Methods

# 1 Self-adjoint ODEs

#### **Fourier Series**

- In the vector space of 'nice enough' periodic functions:  $\mathbb{R} \to \mathbb{R}$  with inner product  $\langle f, g \rangle = \int_{-L}^{L} f(x)g(x)\mathrm{d}x$ , the functions  $\cos\left(\frac{n\pi x}{L}\right)$  and  $\sin\left(\frac{n\pi x}{L}\right)$  form a countable complete orthonormal basis
- Dirichlet conditions:
  - 1. Finitely many bouned discontinuties
  - 2. Finitely many extrema in one period
  - 3. Absolutely integrable (Bounded?)
- Fourier series:

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} \left( a_n \cos\left(\frac{n\pi x}{L}\right) + b_n \sin\left(\frac{n\pi x}{L}\right) \right),$$

where 
$$a_n = \frac{1}{L} \int_{-L}^{L} f(x) \cos\left(\frac{n\pi x}{L}\right) dx$$
,  

$$b_n = \frac{1}{L} \int_{-L}^{L} f(x) \sin\left(\frac{n\pi x}{L}\right) dx$$

- Half-range series
- Fourier series converge to average of left and right limits at discontinuities
- Order of coefficients related to differentiability of f and discontinuties: f(p) discts  $\implies a_n, b_n = O(n^{-(p+1)})$
- Can always integrate, but not always differentiate

Can also formulate using complex exponential as basis over periodic function  $\mathbb{R} \to \mathbb{C}$  with inner product  $\langle f, g \rangle = \int_{-L}^{L} f^*(x) g(x) dx$ :

$$f(x) = \sum_{n=-\infty}^{\infty} c_n e^{\frac{in\pi x}{L}}, \quad \text{where } c_n = \frac{1}{2L} \int_{-L}^{L} f(x) e^{\frac{-in\pi x}{L}} dx$$

• Parseval's Theorem:

$$\int_{-L}^{L} |f|^2 dx = 2L \sum_{n=-\infty}^{\infty} |c_n|^2 = L \left( \frac{a_0^2}{2} + \sum_{n=1}^{\infty} (a_n^2 + b_n^2) \right)$$

## Sturm-Liouville Theory

Background: (DE, LA)

- General solution to DE = Particular + Complementary (*n* linearly indep and spanning), uniqueness by boundary/initial conditions
- On inner product space V, a self-adjoint map (i.e. hermitian matrices for finite dim)  $M: V \to V$  satisfies  $\langle Mu, v \rangle = \langle u, Mv \rangle \, \forall u, v \in V$
- Self-adjoint maps are orthogonally diagonalisable, can use eigenvector basis to solve Mx = b (get coeffs)

S-L Theory: Consider linear differential operators as linear maps between inner product function spaces:

Aim: Solve 2nd order linear ODE:

$$\mathcal{D}y = \alpha y'' + \beta y' + \gamma y = f, \quad a \le x \le b$$

and two boundary conditions

- If  $\alpha \neq 0$ , any eigenvalue problem  $(\mathcal{D}y = \lambda y)$  becomes  $\mathcal{L}y = -(py')' + qy = \lambda wy$ , weight function w with countably many zeroes
- S-L form is self-adjoint iff  $[p(\bar{y_1}'y_2 \bar{y_1}y_2')]_a^b = 0$ , for all  $y_1, y_2$  that satisfy BC, includes many types (e.g. homogeneous Dirichlet, Neumann, mixed)

Good properties of self-adjoint operators:

- If  $\mathcal{L}$  is self-adjoint, get countably infinite, complete basis of orthogonal (orthonormal) simple eigenfunctions with real eigenvalues:  $\mathcal{L}y_n = \lambda_n w y_n$ , orthogonal in the sense  $\langle u, v \rangle_w = \int_a^b w(\bar{u}v) dx$
- Can use basis to solve for particular integral, if forcing has same BC as solution

- Parseval's Theorem for eigenfunction expansion: if  $f(x) = \sum_{n \in \mathbb{N}} a_n y_n$ , then  $\int_a^b w(f(x))^2 dx = \sum_{n \in \mathbb{N}} (a_n)^2 \|y_n\|_w^2$  (norm of  $y_n$  is integral with weight)
- Bessel's inequality, for incomplete set of eigenfunctions
- Least squares approx: Using a finite set of eigenfunctions  $S_N(x) = \sum_{n=1}^N b_n y_n$  to approximate  $f(x) = \sum_{n \in \mathbb{N}} a_n y_n$ : choosing  $b_n = a_n$  minimises mean squared error  $\varepsilon_N = \int_a^b w \left[ f - S_N \right]^2 \mathrm{d}x$

Can solve forced problems Taking product with Green's functions 'inverts' L

# 2 PDEs, Separation of variables

#### Wave equation

1D: 
$$y_{tt} = c^2 y_{xx}$$
  
2D:  $\frac{\partial^2 \phi}{\partial t^2} = c^2 \nabla^2 \phi, c = \sqrt{T/\mu}$ 

- SOV: if y = X(x)T(t) two sides depends on different variables  $\Longrightarrow$  both sides constant.
- If given BC: homogeneous Dirichlet at end points, SOV give eigenvalue problems; solutions must be quantised sine wave, called the normal modes

$$y(x,t) = \sum_{n=1}^{\infty} \left( a_n \cos \left( \frac{n\pi ct}{L} \right) + b_n \sin \left( \frac{n\pi ct}{L} \right) \right) \sin \left( \frac{n\pi x}{L} \right)$$

- IC: initial position and velocity; determine the coefficient of each separable solution using Fourier
- Energy = sum of squares of coeffs, indep of time (use orthogonality)
- Reflection at boundary of changing mass density, three waves: incident, transmitted and reflected; y, y' continuous at boundary at all time; phase shift

#### Drum problem:

• Wave equation in 2D plane polars, with fixed circular boundary conditions  $u(r=1,\theta,t)=0$ 

Suppose solutions separable:  $u = R(r)\Theta(\theta)T(t)$ 

 $\bullet$  Equation for R gives Bessel's equation

$$z^{2} \frac{d^{2}R}{dz^{2}} + z \frac{dR}{dz} + (z^{2} - m^{2})R = 0$$

- Power series solutions: Bessel/Neumann functions  $(J_m(x), Y_m(x))$
- Regularity (boundedness) at r = 0 forces only Bessel function; so solution for space component:

$$X_{mn}(r,\theta) = J_m(j_{mn}r)(A_n\cos(m\theta) + B_n\sin(m\theta))$$

- General space solution = linear combination of  $X_{mn}$ :
- Solution for u:

$$u(r, \theta, t) = \sum_{n=1}^{\infty} J_0(j_{0n}r) \left( A_{0n} \cos(j_{0n}ct) + C_{0n} \sin(j_{0n}ct) \right)$$

$$\sum_{m=1}^{\infty} \sum_{n=1}^{\infty} J_m(j_{mn}r) \left( A_{mn} \cos(m\theta) + B_{mn} \sin(m\theta) \right) \left( \cos(j_{mn}ct) \right)$$

$$+ \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} J_m(j_{mn}r) \left( C_{mn} \cos(m\theta) + D_{mn} \sin(m\theta) \right) \left( \sin(j_{mn}ct) \right)$$

IC determine coefficients by orthogonality

## Diffusion equation

• Flux proportional to diffusion gradient, and opposite direction:

$$\mathbf{q} = -k\nabla\theta$$

- Changes in total heat = (net flux outwards)
- Diffusion equation:

$$\frac{\partial \theta}{\partial t} = D\nabla^2 \theta$$

- Alternative derivation from random walk
- Error function is a solution without BC; good initial approximation for solution with BCs
- SOV: reformulate problem into homogeneous, get transient solution

## Laplace's equation

- Cartesian, Plane polar, Cylindrical polar, Spherical polar
- Spherical polar, assume axisymmetric, equation for  $\Theta$  give Legendre's equation by substitution  $x = \cos \theta \in [-1, 1]$ :

$$-\frac{\mathrm{d}}{\mathrm{d}x}\left((1-x^2)\frac{\mathrm{d}\Theta}{\mathrm{d}x}\right) = \lambda\Theta$$

- Series solution, bounded at endpoints  $\implies$  solutions are polynomials, with  $\lambda_l = l(l+1)$
- Legendre polynomials: Scaling convention  $P_l(1) = 1$ , orthogonal,

$$\int_{-1}^{1} P_n(x) \mathrm{d}x = \frac{2}{2n+1}$$

• Generating function

# 3 Generalised function, Inhom ODE

• Dirac Delta:  $\delta(x)$ : defined by its property

$$\int_{-\infty}^{\infty} \delta(x) f(x) dx = f(0)$$

(sampling property)

• Dirac Delta as a limit of discrete functions

$$\delta_n(x) = \begin{cases} 0 &, |x| > 1/n \\ n/2 &, |x| \le 1/n \end{cases}$$

• Dirac Delta as a limit of continuous Gaussian functions, mean 0, var  $\varepsilon^2$ :

$$\delta_{\varepsilon}(x) = \frac{1}{\varepsilon \sqrt{\pi}} e^{-x^2/\varepsilon^2}$$

- Properties:  $\delta(ax)$ ,  $\delta(f(x))$ , integral of  $\delta'(x)$
- Write Dirac Delta as a Fourier series (Dirac comb): if  $f(x) = \delta(x)$  on -L < x < L, then

$$\delta(x) = \frac{1}{2L} \sum_{n=-\infty}^{\infty} c_n e^{in\pi x/L}$$

• Using Dirac Delta to define Green's function: solution to

$$\mathcal{L}G(x;\xi) = \delta(x-\xi)$$

with homogeneous BC

## 4 Fourier Transform

Fourier Transform:

$$\tilde{f}(k) = \mathcal{F}[f](k) = \int_{-\infty}^{\infty} f(x)e^{-ikx}dx$$

Inverse Fourier Transform:

$$f(x) = \mathcal{F}^{-1}[\tilde{f}](x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{f}(k)e^{ikx}dk$$

(Dual property)

Properties:

- FT of xf(x)
- FT of f':  $\tilde{f}'(k) = ik\tilde{f}(k)$
- Convolution:  $h(x) = (f * g)(x) = \int_{-\infty}^{\infty} f(x u)g(u)du$  iff  $\tilde{h} = \tilde{f}\tilde{g}$  Parseval's Theorem:

$$\int_{-\infty}^{\infty} |f(u)|^2 du = \frac{1}{2\pi} \int_{\infty}^{\infty} |\tilde{f}(k)|^2 dk$$

- FT of trig, Heaviside, Delta
- Express Dirac Delta as a FT:  $delta \longleftrightarrow 1$

#### 5 PDEs

Cauchy problem: include BVP/IVP as auxiliary data (Cauchy data) Well posed problem:

- Solution exists
- Solution is unique
- Solution depends continuously on auxiliary data

#### Method of characteristics

1st order, Change coordinates along characteristics and along boundary curve

#### Classification of 2nd order PDEs

- General form with coefficient functions, compute value of  $b^2 ac$  to classify:
  - ->0: hyperbolic, two real characteristics, e.g. wave
  - -=0: parabolic, one real characteristics, e.g. heat
  - ->0: elliptic, complex characteristics, e.g. Laplace
- Equations for characteristics
- Hyperbolic equation: use two sets of characteristics as coordinates, canonical form for the equation
- d'Alembert's solution for wave equation: change variable to x + ct and x ct, get solution

$$u = f(x + ct) + q(x - ct)$$

Solving PDEs with FT Heat equation: FT wrt x, solve, convolution of IC and source function/fundamental solution/diffusion kernel

$$S_d(x,t) = \frac{1}{\sqrt{4\pi Dt}} e^{\frac{-x^2}{4Dt}}$$

Duhamel's principle:

- (Hyperbolic/parabolic, i.e. wave/heat equation)
- Greens function (forcing is product of deltas);

$$G = H(t - \tau)S_d(x - \xi, t - \tau)$$

• Built in causal relationship

Poisson's equation (forced Laplace's):

- Higher dim delta function (higher dim intergation)
- Solve with delta forcing:

$$\nabla^2 G(\mathbf{r}; \mathbf{r_0}) = \delta(\mathbf{r} - \mathbf{r_0})$$

• Divergence theorem; get formula for free space Green's function (different in 3D and 2D)

## Method of Images

Do something outside of domain to eliminate boundary condition (e.g. another heat source)