

介观物理第三次作业

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Problem I.1 Verify the Kramers-Kronig relation using the explicit solution given in Eq. (12).

$$\bar{\chi}(\omega) = \frac{1}{m} \frac{1}{\omega^2 - \omega_0^2 + i\gamma\omega},$$
$$\bar{\chi}(t - t') = -\frac{1}{m} \int \frac{d\omega}{2\pi} \frac{e^{-i\omega(t-t')}}{\omega^2 - \omega_0^2 + i\gamma\omega}.$$

Proof. Kramers - Kronig 关系为:

$$\chi'(\omega) = \int_{-\infty}^{\infty} \frac{d\omega'}{\pi} \chi''(\omega') P \frac{1}{\omega' - \omega},$$
$$\chi''(\omega) = - \int_{-\infty}^{\infty} \frac{d\omega'}{\pi} \chi'(\omega') P \frac{1}{\omega' - \omega}.$$

其中 $\chi'(\omega)$ 和 $\chi''(\omega)$ 分别为 $\chi(\omega)$ 的实部和虚部。

$$\chi'(\omega) = \frac{1}{m} \frac{\omega_0^2 - \omega^2}{(\omega^2 - \omega_0^2)^2 + \gamma^2 \omega^2};$$
$$\chi''(\omega) = \frac{1}{m} \frac{\gamma \omega}{(\omega^2 - \omega_0^2)^2 + \gamma^2 \omega^2}$$

则可以计算:

$$\begin{aligned} \mathbf{I}_1 &= \int_{-\infty}^{\infty} \frac{d\omega'}{\pi} \frac{1}{m} \frac{\gamma \omega'}{(\omega'^2 - \omega_0^2)^2 + \gamma^2 \omega'^2} P \frac{1}{\omega' - \omega} \\ &= \frac{1}{\pi m} \int_{-\infty}^{\infty} \frac{\gamma \omega'}{(\omega'^2 - \omega_0^2)^2 + \gamma^2 \omega'^2} \left[-\frac{1}{\omega - \omega' + i0^+} - i\pi \delta(\omega - \omega') \right] d\omega' \\ &= -\frac{1}{\pi m} \int_{-\infty}^{\infty} \frac{\gamma \omega'}{(\omega'^2 - \omega_0^2)^2 + \gamma^2 \omega'^2} \frac{1}{\omega - \omega' + i0^+} d\omega' - \frac{i}{m} \frac{\gamma \omega}{(\omega^2 - \omega_0^2)^2 + \gamma^2 \omega^2}. \end{aligned}$$

被积函数记作:

$$\begin{aligned} f(\omega') &= \frac{\gamma \omega'}{(\omega'^2 - \omega_0^2)^2 + \gamma^2 \omega'^2} \frac{1}{\omega - \omega' + i0^+} \\ &= \frac{\gamma \omega'}{(\omega + i0^+ - \omega')(\omega' - \omega_1)(\omega' - \omega_2)(\omega' - \omega_1^*)(\omega' - \omega_2^*)}. \end{aligned}$$

其中

$$\omega_1 = \frac{i\gamma + \sqrt{\Delta}}{2}; \quad \omega_2 = \frac{i\gamma - \sqrt{\Delta}}{2}; \quad \Delta = 4\omega_0^2 - \gamma^2.$$

即被积函数在上半复平面有三个奇点, 分别是 $\omega + i0^+$; ω_1 ; ω_2 . 分别计算其留数可得:

$$\text{Res}[f(\omega'), \omega + i0^+] = -\frac{\gamma \omega}{(\omega^2 - \omega_0^2)^2 + \gamma^2 \omega^2};$$
$$\text{Res}[f(\omega'), \omega_1] = \frac{\gamma \omega_1}{(\omega - \omega_1)(\omega_1 - \omega_2)(\omega_1 - \omega_1^*)(\omega_1 - \omega_2^*)}$$

$$\begin{aligned}
&= \frac{\omega_1}{\omega - \omega_1} \cdot \frac{\gamma}{\sqrt{\Delta}(i\gamma)(i\gamma + \sqrt{\Delta})} \\
&= \frac{1}{2i\sqrt{\Delta}(\omega - \omega_1)}; \\
\mathbf{Res}[f(\omega'), \omega_2] &= \frac{\gamma\omega_2}{(\omega - \omega_2)(\omega_2 - \omega_1)(\omega_2 - \omega_1^*)(\omega_2 - \omega_2^*)} \\
&= \frac{\omega_2}{\omega - \omega_2} \cdot \frac{\gamma}{\sqrt{\Delta}(i\gamma)(-i\gamma + \sqrt{\Delta})} \\
&= \frac{1}{2i\sqrt{\Delta}(\omega - \omega_2)}.
\end{aligned}$$

利用留数定理可得，积分式为：

$$\begin{aligned}
\int_{-\infty}^{\infty} f(\omega') d\omega' &= 2\pi i (\mathbf{Res}[f(\omega'), \omega + i0^+] + \mathbf{Res}[f(\omega'), \omega_1] + \mathbf{Res}[f(\omega'), \omega_2]) \\
&= \frac{2\pi i}{(\omega^2 - \omega_0^2)^2 + \gamma^2\omega^2} \left(-\gamma\omega + \frac{\omega^2 - \omega_0^2 + i\gamma\omega}{2i} \right)
\end{aligned}$$

则有：

$$\begin{aligned}
\mathbf{I}_1 &= -\frac{1}{\pi m} \int_{-\infty}^{\infty} f(\omega') d\omega' - \frac{i}{m} \frac{\gamma\omega}{(\omega^2 - \omega_0^2)^2 + \gamma^2\omega^2} \\
&= \frac{1}{m} \frac{\omega_0^2 - \omega^2}{(\omega^2 - \omega_0^2)^2 + \gamma^2\omega^2} \\
&= \chi'(\omega).
\end{aligned}$$

同样的，考虑

$$\begin{aligned}
\mathbf{I}_2 &= -\int_{-\infty}^{\infty} \frac{d\omega'}{\pi} \frac{1}{m} \frac{\omega_0^2 - \omega'^2}{(\omega'^2 - \omega_0^2)^2 + \gamma^2\omega'^2} P \frac{1}{\omega' - \omega} \\
&= -\frac{1}{\pi m} \int_{-\infty}^{\infty} \frac{\omega_0^2 - \omega'^2}{(\omega'^2 - \omega_0^2)^2 + \gamma^2\omega'^2} \left[-\frac{1}{\omega - \omega' + i0^+} - i\pi\delta(\omega - \omega') \right] d\omega' \\
&= \frac{1}{\pi m} \int_{-\infty}^{\infty} \frac{\omega_0^2 - \omega'^2}{(\omega'^2 - \omega_0^2)^2 + \gamma^2\omega'^2} \frac{1}{\omega - \omega' + i0^+} d\omega' + \frac{i}{m} \frac{\omega_0^2 - \omega^2}{(\omega^2 - \omega_0^2)^2 + \gamma^2\omega^2}.
\end{aligned}$$

被积函数记作：

$$\begin{aligned}
g(\omega') &= \frac{\omega_0^2 - \omega'^2}{(\omega'^2 - \omega_0^2)^2 + \gamma^2\omega'^2} \frac{1}{\omega - \omega' + i0^+} \\
&= \frac{\omega_0^2 - \omega'^2}{(\omega + i0^+ - \omega')(\omega' - \omega_1)(\omega' - \omega_2)(\omega' - \omega_1^*)(\omega' - \omega_2^*)}.
\end{aligned}$$

即被积函数在上半复平面有三个奇点，分别是 $\omega + i0^+$ ； ω_1 ； ω_2 。分别计算其留数可得：

$$\begin{aligned}
\mathbf{Res}[g(\omega'), \omega + i0^+] &= \frac{\omega^2 - \omega_0^2}{(\omega^2 - \omega_0^2)^2 + \gamma^2\omega^2}; \\
\mathbf{Res}[g(\omega'), \omega_1] &= \frac{\omega_0^2 - \omega_1^2}{(\omega - \omega_1)(\omega_1 - \omega_2)(\omega_1 - \omega_1^*)(\omega_1 - \omega_2^*)} \\
&= \frac{\omega_0^2 - \omega_1^2}{\omega - \omega_1} \cdot \frac{1}{\sqrt{\Delta}(i\gamma)(i\gamma + \sqrt{\Delta})}
\end{aligned}$$

$$\begin{aligned}
&= \frac{\omega_2(\omega_0^2 - \omega_1^2)}{\omega - \omega_1} \frac{-1}{2i\gamma\omega_0^2\sqrt{\Delta}}; \\
\mathbf{Res}[g(\omega'), \omega_2] &= \frac{\omega_0^2 - \omega_2^2}{(\omega - \omega_2)(\omega_2 - \omega_1)(\omega_2 - \omega_1^*)(\omega_2 - \omega_2^*)} \\
&= \frac{\omega_0^2 - \omega_2^2}{\omega - \omega_2} \cdot \frac{1}{\sqrt{\Delta}(i\gamma)(-i\gamma + \sqrt{\Delta})} \\
&= \frac{\omega_1(\omega_0^2 - \omega_1^2)}{\omega - \omega_2} \frac{1}{2i\gamma\omega_0^2\sqrt{\Delta}}.
\end{aligned}$$

利用留数定理可得，积分式为：

$$\begin{aligned}
\int_{-\infty}^{\infty} g(\omega') d\omega' &= 2\pi i (\mathbf{Res}[g(\omega'), \omega + i0^+] + \mathbf{Res}[g(\omega'), \omega_1] + \mathbf{Res}[g(\omega'), \omega_2]) \\
&= \frac{2\pi i}{(\omega^2 - \omega_0^2)^2 + \gamma^2\omega^2} \left(-\frac{i\gamma\omega}{2} + \frac{\omega^2 - \omega_0^2}{2} \right)
\end{aligned}$$

则有：

$$\begin{aligned}
\mathbf{I}_2 &= \frac{1}{\pi m} \int_{-\infty}^{\infty} g(\omega') d\omega' + \frac{i}{m} \frac{\omega_0^2 - \omega^2}{(\omega^2 - \omega_0^2)^2 + \gamma^2\omega^2} \\
&= \frac{1}{m} \frac{\gamma\omega}{(\omega^2 - \omega_0^2)^2 + \gamma^2\omega^2} \\
&= \chi''(\omega).
\end{aligned}$$

□

Problem I.2 * Read Kubo's paper entitled "The fluctuation-dissipation theorem".

Problem I.3 Derive the Kubo formula in Eq. (28).

$$\chi(\mathbf{r}, \mathbf{r}'; \omega) = \frac{i}{\hbar} \int_{-\infty}^t dt' e^{i\omega(t-t')} \langle 0 | [\hat{X}(\mathbf{r}, t), \hat{X}(\mathbf{r}', t')] | 0 \rangle.$$

Proof. 响应函数为：

$$\chi(\mathbf{r}, t; \mathbf{r}', t') = \frac{i}{\hbar} \theta(t - t') \langle 0 | [\hat{X}(\mathbf{r}, t), \hat{X}(\mathbf{r}', t')] | 0 \rangle.$$

对于不显含时的哈密顿量，响应函数只与 $t - t'$ 相关，令 $t'' = t - t'$ ，做傅里叶变换可得：

$$\begin{aligned}
\chi(\mathbf{r}, \mathbf{r}'; \omega) &= \int_{-\infty}^{\infty} e^{i\omega t''} \chi(\mathbf{r}, \mathbf{r}'; t'') dt'' \\
&= \frac{i}{\hbar} \int_{-\infty}^{\infty} e^{i\omega(t-t')}\theta(t'') \langle 0 | [\hat{X}(\mathbf{r}, t), \hat{X}(\mathbf{r}', t')] | 0 \rangle dt'' \\
&= \frac{i}{\hbar} \int_0^{\infty} e^{i\omega t''} \langle 0 | [\hat{X}(\mathbf{r}, t), \hat{X}(\mathbf{r}', t')] | 0 \rangle dt'' \\
&= \frac{i}{\hbar} \int_{-\infty}^t e^{i\omega(t-t')} \langle 0 | [\hat{X}(\mathbf{r}, t), \hat{X}(\mathbf{r}', t')] | 0 \rangle dt'.
\end{aligned}$$

即得到 Kubo 公式。

□