

# On background dependent formalisms for general relativity

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## Background dependent formalisms

- General relativity is background independent
  - $g_{ab}$  is defined without reference to an underlying geometric structure
- Why introduce a background?
  - Can choose background with symmetries---potentially useful feature
  - Can build new solutions out of existing ones
  - Uncover new properties of the Einstein field equations (EFE)

# Examples (ordered in increasing generality)

- Kerr-schild perturbation:<sup>\*</sup>  $g_{ab} = \bar{g}_{ab} + \psi k_a k_b$ ,  $k$  null
  - Includes Kerr soln, can linearize vac EFE under some condns.<sup>⊗</sup>
- Extended Kerr-schild:<sup>†</sup>  $g_{ab} = \varphi \bar{g}_{ab} + \psi k_{(a} l_{b)}$ ,  $k, l$  null wrt  $\bar{g}_{ab}$ 
  - Quite general, can reduce nonlinearity in EFE to fifth order.<sup>‡</sup>
- Matrix deformations:<sup>\*†</sup>  $g_{ab} = \bar{g}_{mn} \tau_a{}^m \tau_b{}^n$ 
  - Generalization of KS, extended KS, and tetrad formalism
  - Can write:  $\tau_a{}^m = \tau_a^{(1)n} \tau_n^{(2)p} \dots \tau_q^{(N)m}$  ("Turtles all the way down")

<sup>\*</sup>R. P. Kerr, A. Schild, Gen. Rel. Grav. 41 (10): 2485–2499 (2009).

<sup>†</sup>J. Llosa and D. Soler, CQG 22, 893 (2005); J. Llosa and J. Carot, CQG 26, 055013 (2009)

<sup>\*</sup>S. Capozziello, C. Stornaiolo, Int.J.Geom.Meth.Mod.Phys. 5 (2008) 185-195

<sup>⊗</sup>B. C. Xanthopoulos, J. Math. Phys 19, 1607 (1978)

<sup>‡</sup>A. I. Harte, Phys. Rev. Lett. 113, 261103 (2014)

<sup>†</sup>JCF, S. Carloni, Phys. Rev. D 101, 064002 (2020)\$

# Reference connections

- Introduce reference connection  $\bar{\nabla}$ , then construct:  $W^\sigma{}_{\mu\nu} = \Gamma^\sigma{}_{\mu\nu} - \bar{\Gamma}^\sigma{}_{\mu\nu}$

- Can recover covariant counterpart of  $\Gamma\Gamma$  action:<sup>\*</sup> <sup>⊗</sup>

$$\bar{S}_{EH} = \int_U d^4x \sqrt{-g} \left[ g^{ab} W^c{}_{ad} W^d{}_{cb} - W^a{}_{ab} W^b{}_{cd} g^{cd} + g^{ab} \bar{R}_{ab} + \nabla_a B^a \right]$$

$$B^a := g^{cd} W^a{}_{cd} - g^{ad} W^c{}_{cd}$$

- Explored as early as 1940s<sup>†</sup>, studied by Katz, Bičák, and Lynden-Bell (KBL)<sup>†</sup>
  - Formalism rediscovered at least twice in the past ten years<sup>‡</sup>

<sup>\*</sup>D. Lynden-Bell, J. Katz, and J. Bičák, MNRAS 272, 150 (1995); J. Katz, J. Bičák, and D. Lynden-Bell, Phys. Rev. D 55, 5957 (1997)

<sup>⊗</sup>JCF, S. Chakraborty, Gen. Rel. Grav. 54, 67 (2022)

<sup>†</sup>N. Rosen, Phys. Rev. 57, 147–150 (1940); A. Papapetrou, Proc. Roy. Irish Acad. A 52, 11–23 (1948).

<sup>‡</sup>J. Harada, Phys. Rev. D 101, 024053 (2020), arXiv:2001.06990 [gr-qc]; E. T. Tomboulis, JHEP 09, 145 (2017), arXiv:1708.03977 [hep-th]

# Freedom to choose background: Advantages

- Can choose background with symmetries to get conserved quantities
  - Conserved canonical "energy-momentum" tensor!

For flat background  $\bar{g}_{\mu\nu} = \eta_{\mu\nu}$  and  $\bar{\nabla}_\mu \xi^\nu = 0$ :\*

$$\Pi^\sigma{}_\tau := \sqrt{g/\bar{g}} (2G^\mu{}_\nu - \Theta^\mu{}_\nu), \quad \bar{\nabla}_\sigma (\Pi^\sigma{}_\tau \xi^\tau) = \xi^\tau \bar{\nabla}_\sigma \Pi^\sigma{}_\tau = 0$$

$$\Theta^\alpha{}_\beta := W^{\alpha\sigma\tau} (W_{\sigma\beta\tau} + W_{\tau\beta\sigma}) - W^\sigma{}_{\sigma\tau} (W^\alpha{}_{\beta\tau} + W^\tau{}_{\beta\alpha}) - W^\sigma{}_{\beta\sigma} B^\alpha - \delta^\alpha{}_\beta (R - \nabla_\sigma B^\sigma)$$

- Can we really have our cake and eat it too?
  - Maybe not: Schwarzschild (or PG) with static flat BG:<sup>†</sup>  $\Pi^\sigma{}_\tau = 0$

\*D. Lynden-Bell, J. Katz, and J. Bičák, MNRAS 272, 150 (1995); J. Katz, J. Bičák, and D. Lynden-Bell, Phys. Rev. D 55, 5957 (1997)

<sup>†</sup>JCF, S. Chakraborty, Gen. Rel. Grav. 54, 67 (2022)

<sup>‡</sup>M. Krššák, Phys. Rev. D 110 (2024) 10, 104061

## Freedom is not free: Tradeoffs

- Can choose coordinates for  $\bar{g}_{\mu\nu}$  independently of  $g_{\mu\nu}$ 
  - Pseudotensorial ambiguities  $\Rightarrow$  extra coordinate freedom for  $\bar{g}_{\mu\nu}$
  - Schwarzschild in harmonic coords: "energy" obtained from  $\Pi^{\sigma}_{\tau}$  can either be nice,<sup>†</sup> or negative and divergent<sup>‡</sup>
  - Can appropriate coordinate condition for  $\bar{g}_{\mu\nu}$  fix this?
- Pseudotensor pathologies still present, but tensors nicer than pseudotensors.

*The cake is not edible, but looks appetizing*

<sup>†</sup>N. Nakanishi Prog.Theor.Phys. 75 (1986) 1351,

<sup>‡</sup>JCF, S. Chakraborty, Gen. Rel. Grav. 54, 67 (2022)



# Energy of Schwarzschild in harmonic coords.

- Transform Schwarzschild to Harmonic coordinates:<sup>\*</sup><sup>†</sup>

$$\tilde{r}(r) = c_1(r - m) + c_2 \left\{ (r - m) \ln \left[ 1 - \frac{2m}{r} \right] + 2m \right\}$$

- Integrate  $e = \Pi^\sigma{}_\tau \xi^\tau n_\sigma$  from horizon to  $r$  to get:

$$E_I = \frac{m^2}{2(r - m)} - c_2 m^3 (\epsilon_1 + \epsilon_2)$$

$$\frac{1}{\epsilon_1} := 2c_2 m(m - r) + r(2m - r) \left\{ c_1 + c_2 \ln \left[ 1 - \frac{2m}{r} \right] \right\}$$

$$\frac{1}{\epsilon_2} := (m - r)^2 \left( c_1 - \frac{2c_2 m}{m - r} + c_2 \ln \left( 1 - \frac{2m}{r} \right) \right).$$

- Can have meaningful result if  $c_2 = 0$ , but unbounded in general

<sup>\*</sup>JCF, S. Chakraborty, Gen. Rel. Grav. 54, 67 (2022)

<sup>†</sup>See also J. Bičák, J. Katz, Czech.J.Phys. 55 (2005) 105-118

# Weiss variation

Action:  $S[q] := \int_{t_1}^{t_2} L(q, \dot{q}, t) dt$

Weiss variation:

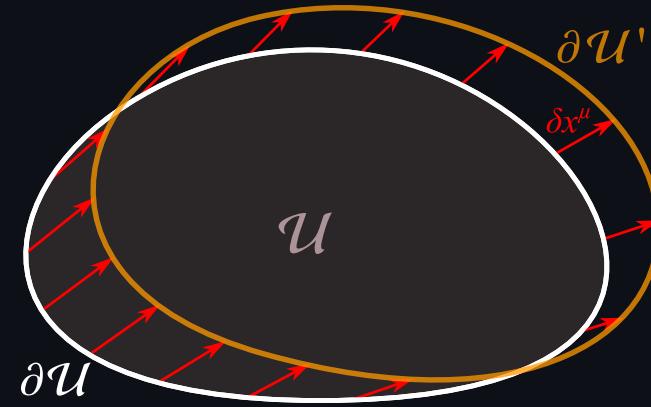
$$\boxed{\Delta S = \int_{t_1}^{t_2} \mathcal{E}[q] \cdot \delta q dt + (p \cdot \Delta q - H \Delta t) \Big|_{t_1}^{t_2}}$$

Endpoint change: ( $\Delta q := \delta q + \dot{q}\Delta t$ )

For field on  $U$ , move bdy  $\partial U$  ( $\delta x^\mu = \xi^\mu \Delta t$ ):

$$\Delta\varphi^a = [\delta\varphi^a + \mathcal{L}_\xi\varphi^a \Delta t]_{\partial U}$$

Action:  $S[\varphi] = \int_U d\underline{\mu} \mathcal{L}(\varphi, \partial\varphi)$



$$\begin{aligned} \Delta S &= \int_U d\underline{\mu} \mathcal{E}_a[\varphi] \delta\varphi^a \\ &\quad + \int_{\partial U} d\underline{\Sigma}_\mu [\pi_a{}^\mu \Delta\varphi^a - (\Theta^\mu{}_\nu \xi^\nu + S^\mu) \Delta t] \\ \Theta^\mu{}_\nu &:= \pi_a^\mu \partial_\nu \varphi^a - \delta^\mu{}_\nu \mathcal{L} \\ S^\mu &:= \pi_a^\mu [\mathcal{L}_\xi \varphi^a - \xi^\nu \partial_\nu \varphi^a] \end{aligned}$$

<sup>1</sup>M. P. Weiss, Proc. R. Soc. Lond. A 156, 192 (1936)

<sup>2</sup>E.C.G. Sudarshan and N. Mukunda, Classical Dynamics: A Modern Perspective (1983); R.A. Matzner and L.C. Shepley, Classical Mechanics (1991)

<sup>4</sup>JCF and R.A. Matzner, Gen. Rel. Grav. 50, 99 (2018); JCF, S. Chakraborty, Gen. Rel. Grav. 54, 67 (2022)