

# 热力学统计物理第四次作业

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1

由定义可知:

$$C_p = \left( \frac{\partial H}{\partial T} \right)_p$$

则有:

$$\left( \frac{\partial C_p}{\partial p} \right)_T = \left( \frac{\partial}{\partial p} \right)_T \left( \frac{\partial}{\partial T} \right)_p H = \left( \frac{\partial}{\partial T} \right)_p \left( \frac{\partial H}{\partial p} \right)_T$$

利用热力学基本方程与麦克斯韦关系可知:

$$\left( \frac{\partial H}{\partial p} \right)_T = T \left( \frac{\partial S}{\partial p} \right)_T + V = -T \left( \frac{\partial V}{\partial T} \right)_p + V$$

则有:

$$\left( \frac{\partial C_p}{\partial p} \right)_T = \left( \frac{\partial}{\partial T} \left( -T \left( \frac{\partial V}{\partial T} \right)_p + V \right) \right)_p = -T \left( \frac{\partial^2 V}{\partial T^2} \right)_p$$

考虑  $(\partial C_V / \partial V)$ :

有定义可知:

$$C_V = \left( \frac{\partial U}{\partial T} \right)_V$$

则有:

$$\left( \frac{\partial C_V}{\partial V} \right)_T = \left( \frac{\partial}{\partial V} \right)_T \left( \frac{\partial}{\partial T} \right)_V U = \left( \frac{\partial}{\partial T} \right)_V \left( \frac{\partial U}{\partial V} \right)_T$$

利用热力学基本方程与麦克斯韦关系可知:

$$\left( \frac{\partial U}{\partial V} \right)_T = T \left( \frac{\partial S}{\partial V} \right)_T - p = T \left( \frac{\partial p}{\partial T} \right)_V - p$$

则有:

$$\left( \frac{\partial C_V}{\partial V} \right)_T = \left( \frac{\partial}{\partial T} \left( T \left( \frac{\partial p}{\partial T} \right)_V - p \right) \right)_V = T \left( \frac{\partial^2 p}{\partial T^2} \right)_V$$

利用热力学基本关系可知:

$$C_V = \left( \frac{\partial U}{\partial T} \right)_V = T \left( \frac{\partial S}{\partial T} \right)_V, \quad C_p = \left( \frac{\partial H}{\partial T} \right)_p = T \left( \frac{\partial S}{\partial T} \right)_p$$

令  $S(T, p) = S(T, V(T, p))$ , 则有:

$$\begin{aligned} dS &= \left( \frac{\partial S}{\partial T} \right)_p dT + \left( \frac{\partial S}{\partial p} \right)_T dp \\ &= \left( \frac{\partial S}{\partial T} \right)_V dT + \left( \frac{\partial S}{\partial V} \right)_T dV \\ &= \left( \frac{\partial S}{\partial T} \right)_V dT + \left( \frac{\partial S}{\partial V} \right)_T \left( \frac{\partial V}{\partial T} \right)_p dT + \left( \frac{\partial S}{\partial V} \right)_T \left( \frac{\partial V}{\partial p} \right)_T dp \end{aligned}$$

则有:

$$\left( \frac{\partial S}{\partial T} \right)_p = \left( \frac{\partial S}{\partial T} \right)_V + \left( \frac{\partial S}{\partial V} \right)_T \left( \frac{\partial V}{\partial T} \right)_p$$

利用麦克斯韦关系可得:

$$\begin{aligned} C_p - C_V &= T \left( \frac{\partial S}{\partial V} \right)_T \left( \frac{\partial V}{\partial T} \right)_p \\ &= T \left( \frac{\partial p}{\partial T} \right)_V \left( \frac{\partial V}{\partial T} \right)_p \\ &= \frac{T \left( \frac{\partial p}{\partial T} \right)_V \left( \frac{\partial V}{\partial T} \right)_p \left( -\frac{1}{V} \left( \frac{\partial V}{\partial p} \right)_T \left( \frac{\partial T}{\partial V} \right)_p \left( \frac{\partial V}{\partial T} \right)_p \right)}{-\frac{1}{V} \left( \frac{\partial V}{\partial p} \right)_T} \\ &= \frac{TV \left( \frac{1}{V} \left( \frac{\partial V}{\partial T} \right)_p \right)^2}{-\frac{1}{V} \left( \frac{\partial V}{\partial p} \right)_T} \end{aligned}$$

其中有:

$$\left( \frac{\partial p}{\partial T} \right)_V \left( \frac{\partial T}{\partial V} \right)_p \left( \frac{\partial V}{\partial p} \right)_T = -1, \quad \kappa_T = -\frac{1}{V} \left( \frac{\partial V}{\partial p} \right)_T, \quad \beta_p = \frac{1}{V} \left( \frac{\partial V}{\partial T} \right)_p$$

则有:

$$C_p - C_V = \frac{TV\beta_p^2}{\kappa_T}$$

3

由 Jacobi 行列式性质可知:

$$\left(\frac{\partial V}{\partial p}\right)_T = \frac{\partial(V, T)}{\partial(p, T)}, \quad \left(\frac{\partial V}{\partial p}\right)_S = \frac{\partial(V, S)}{\partial(p, S)}$$

则有:

$$\frac{\kappa_T}{\kappa_S} = \frac{\frac{1}{V} \left(\frac{\partial V}{\partial p}\right)_T}{\frac{1}{V} \left(\frac{\partial V}{\partial p}\right)_T} = \frac{\frac{\partial(V, T)}{\partial(p, T)}}{\frac{\partial(V, S)}{\partial(p, S)}} = \frac{\frac{\partial(p, S)}{\partial(p, T)}}{\frac{\partial(V, S)}{\partial(V, T)}} = \frac{T \left(\frac{\partial S}{\partial T}\right)_p}{T \left(\frac{\partial S}{\partial T}\right)_V} = \frac{C_p}{C_V} = \gamma$$

其中利用了由热力学基本微分方程推得的:

$$C_V = \left(\frac{\partial U}{\partial T}\right)_V = T \left(\frac{\partial S}{\partial T}\right)_V, \quad C_p = \left(\frac{\partial H}{\partial T}\right)_p = T \left(\frac{\partial S}{\partial T}\right)_p$$

4

由于物质的温度保持不变, 则其内能不发生变化, 由热力学第一定律可知, 磁化过程中释放出的热量为:

$$Q = -\Delta U + W = \int_0^H \mu_0 V H dM = \int_0^H \frac{\mu_0 C V H}{T} dH = \frac{\mu_0 C V H^2}{2T}$$

即磁化过程中释放出的热量为:  $\frac{\mu_0 C V H^2}{2T}$

5

5.1

对  $S = S(T, V)$  做全微分得:

$$dS = \left(\frac{\partial S}{\partial T}\right)_V dT + \left(\frac{\partial S}{\partial V}\right)_T dV$$

由热力学基本微分方程可知:

$$\left(\frac{\partial S}{\partial T}\right)_V = \frac{1}{T} \left(\frac{\partial U}{\partial T}\right)_V = \frac{C_V}{T}$$

3

利用麦克斯韦关系与范德瓦尔斯气体状态方程可知:

$$\left(\frac{\partial S}{\partial V}\right)_T = \left(\frac{\partial p}{\partial T}\right)_V = \frac{\nu R}{V - \nu b}$$

则有:

$$\begin{aligned} S(T, V) &= \int dS = \int_{T_0}^T \frac{C_V}{T} dT + \int \frac{\nu R}{V - \nu b} dV + S_0 \\ &= \int_{T_0}^T \frac{C_V}{T} dT + \nu R \ln(V - \nu b) + S_0 \end{aligned}$$

## 5.2

设  $U = U(T, V)$ , 则有:

$$dU = \left(\frac{\partial U}{\partial T}\right)_V dT + \left(\frac{\partial U}{\partial V}\right)_T dV = C_V dT + \left(\frac{\partial U}{\partial V}\right)_T dV$$

由热力学基本方程以及麦克斯韦关系可知:

$$\left(\frac{\partial U}{\partial V}\right)_T = T \left(\frac{\partial S}{\partial V}\right)_T - p = T \left(\frac{\partial p}{\partial T}\right)_V - p = \frac{\nu RT}{V - \nu b} - \frac{\nu RT}{V - \nu b} + \frac{a\nu^2}{V^2} = \frac{a\nu^2}{V^2}$$

则有:

$$dU = C_V dT + \frac{a\nu^2}{V^2} dV$$

两侧积分可得:

$$U = U_0 + \int_{T_0}^T C_V dT' - \frac{a\nu^2}{V}$$

则自由能为:

$$\begin{aligned} F(T, V) &= U - TS = U_0 + \int_{T_0}^T C_V dT' - \frac{a\nu^2}{V} - T \left( \int_{T_0}^T \frac{C_V}{T'} dT' + \nu R \ln(V - \nu b) + S_0 \right) \\ &= \int_{T_0}^T C_V \left( 1 - \frac{T}{T'} \right) dT' - \frac{a\nu^2}{V} - \nu RT \ln(V - \nu b) + F_0 \end{aligned}$$

## 5.3

由范德瓦尔斯气体方程可知:

$$pV = \left( \frac{\nu RT}{V - \nu b} - \frac{a\nu^2}{V^2} \right) V = \frac{\nu RTV}{V - \nu b} - \frac{a\nu^2}{V}$$

则自由焓为：

$$\begin{aligned} G(T, V) &= F + pV = \int_{T_0}^T C_V \left( 1 - \frac{T}{T'} \right) dT' - \frac{a\nu^2}{V} - \nu RT \ln(V - \nu b) + F_0 + pV \\ &= \int_{T_0}^T C_V \left( 1 - \frac{T}{T'} \right) dT' + \frac{\nu RTV}{V - nb} - \frac{2a\nu^2}{V} - \nu RT \ln(V - \nu b) + G_0 \end{aligned}$$