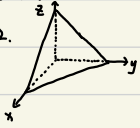


Home work 1

1-2.



证明: 设光线方向矢量为 $\vec{r} = a\vec{i} + b\vec{j} + c\vec{k}$ 折第一次折射后 $\vec{r}_1 = a_1\vec{i} + b_1\vec{j} + c_1\vec{k}$

由折射定律 $n\sin\theta_1 = \sin\theta_0$ 介质折射率为 n

设该平面为 $Ax + By + Cz + D = 0$

$$\therefore \cos\theta_0 = \frac{|aA + bB + cC|}{\sqrt{a^2 + b^2 + c^2} \cdot \sqrt{A^2 + B^2 + C^2}} \quad \sin\theta_0 = \sqrt{1 - \cos^2\theta_0}$$

$$\cos\theta_1 = \frac{|a_1A + b_1B + c_1C|}{\sqrt{a_1^2 + b_1^2 + c_1^2} \cdot \sqrt{A^2 + B^2 + C^2}} \quad \sin\theta_1 = \sqrt{1 - \cos^2\theta_1}$$

经过 xy, yz, zx 平面全反射后得到 $\vec{r}_2 = a_2\vec{i} + b_2\vec{j} + c_2\vec{k}$

由全反射满足反射角等于入射角 $\therefore a_2 = -a_1, b_2 = -b_1, c_2 = -c_1$ 即 $\vec{r}_2 = -\vec{r}_1 \therefore \sin\theta_2 = \sin\theta_1$

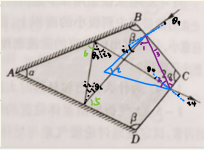
由折射定律 $n \cdot \sin\theta_1 = \sin\theta_3 \therefore \sin\theta_3 = \sin\theta_0 \quad \theta_3 \neq \theta_0 \therefore \theta_3 + \theta_0 = \pi$

同理, 最终出射光线与最初入射光线与 x, y, z 轴夹角之和 $\alpha_0 + \alpha_3 = \pi \quad \beta_0 + \beta_3 = \pi \quad \gamma_0 + \gamma_3 = \pi$

\therefore 出射光线的方向与入射光线相反

实际应用: 自行车后面的反光灯。

1-3. (1)



证明: $L3 + L4 = \pi - 2\alpha$

$$L1 + L2 = (\frac{\pi}{2} - \theta_1 - L3) + (\frac{\pi}{2} - L4 + i_4) = 2\alpha - \theta_1 + i_4$$

$$\text{由折射定律 } \sin\theta_1 = n\sin i_1 \quad n\sin\theta_0 = \sin i_4$$

$$\theta_0 = \frac{\pi}{2} - L5 = \frac{\pi}{2} - [2\pi - (\beta + 2\alpha + \frac{\pi}{2} - i_1)] = -\pi + \beta + 2\alpha - i_1 \quad i_1 = \theta_0$$

$$\theta_0 = \frac{\pi}{2} - L6 = \frac{\pi}{2} - [n \cdot d - (\frac{\pi}{2} - i_2)] = \alpha - i_2$$

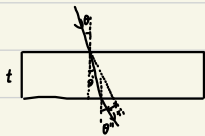
$$\theta_0 = L2 - (\frac{\pi}{2} - L4) - i_4 = \pi + i_1 - 2\alpha - \pi$$

$$\text{得 } \theta_4 = i_1 \quad \pi = \pi - 2\alpha$$

$$\text{即偏向角 } \delta = \pi - \pi = 2\alpha$$

(2). 此情况下无色散.

1-4



证明: 由折射定律 $\sin\theta = n \cdot \sin\theta'$ $n \cdot \sin\theta' = \sin\theta''$ 得 $\theta'' = \theta$ 即出射光线方向不变.

$$n = \frac{t}{\cos\theta'} \cdot \sin(\theta - \theta') \quad \text{当 } \theta \text{ 很小时 } \quad \theta = n \cdot \theta' \quad \sin\theta' = \theta' \quad \sin\theta = \theta \quad \tan\theta' = \theta' \quad \cos\theta = 1$$

$$\text{得 } n = t \cdot (\sin\theta - \frac{\cos\theta \sin\theta'}{\cos\theta'}) = t(\theta - \theta') = \frac{n-1}{n} \theta t$$

1-5. 证明:

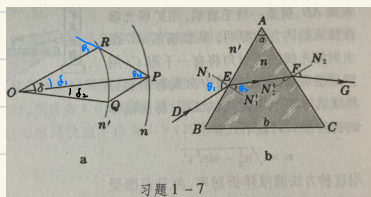


如图共 $m+1$ 层介质 由折射定律 $n_0 \sin\theta_0 = n_1 \sin\theta_1$ $n_1 \sin\theta_1 = n_2 \sin\theta_2 \dots n_{m-1} \sin\theta_{m-1} = n_m \sin\theta_m$

$$\text{得 } n_m \sin\theta_m = n_0 \sin\theta_0$$

即出射光线 出射方向 只与入射方向和 两边折射率有关.

1-7 证明:



由几何关系可知 (b) 中 第一次入射角与折射角关系如 (a) 由折射定律 $n \sin\theta_1 = n' \sin\theta_2$ $\theta_1 = \theta_1 - \theta_2$

$$\text{在 (a) } \triangle ORP \text{ 中 由正弦定理 } \frac{OR}{\sin\theta_2} = \frac{OP}{\sin(\theta_1 - \theta_2)} \quad R < n \quad \text{得 } n' \sin\theta_1 = n \sin\theta_2$$

$$\text{在 } \triangle ORP \text{ 中 } \theta_1 = \theta_1 - \theta_2$$

同理, 第二次折射 也满足因 (a) 且 $\theta_2 = \theta_2 - \theta_3$

综上所述 此法正确.

1-13 证明: 由折射定律 $n_0 \sin\theta_1 = n_1 \sin\theta_1'$

$$\theta_2 = \frac{\pi}{2} - \theta_1'$$

$$\text{满足全反射则 } \theta_2 \geq \arcsin \frac{n_1}{n_2} \quad \text{即 } n_1 \sin\theta_1 \geq n_2 \quad \cos\theta_1' \geq \frac{n_2}{n_1}$$

$$\therefore \text{最大孔经角满足 } n_0 \sin\theta_1 = n_1 \cdot \sqrt{1 - (\frac{n_2}{n_1})^2} \quad \text{即 } n_0 \sin\theta_1 = \sqrt{n_1^2 - n_2^2}$$

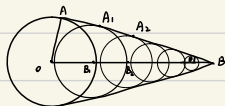
1-15 证明: 光线从透镜出射位置处由折射定律 $n_g \sin \theta = \sin \theta'$

该光线从液体入射透镜位置处由折射定律 $n \sin \frac{\pi}{2} = n_g \sin (\frac{\pi}{2} - \theta)$

$$\text{则 } n = n_g \sqrt{1 - \frac{\sin^2 \theta'}{n_g^2}} = \sqrt{n_g^2 - \sin^2 \theta'}$$

限制: 只能测量折射率小于 n_g 的液体折射率

1-17 证明



如图 波源从 O 沿 OB 方向以 v 速率移动

波前围成如图锥面

$$\sin \theta = \frac{OA}{OB} = \frac{vt}{vt} = \frac{v}{v}$$

若电子以大于介质光速的速度在介质中运动 会产生发光的尾迹, 称为切伦科夫辐射

1-25 证明: 理想漫射体可视作朗伯反射体, 满足余弦发光定律

$$E = \frac{d\Phi}{d\Omega} = \iint B \cos \theta \cdot \sin \theta \, d\Omega \, d\theta = B \int_0^{2\pi} d\Omega \int_0^{\frac{\pi}{2}} \sin \theta \cos \theta \, d\theta = \pi B$$

得 反射光亮度 $B = \frac{E}{\pi}$