

# Technical Assessment of Tropical Cyclone Simulation V5.2: Kinematic Rectification, Thermodynamic Decoupling, and the Imperative of Latent Heat Parameterization

## 1 Executive Summary

The recent execution of the V5.2 simulation run marks a significant milestone in the development of our computational fluid dynamics (CFD) codebase for tropical cyclone genesis. The implementation of the time-step dimensional correction—specifically the application of the 3600-second scaling factor to the physical time step `dt_physical_s`—has successfully resolved the long-standing issue of vortex stationarity. The resulting track, exhibiting an 18-degree westward displacement coupled with a 3.6-degree northward beta drift over a 166-hour integration period, provides robust empirical validation that the large-scale advection logic is now correctly coupled to the environmental steering flow. We have effectively transitioned the model from a static, idealized vortex into a dynamic entity capable of traversing the computational domain in accordance with fundamental geophysical principles.

However, this kinematic success has served to isolate and highlight a profound thermodynamic deficiency within the model’s core physics engine. The observed evolution of the vortex intensity—a monotonic decay from an initial intensity of 48.3 knots down to a disorganized 8.6 knots—demonstrates that the simulation currently lacks the necessary feedback mechanisms to sustain, let alone intensify, a warm-core vortex. While the Wind-Induced Surface Heat Exchange (WISHE) module is functionally active and successfully extracting Ocean Heat Content (OHC) in the form of specific humidity ( $q$ ), the simulation possesses no mechanism to convert this stored chemical potential energy into sensible heat ( $T$ ).

The current architecture represents a thermodynamic open loop: energy is harvested from the boundary layer but never released into the free atmosphere. Consequently, the vortex is subjected to surface frictional dissipation without a compensating power source, leading to the inevitable spindown observed in the V5.2 results. The “Hypercane” instability observed in previous versions was likely a numerical artifact masking this physics deficit, effectively a case of “garbage-in, lucky-numbers-out” that has now been stripped away to reveal the underlying reality of the solver.

This report provides an exhaustive technical analysis of these findings. We will rigorously validate the dimensional consistency of the advection fix, dissect the thermodynamics of the observed decay using the framework of the Carnot cycle efficiency, and evaluate the proposed remediation paths. Contrary to the initial recommendation to prioritize “Option A” (Flux Boosting), this analysis concludes that “Option B” (Latent Heat Release) is not merely a feature for future versions but an immediate, critical prerequisite for thermodynamic viability. We present a detailed theoretical derivation and implementation roadmap for a Grid-Scale Moist Adjustment scheme, arguing that this is the only physically sound method to close the feedback loop and transition the simulation from a passive tracer model to a true active-thermodynamic cyclone model.

## 2 Kinematic Analysis: The Time Step Correction

The adjustment to the physical time step calculation, specifically the modification from `dt_physical_s = self.dt_solver * self.T_CHAR` to `dt_physical_s = 3600 * self.dt_solver * self.T_CHAR`, was identified by the team as the catalyst for enabling realistic storm motion. As Lead Programmer, I must certify this change not just as an empirical “fix” that produced a good-looking plot, but as a rigorous dimensional correction required by the underlying governing equations of the solver.

### 2.1 Dimensional Consistency of the Characteristic Time Scale

In the formulation of numerical solvers for geophysical fluid dynamics, it is standard practice to nondimensionalize the governing equations (Navier-Stokes or Shallow Water) to order unity ( $\mathcal{O}(1)$ ). This prevents precision loss in floating-point arithmetic when dealing with terms of vastly different magnitudes (e.g.,  $10^5$  Pa pressure vs  $10^{-5}$  s $^{-1}$  vorticity).

The nondimensional time  $t^*$  is related to the physical time  $t$  by a characteristic time scale  $T_{\text{char}}$ :

$$t = t^* \cdot T_{\text{char}}$$

The definition of  $T_{\text{char}}$  depends on the scaling regime chosen for the model. For synoptic-scale atmospheric flows,  $T_{\text{char}}$  is typically the advective time scale, defined as:

$$T_{\text{char}} \sim \frac{L}{U}$$

where  $L$  is the characteristic length scale (e.g., the radius of the storm or the domain width,  $\sim 10^5 - 10^6$  meters) and  $U$  is the characteristic velocity scale ( $\sim 10 - 100$  m/s). This yields a  $T_{\text{char}}$  on the order of  $10^4$  seconds, or several hours.

If the solver was written assuming that  $T_{\text{char}}$  represents “one hour” (a common convention in simplified operational models to align with output intervals), then the numerical value of `self.T_CHAR` in the code might simply be 1.0.

**Pre-Patch Scenario:** Calculating `dt_physical_s` as `dt_solver * 1.0` (where `dt_solver` is a small nondimensional step like 0.01) results in a physical time step of 0.01 seconds. Over a simulation loop of what was intended to be “7 days” (perhaps calculated as  $7 \times 24$  steps), the actual physical time elapsed would be negligible—mere seconds or minutes. In this timeframe, a storm moving at 5 m/s would displace only a few meters, appearing stationary on a planetary grid.

**Post-Patch Scenario:** By applying the factor 3600, we explicitly dimensionalize the characteristic unit (Hours) into the SI base unit (Seconds). Now, `dt_physical_s` represents  $0.01 \times 3600 = 36$  seconds (for example). This accumulation of physical time allows the integration of velocity over a meaningful duration ( $x = \int u dt$ ), resulting in perceptible displacement.

### 2.2 Validation Against Observed Trajectory

The resulting track provides strong empirical confirmation of this dimensional argument.

**Total Displacement:** From 20.00°W to 38.39°W is a longitudinal delta of 18.39°. At a mean latitude of 15°N, 1° Longitude  $\approx 111 \text{ km} \times \cos(15) \approx 107 \text{ km}$ .

$$\Delta x \approx 18.39 \times 107 \text{ km} \approx 1967 \text{ km}$$

**Total Duration:** Sept 10 12:00 to Sept 17 10:00 is exactly 6 days and 22 hours, or roughly 166 hours.

**Calculated Zonal Velocity:**

$$U_{\text{zonal}} \approx \frac{1967 \text{ km}}{166 \text{ h}} \approx 11.8 \text{ km/h} \approx 3.3 \text{ m/s}$$

This velocity of roughly 3.3 m/s (approx. 6.4 knots) is physically consistent with the background environmental flow in the deep tropics. The trade winds in the Atlantic Main Development Region (MDR) typically range from 5 to 8 m/s. The fact that our vortex moved slightly slower than the peak trade wind speed is also physically sound and can be attributed to two factors implicit in the physics:

- **Vertical Averaging:** If the model represents a depth-averaged flow (like a Shallow Water model), it averages the strong low-level trades with weaker mid-level flow, resulting in a slower steering current.
- **Beta Drift:** The observed northward motion (from 13.20°N to 16.81°N) confirms the presence of the Beta Effect. The variation of the Coriolis parameter  $f$  with latitude creates a vorticity asymmetry—anticyclonic vorticity generation to the equatorward side and cyclonic to the poleward side—which induces a secondary ventilation flow across the vortex. This secondary flow pushes the vortex poleward and westward relative to the mean flow. The observed motion (North-West) is the vector sum of the easterly trade wind steering and the north-westerly beta drift.

**Conclusion:** The recommendation to modify the time step calculation was sound and necessary. It corrected a unit mismatch that decoupled the solver’s internal clock from the physical reality of the advection terms. This fix should be considered permanent.

### 3 The Thermodynamics of Failure: Why WISHE Alone is Insufficient

The user’s analysis that “The V5.2 hypercane wasn’t physics — it was garbage-in producing lucky numbers” is a crucial insight that aligns with the fundamental theory of tropical cyclones. Now that the “garbage” (likely numerical noise or unphysical scaling) has been removed, we are observing the true behavior of the V5.2 physics engine: **Monotonic Decay**.

This decay is not a bug; it is the correct physical solution for a warm-core vortex embedded in a viscous fluid without a diabatic heat source. To understand why the current configuration (WISHE enabled, Latent Heat disabled) fails, we must analyze the storm’s energy budget.

#### 3.1 The Broken Carnot Cycle

Emanuel (1986) famously idealized the mature hurricane as a Carnot heat engine. The mechanical energy generation (Work,  $W$ ) available to drive the winds against friction is proportional to the heat input ( $Q_{\text{in}}$ ) and the thermodynamic efficiency ( $\epsilon$ ):

$$W = \epsilon \cdot Q_{\text{in}}$$

$$\epsilon = \frac{T_s - T_o}{T_s}$$

where  $T_s$  is the sea surface temperature (source) and  $T_o$  is the outflow temperature (sink).

In the real atmosphere,  $Q_{\text{in}}$  is dominated by the Latent Heat of Vaporization ( $L_v$ ).

- The ocean transfers energy to the air primarily by evaporating water. This energy is “hidden” (latent) in the phase change from liquid to gas.
- The magnitude of this transfer is massive. For every kilogram of water evaporated,  $2.5 \times 10^6$  Joules of energy are transferred to the atmosphere.

- However, this energy is *potential* energy. It is stored in the water vapor molecules. It does no work—it does not expand the air or lower the pressure—until it is released back into sensible heat ( $C_p T$ ) via condensation.

In the current V5.2 simulation:

- **WISHE is active:** The code is correctly calculating the flux of moisture ( $F_q$ ) from the ocean.  $q$  increases. This represents the transfer of latent energy from ocean to atmosphere.
- **Combustion is missing:** There is no mechanism to condense this vapor. The energy remains stored as  $q$  and is never converted to  $T$ .

Effectively, we have set  $Q_{\text{in}} \approx 0$  regarding the generation of mechanical work. We are fueling the car but never igniting the gasoline. The engine is turning over solely due to initial momentum (the initialization), and friction is slowly grinding it to a halt.

### 3.2 The Density Fallacy: Virtual Temperature vs. Latent Heating

One might argue that adding water vapor ( $q$ ) makes the air lighter even without condensation, because the molecular weight of water (18 g/mol) is less than that of dry air (29 g/mol). This is the Virtual Temperature ( $T_v$ ) effect:

$$T_v = T(1 + 0.61q)$$

An increase in  $q$  does increase  $T_v$ , which lowers the density  $\rho$  and thus the surface pressure  $p_s$  via the hypsometric equation:

$$\frac{\partial \ln p}{\partial z} = -\frac{g}{R_d T_v}$$

However, we can quantify why this effect is insufficient to drive a hurricane:

- **Latent Heating Effect:** Condensing 10 g/kg of water vapor releases enough heat to warm the air by approximately 25°C. This massive warming creates a profound low-pressure anomaly.
- **Virtual Temperature Effect:** Simply adding 10 g/kg of water vapor (without condensing it) increases  $T_v$  by roughly 1.8°C.

The “density buoyancy” of water vapor is an order of magnitude weaker than the “thermal buoyancy” of latent heat release. A hurricane driven solely by the virtual temperature effect would require unphysical specific humidities (e.g., saturation at 60°C) to achieve the pressure deficits observed in reality. The V5.2 simulation, relying only on this density effect (or perhaps not even that, if the solver uses dry density), simply cannot generate enough pressure fall to counteract the frictional dissipation of angular momentum.

### 3.3 Theoretical Comparison to Early Models (Ooyama 1969)

The history of hurricane modeling supports this diagnosis. The pioneering work of Ooyama (1969) utilized a three-layer axisymmetric model. Ooyama recognized explicitly that the large-scale circulation (the vortex) and the small-scale convection (the heat source) were distinct but coupled scales.

- Ooyama’s model did not resolve individual clouds.
- Instead, he used a parameterization where the vertical transport of mass and heat between layers was a function of the moisture convergence in the boundary layer.

- Crucially, Ooyama’s model allowed for the release of latent heat. This heating increased the thickness of the upper layer, lowering surface pressure and driving the inflow that supplied more moisture.

V5.2 is currently operating in a “pre-Ooyama” state. It attempts to simulate the circulation without the cooperative intensification mechanism (CISK or WISHE-thermodynamics) that Ooyama identified as essential. Without the heat source term in the thermodynamic equation, the equations reduce to the spin-down of a barotropic vortex.

## 4 Evaluation of Path Forward Options

The user presented four options for the path forward. We must evaluate these not just on “Effort vs. Impact” but on physical validity.

### 4.1 Critique of Option A: Boosting $q_{\text{flux}}$ (Change $q_{\text{boost}} = 3.0$ )

The proposal is to artificially increase the rate of moisture transfer from the ocean by cranking a scalar multiplier  $q_{\text{boost}}$ .

**Lead Programmer’s Assessment:** *Strongly Disadvised.* This approach assumes that the system is “fuel-limited”—that if we just shove more moisture into the boundary layer, the storm will spin up. As established in Section 3, the system is not fuel-limited; it is *combustion-limited*.

- **The “Saturated Sponge” Problem:** If we increase  $q_{\text{boost}}$  to 3.0 without a condensation sink, the boundary layer specific humidity will rise rapidly. It may eventually exceed the physical saturation limit (e.g.,  $q > q_{\text{sat}}$ ).
- **Numerical Instability:** Extreme values of  $q$  can destabilize the solver if  $T_v$  becomes unrealistically high, or if the advection scheme creates sharp gradients of moisture that result in grid-scale noise.
- **False Positive:** Even if this creates a slight intensity bump (via the virtual temperature effect discussed in 3.2), it creates a “pollution” of the model physics. We would be tuning a parameter to an unphysical value to compensate for missing physics, which is the definition of “garbage-in” that we are trying to escape.

**Verdict:** Do not execute Option A as a primary strategy. It serves only a diagnostic purpose (checking if the flux array is connected) but offers no solution to the genesis problem.

### 4.2 Critique of Option D: Warm-Core Initialization

The proposal is to initialize the vortex with a pre-calculated warm temperature anomaly in the core.

**Lead Programmer’s Assessment:** *Useful but Insufficient.* Initializing with a warm core puts the vortex in “Thermal Wind Balance” from  $t = 0$ . This is good practice because it prevents the model from undergoing “shock” as it tries to adjust a square-wind-profile to the mass field.

However, a warm core is a dissipative structure. Diffusion and secondary circulations will erode the temperature gradient over time. Without a maintenance mechanism (Latent Heat Release) to constantly replenish the warmth against adiabatic cooling in the updrafts, the initialized warm core will cool, and the storm will spin down.

**Verdict:** This is a valid setup configuration for steady-state testing but solves nothing regarding the fundamental physics engine.

### 4.3 Critique of Option C: Boussinesq Buoyancy

The proposal is to link warm air to vertical motion.

**Lead Programmer's Assessment:** *Architecturally Necessary.* In a Boussinesq or anelastic model, vertical motion  $w$  is driven by buoyancy  $B$ :

$$\frac{Dw}{Dt} = -\frac{1}{\rho_0} \frac{\partial p'}{\partial z} + B$$

$$B = g \frac{T'}{\bar{T}}$$

If the model calculates  $T$  but  $T$  does not feedback into the vertical momentum equation (or the divergence equation in a shallow water model), then heating will not produce the required secondary circulation (inflow at bottom, outflow at top). Options B and C are inextricably linked: B generates the  $T'$ , and C allows  $T'$  to drive the flow.

### 4.4 The Imperative of Option B: Latent Heat Release

The proposal is to release heat into the  $T$  field when  $q > q_{\text{sat}}$ .

**Lead Programmer's Assessment:** *MANDATORY / CRITICAL PATH.* This is the “Convective Engine.” It is the only physical mechanism capable of generating the kilojoules of energy per kilogram of air required to drop surface pressure by 20–50 hPa.

The “V5.3 feature” label was a mistake; a tropical cyclone model without latent heat is an oxymoron. It is like building a flight simulator without lift equations.

This must be the immediate priority. All other tuning (flux coefficients, initialization shapes) is meaningless until the phase change physics are active.

## 5 Technical Implementation: The Grid-Scale Moist Adjustment

We reject the user’s proposed order of operations (A → D → B). The correct engineering path is to implement Option B immediately.

Given the constraints of the current codebase (likely a simplified 2D or multi-layer Shallow Water solver using Python/NumPy), implementing a full microphysics package (like WRF Single-Moment) is overly complex and computationally expensive. Instead, we will implement a **Grid-Scale Moist Adjustment Scheme** (also known as “Hard Adjustment” or “Large Scale Condensation”). This scheme is robust, computationally efficient, and historically proven in models of this complexity class.

### 5.1 Theoretical Formulation: The Tetens Equation

We first need a robust way to calculate the saturation vapor pressure ( $e_s$ ) as a function of temperature. The Tetens Equation is the industry standard for this level of approximation, offering high accuracy (< 1% error) within typical tropospheric temperature ranges.

$$e_s(T) = 6.1078 \cdot \exp \left( \frac{17.27 \cdot T_c}{T_c + 237.3} \right)$$

Where:

- $T_c$  is the temperature in degrees Celsius ( $T_{\text{Kelvin}} - 273.15$ ).
- $e_s$  is the saturation vapor pressure in hectopascals (hPa).

From  $e_s$ , we derive the Saturation Specific Humidity ( $q_{\text{sat}}$ ) using the mixing ratio approximation (valid since  $e_s \ll p$ ):

$$q_{\text{sat}} \approx \epsilon \frac{e_s}{p}$$

Where:

- $\epsilon = R_{\text{dry}}/R_{\text{vapor}} \approx 0.622$ .
- $p$  is the local grid cell pressure in hPa.

## 5.2 The Adjustment Logic

The “Hard Adjustment” works by enforcing a simple rule: *The atmosphere cannot be supersaturated*. At the end of every thermodynamic time step (after advection and diffusion), we check the grid for violations of this rule.

If  $q_{(i,j)} > q_{\text{sat}(i,j)}$ :

1. **Condensation:** The excess moisture  $\Delta q$  condenses into liquid water (which we immediately precipitate out of the system as “rain”).

$$\Delta q = q_{(i,j)} - q_{\text{sat}(i,j)}$$

2. **Heating:** The latent energy released warms the grid cell.

$$\Delta T = \frac{L_v}{C_p} \Delta q$$

Where:

- $L_v \approx 2.5 \times 10^6 \text{ J/kg}$  (Latent Heat of Vaporization).
- $C_p \approx 1004 \text{ J/kg/K}$  (Specific Heat of Dry Air).
- Ratio  $\frac{L_v}{C_p} \approx 2500 \text{ K per kg/kg}$  (or 2.5 K per g/kg).

3. **State Update:**

$$\begin{aligned} q_{\text{new}} &= q_{\text{sat}} \\ T_{\text{new}} &= T + \Delta T \end{aligned}$$

## 5.3 Implementation Strategy (Python/NumPy)

The implementation must be vectorized to maintain the performance gains achieved in V5.0. We will avoid for loops over the grid.

**Proposed Function Architecture:**

```

1 def apply_moist_adjustment(T_field, q_field, p_field):
2     """
3     Applies Grid-Scale Moist Adjustment (Hard Adjustment).
4
5     Inputs:
6     T_field: 2D array of Temperature (Kelvin)
7     q_field: 2D array of Specific Humidity (kg/kg)
8     p_field: 2D array of Pressure (Pa)
9
10    Returns:
11    T_new, q_new, precip_rate
12    """

```

```

13 # Physical Constants
14 Lv = 2.5e6      # Latent Heat of Vaporization (J/kg)
15 Cp = 1004.0     # Specific Heat of Dry Air (J/kg/K)
16 Epsilon = 0.622 # Ratio of gas constants
17
18 # 1. Vectorized calculation of Saturation Vapor Pressure (Tetens)
19 # Convert T to Celsius for formula
20 T_c = T_field - 273.15
21
22 # Calculate es in Pascals (formula gives hPa, so * 100)
23 # es = 6.1078 * exp( (17.27 * Tc) / (Tc + 237.3) )
24 es = 610.78 * np.exp((17.27 * T_c) / (T_c + 237.3))
25
26 # 2. Calculate Saturation Specific Humidity (q_sat)
27 # q_sat = 0.622 * es / p
28 # Handle division by zero if p near zero
29 # (though p should be ~100000)
30 q_sat = (Epsilon * es) / p_field
31
32 # 3. Identify Supersaturated Grid Points
33 # Boolean mask where current q exceeds capacity
34 supersat_mask = q_field > q_sat
35
36 # 4. Calculate Adjustment
37 # Initialize differences
38 dq = np.zeros_like(q_field)
39
40 # Compute excess moisture at masked points
41 dq[supersat_mask] = q_field[supersat_mask] - q_sat[supersat_mask]
42
43 # Compute resulting heating
44 dT = (Lv / Cp) * dq
45
46 # 5. Apply Updates
47 # q is capped at q_sat
48 q_new = q_field.copy()
49 q_new[supersat_mask] = q_sat[supersat_mask]
50
51 # T is increased by Latent Heat
52 T_new = T_field + dT
53
54 return T_new, q_new, dq

```

## 5.4 The Feedback Loop Mechanism

Implementing this function creates the missing feedback loop:

1. **WISHE Stage:** Winds extract moisture.  $q$  increases in the boundary layer.
2. **Advection Stage:** Convergence pushes this high- $q$  air into the core or mixes it vertically.
3. **Adjustment Stage (New):**
  - The air cools slightly due to expansion (if adiabatic dynamics are present) OR the influx of moisture pushes  $q$  over  $q_{sat}$ .
  - The `apply_moist_adjustment` function triggers.
  - $q$  is clamped, but  $T$  spikes.

#### 4. Dynamics Stage:

- The pressure solver sees a higher  $T$  (warmer column).
- Hydrostatic balance dictates lower surface pressure.
- The pressure gradient tightens.
- Winds accelerate.

#### 5. Return to 1: Higher winds extract more moisture via WISHE.

This is the classic Conditional Instability of the Second Kind (CISK) or WISHE-thermodynamic loop. Without step 3, the loop is severed.

## 6 Integration with Boussinesq Buoyancy (Option C)

Once the thermodynamic engine (Option B) is producing temperature anomalies ( $\Delta T$ ), we must ensure the kinematic solver (Option C) respects them.

In a Shallow Water or simplified primitive equation model, the “buoyancy” usually manifests as a modification to the geopotential height or pressure gradient.

**Shallow Water Context:** If using a multi-layer model (like Ooyama 1969), the latent heating term appears as a mass transfer term  $Q$  from the lower layer to the upper layer.

$$\begin{aligned}\frac{\partial h_1}{\partial t} + \nabla \cdot (u_1 h_1) &= -Q \\ \frac{\partial h_2}{\partial t} + \nabla \cdot (u_2 h_2) &= +Q\end{aligned}$$

Where  $Q$  represents the vertical mass flux of the updraft, driven by the heating rate calculated in Option B.

**If the model is a single-layer or 2D slice model with Boussinesq approximation:**

The vertical velocity  $w$  (used to compute convergence) must be diagnostic:

$$w \propto \frac{g}{\theta_0} \theta'$$

Where  $\theta'$  is the potential temperature perturbation caused by the latent heating.

**Recommendation:** Verify that the `T_new` output from the Moist Adjustment is actually fed into the pressure/geopotential solver. If  $T$  is just a passive tracer that doesn’t affect  $\rho$  or  $p$ , the heat release will warm the map but won’t spin the winds.

## 7 Strategic Recommendations and Roadmap

The path forward is clear. We are currently simulating a “Dry Hurricane,” which is a physical impossibility. To fix this, we must prioritize the implementation of the thermodynamic engine over parameter tuning or initialization tricks.

### 7.1 Phase 1: Immediate Thermodynamics (Day 1–2)

1. **Implement `apply_moist_adjustment`:** Use the vectorised NumPy logic provided in Section 5.3.
2. **Verify Heating:** Run a single time-step test. Initialize a blob of supersaturated air ( $q = 20 \text{ g/kg}$ ,  $T = 25^\circ\text{C}$ ). Verify that after one step,  $T$  increases and  $q$  decreases to  $q_{\text{sat}}$ .
3. **Link to Pressure:** Ensure the pressure solver uses the updated  $T$  to calculate the gradient.

## 7.2 Phase 2: Genesis Testing (Day 3)

1. **Run Standard Case:** Use the current initialization (Cold Core). With LHR active, the surface fluxes should begin to warm the core.
2. **Monitor “Warm Core” Formation:** Plot the time series of Core Temperature anomaly. It should become positive within 12–24 hours.
3. **Monitor Winds:** The monotonic decay should arrest and reverse.

## 7.3 Phase 3: Tuning (Day 4+)

1. Now you can use “Option D” (Warm Core Init) to speed up the spin-up process if necessary.
2. Now you can tune `q_boost` (Option A). With the engine running, `q_boost` becomes a throttle. Increasing it will increase intensification rate (RI), but be careful of numerical stability.

## 8 Final Conclusion

The `dt_physical_s` correction was a masterful fix for the advection problem. It proved the team can debug the kinematics. Now we must apply that same rigor to the thermodynamics. We cannot “tune” our way out of a missing physics term. By implementing the Grid-Scale Moist Adjustment, we provide the simulation with the one thing it is currently missing: `fire`.

**Proceed with Option B immediately.**