

高等热力学与统计物理第四次作业

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1 计算下列系统内电子气体的 Fermi 能, Fermi 温度和 Fermi 速度

- (1) 室温下的金属钠: 密度为 0.97g/cm^3 , 每个原子贡献一个传导电子, 假设他们的能量和动量关系为 $E = \frac{p^2}{2m}$, 即忽略晶格场对电子运动的影响。

在 0K 下, 费米狄拉克分布为:

$$f = \begin{cases} 1, & E < \epsilon_F, \\ 0, & E > \epsilon_F. \end{cases}$$

则粒子数 N 为:

$$\begin{aligned} N &= 2 \left(\frac{L}{2\pi} \right)^3 \iiint dk_x dk_y dk_z \\ &= 2 \frac{V}{(2\pi)^3} \int 4\pi k^2 dk \\ &= \frac{V}{2\pi^2} \frac{(2m)^{3/2}}{\hbar^3} \int_0^{\epsilon_F} \sqrt{E} dE \\ &= \frac{V}{3\pi^2} \frac{(2m)^{3/2}}{\hbar^3} \epsilon_F^{3/2}. \end{aligned}$$

其中粒子数 $N = nN_A$, 体积 $V = \frac{nM}{\rho}$, n 为物质的量, N_A 为阿伏伽德罗常数, M 为钠的摩尔质量。则有:

$$\epsilon_F = \frac{1}{2m} \left(\frac{3\pi^2 \hbar^3 N_A \rho}{M} \right)^{2/3}.$$

代入数据可得费米能量为:

$$\epsilon_F = 5.048 \times 10^{-19} \text{J} = 3.15 \text{eV}.$$

费米温度为:

$$T_F = \frac{\epsilon_F}{k_B} = 36562 \text{K}.$$

费米速度为:

$$v_F = \sqrt{\frac{2\epsilon_F}{m}} = 1.05 \times 10^6 \text{m/s}.$$

- (2) 天狼星的伴星 (白矮星): 其质量约为太阳质量的 0.98 倍, 半径约为太阳半径的 0.0084 倍。假设星体全部由氦构成。

对于白矮星而言，密度为：

$$\rho = \frac{0.98m_{sun}}{\frac{4}{3}\pi(0.0084r_{sun})^3} = 2.315 \times 10^9 kg/m^3.$$

每一个氮原子贡献 2 个电子，考虑相对论条件粒子数满足：

$$N = \frac{2V}{\pi^2} \int_0^{k_F} k^2 dk = \frac{2V}{3\pi^2} k_F^3.$$

则有：

$$k_F = \sqrt[3]{\frac{6\pi^2 N_A \rho}{M}} = 2.743 \times 10^{12} m^{-1}.$$

费米能量为：

$$\epsilon_F = \sqrt{p^2 c^2 + m^2 c^4} - mc^2 = 3.739 \times 10^{-14} J = 2.33 \times 10^5 eV.$$

费米温度为：

$$T_F = \frac{\epsilon_F}{k_B} = 2.708 \times 10^9 K.$$

费米速度为：

$$v_F = \sqrt{\frac{p^2 c^2}{p^2 + m^2 c^2}} = 2.180 \times 10^8 m/s.$$

2 证明非相对论简并电子气体的热力学函数是

$$\begin{aligned} G &= N\mu = N\epsilon_F \left[1 - \frac{1}{12}\pi^2 \left(\frac{kT}{\epsilon_F} \right)^2 - \frac{1}{80}\pi^4 \left(\frac{kT}{\epsilon_F} \right)^4 \dots \right] \\ E &= \frac{3}{5}N\epsilon_F \left[1 + \frac{5}{12}\pi^2 \left(\frac{kT}{\epsilon_F} \right)^2 - \frac{1}{16}\pi^4 \left(\frac{kT}{\epsilon_F} \right)^4 \dots \right] \\ C_V &= \frac{1}{2}N\pi^2 \frac{k^2 T}{\epsilon_F} \left[1 - \frac{3}{10}\pi^2 \left(\frac{kT}{\epsilon_F} \right)^2 \dots \right] \\ S &= \frac{1}{2}N\pi^2 \frac{k^2 T}{\epsilon_F} \left[1 - \frac{1}{10}\pi^2 \left(\frac{kT}{\epsilon_F} \right)^2 \dots \right] \end{aligned}$$

其中 $N = \frac{V}{3\pi^2} \left(\frac{2m\epsilon_F}{\hbar} \right)^{3/2}$ 为电子总数，... 代表 $\frac{kT}{\epsilon_F}$ 的更高阶项。

粒子数 N 满足：

$$\begin{aligned} N &= \frac{V(2m)^{3/2}}{2\pi^2 \hbar^3} \int_0^\infty \frac{\sqrt{\epsilon}}{e^{\beta(\epsilon-\mu)} + 1} d\epsilon \\ &= \frac{V(2m)^{3/2}}{2\pi^2 \hbar^3} \int_0^\infty \frac{2\beta}{3} \epsilon^{3/2} \frac{e^{\beta(\epsilon-\mu)}}{(e^{\beta(\epsilon-\mu)} + 1)^2} d\epsilon. \end{aligned}$$

内能为:

$$\begin{aligned} E &= \frac{V(2m)^{3/2}}{2\pi^2\hbar^3} \int_0^\infty \frac{\epsilon^{3/2}}{e^{\beta(\epsilon-\mu)} + 1} d\epsilon \\ &= \frac{V(2m)^{3/2}}{2\pi^2\hbar^3} \int_0^\infty \frac{2\beta}{5} \epsilon^{5/2} \frac{e^{\beta(\epsilon-\mu)}}{(e^{\beta(\epsilon-\mu)} + 1)^2} d\epsilon. \end{aligned}$$

令 $\xi = \beta(\epsilon - \mu)$, 则积分化为:

$$\begin{aligned} N &= \frac{V(2m)^{3/2}}{3\pi^2\hbar^3} \int_{-\beta\mu}^\infty (kT\xi + \mu)^{3/2} \frac{e^\xi}{(e^\xi + 1)^2} d\xi. \\ E &= \frac{V(2m)^{3/2}}{5\pi^2\hbar^3} \int_{-\beta\mu}^\infty (kT\xi + \mu)^{5/2} \frac{e^\xi}{(e^\xi + 1)^2} d\xi. \end{aligned}$$

令

$$I = \int_{-\beta\mu}^\infty (kT\xi + \mu)^{3/2} \frac{e^\xi}{(e^\xi + 1)^2} d\xi \approx \int_{-\infty}^\infty (kT\xi + \mu)^{3/2} \frac{e^\xi}{(e^\xi + 1)^2} d\xi.$$

对 $(kT\xi + \mu)^{3/2}$ 做泰勒展开得:

$$(kT\xi + \mu)^{3/2} = \mu^{3/2} + \frac{3}{2}kT\mu^{1/2}\xi + \frac{3}{8}(kT)^2\mu^{-1/2}\xi^2 - \frac{1}{16}(kT)^3\mu^{-3/2}\xi^3 + \frac{3}{128}(kT)^4\mu^{-5/2}\xi^4 + \mathcal{O}(\xi^4).$$

对 $(kT\xi + \mu)^{5/2}$ 做泰勒展开可得:

$$(kT\xi + \mu)^{5/2} = \mu^{5/2} + \frac{5}{2}kT\mu^{3/2}\xi + \frac{15}{8}(kT)^2\mu^{1/2}\xi^2 + \frac{5}{16}(kT)^3\mu^{-1/2}\xi^3 - \frac{5}{128}(kT)^4\mu^{-3/2}\xi^4 + \mathcal{O}(\xi^4).$$

由于 $\frac{e^\xi}{(e^\xi + 1)^2}$ 是关于 ξ 的偶函数, 则泰勒展开只有偶次幂的积分不为零. 令 $c_n = \int_0^\infty \frac{\xi^{2n} e^\xi}{(e^\xi + 1)^2} d\xi$. 当 $n = 0$ 时, $c_0 = \frac{1}{2}$. 当 $n > 0$ 时, 有:

$$\frac{e^\xi}{(e^\xi + 1)^2} = \frac{e^{-\xi}}{(e^{-\xi} + 1)^2} = \sum_{l=1}^{\infty} (-1)^{l-1} l e^{-l\xi}.$$

则有:

$$c_n = \sum_{l=1}^{\infty} (-1)^{l-1} l \int_0^\infty \xi^{2n} e^{-l\xi} d\xi = (2n)! \sum_{l=1}^{\infty} \frac{(-1)^{l-1}}{l^{2n}}.$$

则有:

$$\begin{aligned} c_1 &= 2 \left(\sum_{l=odd} \frac{1}{l^2} - \sum_{l=even} \frac{1}{l^2} \right) = 2 \left(\sum_{l=1}^{\infty} \frac{1}{l^2} - \frac{2}{2^2} \sum_{l=1}^{\infty} \frac{1}{l^2} \right) = \frac{\pi^2}{6}. \\ c_2 &= 24 \left(\sum_{l=odd} \frac{1}{l^4} - \sum_{l=even} \frac{1}{l^4} \right) = 24 \left(\sum_{l=1}^{\infty} \frac{1}{l^4} - \frac{2}{16} \sum_{l=1}^{\infty} \frac{1}{l^4} \right) = \frac{7\pi^4}{30}. \end{aligned}$$

则有:

$$\begin{aligned} N &= \frac{V(2m)^{3/2}}{3\pi^2\hbar^3} \mu^{3/2} \left[1 + \frac{\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 + \frac{7\pi^4}{640} \left(\frac{kT}{\mu} \right)^4 + \dots \right]. \\ E &= \frac{V(2m)^{3/2}}{5\pi^2\hbar^3} \mu^{5/2} \left[1 + \frac{5\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 - \frac{7\pi^4}{384} \left(\frac{kT}{\mu} \right)^4 \dots \right]. \end{aligned}$$

则有:

$$E = \frac{3}{5}N\mu \left[1 + \frac{\pi^2}{2} \left(\frac{kT}{\mu} \right)^2 - \frac{11\pi^4}{120} \left(\frac{kT}{\mu} \right)^4 \right].$$

对比

$$N = \frac{V(2m)^{3/2}}{3\pi^2\hbar^3} \epsilon_F^{3/2}.$$

有:

$$\epsilon_F = \mu \left[1 + \frac{\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 + \frac{7\pi^4}{640} \left(\frac{kT}{\mu} \right)^4 + \dots \right]^{2/3}$$

泰勒展开可得:

$$\epsilon_F = \mu \left[1 + \frac{\pi^2}{12} \left(\frac{kT}{\mu} \right)^2 + \frac{\pi^4}{180} \left(\frac{kT}{\mu} \right)^4 \right]$$

利用级数反演可得:

$$\mu = \epsilon_F \left[1 - \frac{\pi^2}{12} \left(\frac{kT}{\epsilon_F} \right)^2 - \frac{\pi^4}{80} \left(\frac{kT}{\epsilon_F} \right)^4 \dots \right]$$

则 Gibbs 势为:

$$G = N\mu = N\epsilon_F \left[1 - \frac{\pi^2}{12} \left(\frac{kT}{\epsilon_F} \right)^2 - \frac{\pi^4}{80} \left(\frac{kT}{\epsilon_F} \right)^4 \dots \right].$$

内能为:

$$E = \frac{3}{5}N\epsilon_F \left[1 + \frac{5\pi^2}{12} \left(\frac{kT}{\epsilon_F} \right)^2 - \frac{\pi^4}{16} \left(\frac{kT}{\epsilon_F} \right)^4 \dots \right]$$

则等容热容为:

$$C_V = \left(\frac{\partial E}{\partial T} \right)_V = \frac{1}{2}N\pi^2 \frac{k^2 T}{\epsilon_F} \left[1 - \frac{3}{10}\pi^2 \left(\frac{kT}{\epsilon_F} \right)^2 + \dots \right]$$

在非相对论条件下有:

$$PV = \frac{2}{3}E.$$

则气体自由能为:

$$F = G - PV = G - \frac{2}{3}E = N\epsilon_F \left[\frac{3}{5} - \frac{\pi^2}{4} \left(\frac{kT}{\epsilon_F} \right)^2 + \frac{\pi^4}{80} \left(\frac{kT}{\epsilon_F} \right)^4 + \dots \right]$$

则熵为:

$$S = - \left(\frac{\partial F}{\partial T} \right)_V = N\epsilon_F \left(\frac{\pi^2 k^2 T}{2 \epsilon_F^2} - \frac{\pi^4 k^4 T^3}{20 \epsilon_F^4} + \dots \right).$$