

## 高等热力学与统计物理第九次作业

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1. 一稀薄气体处于外力场内，相应的势能为  $V(\vec{r})$ 。假设  $V(\vec{r})$  在分子相互作用力程范围内的变化很小，求出 Boltzmann 方程的近似静态解并用平均数密度平均动能定义的温度表示所得解。

处于势场  $V(\vec{r})$  中受力为：

$$\vec{F}(\vec{r}) = -\vec{\nabla}V(\vec{r}).$$

则 Boltzmann 方程可以写为：

$$\begin{aligned} & \frac{\partial}{\partial t} f(\vec{r}, \vec{p}; t) + \left( \frac{\vec{p}}{m} \cdot \vec{\nabla}_{\vec{r}} + \vec{F} \cdot \vec{\nabla}_{\vec{p}} \right) f(\vec{r}, \vec{p}; t) \\ &= \int d^3\vec{p}_2 d^3\vec{p}'_1 d^3\vec{p}'_2 |T_{fi}|^2 \delta^4(P_f - P_i) [f(\vec{r}, \vec{p}'_1; t) f(\vec{r}, \vec{p}'_2; t) - f(\vec{r}, \vec{p}; t) f(\vec{r}, \vec{p}_2; t)]. \end{aligned}$$

假设 Boltzmann 方程的试探解为：

$$f(\vec{r}, \vec{p}; t) = C \rho e^{-A(\vec{p} - \vec{p}_0)^2}$$

则坐标空间中粒子数密度为：

$$\rho = \int d^3\vec{p} f(\vec{r}, \vec{p}; t) = C \rho \int d^3\vec{p} e^{-A(\vec{p} - \vec{p}_0)^2} = C \rho \left( \frac{\pi}{A} \right)^{3/2}.$$

则有系数间关系为：

$$C = \left( \frac{A}{\pi} \right)^{3/2}$$

则考虑动量空间中粒子数密度为： $g(\vec{p}; t) = \frac{f(\vec{r}, \vec{p}; t)}{\rho(\vec{r}; t)} = \left( \frac{A}{\pi} \right)^{3/2} e^{-A(\vec{p} - \vec{p}_0)^2}$ 。则平均动量为：

$$\langle \vec{p} \rangle = \int d^3\vec{p} \vec{p} \left( \frac{A}{\pi} \right)^{3/2} e^{-A(\vec{p} - \vec{p}_0)^2} = \left( \frac{A}{\pi} \right)^{3/2} \int d^3\vec{p} (\vec{p} + \vec{p}_0) e^{-A\vec{p}^2} = \vec{p}_0.$$

平均动量为 0 条件即为  $\vec{p}_0 = 0$ 。此时平均动能为：

$$\epsilon \equiv \left\langle \frac{\vec{p}^2}{2m} \right\rangle = \left( \frac{A}{\pi} \right)^{3/2} \frac{1}{2m} \int d^3\vec{p} \vec{p}^2 e^{-A\vec{p}^2} = \frac{3}{4mA}.$$

若定义平均动能对应的温度为  $T_0 = \frac{2\epsilon}{3k}$ ，其中  $k$  为 Boltzmann 常数，则系数可以写为：

$$A = \frac{1}{2mkT_0}.$$

对于该试探解，Boltzmann 方程右侧为 0，则有：

$$\frac{\partial}{\partial t} \rho(\vec{r}; t) + \frac{\vec{p}}{m} \cdot \vec{\nabla}_{\vec{r}} \rho(\vec{r}; t) - \vec{\nabla}_{\vec{r}} V(\vec{r}) \cdot (-2A\vec{p}) \rho(\vec{r}; t) = 0.$$

通解为:

$$\rho(\vec{r}) = C_1 e^{-2mAV(\vec{r})}.$$

其中  $C_1$  满足:

$$\int d^3\vec{r} \rho(\vec{r}) = N.$$

$N$  为气体分子总数。则 Boltzmann 方程近似静态解为:

$$f(\vec{r}, \vec{p}) = \frac{C_1}{(2\pi mkT_0)^{3/2}} \exp\left(-\frac{V(\vec{r})}{kT_0} - \frac{\vec{p}^2}{2mkT_0}\right).$$

2. 写下一个均匀且无外力作用的气体的 Boltzmann 方程并证明下列 Boltzmann H-定理:

$$\frac{dH}{dt} \leq 0$$

其中  $H \equiv \int d^3\vec{p} f(\vec{p}, t) \ln f(\vec{p}, t)$ .

对于均匀且无外力作用的气体, 有:

$$\vec{\nabla}_{\vec{r}} f = 0, \vec{F} = 0.$$

则 Boltzmann 方程可以写为:

$$\frac{\partial}{\partial t} f(\vec{p}; t) = \int d^3\vec{p}_2 d^3\vec{p}_1 d^3\vec{p}_2 |T_{fi}|^2 \delta^4(P_f - P_i) [f(\vec{p}_1'; t) f(\vec{p}_2'; t) - f(\vec{p}; t) f(\vec{p}_2; t)].$$

可以计算:

$$\begin{aligned} \frac{dH}{dt} &= \int d^3\vec{p} \frac{\partial}{\partial t} (f(\vec{p}, t) \ln f(\vec{p}, t)) \\ &= \int d^3\vec{p} \frac{\partial}{\partial t} f(\vec{p}; t) + \int d^3\vec{p} \ln f(\vec{p}; t) \frac{\partial}{\partial t} f(\vec{p}; t) \end{aligned}$$

考虑等式右侧第一项则有:

$$\int d^3\vec{p} \frac{\partial}{\partial t} f(\vec{p}; t) = \frac{d}{dt} \int d^3\vec{p} f(\vec{p}; t) = \frac{d\rho}{dt} = 0.$$

其中  $\rho$  为实空间中粒子数密度, 由于均匀且无外力作用, 则有:  $\frac{d\rho}{dt} = 0$ . 将 Boltzmann 方程带入上式可得:

$$\frac{dH}{dt} = \int d^3\vec{p} d^3\vec{p}_2 d^3\vec{p}_1 d^3\vec{p}_2 |T_{fi}|^2 \delta^4(P_f - P_i) [f(\vec{p}_1'; t) f(\vec{p}_2'; t) - f(\vec{p}; t) f(\vec{p}_2; t)] \ln f(\vec{p}; t).$$

交换  $\vec{p}$  和  $\vec{p}_2$  以及  $\vec{p}_1$  和  $\vec{p}_2'$  可得:

$$\frac{dH}{dt} = \int d^3\vec{p} d^3\vec{p}_2 d^3\vec{p}_1 d^3\vec{p}_2 |T_{fi}|^2 \delta^4(P_f - P_i) [f(\vec{p}_1'; t) f(\vec{p}_2'; t) - f(\vec{p}; t) f(\vec{p}_2; t)] \ln f(\vec{p}_2; t).$$

对上面两式利用时间反演不变性可得:

$$\frac{dH}{dt} = \int d^3\vec{p} d^3\vec{p}_2 d^3\vec{p}_1 d^3\vec{p}_2 |T_{fi}|^2 \delta^4(P_f - P_i) [f(\vec{p}; t) f(\vec{p}_2; t) - f(\vec{p}_1'; t) f(\vec{p}_2'; t)] \ln f(\vec{p}_1'; t).$$

$$\frac{dH}{dt} = \int d^3\vec{p} d^3\vec{p}_2 d^3\vec{p}_1 d^3\vec{p}_2 |T_{fi}|^2 \delta^4(P_f - P_i) [f(\vec{p}; t) f(\vec{p}_2; t) - f(\vec{p}_1'; t) f(\vec{p}_2'; t)] \ln f(\vec{p}_2'; t).$$

则有：

$$\begin{aligned} \frac{dH}{dt} = & \frac{1}{4} \int d^3\vec{p} d^3\vec{p}_2 d^3\vec{p}_1' d^3\vec{p}_2' |T_{fi}|^2 \delta^4(P_f - P_i) \\ & \times [f(\vec{p}_1'; t) f(\vec{p}_2'; t) - f(\vec{p}; t) f(\vec{p}_2; t)] \ln \frac{f(\vec{p}; t) f(\vec{p}_2; t)}{f(\vec{p}_1'; t) f(\vec{p}_2'; t)}. \end{aligned}$$

注意到如果  $f(\vec{p}_1'; t) f(\vec{p}_2'; t) - f(\vec{p}; t) f(\vec{p}_2; t) > 0$ ，则有：  $\ln \frac{f(\vec{p}; t) f(\vec{p}_2; t)}{f(\vec{p}_1'; t) f(\vec{p}_2'; t)} < 0$ ;

如果  $f(\vec{p}_1'; t) f(\vec{p}_2'; t) - f(\vec{p}; t) f(\vec{p}_2; t) < 0$ ，则有：  $\ln \frac{f(\vec{p}; t) f(\vec{p}_2; t)}{f(\vec{p}_1'; t) f(\vec{p}_2'; t)} > 0$ ;

如果  $f(\vec{p}_1'; t) f(\vec{p}_2'; t) - f(\vec{p}; t) f(\vec{p}_2; t) = 0$ ，则有：  $\ln \frac{f(\vec{p}; t) f(\vec{p}_2; t)}{f(\vec{p}_1'; t) f(\vec{p}_2'; t)} = 0$ ;

综上所述，有

$$\frac{dH}{dt} \leq 0.$$