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Chapter 1

Introduction and Theory Overview

1.1 Introduction

This thesis presents the analysis details and the results of the search for heavy resonances decaying into a Z boson and a Higgs boson (h) at the center-of-mass energy of 8 TeV, using 19.7 fb^{-1} p-p collision data. In turn, the Z boson is identified through its leptonic decays (Leptons often refer to e and μ only in experiments. $l = e, \mu$). The Higgs boson h is expected to hadronically decay into a pair of b-quarks. The investigated final states consist of two charged leptons which are identified in the detector and limit the presence of the background, and two b-quarks from the hadronic Higgs decay which collects the largest possible fraction of Higgs events.

This thesis is organised as follows. In the latter part of this chapter, the heavy resonances model is introduced, including the expected cross section and the specification of model parameters. In chapter 2, the LHC and the CMS experiment are introduced, including the information of each sub-detector and the trigger system of the CMS. The details of the analysis are shown in chapter 3. This chapter reveals the way to reconstruct physical objects in CMS. By adding some proper kinematic selections on those physics objects, the interested event in data collected by the CMS detector can be selected. Moreover, this chapter shows the comparison between data and simulation. In the last chapter, the results of the search and the conclusion are shown.

1.2 Theory Overview

Although the Higgs boson discovery by the ATLAS and CMS collaborations [1–3] imposes strong constraints on theories beyond the Standard Model (SM), the extreme fine tuning in quantum corrections required to have a light fundamental Higgs boson with mass close to 125 GeV [4–7] suggests that the Standard Model may be incomplete, and not valid beyond a scale of a few TeV. Various dynamical electroweak symmetry breaking scenarios which attempt to solve this naturalness problem, such as Minimal Walking Technicolor [8], Little Higgs [9–11], or composite Higgs models [12–14] predict the existence of new resonances decaying to a vector boson plus a Higgs boson.

1.2.1 Heavy Vector Triplet Model

Resonant searches are typically not sensitive to all the details and the free parameters of the underlying model, but only to those parameters or combinations of parameters that control the mass of the resonance and the interactions involved in its production and decay. Therefore one can employ a simplified description of the resonance defined by a phenomenological Lagrangian where only the relevant couplings and mass parameters are retained. This model-independent strategy applies a Heavy Vector Triplet (HVT) [15] to the Standard Model group and reproduces a large class of explicit models. In Eq. (1.1), the mathematical form of the simplified Lagrangian is defined, where V_ν^a , $a = 1, 2, 3$, is a real vector with vanishing hypercharge in the adjoint representation of $SU(2)_L$, it describes one charged and one neutral heavy spin-1 particle with charge eigenstate fields, and $D_{[\mu}V_{\nu]}^a$ represents the covariant derivative.

$$\begin{aligned} \mathcal{L}_V = & -\frac{1}{4}D_{[\mu}V_{\nu]}^a D^{[\mu}V^{\nu]a} + \frac{m_V^2}{2}V_\mu^a V^{\mu a} \\ & + ig_V c_H V_\mu^a H^\dagger \tau^a \overleftrightarrow{D}^\mu H + \frac{g^2}{g_V} c_F V_\mu^a \sum_f \bar{f}_L \gamma^\mu \tau^a f_L \\ & + \frac{g_V}{2} c_{VVV} \epsilon_{abc} V_\mu^a V_\nu^b D^{[\mu}V^{\nu]c} + \text{quadrilinear terms} \end{aligned} \quad (1.1)$$

$$V_\mu^\pm = \frac{V_\mu^1 \mp iV_\mu^2}{\sqrt{2}}, \quad V_\mu^0 = V_\mu^3 \quad (1.2)$$

$$D_{[\mu}V_{\nu]}^a = D_\mu V_\nu^a - D_\nu V_\mu^a, \quad D_\mu V_\nu^a = \partial_\mu V_\nu^a + g\epsilon^{abc}W_\mu^b V_\nu^c \quad (1.3)$$

In these models, new heavy vector bosons (V^\pm, V^0) that couple to the Higgs and SM gauge bosons with the parameters g_V and c_H and to the fermions via the combination $(g^2/g_V)c_F$. The parameter g_V represents the strength of the new vector boson interaction, while c_H and c_F represent the couplings to the Higgs and the fermions respectively, and are expected to be of order unity in most models.

1.2.2 Basic Phenomenology

Masses and Mixings

After electro-weak symmetry breaking (EWSB), the only massless state is the photon which can be identified as the gauge field associated with the unbroken $U(1)_{em}$. The two other neutral mass eigenstates are the SM Z boson and one heavy vector of mass M_0 which are obtained by diagonalizing the mass matrix of the (Z, V^0) system by a rotation with angle θ_N

$$\begin{pmatrix} Z \\ V^0 \end{pmatrix} \rightarrow \begin{pmatrix} \cos \theta_N & \sin \theta_N \\ -\sin \theta_N & \cos \theta_N \end{pmatrix} \begin{pmatrix} Z \\ V^0 \end{pmatrix}. \quad (1.4)$$

The mass matrix is

$$\mathcal{M}_N^2 = \begin{pmatrix} \hat{m}_Z^2 & c_H \xi \hat{m}_Z \hat{m}_V \\ c_H \xi \hat{m}_Z \hat{m}_V & \hat{m}_V^2 \end{pmatrix}, \quad \text{where} \quad \begin{cases} \hat{m}_Z = \frac{e\hat{v}}{2 \sin \theta_W \cos \theta_W} \\ \hat{m}_V^2 = m_V^2 + g_V^2 c_{VVHH} \hat{v}^2 \\ \xi = \frac{g_V \hat{v}}{2 \hat{m}_V} \end{cases} \quad (1.5)$$

In the above equations \hat{v} denotes the Vacuum Expectation Value (VEV) defined by $\langle H^\dagger H \rangle = \hat{v}^2/2$, and one should know the masses \hat{m}_Z and \hat{m}_V do not coincide with the physical Z boson and new resonance masses of this model, although they do in the approximations later. The mass eigenvalues and the rotation angles are easily obtained by inverting the relations

$$\begin{aligned} Tr[\mathcal{M}_N^2] &= \hat{m}_Z^2 + \hat{m}_V^2 = m_Z^2 + M_0^2 , \\ Det[\mathcal{M}_N^2] &= \hat{m}_Z^2 \hat{m}_V^2 (1 - c_H^2 \xi^2) = m_Z^2 M_0^2 , \\ \tan 2\theta_N &= \frac{2c_H \xi \hat{m}_Z \hat{m}_V}{\hat{m}_V^2 - \hat{m}_Z^2} . \end{aligned} \quad (1.6)$$

Notice that the tangent can be uniquely inverted because the angle θ_N is in the range $[-\pi/4, \pi/4]$ in the parameter region we will be interested in, where $\hat{m}_Z < \hat{m}_V$.

The situation is similar in the charged vector where the mass matrix of the (W^\pm, V^\pm) system reads

$$\mathcal{M}_C^2 = \begin{pmatrix} \hat{m}_W^2 & c_H \xi \hat{m}_W \hat{m}_V \\ c_H \xi \hat{m}_W \hat{m}_V & \hat{m}_V^2 \end{pmatrix} , \text{ where } \hat{m}_W = \frac{e\hat{v}}{2 \sin \theta_W} = \cos \theta_W \hat{m}_Z , \quad (1.7)$$

and it is diagonalized by

$$\begin{aligned} Tr[\mathcal{M}_C^2] &= \hat{m}_W^2 + \hat{m}_V^2 = m_W^2 + M_\pm^2 , \\ Det[\mathcal{M}_C^2] &= \hat{m}_W^2 \hat{m}_V^2 (1 - c_H^2 \xi^2) = m_W^2 M_\pm^2 , \\ \tan 2\theta_C &= \frac{2c_H \xi \hat{m}_W \hat{m}_V}{\hat{m}_V^2 - \hat{m}_W^2} . \end{aligned} \quad (1.8)$$

By checking Eq. (1.5) and Eq. (1.7), the charged and neutral mass matrices are connected by custodial symmetry, which can be shown in full generality to imply

$$\mathcal{M}_C^2 = \begin{pmatrix} \cos \theta_W & 0 \\ 0 & 1 \end{pmatrix} \mathcal{M}_N^2 \begin{pmatrix} \cos \theta_W & 0 \\ 0 & 1 \end{pmatrix} . \quad (1.9)$$

By taking the determinant of the above equation, or equivalently by comparing the charged and neutral determinants in Eq. (1.6) and Eq. (1.8), we obtain a generalized custodial relation among the physical masses

$$m_W^2 M_{\pm}^2 = \cos^2 \theta_W m_Z^2 M_0^2 . \quad (1.10)$$

From the simple formula above, we can start to identify the physically reasonable region of the parameter space in this model. We aim at describing new vectors with masses at or above the TeV scale, but we also want the SM masses $m_{W,Z} \sim 100$ GeV to be reproduced. Therefore we require a hierarchy in the mass relation of SM Z and W bosons versus the new vectors.

$$\frac{\hat{m}_{W,Z}}{\hat{m}_V} \sim \frac{m_{W,Z}}{M_{\pm,0}} \leq 10^{-1} \ll 1 \quad (1.11)$$

In the limit of Eq. (1.11) we obtain simple approximation for m_W and m_Z

$$\begin{aligned} m_Z^2 &= \hat{m}_Z^2 (1 - c_H^2 \xi^2) (1 + \mathcal{O}(\hat{m}_Z^2 / \hat{m}_V^2)) , \\ m_W^2 &= \hat{m}_W^2 (1 - c_H^2 \xi^2) (1 + \mathcal{O}(\hat{m}_W^2 / \hat{m}_V^2)) . \end{aligned} \quad (1.12)$$

The parameter ξ can be either very small or of order unity. Both cases are realized in explicit models. While $\xi \ll 1$ is the most common situation, $\xi \sim 1$ only occurs in strongly coupled scenarios at very large g_V . In these approximations, SM tree-level experimental observation can be reproduced to percent accuracy.

Since $\hat{m}_W = \cos \theta_W \hat{m}_Z$, the W - Z mass ratio is thus given by

$$\frac{m_W^2}{m_Z^2} \simeq \cos^2 \theta_W . \quad (1.13)$$

Eq. (1.13) has one important implication on the masses of the new vectors. When combined with the custodial relation Eq. (1.10), it tells us that the charged and

neutral V s are practically degenerate

$$M_{\pm}^2 = M_0^2(1 + \mathcal{O}(\%)) , \quad (1.14)$$

In the following, when working at the leading order in the limit Eq. (1.11), we can ignore the mass splitting and denote the mass of the charged and the neutral states collectively as M_V . It is easy to check that in that limit $M_V = \hat{m}_V$.

Decay Widths

1.2.3 Explicit Models

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