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¹ Chapter 1

² Introduction and Theory Overview

³ 1.1 Introduction

⁴ This thesis presents the analysis details and the results of the search for heavy reso-
⁵ nances decaying into a Z boson and a Higgs boson (h) at the center-of-mass energy of
⁶ 8 TeV, using 19.7 fb^{-1} p-p collision data. In turn, the Z boson is identified through
⁷ its leptonic decays (leptons often refer to e and μ only in experiments. $l = e, \mu$). The
⁸ Higgs boson h is expected to hadronically decay into a pair of b-quarks. The investi-
⁹ gated final states consist of two charged leptons which are identified in the detector
¹⁰ and limit the presence of the background, and two b-quarks from the hadronic Higgs
¹¹ decay which collects the largest possible fraction of Higgs events.

¹² This thesis is organised as follows. In the latter part of this chapter, the model that
¹³ predicts heavy resonances is introduced, including the expected cross section and the
¹⁴ specification of model parameters. In chapter 2, the LHC and the CMS experiment
¹⁵ are described, including the information of each sub-detector and the trigger system
¹⁶ of the CMS. The details of the analysis are shown in chapter 3. This chapter reveals
¹⁷ the way to reconstruct physical objects in CMS. By adding some proper kinematic
¹⁸ selections on those physics objects, the interested events in data collected by the CMS
¹⁹ detector can be selected. Moreover, this chapter shows the comparison between data
²⁰ and simulation. In the last chapter, the results of the search and the conclusion are
²¹ presented.

22 1.2 Theory Overview

23 Although the Higgs boson discovered by the ATLAS and CMS collaborations [1–3]
 24 imposes strong constraints on theories beyond the Standard Model(SM), the extreme
 25 fine tuning in quantum corrections required to have a light fundamental Higgs boson
 26 with mass close to 125 GeV [4–7] suggests that the Standard Model may be incom-
 27 plete, and not valid beyond a scale of a few TeV. Various dynamical electroweak
 28 symmetry breaking scenarios which attempt to solve this naturalness problem, such
 29 as Minimal Walking Technicolor [8], Little Higgs [9–11], or composite Higgs mod-
 30 els [12–14] predict the existence of new resonances decaying to a vector boson plus a
 31 Higgs boson.

32 1.2.1 Heavy Vector Triplet Model

33 Resonance searches are typically not sensitive to all the details and the free parameters
 34 of the underlying model, but only to those parameters or combinations of parameters
 35 that control the mass of the resonance and the interactions involved in its production
 36 and decay. Therefore, one can employ a simplified description of the resonance de-
 37 fined by a phenomenological Lagrangian where only the relevant couplings and mass
 38 parameters are retained. This model-independent strategy applies a Heavy Vector
 39 Triplet (HVT) [15] to the Standard Model group and reproduces a large class of ex-
 40 plicit models. In Eq. (1.1), the mathematical form of the simplified Lagrangian is
 41 defined, where V_ν^a , $a = 1,2,3$, is a real vector with vanishing hypercharge in the ad-
 42 joint representation of $SU(2)_L$, it describes one charged and one neutral heavy spin-1
 43 particle with charge eigenstate fields, and $D_{[\mu} V_{\nu]}^a$ represents the covariant derivative.

$$\begin{aligned} \mathcal{L}_V = & -\frac{1}{4} D_{[\mu} V_{\nu]}^a D^{[\mu} V^{\nu]}_a + \frac{m_V^2}{2} V_\mu^a V^{\mu a} \\ & + ig_V c_H V_\mu^a H^\dagger \tau^a \overset{\leftrightarrow}{D}^\mu H + \frac{g^2}{g_V} c_F V_\mu^a \sum_f \bar{f}_L \gamma^\mu \tau^a f_L \\ & + \frac{g_V}{2} c_{VVV} \epsilon_{abc} V_\mu^a V_\nu^b D^{[\mu} V^{\nu]}_c + \text{quadrilinear terms} \end{aligned} \quad (1.1)$$

$$V_\mu^\pm = \frac{V_\mu^1 \mp i V_\mu^2}{\sqrt{2}}, \quad V_\mu^0 = V_\mu^3 \quad (1.2)$$

$$D_{[\mu} V_{\nu]}^a = D_\mu V_\nu^a - D_\nu V_\mu^a, \quad D_\mu V_\nu^a = \partial_\mu V_\nu^a + g \epsilon^{abc} W_\mu^b V_\nu^c \quad (1.3)$$

$$H = \begin{pmatrix} \phi^+ \\ \phi^- \end{pmatrix} = \begin{pmatrix} \frac{1}{\sqrt{2}}(\phi_1 + i\phi_2) \\ \frac{1}{\sqrt{2}}(\phi_3 + i\phi_4) \end{pmatrix} \quad (1.4)$$

⁴⁴

⁴⁵ In these models, new heavy vector bosons (V^\pm, V^0) that couple to the SM Higgs
⁴⁶ doublet (Eq. 1.4) and SM gauge bosons with the parameters c_H and g_V and to the
⁴⁷ fermions via the combination $(g^2/g_V)c_F$. The parameter g_V represents the strength
⁴⁸ of the new vector boson interaction, while c_H and c_F represent the couplings to the
⁴⁹ Higgs and the fermions respectively, and are expected to be of the order of unity in
⁵⁰ most models.

⁵¹ 1.2.2 Basic Phenomenology

⁵² Masses and Mixings

After electro-weak symmetry breaking (EWSB), the only massless state is photon, which can be identified as the gauge field associated with the unbroken $U(1)_{em}$. The two other neutral mass eigenstates are the SM Z boson and one heavy vector of mass M_0 which are obtained by diagonalizing the mass matrix of the (Z, V^0) system by a rotation with angle θ_N

$$\begin{pmatrix} Z \\ V^0 \end{pmatrix} \rightarrow \begin{pmatrix} \cos \theta_N & \sin \theta_N \\ -\sin \theta_N & \cos \theta_N \end{pmatrix} \begin{pmatrix} Z \\ V^0 \end{pmatrix}. \quad (1.5)$$

The mass matrix is

$$\mathcal{M}_N^2 = \begin{pmatrix} \hat{m}_Z^2 & c_H \xi \hat{m}_Z \hat{m}_V \\ c_H \xi \hat{m}_Z \hat{m}_V & \hat{m}_V^2 \end{pmatrix}, \text{ where } \begin{cases} \hat{m}_Z = \frac{e\hat{v}}{2\sin\theta_W \cos\theta_W} \\ \hat{m}_V^2 = m_V^2 + g_V^2 c_{VVHH} \hat{v}^2 \\ \xi = \frac{g_V \hat{v}}{2\hat{m}_V} \end{cases}. \quad (1.6)$$

In the above equations \hat{v} denotes the Vacuum Expectation Value (VEV) defined by $\langle H^\dagger H \rangle = \hat{v}^2/2$, and one should know the masses \hat{m}_Z and \hat{m}_V do not coincide with the physical Z boson and the masses of the new resonances of this model, although they do in the approximations later (Eq. 1.12). The mass eigenvalues and the rotation angles are easily obtained by inverting the relations

$$\begin{aligned} Tr[\mathcal{M}_N^2] &= \hat{m}_Z^2 + \hat{m}_V^2 = m_Z^2 + M_0^2, \\ Det[\mathcal{M}_N^2] &= \hat{m}_Z^2 \hat{m}_V^2 (1 - c_H^2 \xi^2) = m_Z^2 M_0^2, \\ \tan 2\theta_N &= \frac{2c_H \xi \hat{m}_Z \hat{m}_V}{\hat{m}_V^2 - \hat{m}_Z^2}. \end{aligned} \quad (1.7)$$

- ⁵³ Note that M_0 represents the real mass eigenvalue of the neutral heavy vector boson.
- ⁵⁴ Moreover, since we assume $\hat{m}_V > \hat{m}_Z$, the only variable controls the sign of the
- ⁵⁵ tangent function is c_H , which is model dependent. Once the sign of c_H is determined,
- ⁵⁶ the tangent can be uniquely inverted.

The situation is similar in the charged vector mass matrix of (W^\pm, V^\pm) system, and M_\pm denotes the real mass eigenvalue of charged states as well.

$$\mathcal{M}_C^2 = \begin{pmatrix} \hat{m}_W^2 & c_H \xi \hat{m}_W \hat{m}_V \\ c_H \xi \hat{m}_W \hat{m}_V & \hat{m}_V^2 \end{pmatrix}, \text{ where } \hat{m}_W = \frac{e\hat{v}}{2\sin\theta_W} = \cos\theta_W \hat{m}_Z, \quad (1.8)$$

where it is diagonalized by

$$\begin{aligned} Tr[\mathcal{M}_C^2] &= \hat{m}_W^2 + \hat{m}_V^2 = m_W^2 + M_\pm^2 , \\ Det[\mathcal{M}_C^2] &= \hat{m}_W^2 \hat{m}_V^2 (1 - c_H^2 \xi^2) = m_W^2 M_\pm^2 , \\ \tan 2\theta_C &= \frac{2c_H \xi \hat{m}_W \hat{m}_V}{\hat{m}_V^2 - \hat{m}_W^2} . \end{aligned} \quad (1.9)$$

By checking Eq. (1.6) and Eq. (1.8), the charged and neutral mass matrices are connected by custodial symmetry, which can be shown in full generality to imply

$$\mathcal{M}_C^2 = \begin{pmatrix} \cos \theta_W & 0 \\ 0 & 1 \end{pmatrix} \mathcal{M}_N^2 \begin{pmatrix} \cos \theta_W & 0 \\ 0 & 1 \end{pmatrix} . \quad (1.10)$$

By taking the determinant of the above equation, or equivalently by comparing the charged and neutral determinants in Eq. (1.7) and Eq. (1.9), we obtain a generalized custodial relation among the physical masses

$$m_W^2 M_\pm^2 = \cos^2 \theta_W m_Z^2 M_0^2 . \quad (1.11)$$

From the simple formula above, we can start to identify the physically reasonable region of the parameter space in this model. We aim at describing new vectors with masses at or above the TeV scale, but we also want the SM masses $m_{W,Z} \sim 100$ GeV to be reproduced. Therefore we require a hierarchy in the mass relation of SM Z and W bosons versus the new vectors.

$$\frac{\hat{m}_{W,Z}}{\hat{m}_V} \sim \frac{m_{W,Z}}{M_{\pm,0}} \leq 10^{-1} \ll 1 \quad (1.12)$$

Use the limit above, we can expand the determinant formulas both in Eq. (1.7) and Eq. (1.9) to obtain simple approximations for m_W and m_Z

$$\begin{aligned} m_Z^2 &= \hat{m}_Z^2 (1 - c_H^2 \xi^2) (1 + \mathcal{O}(\hat{m}_Z^2 / \hat{m}_V^2)) , \\ m_W^2 &= \hat{m}_W^2 (1 - c_H^2 \xi^2) (1 + \mathcal{O}(\hat{m}_W^2 / \hat{m}_V^2)) . \end{aligned} \quad (1.13)$$

The parameter ξ can be either very small or of order unity. Both cases are realized in explicit models. While $\xi \ll 1$ is the most common situation, $\xi \sim 1$ only occurs in strongly coupled scenarios at very large g_V . In these approximations, SM tree-level experimental observation can be reproduced to percent accuracy.

Since $\hat{m}_W = \cos \theta_W \hat{m}_Z$, the W - Z mass ratio is thus given by

$$\frac{m_W^2}{m_Z^2} \simeq \cos^2 \theta_W . \quad (1.14)$$

Eq. (1.14) has one important implication on the masses of the new vectors. When combined with the custodial relation Eq. (1.11), it tells us that the charged and neutral V s are practically degenerate

$$M_{\pm}^2 = M_0^2(1 + \mathcal{O}(\%)) , \quad (1.15)$$

⁵⁷ In the following, when working at the leading order in the limit Eq. (1.12), we can
⁵⁸ ignore the mass splitting and denote the mass of the charged and the neutral states
⁵⁹ collectively as M_V . It is easy to check that in that limit $M_V = \hat{m}_V$.

⁶⁰ Decay Widths

Because of the hierarchy in the mass matrices, the mixing angles are naturally small. By looking at Eqs. (1.7), (1.9) and (1.12) we can estimate

$$\theta_{N,C} \simeq c_H \xi \frac{\hat{m}_{W,Z}}{\hat{m}_V} \leq 10^{-1} , \quad (1.16)$$

and after rotating to the mass basis, the coupling of the neutral and charged resonances to left- and right-handed fermion chiralities can be written in a compact form

for each fermion species $F = \{l, q, 3\}$.

$$\begin{cases} g_L^N = \frac{g^2}{g_V} \frac{c_F}{2} \cos \theta_C + (g_L^Z)_{SM} \sin \theta_N \simeq \frac{g^2}{g_V} \frac{c_F}{2} , \\ g_R^N = (g_R^Z)_{SM} \sin \theta_N \simeq 0 \\ g_L^C = \frac{g^2}{g_V} \frac{c_F}{2} \cos \theta_C + (g_L^W)_{SM} \sin \theta_C \simeq \frac{g^2}{g_V} \frac{c_F}{2} , \\ g_R^C = 0 \end{cases} \quad (1.17)$$

In the above equation $(g_{L,R}^{W,Z})_{SM}$ denote the ordinary SM W and Z couplings (with the normalization given by $g_L^W = g/\sqrt{2}$).

Given that the rotation angles are small, the couplings further simplify, as also shown in the equation. We could see that V interact mainly with left-handed chiralities and that all the couplings for each fermion species are controlled by the parameter combination $g^2/g_V c_F$. This gives tight correlations among different channels

$$\Gamma_{V_\pm \rightarrow f\bar{f}'} \simeq 2\Gamma_{V_0 \rightarrow f\bar{f}'} \simeq N_C[f] \left(\frac{g^2 c_F}{g_V}\right)^2 \frac{M_V}{48\pi}, \quad (1.18)$$

where $N_C[f]$ is the number of colors (3 for the di-quark and 1 for the dilepton decays). The parameters $c_F = \{c_l, c_q, c_3\}$ control the relative BRs to leptons, light quarks and the third family.

In the case of di-boson decay width

$$\begin{aligned} \Gamma_{V_0 \rightarrow W_L^+ W_L^-} &\simeq \Gamma_{V_\pm \rightarrow W_L^\pm Z_L} \simeq \frac{g_V^2 c_H^2 M_V}{192\pi} \frac{(1 + c_H c_{VVV} \xi^2)^2}{(1 - c_H^2 \xi^2)^2} = \frac{g_V^2 c_H^2 M_V}{192\pi} [1 + \mathcal{O}(\xi^2)], \\ \Gamma_{V_0 \rightarrow Z_L h} &\simeq \Gamma_{V_\pm \rightarrow W_L^\pm h} \simeq \frac{g_V^2 c_H^2 M_V}{192\pi} \frac{(1 - 4c_{VVH} \xi^2)^2}{1 - c_H^2 \xi^2} = \frac{g_V^2 c_H^2 M_V}{192\pi} [1 + \mathcal{O}(\xi^2)]. \end{aligned} \quad (1.19)$$

⁶¹ Note that Eq. (1.19) is derived in the Equivalent Gauge [16] because the decay to transverse SM vectors is highly suppressed while to the longitudinal parts grows with the energy of the process, therefore the Unitary Gauge which is used in the original

⁶⁴ Lagrangian is instead useful. The channels that are not shown in the above equations
⁶⁵ are either forbidden or suppressed like the decays to transverse polarizations.

⁶⁶ From this section, a very simple picture emerges. At small ξ , all the decay widths
⁶⁷ are fixed with a given resonance mass M_V and the couplings $\{g^2 c_F/g_V, g_V c_H\}$ which
⁶⁸ control the BRs in all relevant channels. Parameters c_{VVV} , c_{VHH} and c_{VW} are
⁶⁹ basically irrelevant. Thus, the basic phenomenology of this model is well described
⁷⁰ by a good approximation.

⁷¹ 1.2.3 Explicit Models

Now the general picture is clear, we can get exact values of the widths and BRs from explicit models. Consider two benchmark models, A and B, which correspond to two explicit models describing the heavy vectors in Refs. [17] and [12] respectively. All the c parameters are fixed to specific values in these models and the only free parameters are the resonance mass M_V and coupling g_V . Moreover, model A is inspired by weakly coupled extensions of the SM gauge group while model B is by strongly coupled scenarios of EWSB, *i.e.* Composite Higgs models, we will consider them in different regions of g_V , relatively small $g_V \leq 3$ and relatively large $g_V \geq 3$.

Figure 1-1 shows the BRs as functions of the mass in model A and B. As expected from the previous discussion and according to Refs. [17], model A predicts

$$c_H \simeq -g^2/g_V^2, c_F \simeq 1, \\ g_V^2 c_H^2 \simeq g^4 c_F^2/g_V^2 \simeq g^2/g_V. \quad (1.20)$$

Therefore Eq. (1.18) and (1.19) can be determined in the following form for V_0 in model A ($g_V = 1$),

$$\Gamma_{V_0 \rightarrow f\bar{f}'} \simeq N_c[f] \frac{g^4 M_V}{96\pi} \\ \Gamma_{V_0 \rightarrow W^+W^-} \simeq \Gamma_{V_0 \rightarrow Zh} \simeq \frac{g^4 M_V}{192\pi}. \quad (1.21)$$

One can easily check either from the plot or the equation, a factor of two difference comparing the BRs between fermions and bosons. Due to the color factor, leptons and quarks also have a difference by a factor of three. Since the c_F term is universal both in A and B. The total width in model A decreases with increasing g_V because of the overall suppression (g^2/g_V) in Eq. (1.20).

On the contrary, in model B the c_H term is unsuppressed

$$c_H \simeq c_F \simeq 1 , \\ g_V^2 c_H^2 \simeq g_V^2 , g^2 c_{c_F} / g_{g_V} \simeq g^2 / g_V . \quad (1.22)$$

Thus the determinate V_0 decay widths for model B ($g_V = 3$) are

$$\Gamma_{V_0 \rightarrow f\bar{f}'} \simeq N_c[f] \frac{g^4 M_V}{342\pi} \\ \Gamma_{V_0 \rightarrow W^+W^-} \simeq \Gamma_{V_0 \rightarrow Zh} \simeq \frac{3M_V}{64\pi} . \quad (1.23)$$

- ⁷² For model B _{$g_V=3$} the dominant BRs are into di-bosons and the fermionic decays are
- ⁷³ extremely suppressed. Moreover, the total width increases with increasing g_V since it
- ⁷⁴ is dominated by the di-boson width which grows with g_V as expected from Eq. (1.22).
- ⁷⁵ This model B is particularly interesting for the present search, since it predicts signal
- ⁷⁶ cross sections of the order of fb [15] [18] [Fig. 1-2], branching ratios to vector bosons
- ⁷⁷ close to unity, and thus being accessible at the LHC. In the latter chapters, the mass
- ⁷⁸ eigenstate of the neutral heavy vector boson in model B scenario refers to the Z'
- ⁷⁹ particle, which is the search target of this thesis.

⁸⁰

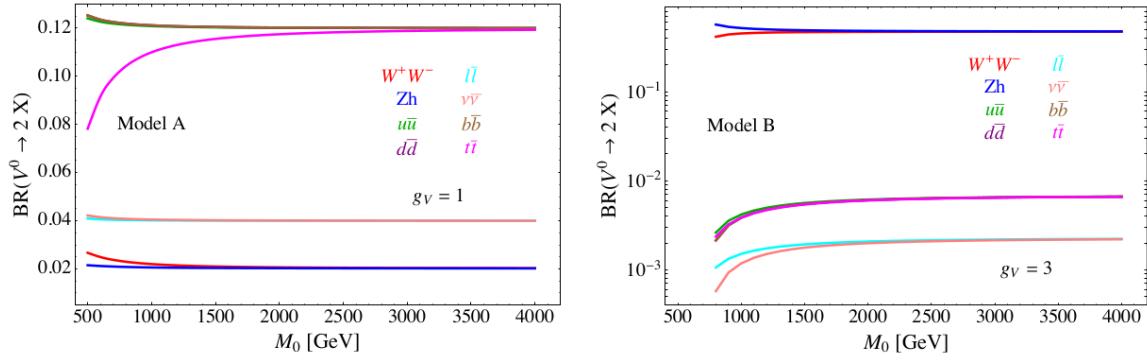


Figure 1-1: Branching ratios as a function of the resonance mass for the HVT benchmark model A(left) and model B(right).

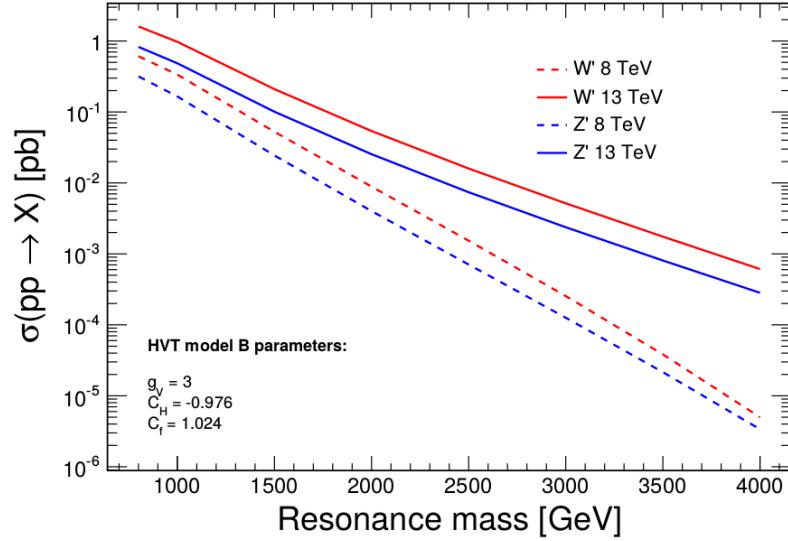


Figure 1-2: Theoretical production cross-section as a function of new resonance particles for HVT model B benchmark. Dash lines are 8 TeV predictions while solid lines are 13 TeV predictions.

81 **Chapter 2**

82 **The LHC and the CMS Detector**

83 This thesis is based on the data collected by the Compact Muon Solenid (CMS)
84 detector at the Large Hadron Collider (LHC). CMS is one of the two largest detectors
85 built at the LHC. This chapter will briefly introduce the LHC and the CMS detector.

86 **2.1 Large Hadron Collider**

87 The LHC is the world's most powerful hadron collider and the largest experimental
88 facility ever. It was built by the European Organization for Nuclear Research (CERN)
89 between 1998 and 2008 in collaboration with over 10,000 scientists and engineers from
90 over 100 countries, as well as hundreds of universities and laboratories. It lies in a
91 tunnel of 27 km in circumference, as deep as 175 m beneath the France–Switzerland
92 border near Geneva. The designed maximum collision energy and highest luminosity
93 of the LHC are 14 TeV and $10^{-34}\text{cm}^{-2}\text{s}^{-1}$, respectively.

94 Other accelerators that had been originally built at CERN for previous experi-
95 ments serve as an injection chain for the LHC now (Fig. 2-2). The proton beam starts
96 from LINAC, a small linear accelerator, where the energy of protons first reaches at
97 50 MeV. The proton beam then passes through a booster and goes to the PS, where
98 it is accelerated up to 25 GeV. After that, it reaches 450 GeV in the SPS. The beam
99 is finally injected in the LHC ring from the SPS, and it had been accelerated up to
100 4 TeV in 2012. In early 2015, the proton beam had been accelerated to 6.5 TeV, a

101 value near its designed energy, before undergoing collision.

102 There are four collision points at the LHC, corresponding to four main experiments, CMS, ATLAS, LHCb and ALICE. The ALICE experiment is optimized to
103 study heavy-ion (Pb-Pb nuclei) collisions and focusing on the physics of strongly
104 interacting matter at extreme energy densities. LHCb is a specialized b-physics ex-
105 periment, measuring the parameters of CP violation in the interactions of b-hadrons.
106 Such studies can help to explain the matter-antimatter asymmetry of the universe.
107 Last, CMS and ATLAS are two general purpose detectors. The aims of these two
108 experiments are investigating a wide range of physics, including the search for the
109 beyond standard model particles, extra dimensions, and dark matter.
110



Figure 2-1: Overview of the LHC and relative location of the detectors.

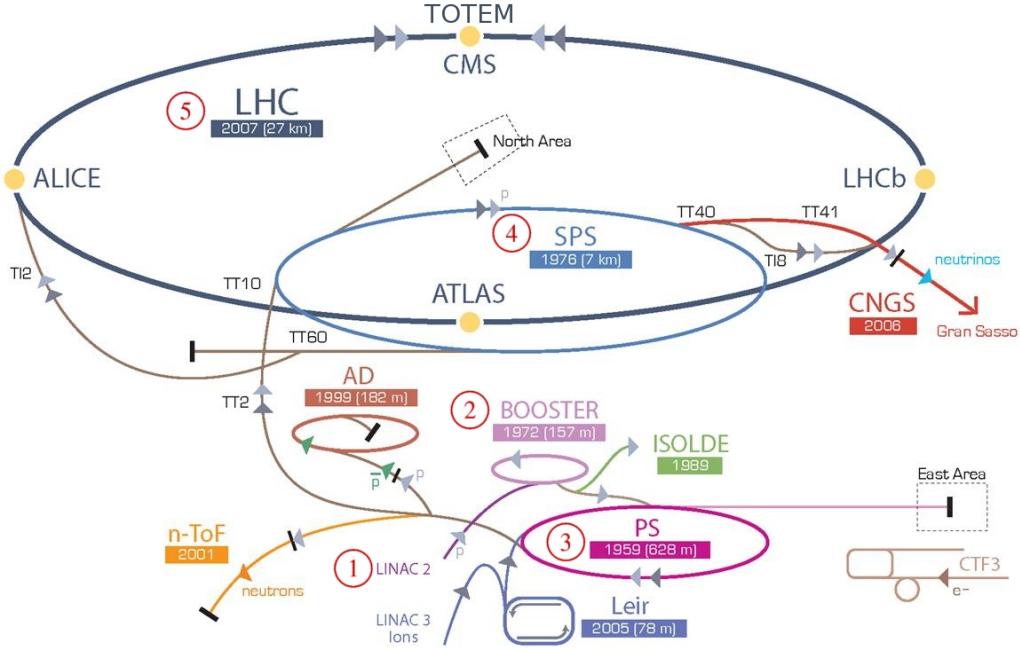


Figure 2-2: CERN accelerator complex.

2.2 Compact Muon Solenoid

The Compact Muon Solenoid (CMS) is designed to cope very high rate of interactions expected to take place at the high LHC luminosity. It has the typical structure of detectors at hadron colliders: a central region (*barrel*) enclosed by two disks (*endcaps*). The structure of CMS can be seen in Fig. 2-3.

Solenoid and Sub-detectors

CMS features a powerful superconducting coil, generating a solenoidal magnetic field around 3.8 Tesla in a large volume which hosts different sub-detectors. The magnetic field lines close through steel yoke in the outer region. The distinct sub-detectors are designed in order to obtain the highest possible resolution and the largest acceptance for every kind of particles.

The innermost layer is a silicon-based tracker. Surrounding it is a scintillating crystal electromagnetic calorimeter (ECAL), which is itself surrounded with a sampling calorimeter for hadrons (HCAL). The tracker and the calorimeters are compact

125 enough to fit inside the CMS Solenoid. Outside the magnet are the large muon
126 chambers.

CMS Detector

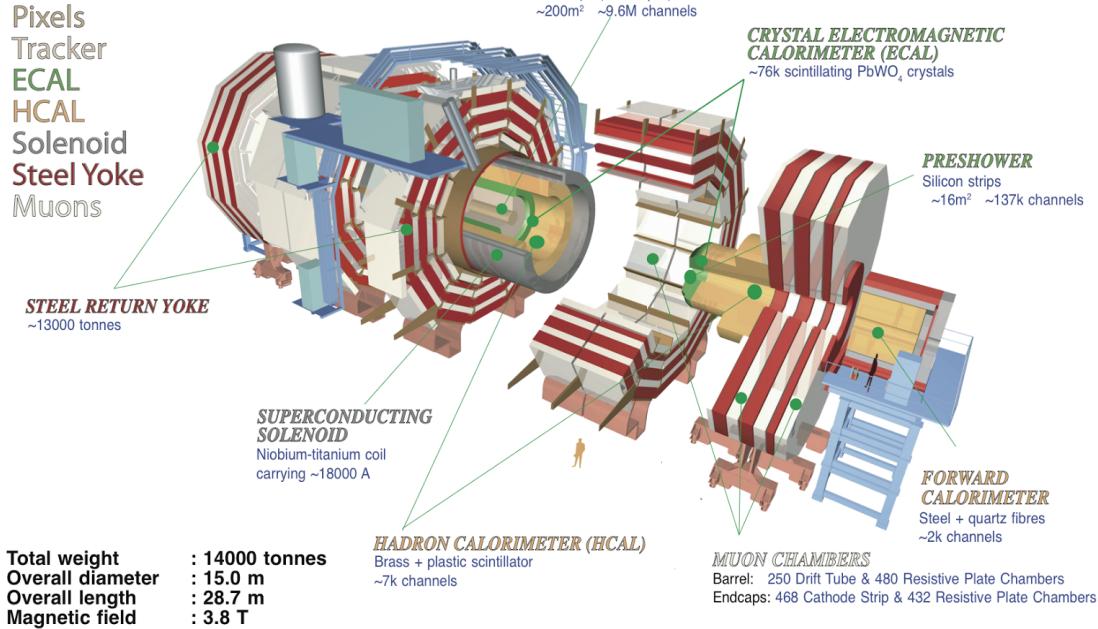


Figure 2-3: Structure overview of the CMS detector.

127 Coordinate System

128 The CMS coordinate system is oriented such that the x -axis points to the center of
129 the LHC ring, the y -axis points vertically upward and the z -axis is in the direction
130 of the beam. The azimuthal angle ϕ is measured from the x -axis in the $x - y$ plane
131 and the radial coordinate in this plane is denoted by r . The polar angle θ is defined
132 in the $r - z$ plane, while the pseudo-rapidity $\eta = -\ln \tan(\theta/2)$. The momentum
133 component transverse to the beam direction, denoted by p_T , is computed from the x -
134 and y -components, and the transverse energy is defined as $E_T = E \sin \theta$.

¹³⁵ 2.2.1 Tracker

¹³⁶ Tracker is the most inner part of CMS that records the productions of collisions in the
¹³⁷ first place. It traces the charged particles' trajectories. Physicists can reconstruct the
¹³⁸ vertices of the interaction and the momentum of charged particles by linking tracks
¹³⁹ to the collider's pipe and measuring the curves of particles under magnetic field.

¹⁴⁰ The tracking system is composed of two kinds of detector, the pixel detector and
¹⁴¹ silicon strip detector. The pixel detector is built from three barrel layers at $r = 44$,
73, 102 mm, and two endcap disks on each side at $z = \pm 345, \pm 465$ mm.

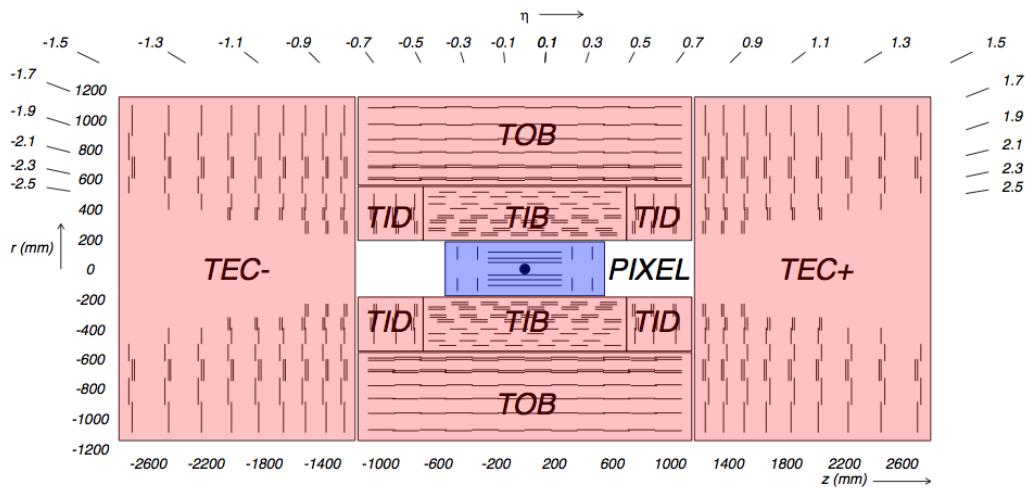


Figure 2-4: Schematic layout of tracker.

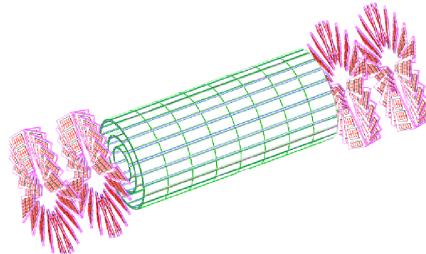


Figure 2-5: The pixel detector inside tracker.

¹⁴²

¹⁴³ The pixel detector, though about the size of a shoebox, consists of 1440 segmented
¹⁴⁴ silicon sensor modules with total 66 million readout channels. Charge carriers are
¹⁴⁵ distributed over several pixels. The analog pulse height information can be used

146 to calculate the center of certain charge distribution which could improve the hit
147 information. The spatial resolution is measured to be about $10 \mu\text{m}$ for the $r - \phi$
148 plane or about $20 \mu\text{m}$ for z direction measurement.

149 Outside the pixel detector, there comes the silicon strip detector. The barrel
150 region of silicon strip detector is divided into two parts, the Tracker Inner Barrel
151 (TIB) and the Tracker Outer Barrel (TOB). The former is composed of four layers
152 of silicon sensors with a thickness of $320 \mu\text{m}$ and of strip pitches varying from 80
153 to $120 \mu\text{m}$. The TOB is made of six layers. In this kind of sub-detector, thicker
154 silicon sensors ($500 \mu\text{m}$) are employed, while the strip pitch varies from 120 to 180
155 μm . The endcap region ($|\eta| > 1.6$) is covered by the Tracker Inner Disks (TID)
156 and the Tracker End Cap (TEC). The entire silicon strip detector is comprised of
157 15200 high-sensitivity modules consisting of detecting unit, supporting structure and
158 readout electronic system.

159 2.2.2 ECAL

160 The Electromagnetic Calorimeter (ECAL) measures the energy of photons, electrons
161 and positrons. It it is placed just outside the tracker, but still inside the solenoid.
162 ECAL is made of 74848 lead-tungstate (PbWO_4) crystals. This material is charac-
163 terized by a high density (8.28 g/cm^3), which gives the crystals a very compact form
164 and makes them particularly suitable to be placed inside the magnetic coil. Another
165 reason, this material has also a fast temporal response ($\sim 10 \text{ ns}$) and its radiation
166 length (X_0) of 0.89 cm give ECAL the possibility to fully contain the expansion of
167 the electromagnetic shower.

168 The arrangement of ECAL is shown in Fig. (2-6). The barrel crystals have a front
169 face area of $2.2 \times 2.2 \text{ cm}^2$ and a length of 23 cm. They are positioned at $r = 1.29$
170 m in pseudo-rapidity region $0 < |\eta| < 1.479$. The crystals in the endcaps have a
171 $2.47 \times 2.47 \text{ cm}^2$ front face, a 22 cm length and they are positioned at $z = 3.17 \text{ m}$ in
172 $1.479 < |\eta| < 3.0$. A Preshower detector is placed in front of the endcaps crystals.
173 The active elements of Preshower are two planes of silicon strips with a pitch of 1.9
174 mm, which lie behind disks of lead absorber at depths of $2X_0$ and $3X_0$. It allows the

175 rejection of photon pairs from π^0 decays and improves the estimation of the direction
176 of photons, to enhance the measurement of the two-photon invariant mass.

The energy resolution of the ECAL is given by three different contributions [19]
(E in GeV),

$$\frac{\sigma_E}{E} = \frac{2.8\%}{\sqrt{E}} \oplus \frac{12\%}{E} \oplus 0.3\% \quad (2.1)$$

177 where the first term is statistical in nature, it also contains fluctuation in showering
178 and in the amplification through photodiodes, the second one considers electronic
179 noise and pile-up, the last term is mainly due to the calibration.

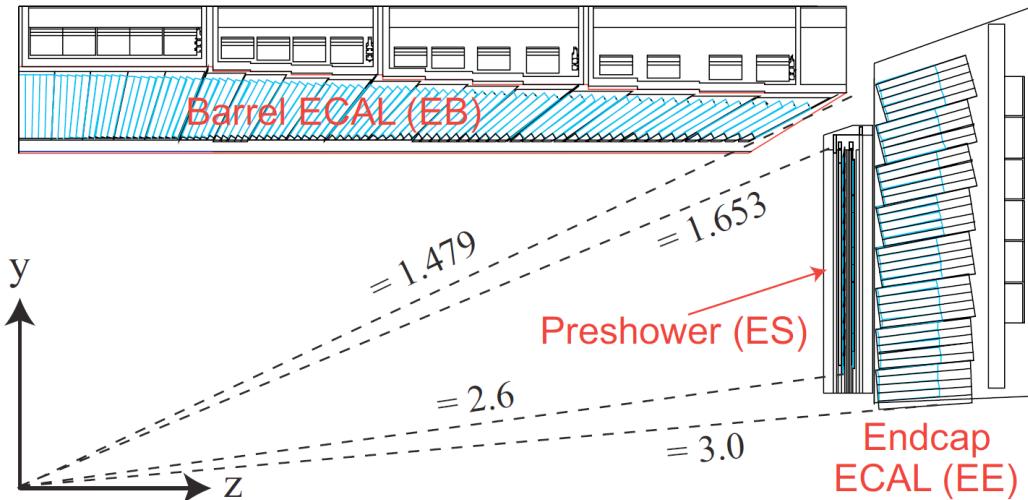


Figure 2-6: Schematic layout of the CMS ECAL.

180 2.2.3 HCAL

181 The hadron calorimeter (HCAL) is placed mainly between ECAL and the magnet
182 coil. It measures the energy of hadrons and mesons. Additionally it provides indirect
183 measurement of the presence of non-interacting, uncharged particles such as neutrinos.
184 The design is strongly influenced by these aims, hence an important requirement is the
185 high hermeticity (the ability to capture every particle emerging from the collisions).
186 This means the detector must cover the biggest possible portion of the solid angle.

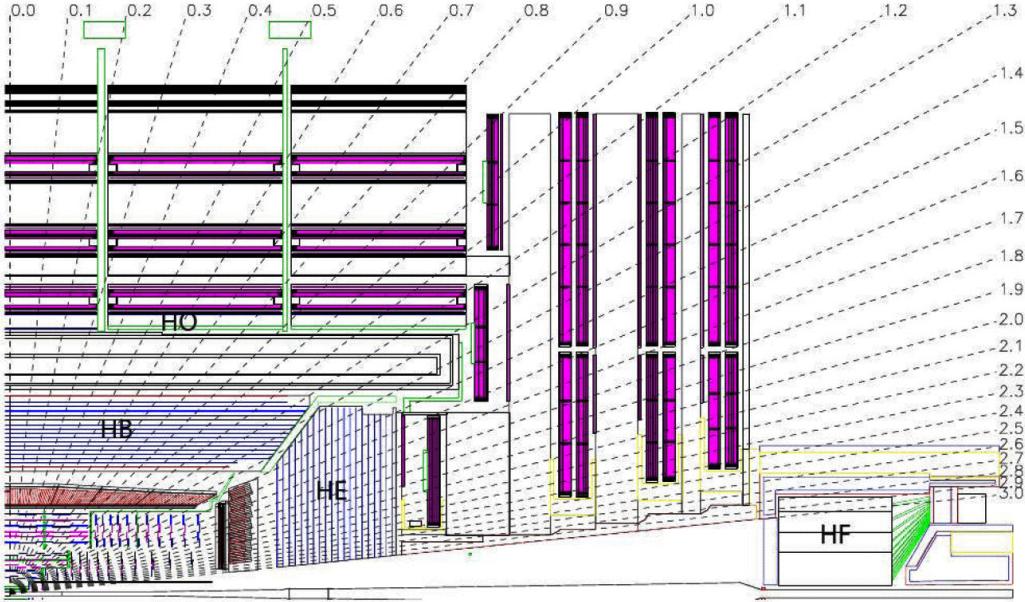


Figure 2-7: Longitudinal view of one quarter of CMS and the location of HB (hadron barrel calorimeter), HE (hadron endcap calorimeter), HF (hadron forward calorimeter) and HO (hadron outer calorimeter)

187 For this reason, a hadron forward calorimeter is required, which is placed outside
 188 the magnet return yokes, with a total coverage of $3 < |\eta| < 5.31$ at 11 m from the
 189 interaction point. Moreover, an outer hadronic calorimeter is placed in the first muon
 190 absorber layer in order to enhance the containment of high energy jets in the central
 191 region of the detector.

192 HCAL is a sampling calorimeter, whose active elements are plastic scintillators
 193 interleaved with brass absorber plates and read out by wavelength shifting fibers.
 194 Brass has been chosen as absorber material for its short interaction length and because
 195 it is non-magnetic. The thickness of the absorber layers is between 60 mm in the barrel
 196 and 80 mm in the endcaps. The barrel has 5.46 interaction lengths at $\eta = 0$ and 10.82
 197 at $\eta = 1.3$, while the endcaps have an average of 11 interaction lengths [20].

The HCAL energy resolution (E in GeV and measured by pion) [21] is

$$\frac{\sigma_E}{E} \simeq \frac{a}{\sqrt{E}} \oplus 5\% \quad (2.2)$$

198 where $a \simeq 65\%$ in the barrel, $a \simeq 85\%$ in the endcaps and $a \simeq 100\%$ in the HF.

199 **2.2.4 Muon Chamber**

200 The efficient detection of muons has primary importance, as muons represent a clear
201 signature for a large number of processes. Muons can penetrate several meters of
202 iron without interacting. Unlike most particles, they are not stopped by any of
203 calorimeters in CMS. Therefore, chambers to detect muons are placed at the very
204 edge of the experiment where they are the only particles likely to register a signal.

205 The muon system fulfills three purposes, muon identification, momentum mea-
206 surement and triggering. Three different types of gaseous detectors are used for CMS
207 muon system depending on the requirements.

208 **Drift Tube**

209 The drift tube (DT) system measures muon positions in the barrel part of the detector.
210 Each DT chamber, on average $2\text{ m} \times 2.5\text{ m}$ in size, consists of 12 aluminium layers,
211 arranged in three groups of four segmentations, each with up to 60 4-cm-wide tubes
212 that contain a stretched wire within each gas volume. The middle group measures
213 the coordinate along the direction parallel to the beam and the two outside groups
214 measure the perpendicular coordinate.

215 **Cathod Strip Chamber**

216 In the two endcaps, where the muon flux and the residual inhomogeneous magnetic
217 field are higher, cathode strip chambers (CSC) are used. CSC is composed of anode
218 wires and cathod strips in the gas volume. The chambers are arranged in 4 disks
219 perpendicular to the beam, and in concentric rings (3 rings in the innermost station,
220 2 in the others) in each of the endcaps.

221 **Resistive Plate Chambers**

222 Resistive plate chambers (RPC) are fast gaseous detectors that provide a muon trigger
223 system parallel with DTs and CSCs. Each RPC consists of two parallel plates, a

²²⁴ positively charged anode and a negatively charged cathode, both made of a very high
²²⁵ resistivity plastic material and separated by a gas volume.

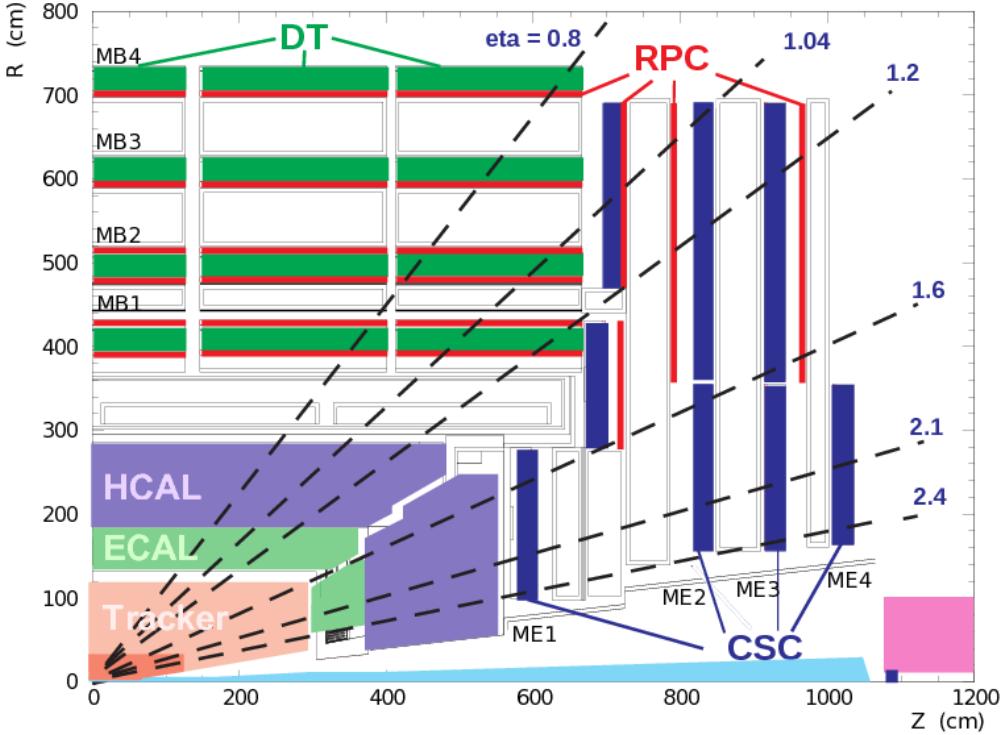


Figure 2-8: Slice view of one quarter of muon chamber system.

²²⁶ 2.2.5 Trigger System

²²⁷ To have a good chance of producing rare particles, a very large number of collisions
²²⁸ is required (LHC proton bunches collide every 25 ns). Most collision events in the
²²⁹ detector are soft and do not produce interesting events. The amount of data from
²³⁰ each crossing is approximately 1 megabytes, which at the 40 MHz crossing rate would
²³¹ result in 40 terabytes of data a second, an amount that the experiment cannot store.
²³² The task of the trigger system is to reduce the storage rate while keeping a high
²³³ efficiency on the potentially interesting events. In CMS the input rate is reduced in
²³⁴ two steps, Level-1 Trigger (L1T) and High-Level Trigger (HLT).

235 **Level-1 Trigger**

236 After the Level-1 Trigger selection, the event recording frequency is decreased to
237 100kHz [22], which is much smaller than the collision rate. The L1T objects are
238 particles (such as photons, muons and electrons), jet candidates, global transverse
239 energy and missing transverse energy. Level-1 Trigger just chooses the event with E_T
240 and P_T higher than the thresholds.

241 **High Level Trigger**

242 High Level Trigger is behind the readout buffers after Level-1 Trigger. It reduces
243 the data output rate to 100Hz by using all the information from CMS including the
244 sub-detectors. The reconstruction algorithms are the same as the off-line analysis.
245 However, the triggering procedure doesn't need maximal precision, therefore these
246 algorithms are modified to be faster even with lower resolution.

247

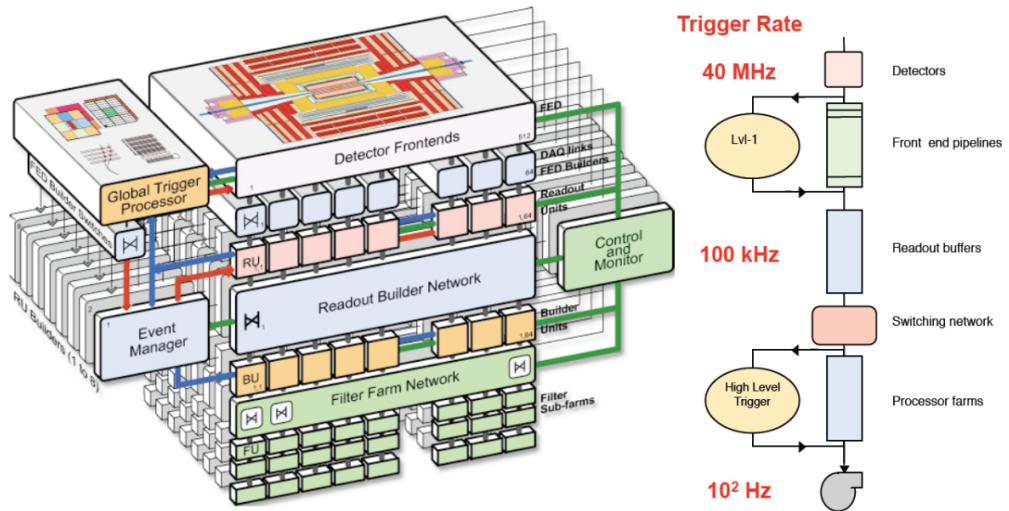


Figure 2-9: CMS triggering and data acquisition architecture.

248 **Chapter 3**

249 **Analysis Procedures**

250 In this chapter, the analysis procedures of the search for Z' decaying into $Z h$ in $llbb$
251 final state are reported. The data sets and Monte Carlo (MC) samples we used in this
252 analysis will be indicated. Physics objects reconstruction and event selections are also
253 introduced. Moreover, background yields and the effects of systematic uncertainties
254 will be demonstrated in the end of this chapter.

255 **3.1 Monte Carlo Samples and Data sets**

256 **3.1.1 Signal MC**

257 As introduced in section 1.2.3, the signal hypothesis is HVT model B benchmark.
258 The heavy resonance (Z') is tested using a wide set of masses from 800 GeV to
259 2000 GeV, one masspoint every 100 GeV (Table 3.1). The signal is simulated by
260 MadGraph5_aMC@NLO [23] in LO mode, as a narrow spin-1 neutral resonance and
261 is forced to decay in the $Z' \rightarrow Z h \rightarrow llqq$ channel. Showering and hadronization are
262 performed with PYTHIA6 [24].

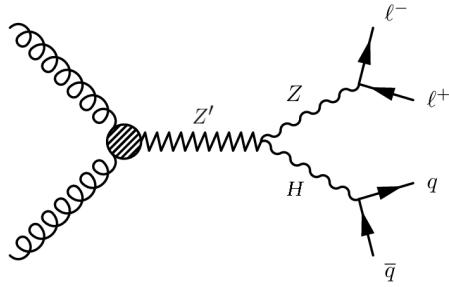


Figure 3-1: Feynman diagram for $Z' \rightarrow Zh \rightarrow 2l2q$.

Sample	Number of Processed Events	$\sigma_{LO}(\text{pb})$
ZPrime_ZH_lljj_M800-MADGRAPH	10710	0.00685367
ZPrime_ZH_lljj_M900-MADGRAPH	10209	0.00485861
ZPrime_ZH_lljj_M1000-MADGRAPH	19997	0.003263
ZPrime_ZH_lljj_M1100-MADGRAPH	9370	0.00217483
ZPrime_ZH_lljj_M1200-MADGRAPH	10710	0.00145484
ZPrime_ZH_lljj_M1300-MADGRAPH	9369	0.000979745
ZPrime_ZH_lljj_M1400-MADGRAPH	10497	0.000664783
ZPrime_ZH_lljj_M1500-MADGRAPH	19999	0.000454339
ZPrime_ZH_lljj_M1600-MADGRAPH	8950	0.000312541
ZPrime_ZH_lljj_M1700-MADGRAPH	9369	0.000216282
ZPrime_ZH_lljj_M1800-MADGRAPH	10708	0.000150398
ZPrime_ZH_lljj_M1900-MADGRAPH	10498	0.000105039
ZPrime_ZH_lljj_M2000-MADGRAPH	19999	7.36377e-05

Table 3.1: Signal samples used in the analysis.

263 **3.1.2 Background MC**

264 Since we are looking for new resonances decaying in semi-leptonic final state, the
 265 background samples of this analysis are originated by all SM events with two leptons
 266 and at least one jet as final state. The dominant background contribution is the
 267 production of Z boson with jets. This Z+jets sample is produced by MADGRAPH.
 268 In the matrix element level, the Z boson is forced decaying into two leptons, and
 269 further this sample is divided into two samples depending on the Z p_T , higher than
 270 100 GeV or between 70 and 100 GeV. The contribution of events with Z p_T less than
 271 70 GeV is negligible due to further cut on the objects p_T in the selection criteria.

272 The second dominant source of background is $t\bar{t}$ production. Both of the two top
 273 quarks decay into all leptonic final state (top decays into a W boson and a b quark
 274 first) which gives two leptons, neutrinos and two b-jets. This sample is generated by
 275 POWHEG [25].

276 Other sources of background considered are SM di-boson productions (WW, WZ
 277 and ZZ) generated by PYTHIA6. All the background samples are required to pass
 278 phase-space cuts, $p_T^l > 60$ GeV and $60 < M_{ll} < 120$ GeV. Related statistics are reported
 279 in Table 3.2.

Sample	Number of Processed Events	$\sigma_{NLO}(\text{pb})$
DYJetsToLL_PtZ-70To100	11764538	63.5
DYJetsToLL_PtZ-100	12511326	39.4
TTTo2L2Nu2B	10783509	25.8
WW	7759752	$56.0 \pm 2.3 (\pm 0.3)$
WZ	9910267	22.4
ZZ	9769891	$7.6 \pm 0.3 (\pm 0.3)$

Table 3.2: Background samples used in the analysis.

280 **3.1.3 Data Samples**

281 In this analysis, the full CMS data collected in 2012 is used, corresponding to the
 282 integrated luminosity of 19.7 fb^{-1} at center-of-mass energy $\sqrt{s} = 8$ TeV. For each
 283 lepton channel, there are four datasets. All datasets are collected with a double

muon or a double electron trigger, as explained in detail in the next section. The trigger algorithm employed for the electron samples doesn't use any information from the tracker but only the energy deposite in the ECAL. This expedient is implemented in order to avoid any possible inefficiencies due to the presence of two tracks very close to each other when the Z is highly boosted and its decay products are very collimated. Such a trigger is contained in the Photon/DoublePhotonHighPt dataset.
 The full dataset names are listed in Table 3.3.

AOD Sample	Luminosity (pb^{-1})
DoubleMu/Run2012A-22Jan2013-v1	876.225
DoubleMuParked/Run2012B-22Jan2013-v1	4409
DoubleMuParked/Run2012C-22Jan2013-v1	7017
DoubleMuParked/Run2012D-22Jan2013-v1	7369
Photon/Run2012A-22Jan2013-v1	876.225
DoublePhotonHighPt/Run2012B-22Jan2013-v1	4412
DoublePhotonHighPt/Run2012C-22Jan2013-v1	7055
DoublePhotonHighPt/Run2012D-22Jan2013-v1	7369

Table 3.3: Data sets used in this analysis.

3.2 Trigger

Since the final state contains two same flavour leptons and at least one jet, we perform this analysis on the DoubleMu and Photon/DoublePhotonHighPt datasets. The first dataset is triggered by two muons, the second one is triggered by two eletrons. These triggers are:

- HLT_Mu22_TkMu8* (for DoubleMu datasets)
- HLT_DoubleEle33_*(for Photon/DoublePhontonHighPt datasets)

The muon trigger has a double p_T threshold, requires leading muon p_T greater than 22 GeV and sub-leading muon p_T greater than 8 GeV. Differently, the double electron trigger requires a higher threshold of 33 GeV to electrons. The trigger efficiencies are close to 1 in both cases.

302 **3.3 Physics Objects**

303 **3.3.1 Muon**

304 **Reconstruction**

305 The muon reconstruction algorithm at CMS takes advantage of the redundancy of
306 detection methods. Muon tracks are first reconstructed independently in the inner
307 tracker (tracker track) and in the muon system (standalone track). Based on these
308 objects, two reconstruction approaches are used [26]:

309 • *Globol Muon* (outside-in): Starting from a standalone track, this algorithm
310 finds a best tracker track to match the standalone track. Then, the fit of the
311 track is repeated using the hits both in the tracker and in the muon system [27].
312 The resulting object is called a *Global Muon*. At large transverse momentum
313 ($p_t > 200$ GeV), the global muon fit can improve the momentum resolution
314 compared to the tracker only fit.

315 • *Tracker Muon* (inside-out): A tracker muon is reconstructed by an opposite
316 direction from a global muon. In this approach, all tracker tracks with $p_T >$
317 0.5 GeV and the total momentum $p > 2.5$ GeV are considered as possible
318 muon candidates. The extrapolation to the muon system takes into account the
319 magnetic field, average expected energy losses, and multiple scattering in the
320 detector material. If at least one muon segment matches the extrapolated track,
321 the corresponding track track qualifies as a *Tracker Muon*. This algorithm is
322 useful for low- p_T muons that are not fully penetrate the muon system, and
323 therefore only register a few hits

324 If no match is found when extrapolating outside-in, the standalone track is stored
325 as a *Stanalone Muon*. This happens only for less than 1% of the muons produced in
326 a collison, and the reconstruction efficiency is about 99% for the muon which carries
327 enough high momentum within detector coverage [26].

328 **Identification**

329 We use both tracker muons and global muons in this analysis. To identify muons
330 from the signal, the muons must pass one of these two off-line selections, high- p_T
331 muon ID or tracker-based muon ID [28]. The requirements are listed as follows:

332

333 High- p_T muon ID

334 • Muon identified as a *Global Muon*.

335 • Number of muon hits in the global track > 0 .

336 • Number of matched muon stations > 1 .

337 • Number of pixel hits > 0 .

338 • Number of tracker layer with hits > 8 .

339 • Transverse impact parameter $d_{xy} < 0.2$ cm.

340 • Longitudinal impact parameter $d_z < 0.5$ cm.

341 • Relative error on the track transverse momentum $\sigma_{p_T}/p_T < 0.3$.

342

343 In the tracker-based muon ID, the muon has to be identified as a *Tracker Muon*,
344 and the requirement of muon hits in the global track is removed. Other requirements
345 are the same.

346 An additional useful variable for lepton identification is the isolation. It is defined
347 as the scalar sum of the p_T of the reconstructed objects within a cone (typical size
348 is $\Delta R = 0.3$) space around the lepton track but excluding the p_T of the lepton
349 itself. Moreover, the relative isolation is defined as isolation divided by the lepton
350 p_T ($I_{rel} = Iso/p_T^{lept}$). The relative isolation is more frequently used in the modern
351 analysis.

352 In this analysis, a modified isolation criteria is used. The two muons originated
353 from boosted Z decay are close to each other, and consequently the presence of another

354 muon in the isolation cone could break the function of this variable. In order to solve
 355 this problem we use a tracker-based isolation relative quantity, explicitely removing
 356 the momentum flow from any other muon passing our muon selection. Moreover, a
 357 tracker-based isolation is well motivated also by two additional aspects: it is more
 358 independent of pile up (pile-up tracks tipically do not match the primary vertex)
 359 and does not include possible muon radiation. Finally, the modified requirement is
 360 $I_{rel}^{mod} < 0.1$.

361

Variable	High- p_T	Tracker-based
Muon type	Global muon	Tracker muon
Muon hits in global track	≥ 1	-
Muon stations matched	≥ 2	≥ 2
d_{xy}	< 0.2 cm	< 0.2 cm
d_z	< 0.5 cm	< 0.5 cm
Pixel hits	≥ 1	≥ 1
Tracker layers	≥ 8	≥ 8
σ_{p_T}/p_T	< 0.3	< 0.3
I_{rel}^{mod}	< 0.1	< 0.1

Table 3.4: Summary of the muon ID selection criteria.

362 **3.3.2 Electron**

363 **Reconstruction**

364 Electrons are reconstructed from energy deposits in the ECAL matched to tracks
365 reconstructed in the silicon tracker. They have less distinguishable signatures than
366 muons in the detector due to the jet induced background in ECAL. Therefore, to
367 reconstruct an election, it's essential to find a brilliant way to match the infomation
368 acquired in both tracker and ECAL.

369 **Clustering**

370 For every single hit from an electron in ECAL, approximately 94% of energy is con-
371 tained in 3×3 crystals and 97% in 5×5 crystals. To form a cluster, crystals with
372 $E_T > 1$ GeV is picked as seeds. Then starting from seed crystal, dominos of size 1×3
373 or 1×5 are created in $\eta - \phi$ plane. If the energy of these dominos exceed 0.1 GeV
374 threshold, thus added up the energy of dominos along the ϕ direction.

375 Moreover, the effects of conversion process and bremsstrahlung radiation must be
376 taken into account to reconstruct the electron energy more precisely. The concept
377 of electron superclustering is to collect the energy of photons from bremsstrahlung
378 radiation along the ϕ direction with fixed η width of the seed crystal. The technical
379 details of the supercluster can be found in [29].

380 **Tracking**

381 There are several steps to reconstruct electron tracks in CMS experiment. The hit
382 on the pixel and the tracker are reconstructed in the first step. The second step is
383 to find the seed of tracks by matching at least two hits in pixel detector. Lst step
384 is to fit the trajetory starting from the seed. To perform this fit, the Gaussian-sum
385 filter (GSF) algorithm is used [30]. In GSF algorithm, the energy loss probability
386 density function is constructed by multiple Gaussian functions. Finally, the electrons
387 are built by matching the superclusters to the GSF tracks. Reconstruction efficiencies

388 for electrons $E_T > 20$ GeV are generally above 95% in EB and 90% in EE [31].

389 Identification

390 The electron identification used in this analysis is based on HEEPv4.1 [32]. As what
391 we did for the muon id, we need to modify the isolation definition again because of
392 the small ΔR between two electrons. The requirements are listed as follows:

393

394 Modified HEEPv4.1 electron ID

- 395 • Transversal supercluster energy > 35 GeV.
- 396 • Pseudorapidity of supercluster $\eta_{SC} < 1.442$ for barrel electrons, or $1.56 < \eta_{SC} <$
397 2.5 for endcap electrons.
- 398 • Have either $E_{2\times 5}/E_{5\times 5} > 0.94$ or $E_{1\times 5}/E_{5\times 5} > 0.83$.

399 • Ratio of HCAL deposit energy to ECAL deposit energy (Fig. 3-2) smaller than
400 0.5.

- 401 • Number of inner layer lost hits smaller than 2.
- 402 • Have $d_{xy} < 0.02$ cm for barrel electrons, or $d_{xy} < 0.05$ cm for endcap electrons.

403

404 As what we did for the muon id, by the same reason we need to modify the
405 isolation definition again. In this case there are three isolation variables have to be
406 changed.

- 407 • Modified track isolation is required to lower than 5 GeV. This variable is defined
408 as the scalar p_T sum of the tracks within a $\Delta R = 0.3$ cone around the electron,
409 and excluding the p_T of another electron which passes the above selections and
410 its track is inside the cone.

- 411 • The electromagnetic calorimeter isolation I_{ECAL} is defined as the scalar sum of
412 E_T of the crystals in a $\Delta R = 0.3$ cone around the particular electron (an inner

area of full-width 3 crystal), excluding a 4 crystals ΔR cone around any other electron. The dimension of the ECAL crystals corresponds roughly at $\Delta R \sim 0.01$ to 0.02 . The threshold of I_{ECAL} is varying with the electron transverse energy.

- The hadronic calorimeter isolation I_{HCAL} is defined as the scalar sum of E_T of the HCAL caloTowers with a center in a $\Delta R = 0.3$ cone around the electron, excluding those lying within $\Delta R = 0.15$ of the electron itself and of any other electron [33]. The threshold of this variable also varies with the electron transverse energy.

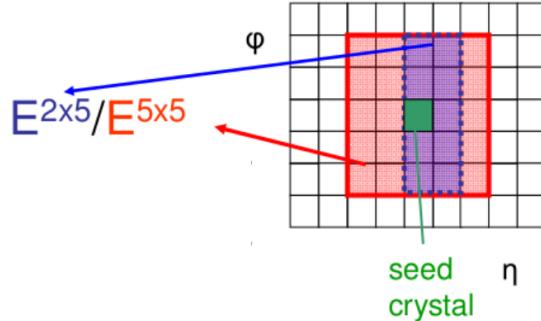


Figure 3-2: Example of the $E_{2\times 5}/E_{5\times 5}$ computation in the ECAL crystals. $E_{i\times j}$ is the energy contained in a $i \times j$ block around the seed crystal (defined as the highest deposit of the energy of the cluster).

Variable	Barrel	Endcap
E_T	$> 35 \text{ GeV}$	$> 35 \text{ GeV}$
$ \eta_{SC} $	$ \eta_{SC} < 1.442$	$1.56 < \eta_{SC} < 2.5$
H/E	< 0.05	< 0.05
$E_{2\times 5}/E_{5\times 5}$	> 0.94 or $E_{1\times 5}/E_{5\times 5} > 0.83$	-
Inner Layer Lost Hits	$<= 1$	$<= 1$
$ d_{xy} $	$< 0.02 \text{ cm}$	$< 0.05 \text{ cm}$
Iso_{Trk}^{mod}	$< 5 \text{ GeV}$	$< 5 \text{ GeV}$
$Iso_{Ecal,Hcal}^{mod}$	$< 2 \text{ GeV} + 0.03E_T$	$< 2.5 \text{ GeV for } E_T < 50 \text{ GeV}$ $< 2.5 \text{ GeV} + 0.03E_T \text{ for } E_T > 50 \text{ GeV}$

Table 3.5: Summary of the modified HEEPv4.1 electron ID.

422 **3.3.3 Jet**

423 Since gluons and quarks cannot exist in free state due to the color confinement
424 [34], they fragment into hadrons. The result of this fragmentation (or called the
425 "hadronization") is a jet of particles depositing energy in the detectors.

426 **Reconstruction**

427 Jet clustering algorithms are among the main tools for analyzing data from hadronic
428 collisions. In this section, a extensively used clustering algorithm called "sequential
429 recombination algorithm" [35,36] will be introduced. At first, events are reconstructed
430 using the particle flow algorithm (PF) [37], which attempts to identify and measure all
431 the stable particles in a collision by combining information from all the sub-detectors.
432 The PF algorithm categorizes all particles into five types: muons, electrons, photons,
433 charged and neutral hadrons. The resulting particle flow candidates are passed to
434 the anti- k_T (AK) [35] or Cambridge/Aachen (CA) [36,38] jet clustering algorithms to
435 create jets (AK and CA are two commonly used branches of sequential recombination
436 algorithms).

437 The jet clustering algorithms are implemented as follows:

$$d_{iB} = p_{T,i}^{2p}$$
$$d_{ij} = \min(p_{T,i}^{2p}, p_{T,j}^{2p}) \frac{\Delta R_{i,j}^2}{R^2} \quad (3.1)$$

438 In Eq. (3.1), the pair of objects (i , j) denote two input candidate particles to
439 the algorithm. $\Delta R_{i,j}$ is the spatial distance between the two objects, and R is the
440 typical cone size of the jet. The parameter p corresponds to different catagories of
441 algorithms and will be introduced in the later paragraph.

442 The clustering proceeds by comparing the value of the two quantities, d_{iB} and d_{ij} .
443 If $d_{ij} > d_{iB}$, the algorithm will look for another possible combination of (i , j). If
444 $d_{ij} < d_{iB}$, the 4-momentum of object i will merge into j, and then the object j forms a

445 pseudo-jet, but the algorithm will not stop looking for a new object until all particles
446 are clustered into jet.

447 The parameter p can be chosen as follows:

448 • $p = 0$: Cambridge/Aachen algorithm;

449 • $p = -1$: Anti- k_T algorithm;

450 The difference between CA/AK algorithms is clear, CA algorithm is only consid-
451 ering the spatial distance of the two candidate objects (if $\Delta R_{ij} < R$, merged). In the
452 case of anti- k_T algorithm, p_T of the object presents as a weight for the distance.

453 The anti- k_T algorithm have better energy calibration, while the CA algorithm
454 was found to be more efficient at finding hard subjets within the jets in simulation-
455 based studies [39]. Therefore the CA jets are used in this analysis to identify events
456 containing hadronically-decaying Higgs bosons.

457 Identification

458 As discussed in the previous section, this analysis use the Cambridge/Aachen algo-
459 rithm with a cone radius of $R = 0.8$ for the identification of jets (CA8 jets). Further-
460 more, jets are required to pass the following loose identification criteria [40, 41]:

461

462 Loose JetID ($> 99\%$ efficiency)

463 • Muon energy fraction smaller than 0.99

464 • Photon energy fraction smaller than 0.99

465 • Charged electromagnetic energy fraction smaller than 0.99

466 • Neutral hadron energy fraction smaller than 0.99

467 • Charged hadron energy fraction larger than 0

468 • Number of constituent particles larger than 1

469 For the "charged" variables, $\eta < 2.4$ requirement are also applied, since there is
470 no track coverage outside of this region, whereas the "neutral" variables extend to
471 the whole η region.

472 3.3.4 Jet Grooming Algorithms and Pruned Jet Mass

473 The jet mass is the main observable in distinguishing a boson-jet from a QCD jet. Jet
474 grooming aims to suppress underlying events and pile-up radiation from the target jet,
475 and improves the discrimination by pushing the jet mass for QCD jets towards lower
476 values while maintaining the jet mass for boson-jets around the boson-mass [42, 43].

477 Typically three different grooming algorithms are considered, trimming [44], fil-
478 tering [45] or pruning [46, 47]. Studies of these different grooming methods in CMS
479 are presented in [48].

480 • **Trimming algorithm:** Trimming is a technique that ignores subjets below a
481 minimum p_T fraction threshold within the original jet. Trimming reclusters the
482 jet's constituents with a smaller radius R_{sub} . Then accepts only the subjets
483 that have $p_{T,frac} > f_{cut}$, where the $p_{T,frac}$ is defined as subjet p_T divided by the
484 original jet p_T , and f_{cut} is typically taken proportional to H_T , the scalar sum of
485 the p_T of all jet reconstructed in the event.

486 • **Filtering algorithm:** Filtering reclusters jet constituents with smaller radius
487 R_{filt} , and sorts the subjets in order of their p_T . The final jet is thus redefined as
488 the sum of the four-momentum of the n hardest subjets. By default the factor
489 is n=3, but it's not a fixed number, n depends on the analysis.

490 • **Pruning algorithm:** The idea is to take a jet of interest and then to reclus-
491 ter it using a sequential clustering algorithm for vetoing soft and large-angle
492 recombinations between pseudojets i and j.

493 Clustering proceeds as explained in the previous section, but it is vetoed if the
494 candidates are too far away in ΔR .

$$\text{veto if } \Delta R_{ij} > r_{cut} \times 2m/p_T \quad (3.2)$$

And the energy sharing is too asymmetric.

$$\text{veto if } z_{ij} = \frac{\min(p_{T,i}, p_{T,j})}{p_{T,i+j}} < z_{cut} \quad (3.3)$$

Where z_{cut} and r_{cut} are parameters of the algorithm (default value: $z_{cut} = 0.1$, $r_{cut} = 0.5$). If both these conditions are satisfied the softer one of the two candidate is not considered.

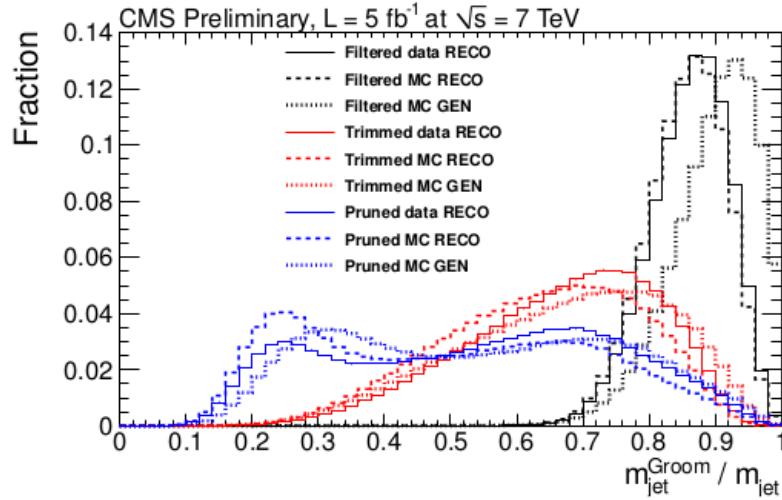


Figure 3-3: Comparison of the jet mass in generic QCD events from the groomed jets divided by the jet mass of matched ungroomed jets for the three grooming techniques, for both data and the PYTHIA 6 Monte Carlo. [49] Events are collected with a single jet trigger.

As shown in Fig. 3-3, the filtering algorithm is the least aggressive grooming technique, with groomed jet masses close to the original case. The trimming algorithm is moderately aggressive which produces a much wider final mass distribution. Pruning is the most aggressive tecnicue and a bimodal distribution begins to appear.

In cases where the pruned jet mass is small, jets usually have most of their energy configured in core components with little gluon radiation, which leads to narrow jets.

499 Instead, when the pruned jet mass is large, the jets are split more symmetrically and
500 more similar with the boson-jet structure. In this analysis we use the pruned jet mass
501 because of its capability to improve the jet mass resolution and background rejection.

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