

第 1 题得分：_____。(i) 证明下列旋量波函数的表达式：

$$u_{\vec{p},s} = \eta \begin{pmatrix} \sqrt{\frac{E_p+m}{2E_p}} \varphi_s \\ 2s \sqrt{\frac{E_p-m}{2E_p}} \varphi_s \end{pmatrix} \quad v_{-\vec{p},s} = \eta' \begin{pmatrix} -2s \sqrt{\frac{E_p-m}{2E_p}} \varphi_s \\ \sqrt{\frac{E_p+m}{2E_p}} \varphi_s \end{pmatrix}$$

其中 $E_p = \sqrt{\vec{p}^2 + m^2}$, η, η' 为任意相因子, $|\eta| = |\eta'| = 1$, 2×1 列矩阵 φ_s 为 $\vec{\tau} \cdot \hat{p}$ 的本征矢量, 即

$$\vec{\sigma} \cdot \hat{p} \varphi_s = 2s \varphi_s$$

(ii) 讨论上述旋量波函数的非相对论极限。

解: (i) 旋量波函数 $u_{\vec{p},s}$ 满足

$$\begin{cases} (\vec{\alpha} \cdot \vec{p} + \beta m) u_{\vec{p},s} = E_p u_{\vec{p},s}, \\ \vec{\sigma} \cdot \hat{p} u_{\vec{p},s} = 2s u_{\vec{p},s}. \end{cases}$$

其中 $\vec{\alpha} = \rho_1 \vec{\sigma} = \begin{pmatrix} 0 & \vec{\tau} \\ \vec{\tau} & 0 \end{pmatrix}$, $\beta = \begin{pmatrix} I & 0 \\ 0 & -I \end{pmatrix}$. 令 $u_{\vec{p},s} = \begin{pmatrix} u_1 \\ u_2 \end{pmatrix}$, 其中 u_1, u_2 分别表示两个 2×1 列矩阵。则上述方程可以写为：

$$\begin{pmatrix} m & \vec{\tau} \cdot \vec{p} \\ \vec{\tau} \cdot \vec{p} & -m \end{pmatrix} \begin{pmatrix} u_1 \\ u_2 \end{pmatrix} = E_p \begin{pmatrix} u_1 \\ u_2 \end{pmatrix}; \quad \begin{pmatrix} \vec{\tau} \cdot \hat{p} & 0 \\ 0 & \vec{\tau} \cdot \hat{p} \end{pmatrix} \begin{pmatrix} u_1 \\ u_2 \end{pmatrix} = 2s \begin{pmatrix} u_1 \\ u_2 \end{pmatrix}.$$

可以验证：

$$\begin{aligned} \eta \left(m \sqrt{\frac{E_p+m}{2E_p}} + (2s)^2 |\vec{p}| \sqrt{\frac{E_p-m}{2E_p}} \right) \varphi_s &= \eta \sqrt{\frac{E_p+m}{2E_p}} \left(m + |\vec{p}| \frac{E_p-m}{\sqrt{E_p^2-m^2}} \right) \varphi_s = \eta \sqrt{\frac{E_p+m}{2E_p}} \varphi_s; \\ \eta \left(2s |\vec{p}| \sqrt{\frac{E_p+m}{2E_p}} - m 2s \sqrt{\frac{E_p-m}{2E_p}} \right) \varphi_s &= \eta (2s) \sqrt{\frac{E_p-m}{2E_p}} \left(|\vec{p}| \frac{E_p+m}{\sqrt{E_p^2-m^2}} - m \right) \varphi_s = \eta 2s \sqrt{\frac{E_p-m}{2E_p}} \varphi_s. \end{aligned}$$

其中, 利用了 $\sqrt{E_p^2-m^2} = |\vec{p}|$ 和 $2s = \pm 1, 4s^2 = 1$ 。

同理, 对于 $v_{-\vec{p},s}$, Dirac Hamiltonian 为:

$$\begin{cases} (\vec{\alpha} \cdot \vec{p} + \beta m) v_{-\vec{p},s} = -E_p v_{-\vec{p},s}, \\ \vec{\sigma} \cdot \hat{p} v_{-\vec{p},s} = 2s v_{-\vec{p},s}. \end{cases}$$

可以验证：

$$\begin{aligned} \eta' \left(-2sm \sqrt{\frac{E_p-m}{2E_p}} + 2s |\vec{p}| \sqrt{\frac{E_p+m}{2E_p}} \right) \varphi_s &= \eta' (-2s) \sqrt{\frac{E_p-m}{2E_p}} \left(m - |\vec{p}| \frac{E_p+m}{\sqrt{E_p^2-m^2}} \right) \varphi_s = -E_p \eta' (-2s) \sqrt{\frac{E_p-m}{2E_p}} \varphi_s; \\ \eta' \left(-(2s)^2 |\vec{p}| \sqrt{\frac{E_p-m}{2E_p}} - m \sqrt{\frac{E_p+m}{2E_p}} \right) \varphi_s &= -\eta' \sqrt{\frac{E_p+m}{2E_p}} \left(|\vec{p}| \frac{E_p-m}{\sqrt{E_p^2-m^2}} + m \right) \varphi_s = -\eta' \sqrt{\frac{E_p+m}{2E_p}} \varphi_s. \end{aligned}$$

(ii) 对于非相对论极限, 有 $m \gg |\vec{p}|$, 即有:

$$\frac{E_p}{m} = \sqrt{1 + \frac{\vec{p}^2}{m^2}} = 1 + \frac{\vec{p}^2}{2m^2} + \mathcal{O}(\vec{p}^2/m^2).$$

则有：

$$\begin{cases} \frac{E_p+m}{2E_p} = \frac{1}{2} + \frac{1}{2(E_p/m)} = \frac{1}{2} + \frac{1}{2} \left(1 - \frac{\vec{p}^2}{2m^2} \right) = 1 - \frac{\vec{p}^2}{4m^2}; \\ \frac{E_p-m}{2E_p} = \frac{1}{2} - \frac{1}{2(E_p/m)} = \frac{1}{2} - \frac{1}{2} \left(1 - \frac{\vec{p}^2}{2m^2} \right) = \frac{\vec{p}^2}{4m^2}. \end{cases}$$

进而有：

$$\begin{cases} \sqrt{\frac{E_p + m}{2E_p}} = \sqrt{1 - \frac{\vec{p}^2}{4m^2}} = 1 - \frac{\vec{p}^2}{8m^2}; \\ \sqrt{\frac{E_p - m}{2E_p}} = \frac{|\vec{p}|}{2m}. \end{cases}$$

则上述旋量波函数的非相对论极限为：

$$u_{\vec{p},s} = \eta \begin{pmatrix} \sqrt{1 - \frac{\vec{p}^2}{4m^2}} \varphi_s \\ s \frac{|\vec{p}|}{m} \varphi_s \end{pmatrix} \quad v_{-\vec{p},s} = \eta' \begin{pmatrix} -s \frac{|\vec{p}|}{m} \varphi_s \\ \sqrt{1 - \frac{\vec{p}^2}{4m^2}} \varphi_s \end{pmatrix}$$

□

第 2 题得分：_____. 证明

$$\begin{aligned} \sum_s u_{\vec{p},s} u_{\vec{p},s}^\dagger \beta &= \frac{\not{p} + m}{2p_0} \\ \sum_s v_{\vec{p},s} v_{\vec{p},s}^\dagger \beta &= \frac{\not{p} - m}{2p_0} \end{aligned}$$

其中 $\not{p} = -i\gamma_\mu p_\mu$, $p_4 = ip_0 = i\sqrt{\vec{p}^2 + m^2}$.

解：可以计算得：

$$u_{\vec{p},s} u_{\vec{p},s}^\dagger \beta = \begin{pmatrix} \sqrt{\frac{E_p+m}{2E_p}} \varphi_s \\ 2s \sqrt{\frac{E_p-m}{2E_p}} \varphi_s \end{pmatrix} \begin{pmatrix} \sqrt{\frac{E_p+m}{2E_p}} \varphi_s^\dagger & 2s \sqrt{\frac{E_p-m}{2E_p}} \varphi_s^\dagger \end{pmatrix} \begin{pmatrix} I & 0 \\ 0 & -I \end{pmatrix} = \begin{pmatrix} \frac{E_p+m}{2E_p} \varphi_s \varphi_s^\dagger & -s \frac{|\vec{p}|}{E_p} \varphi_s \varphi_s^\dagger \\ s \frac{|\vec{p}|}{E_p} \varphi_s \varphi_s^\dagger & -\frac{E_p-m}{2E_p} \varphi_s \varphi_s^\dagger \end{pmatrix}$$

对于等号右侧，可以计算：

$$\frac{\not{p} + m}{2p_0} = \frac{1}{2E_p} (-i\vec{\gamma} \cdot \vec{p} - i\gamma_4 p_4 + m) = \frac{1}{2E_p} \begin{pmatrix} E_p + m & -\vec{\tau} \cdot \vec{p} \\ \vec{\tau} \cdot \vec{p} & m - E_p \end{pmatrix}$$

注意到：

$$\sum_s s |\vec{p}| \varphi_s \varphi_s^\dagger = \sum_s \tau \cdot \vec{p} \varphi_s \varphi_s^\dagger = \tau \cdot \vec{p} \sum_s \varphi_s \varphi_s^\dagger = \tau \cdot \vec{p}.$$

其中，利用了 $\sum_s \varphi_s \varphi_s^\dagger = I$ 。则有：

$$\sum_s u_{\vec{p},s} u_{\vec{p},s}^\dagger \beta = \frac{1}{2E_p} \begin{pmatrix} (E_p + m) \sum_s \varphi_s \varphi_s^\dagger & \sum_s -2s |\vec{p}| \varphi_s \varphi_s^\dagger \\ \sum_s 2s |\vec{p}| \varphi_s \varphi_s^\dagger & (m - E_p) \sum_s \varphi_s \varphi_s^\dagger \end{pmatrix} = \frac{1}{2E_p} \begin{pmatrix} E_p + m & -\tau \cdot \vec{p} \\ \tau \cdot \vec{p} & m - E_p \end{pmatrix} = \frac{\not{p} + m}{2p_0}$$

同理，对于第二个方程，可以计算：

$$v_{\vec{p},s} v_{\vec{p},s}^\dagger \beta = \begin{pmatrix} -2s \sqrt{\frac{E_p-m}{2E_p}} \varphi_s \\ \sqrt{\frac{E_p+m}{2E_p}} \varphi_s \end{pmatrix} \begin{pmatrix} -2s \sqrt{\frac{E_p-m}{2E_p}} \varphi_s^\dagger & \sqrt{\frac{E_p+m}{2E_p}} \varphi_s^\dagger \end{pmatrix} \begin{pmatrix} I & 0 \\ 0 & -I \end{pmatrix} = \begin{pmatrix} \frac{E_p-m}{2E_p} \varphi_s \varphi_s^\dagger & s \frac{|\vec{p}|}{E_p} \varphi_s \varphi_s^\dagger \\ -s \frac{|\vec{p}|}{E_p} \varphi_s \varphi_s^\dagger & -\frac{E_p+m}{2E_p} \varphi_s \varphi_s^\dagger \end{pmatrix}$$

求和得到：

$$\sum_s v_{\vec{p},s} v_{\vec{p},s}^\dagger \beta = \frac{1}{2E_p} \begin{pmatrix} (E_p - m) \sum_s \varphi_s \varphi_s^\dagger & \sum_s 2s |\vec{p}| \varphi_s \varphi_s^\dagger \\ \sum_s -2s |\vec{p}| \varphi_s \varphi_s^\dagger & -(E_p + m) \sum_s \varphi_s \varphi_s^\dagger \end{pmatrix} = \frac{1}{2E_p} \begin{pmatrix} E_p - m & \tau \cdot \vec{p} \\ -\tau \cdot \vec{p} & -E_p - m \end{pmatrix}$$

对于等号右侧，可以计算：

$$\frac{\not{p} - m}{2p_0} = \frac{1}{2E_p} (-i\vec{\gamma} \cdot \vec{p} - i\gamma_4 p_4 - m) = \frac{1}{2E_p} \begin{pmatrix} E_p - m & \vec{\tau} \cdot \vec{p} \\ -\vec{\tau} \cdot \vec{p} & -E_p - m \end{pmatrix} = \sum_s v_{\vec{p},s} v_{\vec{p},s}^\dagger \beta.$$

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