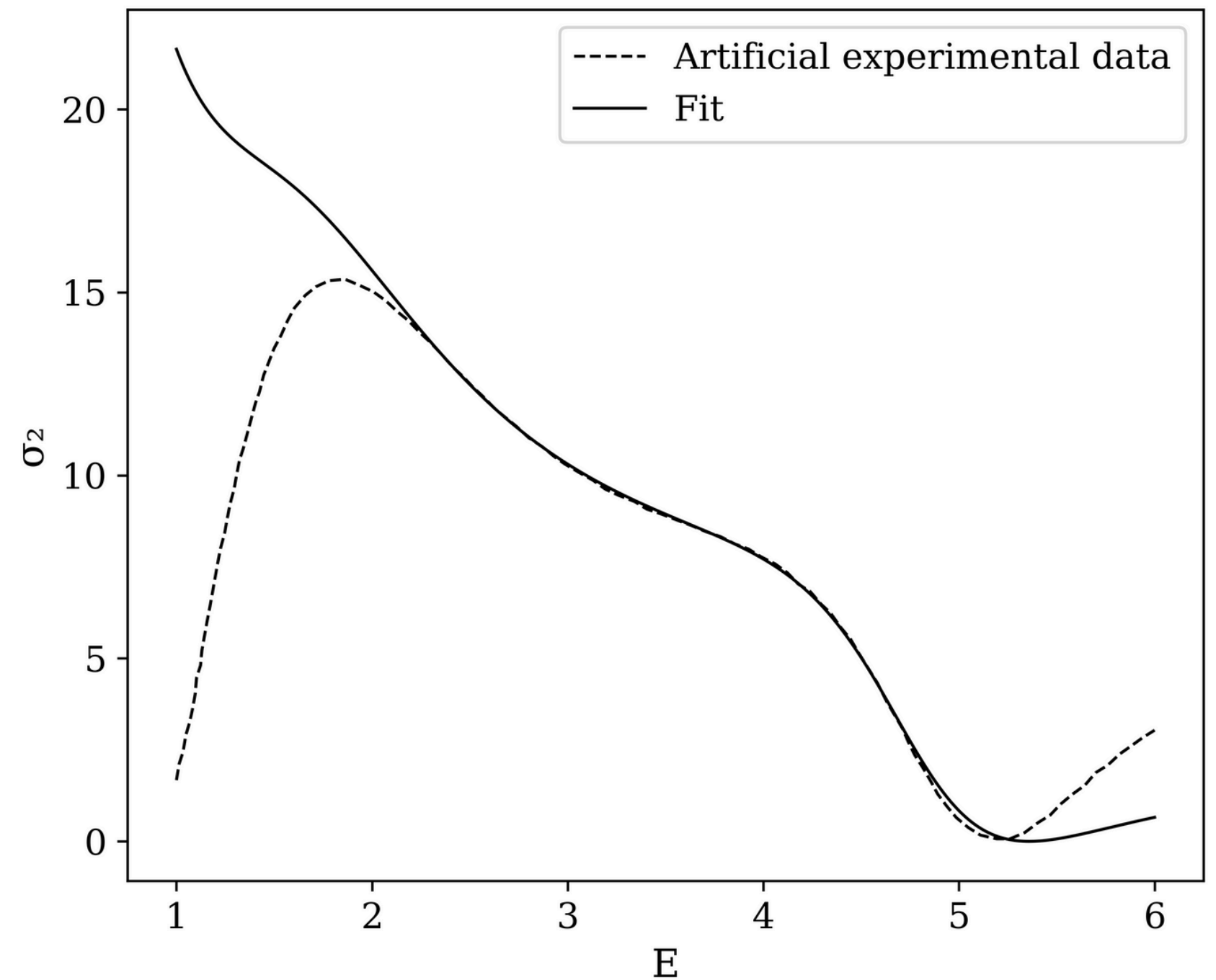
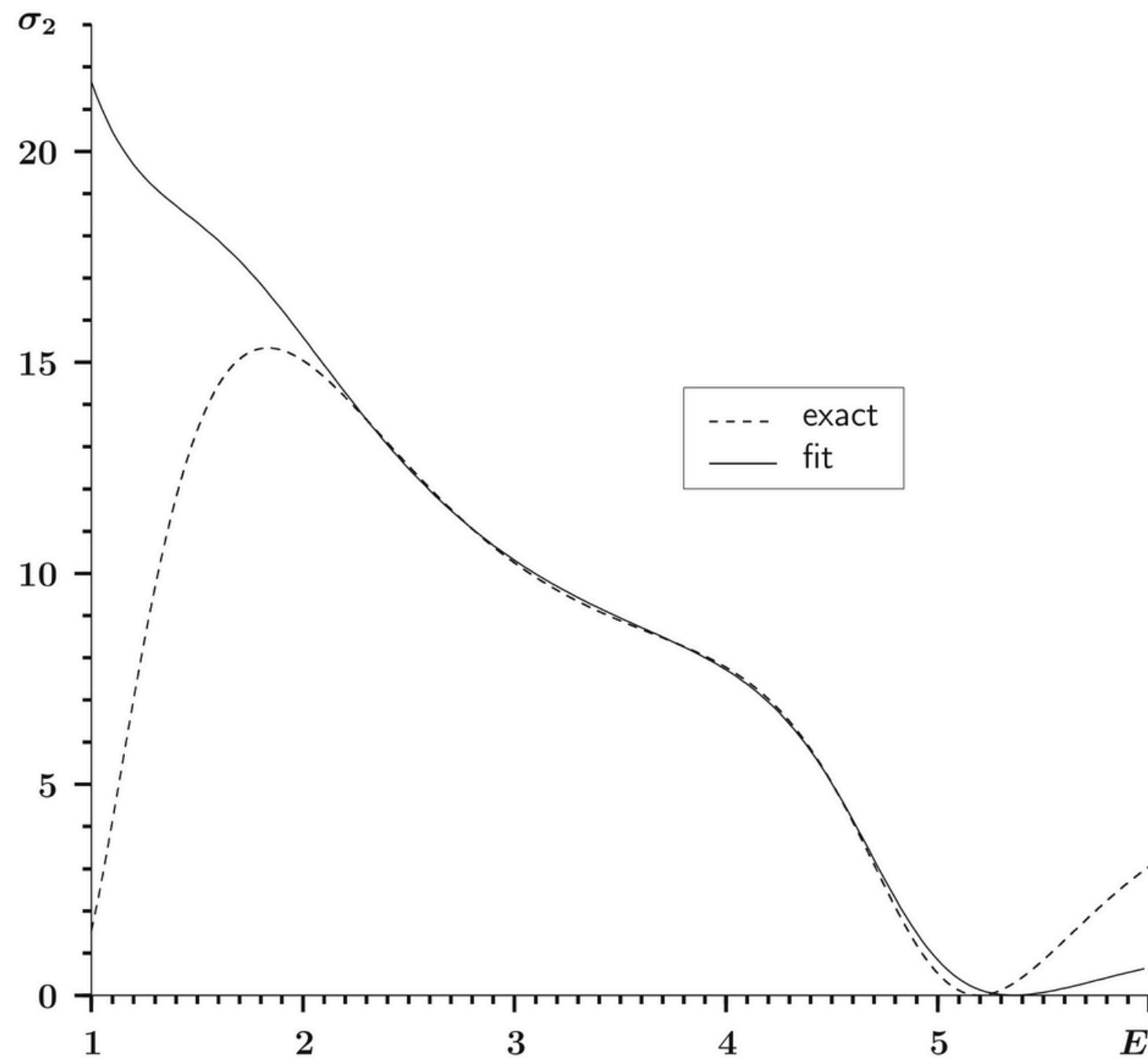
Results

Partial cross-sections



Results

Partial cross-sections



Results

Resonance parameters





ℓ	Exact		Paper fit		Our fit	
	E_r	Γ	E_r	Γ	E_r	Γ
0	1.7805	9.5719×10^{-5}	1.7805	9.6022×10^{-5}	1.7805	9.6022×10^{-5}
	4.1015	1.1563	4.1001	1.1265	4.0918	1.1242
1	3.8480	0.2754	3.8541	0.2724	3.8542	0.2736
2	4.9005	1.5675	4.8029	1.4238	4.7940	1.4365

Results

Resonance parameters

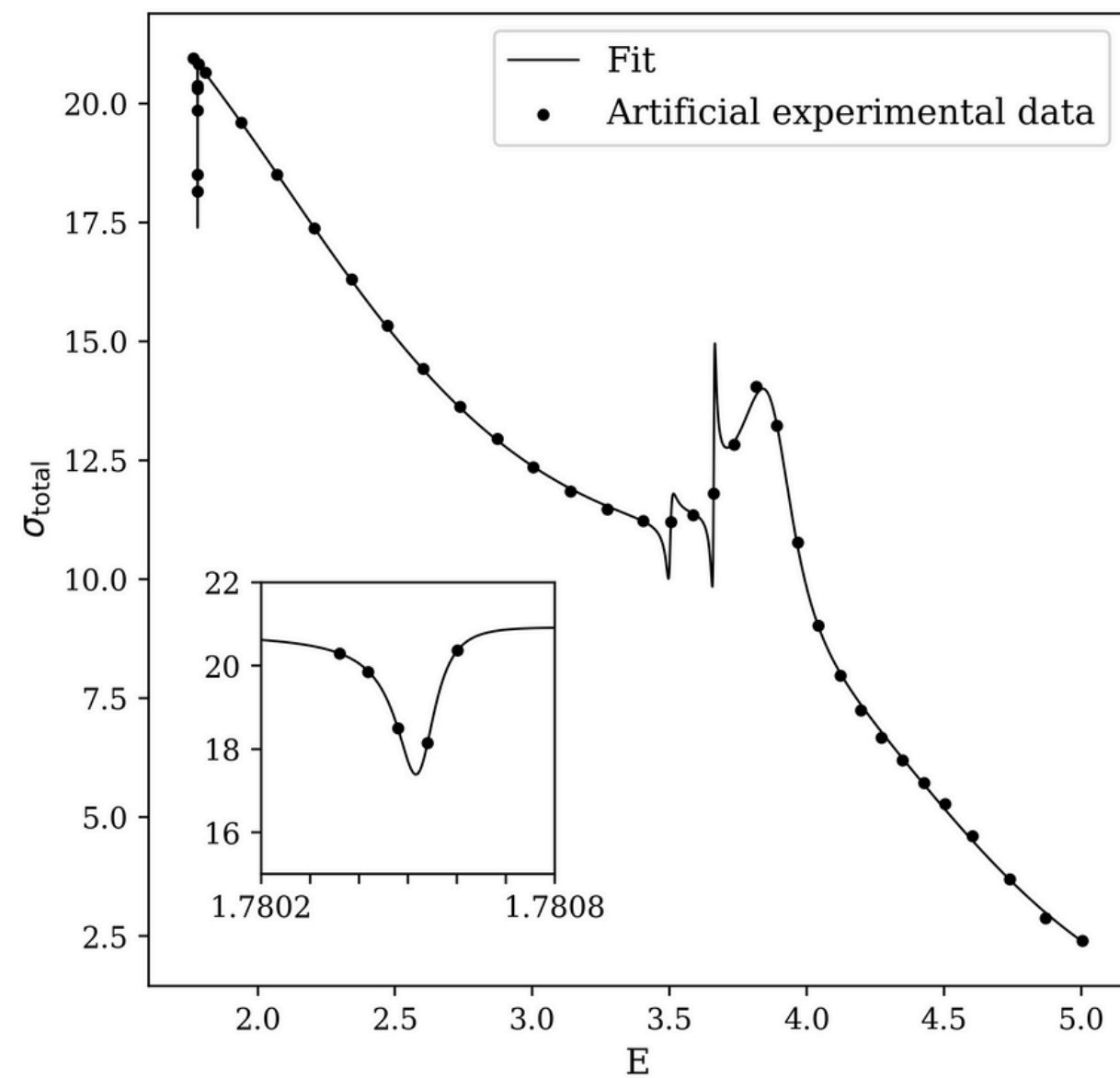
ℓ	Reported errors (%)			
	Paper fit		Our fit	
	E_r	Γ	E_r	Γ
0	6.5×10^{-6}	0.32	6.5×10^{-6}	0.32
	0.03	2.66	0.24	2.86
1	0.16	1.08	0.16	0.64
2	1.99	9.18	2.17	8.36

Conclusions

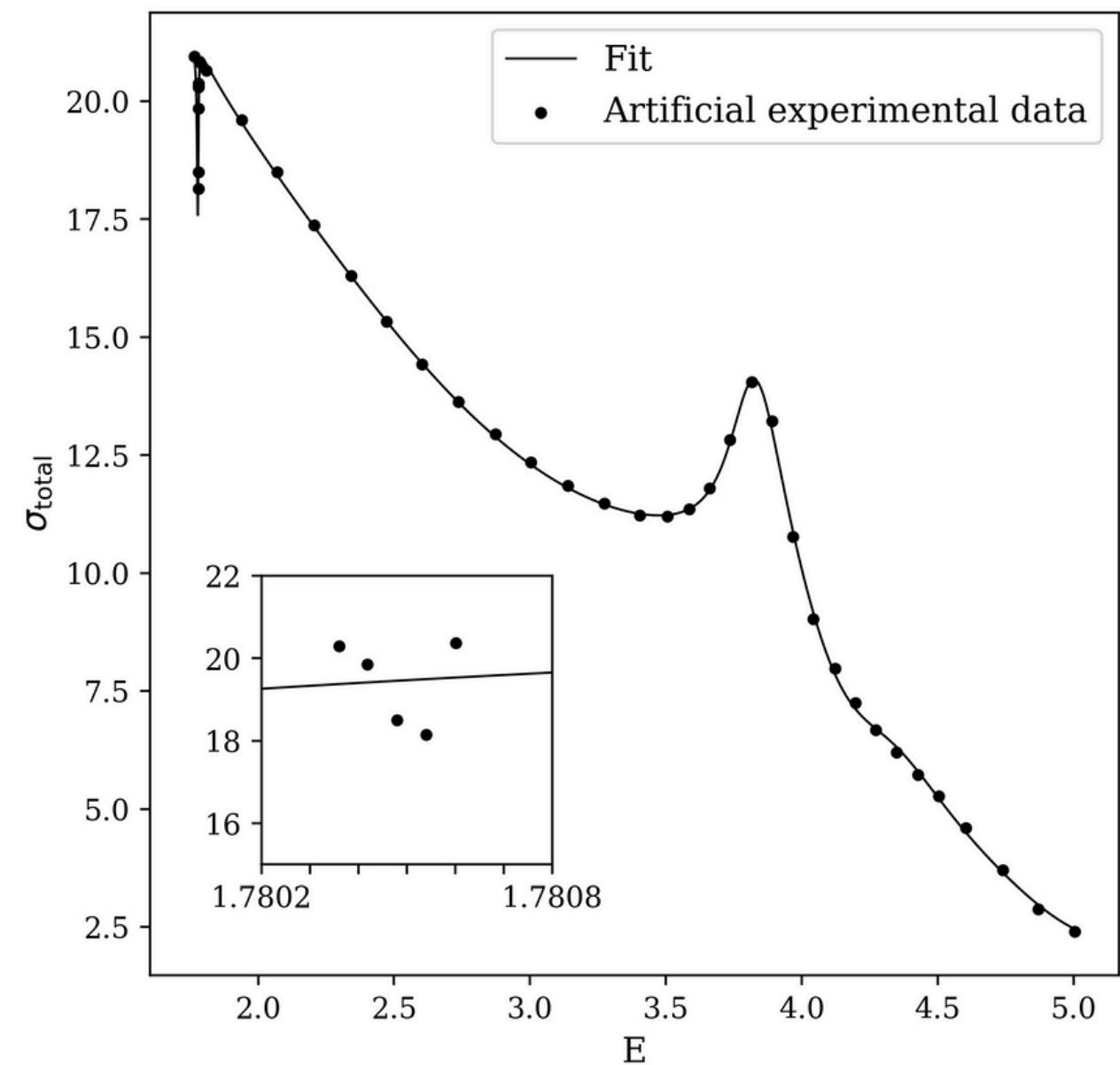
-  We successfully reproduced the method described in the article, as well as the associated **fitting curves** and **resonance parameters**
 -  MINUIT was actually pretty efficient at fitting many parameters
-  The theoretical curve was not successfully reproduced, mainly due to **computational limitations** during the implementation
-  The method has been shown to **outperform the original approaches**. However, its performance is still limited by the **high number of fitting parameters** and its **sensitivity** to specific parameter choices

Some interesting result

Changing β_{10} from 10 to 15



Changing E_{10} from 1.78 to 1.7

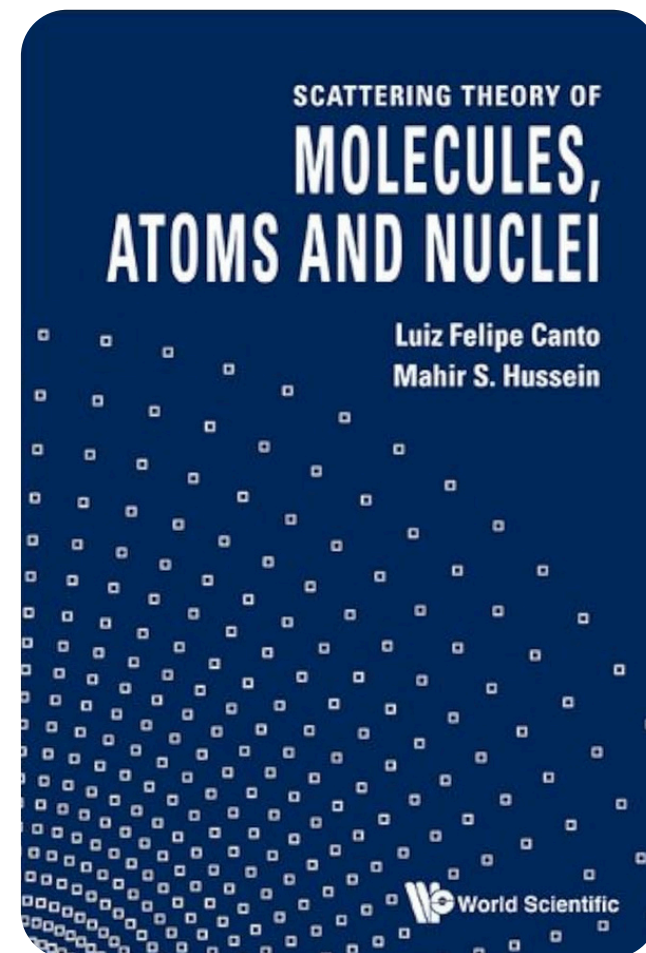


Undiscussed statements

The poles of the S-matrix contain the information about the resonance

The S-matrix may fail to produce physical resonances when it is obtained from models based on polynomial sums

Section 8.1.4



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On beautiful analytic structure of the S-matrix

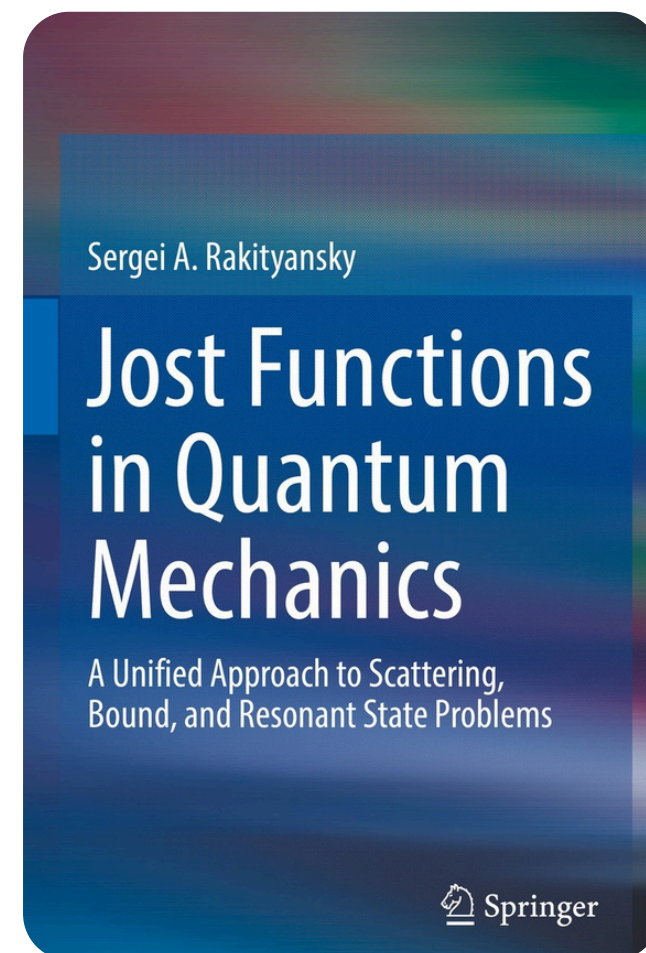
Alexander Moroz and Andrey E Miroshnichenko

Undiscussed statements

The S-matrix is the ratio of Jost functions

The Jost functions can be expressed in terms of the R-matrix

Section 4.6



The Jost function and Siegert pseudostates from R -matrix calculations at complex wavenumbers

P. Vaandrager, J. Dohet-Eraly and J-M. Sparenberg

*Nuclear Physics and Quantum Physics, C.P. 229,
Université libre de Bruxelles (ULB), B 1050 Brussels, Belgium*

June 22, 2023

Undiscussed statements

The boundary parameter can be set to zero

Section 3.4

IOP PUBLISHING

Rep. Prog. Phys. **73** (2010) 036301 (44pp)

REPORTS ON PROGRESS IN PHYSICS

[doi:10.1088/0034-4885/73/3/036301](https://doi.org/10.1088/0034-4885/73/3/036301)

The *R*-matrix theory

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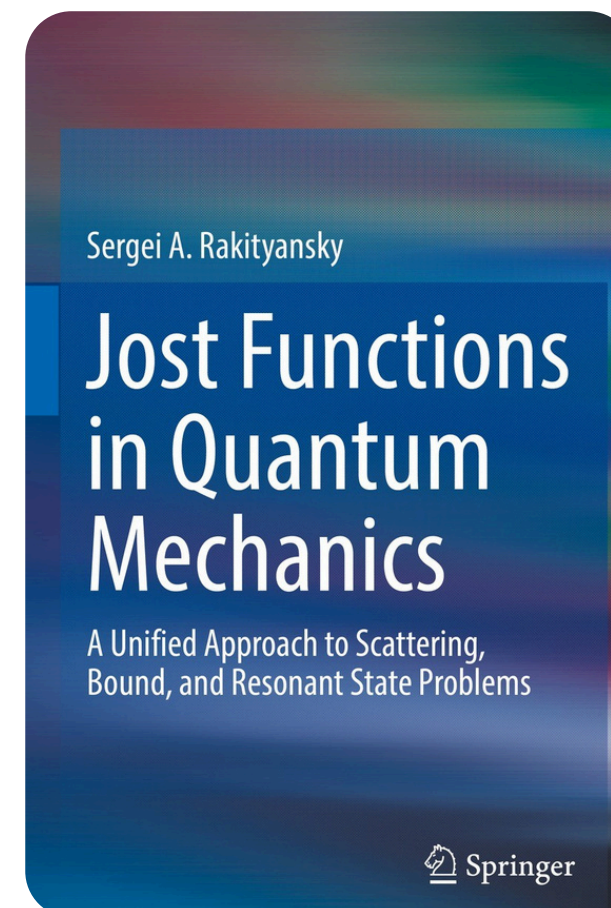
Final comment

Poles of the S-matrix related to
resonance parameters



Complex calculus basis

Chapter 4



Thank you!

