

Problem Set 1

1. (a) Assume: $\hat{e}_i, \hat{e}_j, \hat{e}_k$ are unit vectors in the direction of i, j and k , in Cartesian coordinates.

$\hat{e}_l, \hat{e}_m, \hat{e}_n$ are unit vectors in the direction of l, m and n , in Cartesian coordinates.

Then we have: $\hat{e}_k \epsilon_{ijk} = \hat{e}_i \times \hat{e}_j$, $\hat{e}_k \cdot \hat{e}_k \epsilon_{ijk} = \epsilon_{ijk} = \hat{e}_k \cdot (\hat{e}_i \times \hat{e}_j)$

$$\hat{e}_n \epsilon_{lmn} = \hat{e}_l \times \hat{e}_m, \quad \hat{e}_n \cdot \hat{e}_n \epsilon_{lmn} = \epsilon_{lmn} = \hat{e}_n \cdot (\hat{e}_l \times \hat{e}_m)$$

$$\epsilon_{ijk} \epsilon_{lmn} = [\hat{e}_k \cdot (\hat{e}_i \times \hat{e}_j)] [\hat{e}_n \cdot (\hat{e}_l \times \hat{e}_m)]$$

According to $\det A \det B = \det(AB)$

A and B are 3×3 matrix

$$= \begin{vmatrix} \hat{e}_k \cdot \hat{e}_n & \hat{e}_k \cdot \hat{e}_l & \hat{e}_k \cdot \hat{e}_m \\ \hat{e}_i \cdot \hat{e}_n & \hat{e}_i \cdot \hat{e}_l & \hat{e}_i \cdot \hat{e}_m \\ \hat{e}_j \cdot \hat{e}_n & \hat{e}_j \cdot \hat{e}_l & \hat{e}_j \cdot \hat{e}_m \end{vmatrix}$$

Assume $l \rightarrow i$, and m, n are not the same as j, k

$$\text{then } \epsilon_{ijk} \epsilon_{imn} = \begin{vmatrix} \delta_{kn} & 0 & \delta_{km} \\ 0 & 1 & 0 \\ \delta_{jn} & 0 & \delta_{jm} \end{vmatrix} = \delta_{jm} \delta_{kn} - \delta_{jn} \delta_{km}$$

$$\text{So } \epsilon_{ijk} \epsilon_{imn} = \delta_{jm} \delta_{kn} - \delta_{jn} \delta_{km}$$

$$(b) (i) [\vec{A} \times (\vec{B} \times \vec{C})]_i = \sum \epsilon_{ijk} A_j (\vec{B} \times \vec{C})_k = \sum \epsilon_{ijk} A_j \epsilon_{kmn} B_m C_n = \sum \epsilon_{kij} \epsilon_{kmn} A_j B_m C_n$$

$$= \sum (\delta_{im} \delta_{jn} - \delta_{in} \delta_{jm}) A_j B_m C_n$$

$$[\vec{B} (\vec{A} \cdot \vec{C}) - \vec{C} (\vec{A} \cdot \vec{B})]_i = [\sum \delta_{jn} A_j C_n \vec{B} - (\sum \delta_{jm} A_j B_m) \vec{C}]_i = \sum \delta_{im} \delta_{jn} A_j C_n B_m - \sum \delta_{jm} \delta_{in} A_j B_m C_n$$

$$= \sum (\delta_{im} \delta_{jn} - \delta_{in} \delta_{jm}) A_j B_m C_n$$

$$\text{So } \vec{A} \times (\vec{B} \times \vec{C}) = \vec{B} (\vec{A} \cdot \vec{C}) - \vec{C} (\vec{A} \cdot \vec{B})$$

$$(ii) \vec{\nabla} \cdot (\vec{A} \times \vec{B}) = \partial_i \epsilon_{ijk} A_j B_k = B_k \epsilon_{ijk} \partial_i A_j + A_j \epsilon_{ijk} \partial_i B_k = B_k \epsilon_{kij} \partial_i A_j - A_j \epsilon_{jik} \partial_i B_k$$

$$= \vec{B} \cdot (\vec{\nabla} \times \vec{A}) - \vec{A} \cdot (\vec{\nabla} \times \vec{B})$$

$$\text{So } \nabla \cdot (\vec{A} \times \vec{B}) = \vec{B} \cdot (\nabla \times \vec{A}) - \vec{A} \cdot (\nabla \times \vec{B})$$

$$(iii) [\vec{A} \times (\nabla \times \vec{B})]_i = \epsilon_{ijk} A_j (\nabla \times \vec{B})_k = \epsilon_{kij} A_j \epsilon_{kmn} \partial_m B_n = (\delta_{im} \delta_{jn} - \delta_{in} \delta_{jm}) A_j \partial_m B_n = A_j \partial_i B_j - A_j \partial_j B_i$$

$$[\nabla (\vec{A} \cdot \vec{B}) - \vec{B} \times (\nabla \times \vec{A}) - (\vec{A} \cdot \nabla) \vec{B} - (\nabla \cdot \vec{B}) \vec{A}]_i = \partial_i A_i B_i - \epsilon_{ijk} \epsilon_{kmn} B_j \partial_m A_n - B_i \partial_j A_j - A_i \partial_j B_j$$

$$= A_j \partial_i B_j - A_j \partial_j B_i$$

$$\text{So } \vec{A} \times (\nabla \times \vec{B}) = \nabla (\vec{A} \cdot \vec{B}) - \vec{B} \times (\nabla \times \vec{A}) - (\vec{A} \cdot \nabla) \vec{B} - (\nabla \cdot \vec{B}) \vec{A}$$

$$(iv) [\nabla \times (\vec{A} \times \vec{B})]_i = \epsilon_{ijk} \epsilon_{lmn} \partial_j (A_l B_m) = (\delta_{jm} \delta_{in} - \delta_{in} \delta_{jm}) (B_n \partial_j A_m + A_m \partial_j B_n) \cdot (-1)$$

$$= B_j \partial_j A_i - B_i \partial_j A_j + A_i \partial_j B_j - A_j \partial_j B_i$$

$$= [(\vec{B} \cdot \nabla) \vec{A} - (\vec{A} \cdot \nabla) \vec{B} + \vec{A} (\nabla \cdot \vec{B}) - \vec{B} (\nabla \cdot \vec{A})]_i$$

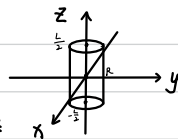
$$\begin{aligned}
 2. (a). \quad \delta(x^2 - a^2) &= \frac{1}{2|a|} [\delta(x+a) + \delta(x-a)] \\
 \text{So } \int_{-\infty}^{+\infty} \phi(x) \delta(x^2 - a^2) dx &= \int_{-\infty}^{+\infty} \frac{1}{2|a|} \phi(x) [\delta(x+a) + \delta(x-a)] dx \\
 &= \frac{1}{2|a|} [\phi(-a) + \phi(a)]
 \end{aligned}$$

(b). The average surface charge density $\sigma_0 = \frac{Q}{2\pi R^2 + 2\pi RL}$

$$\sigma = \sigma_0 \cdot 2\pi r z \cdot \delta(r-R) [\delta(z-\frac{L}{2}) - \delta(z+\frac{L}{2})] + \sigma_0 \cdot \pi r^2 \delta(r-R) [\delta(z-\frac{L}{2}) + \delta(z+\frac{L}{2})]$$

$$Q' = \iint \sigma \, d\mathbf{r} \, dz = \sigma_0 \iint \{ 2\pi r z \delta(r-R) [\delta(z-\frac{L}{2}) - \delta(z+\frac{L}{2})] + \pi r^2 \delta(r-R) [\delta(z-\frac{L}{2}) + \delta(z+\frac{L}{2})] \} \, dr \, dz$$

$$= \sigma_0 [2\pi R [\frac{L}{2} - (-\frac{L}{2})] + \pi R^2 \cdot (1+1)] = \sigma_0 \cdot (2\pi RL + 2\pi R^2) = Q$$



Which shows this yields the proper total charge when integrated over all space.

(c). $\rho(\vec{r}) = A \delta(\vec{r} \cdot \vec{r}_0 - |\vec{r}_0|^2)$ The charge distribution given by the form is there is a point charge at position \vec{r}_0 .

The unit of A is a line charge density.

Because the unit of $\rho(\vec{r})$ is a volume charge density and the unit of $\delta(\vec{r} \cdot \vec{r}_0 - |\vec{r}_0|^2)$ is meter^{-2} so the unit of A is a line charge density.

3. (a) Since the rotational symmetry of the disk, the direction of the electric field at $(0,0,z)$ is along z axis

According to Coulomb's Law

$$d\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{\sigma \cdot 2\pi r dr}{r^2 + z^2} \cdot \cos\theta = \frac{\sigma z}{2\epsilon_0} \frac{r}{(r^2 + z^2)^{3/2}} dr$$

$$\text{So } \vec{E} = \int_0^R d\vec{E} = \frac{\sigma z}{2\epsilon_0} \int_0^R \frac{r}{(r^2 + z^2)^{3/2}} dr = \frac{\sigma}{2\epsilon_0} \left(1 - \frac{z}{\sqrt{R^2 + z^2}}\right)$$

So the electric field at $(0,0,z)$ is $\vec{E} = \frac{\sigma}{2\epsilon_0} \left(1 - \frac{z}{\sqrt{R^2 + z^2}}\right)$

- (b) If $z \gg R$, then $\frac{R^2}{z^2} \rightarrow 0$ $1 - \frac{z}{\sqrt{R^2 + z^2}} = 1 - \left(1 + \frac{R^2}{z^2}\right)^{-1/2} \approx \frac{R^2}{2z^2}$

$$\text{So } \vec{E} = \frac{\sigma R^2}{4\epsilon_0} \cdot \frac{1}{z^2} \hat{e}_z$$

The charged circular disk seems like a point charge.

- (c) If $z \ll R$ then $\frac{z^2}{R^2} \rightarrow 0$ $1 - \frac{z}{\sqrt{R^2 + z^2}} \approx 1$

$$\text{So } \vec{E} = \frac{\sigma}{2\epsilon_0} \hat{e}_z$$

The charged circular disk seems like an infinite disk.

4. (a) Since it is a uniformly charged, the surface charge $\sigma = \frac{Q}{2\omega L}$

According to symmetry, the electric field at any position $x > 0$ on the positive x -axis.

$$\begin{aligned} E &= \iint_S \frac{1}{4\pi\epsilon_0} \frac{\sigma}{y^2 + (x-x_0)^2} \cdot \cos\theta \, dx \, dy = \iint_S \frac{\sigma}{4\pi\epsilon_0} \frac{x_0 - x}{[y^2 + (x-x_0)^2]^{3/2}} \, dx \, dy = \frac{\sigma}{4\pi\epsilon_0} \int_{-\frac{L}{2}}^{\frac{L}{2}} dy \int_{-\omega}^{\omega} \frac{x_0 - x}{[y^2 + (x-x_0)^2]^{3/2}} \, dx \\ &= \frac{\sigma}{4\pi\epsilon_0} \int_{-\frac{L}{2}}^{\frac{L}{2}} \left[\frac{1}{\sqrt{y^2 + (\omega-x_0)^2}} - \frac{1}{\sqrt{y^2 + (\omega+x_0)^2}} \right] dy = \frac{\sigma}{4\pi\epsilon_0} \left[2 \operatorname{arcsinh} \frac{L}{2|\omega-x_0|} - 2 \operatorname{arcsinh} \frac{L}{2(\omega+x_0)} \right] \\ &= \frac{\sigma}{2\pi\epsilon_0} \left[\operatorname{arcsinh} \frac{L}{2|\omega-x_0|} - \operatorname{arcsinh} \frac{L}{2(\omega+x_0)} \right] \end{aligned}$$

So the electric field at any position $x > 0$ on the positive x -axis,

for example, $(x_0, 0)$ $\vec{E}(x_0, 0) = \frac{1}{4\pi\epsilon_0} \frac{Q}{\omega L} \left[\operatorname{arcsinh} \frac{L}{2|\omega-x_0|} - \operatorname{arcsinh} \frac{L}{2(\omega+x_0)} \right] \hat{i}$

- (b). (i) At the origin, we have $x_0 = 0$.

Since the symmetry of the rectangular charged sheet, we can easily get

$$\vec{E}(0,0) = 0$$

My solution also says that $\vec{E}(0,0) = 0$

So my result behaves exactly as expand at the origin.

- (ii) If $x \gg L, \omega$, we have $\frac{\omega}{x} \rightarrow 0$ $\frac{L}{2(\omega+x)} \rightarrow 0$ $\frac{L}{2(x-\omega)} \rightarrow 0$

$$\operatorname{arcsinh} \frac{L}{2(\omega+x)} \approx \frac{L}{2(\omega+x)} = \frac{L}{2x} \cdot \frac{1}{1 + \frac{\omega}{x}} \approx \frac{L}{2x} \left(1 - \frac{\omega}{x}\right)$$

$$\operatorname{arcsinh} \frac{L}{2(x-w)} \approx \frac{L}{2(x-w)} = \frac{L}{2x} \cdot \frac{1}{1 - \frac{w}{x}} \approx \frac{L}{2x} \cdot (1 + \frac{w}{x})$$

$$\text{So } \operatorname{arcsinh} \frac{L}{2(x-w)} - \operatorname{arcsinh} \frac{L}{2(x+w)} \approx \frac{L}{2x} \cdot [(1 + \frac{w}{x}) - (1 - \frac{w}{x})] = \frac{Lw}{x^2}$$

So $\vec{E}(x,0) = \frac{1}{4\pi\epsilon_0} \frac{Q}{x^2}$ and this is the leading behavior, which shows that the rectangular charged sheet seems like a point charge.

(iii). If $x \gg L \gg w$, $\frac{L}{x+w} \rightarrow 0$, $\frac{L}{x-w} \rightarrow 0$, $\frac{w}{x} \rightarrow 0$

$$\operatorname{arcsinh} \frac{L}{2(x+w)} \approx \frac{L}{2(x+w)} - \frac{1}{6} \frac{L^3}{8(x+w)^3} = \frac{L}{2x} \cdot \frac{1}{1 + \frac{w}{x}} - \frac{L^3}{48x^3} \cdot \frac{1}{(1 + \frac{w}{x})^3}$$

$$\approx \frac{L}{2x} (1 - \frac{w}{x} + \frac{w^2}{x^2} - \frac{w^3}{x^3}) - \frac{L^3}{48x^3} (1 - 3\frac{w}{x} + 6\frac{w^2}{x^2} - 10\frac{w^3}{x^3})$$

$$\approx \frac{L}{2x} - \frac{Lw}{2x^2} - (\frac{L^3}{48} - \frac{Lw^2}{2}) \frac{1}{x^3} + (\frac{L^3w}{16} - \frac{Lw^3}{2}) \frac{1}{x^4}$$

$$\operatorname{arcsinh} \frac{L}{2(x-w)} \approx \frac{L}{2(x-w)} - \frac{1}{6} \frac{L^3}{8(x-w)^3} = \frac{L}{2x} \cdot \frac{1}{1 - \frac{w}{x}} - \frac{L^3}{48x^3} \cdot \frac{1}{(1 - \frac{w}{x})^3}$$

$$\approx \frac{L}{2x} (1 + \frac{w}{x} + \frac{w^2}{x^2} + \frac{w^3}{x^3}) - \frac{L^3}{48x^3} (1 + 3\frac{w}{x} + 6\frac{w^2}{x^2} + 10\frac{w^3}{x^3})$$

$$\approx \frac{L}{2x} + \frac{Lw}{2x^2} - (\frac{L^3}{48} - \frac{Lw^2}{2}) \frac{1}{x^3} - (\frac{L^3w}{16} - \frac{Lw^3}{2}) \frac{1}{x^4}$$

$$\text{So } \operatorname{arcsinh} \frac{L}{2(x-w)} - \operatorname{arcsinh} \frac{L}{2(x+w)} \approx \frac{Lw}{x^2} - (\frac{L^3w}{8} - Lw^3) \frac{1}{x^4} \approx \frac{Lw}{x^2} - \frac{L^3w}{8x^4}$$

$$\text{So } \vec{E}(x,0) = \frac{Q}{4\pi\epsilon_0} \cdot (\frac{1}{x^2} - \frac{L^2}{8x^4})$$

The first correction to the leading behavior is $-\frac{Q}{32\pi\epsilon_0} \frac{L^2}{x^4}$

If there is a line of charge of length L and the same total charge as the sheet, since the symmetry, the electrical field on x -axis is along x -axis.

$$dE = \frac{1}{4\pi\epsilon_0} \frac{\lambda dy}{x^2 + y^2} \cos\theta = \frac{1}{4\pi\epsilon_0} \frac{\lambda x}{(x^2 + y^2)^{3/2}} dy \quad \lambda = \frac{Q}{L}$$

$$E = \int_{-\frac{L}{2}}^{\frac{L}{2}} dE = \frac{\lambda}{4\pi\epsilon_0} \int_{-\frac{L}{2}}^{\frac{L}{2}} \frac{x}{(x^2 + y^2)^{3/2}} dy = \frac{\lambda}{4\pi\epsilon_0} \int_{-\frac{L}{2}}^{\frac{L}{2}} \frac{1}{x} \cdot \frac{1}{(1 + \frac{y^2}{x^2})^{3/2}} d\frac{y}{x} = \frac{\lambda}{4\pi\epsilon_0} \cdot \frac{1}{x} \cdot \frac{2L}{\sqrt{4x^2 + L^2}}$$

$$\text{If } x \gg L \quad \frac{L}{x} \rightarrow 0 \quad E = \frac{Q}{4\pi\epsilon_0} \cdot \frac{1}{x^2} \cdot \frac{1}{\sqrt{1 + (\frac{L}{2x})^2}} \approx \frac{1}{4\pi\epsilon_0} \frac{Q}{x^2} \cdot [1 - \frac{1}{2} (\frac{L}{2x})^2] = \frac{1}{4\pi\epsilon_0} \frac{Q}{x^2} (1 - \frac{L^2}{8x^2})$$

The first correction is also $-\frac{1}{32\pi\epsilon_0} \frac{QL^2}{x^4}$ and this is the same as the first correction of my solution.