## Why pseudobands works?

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## A review of GW

## **GW** in momentum space

$$-\,\mathrm{i}\,\Sigma\approx\,GW^{\mathrm{RPA}}=\underbrace{\phantom{+}\cdots\phantom{+}}_{\text{viq+G}}\,.\,\,\mathrm{In}\,\,\mathbf{k}\,\,\mathrm{space},\,\,\mathit{M}_{\mathit{nn'}}(\mathbf{k},\mathbf{q},\mathbf{G})=\underbrace{\langle\mathit{n}\mathbf{k}+\mathbf{q}|\mathrm{e}^{\mathrm{i}(\mathbf{q}+\mathbf{G})\cdot\mathbf{r}}|\mathit{n'}\mathbf{k}\rangle}_{\text{quasiparticle approx.}};\,\,\mathit{v}(\mathbf{q}+\mathbf{G})=\frac{4\pi}{V|\mathbf{q}+\mathbf{G}|^2}.$$

In epsilon  $\epsilon_{\mathbf{G}\mathbf{G}'}(\mathbf{q},\omega) = \delta_{\mathbf{G}\mathbf{G}'} - v(q+\mathbf{G})\chi_{\mathbf{G}\mathbf{G}'}(\mathbf{q},\omega)$ ,

$$\chi^{\text{r/a}}_{\mathbf{G}\mathbf{G}'}(\mathbf{q},\omega) = \sum_{\mathbf{k}} \sum_{n}^{\text{occ}} \sum_{n'}^{\text{emp}} M_{nn'}(\mathbf{k}\mathbf{q}\mathbf{G}) M_{nn'}^*(\mathbf{k},\mathbf{q},\mathbf{G}') \left( \frac{1}{\omega + E_{n,\mathbf{k}+\mathbf{q}} - E_{n'\mathbf{k}} \pm \mathrm{i}\, 0^+} + \frac{1}{-\omega + E_{n,\mathbf{k}+\mathbf{q}} - E_{n'\mathbf{k}} \mp \mathrm{i}\, 0^+} \right).$$

In sigma  $\Sigma^{\int_{contour}}\Sigma^{CH}+\Sigma^{SX}$ ,

$$\langle n\mathbf{k}|\Sigma^{\mathsf{SX}}(\omega)|n'\mathbf{k}\rangle = -\sum_{n'}^{\mathsf{CC}}\sum_{\mathbf{CG}'}M_{n''n}^*(\mathbf{k}, -\mathbf{q}, -\mathbf{G})M_{n''n'}(\mathbf{k}, -\mathbf{q}, -\mathbf{G}')\epsilon_{\mathbf{GG}'}^{-1}(\mathbf{q}, \omega - E_{n'', \mathbf{k} - \mathbf{q}})\nu(\mathbf{q} + \mathbf{G}')$$

$$\label{eq:sigma_loss} \left\langle \textit{n} \textbf{k} \middle| \Sigma^{\text{CH}}(\omega) \middle| \textit{n}' \textbf{k} \right\rangle = \frac{i}{2\pi} \sum_{\textit{n}''} \sum_{\textit{q} \textbf{G} \textit{G}'} \textit{M}^*_{\textit{n}'',\textit{n}}(\textbf{k}, -\textbf{q}, -\textbf{G}) \textit{M}_{\textit{n}'',\textit{n}'}(\textbf{k}, -\textbf{q}, -\textbf{G}') \int_0^\infty \frac{\left[ \epsilon^r(\textbf{q}, \omega') - \epsilon^a(\textbf{q}, \omega') \right]_{\textbf{G} \textbf{G}'}^{-1} \mathrm{d}\omega'}{\omega - \textit{E}_{\textit{n} \textbf{k}} - \omega' + i \, 0^+ \, \text{sgn}(\textit{E}_{\textit{n} \textbf{k}})} \textit{v}(\textbf{q} + \textbf{G}').$$

Infinite  $\sum_{nn'}^{CH}$  in  $\chi$  and  $\Sigma_{nn'}^{CH}$ : A major bottleneck of GW in momentum space

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## The problem of summing over many empty bands

One further simplification trick: pseudobands. For  $E_{nk} \gg E_F$ :

$$\begin{split} \{\phi_{n\mathbf{k}}\}_{\text{adjacent energy block }n_{b}} &\to \sum_{n \in \text{block }n_{b}} \phi_{n\mathbf{k}}, \\ \{E_{n\mathbf{k}}\}_{\text{adjacent energy block }n_{b}} &\to \frac{1}{|n_{b}|} \sum_{n \in \text{block }n_{b}} E_{n\mathbf{k}}. \end{split}$$

4000 bands  $\rightarrow \sim$  400 to  $\sim$  1000 bands.

...but is this snake oil? It makes no sense!!!

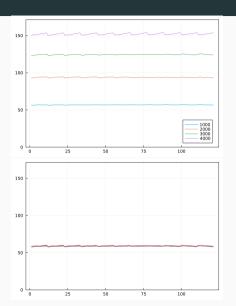
pseudobands, and when it works

## Understanding pseudobands: energy averaging

#### Flatness of high energy bands

High-energy DFT bands of WTe<sub>2</sub> monolayer (ONCV-SG15, 120 electrons; 80 Ry cutoff,  $20 \times 20 \times 1$  grid). x axis =  $\mathbf{k}$  index in irreducible 1BZ; fastest varying coordinate =  $k_y$ 

 $lap{r}{r} \sim 1000$ : bands are generally flat

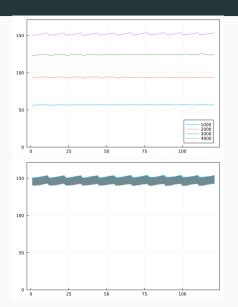


## Understanding pseudobands: energy averaging

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**•** High-lying bands are more dispersive but since they are high,  $1/(E_{\rm emp}-E_{\rm occ})$  is still flat

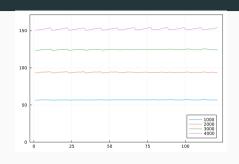


## Understanding pseudobands: energy averaging

#### Flatness of high energy bands

- $\chi \sim MM^* \times f_1(E_{\text{emp}} E_{\text{occ}});$
- $\Sigma^{CH} \sim MM^* \times f_2(\omega E)$ ,  $\omega \sim E_F$ .

For  $E_{n\mathbf{k}}\gg E_{\mathrm{F}}$  ( $\omega$ ):  $f_{1,2}\sim$  const. for all  $\mathbf{k}$ .



## Thus in both $\chi$ and $\Sigma^{CH}$ :

$$\sum_{n''}^{\text{high emp.}} M_{n''n}^* M_{n''n'} \times f(E_{n''k}) = \sum_{\text{block } n_b}^{\text{high emp.}} f(\bar{E}_{n_b}) \sum_{n''}^{\text{block } n_b} M_{n''n}^* M_{n''n'}.$$

## Understanding pseudobands: wave function averaging?

For pseudobands to work, we need

$$\underbrace{\sum_{n'' \in \text{block } n_b} M_{n''n}^* M_{n''n'}}_{\text{normal}} \sim M_{\text{averaged band},n}^* M_{\text{averaged band},n} M_{\text{averaged band},n'}$$

$$= \underbrace{\sum_{n''_1, n''_2 \in \text{block } n_b} M_{n''_1n'} M_{n''_2n}^*}_{\text{pseudobands}}.$$
(1)

The question: is  $[M_{n''_1n'}M^*_{n''_2n}]_{n''_1n''_2}$  diagonal? Easy to verify: it's not.

Then in which case is (1) correct in some sense?

## Technical issues in calculating

 $M_{nn'}(\mathbf{k}, \mathbf{q}, \mathbf{G})$ 

## The structure of $\phi_{n\mathbf{k}}$

Plane wave basis In BerkeleyGW WFN.h5,

$$\phi_{n\mathbf{k}}(\mathbf{r},\sigma) = \frac{1}{\sqrt{V}} \sum_{\mathbf{G}} e^{i(\mathbf{k}+\mathbf{G})\cdot\mathbf{r}} c_{n\mathbf{k},\mathbf{G}\sigma}.$$

Thus

$$M_{nn'}(\mathbf{k}, \mathbf{q}, \mathbf{G}) = \langle n\mathbf{k} + \mathbf{q} | e^{i(\mathbf{q} + \mathbf{G}) \cdot \mathbf{r}} | n'\mathbf{k} \rangle = \sum_{\mathbf{G}', \sigma} c_{n\mathbf{k} + \mathbf{q}, \mathbf{G} + \mathbf{G}'\sigma} c_{n\mathbf{k}, \mathbf{G}'\sigma}.$$

Cutoff Each k has its own G-grid ( $\sim$  30000 vectors for 80 Ry).

#### Procedure

#### Input

- indices of **k**, **q** in **k**-grid;
- index of **G** in **G**-grid of **k** (expect a **G** in *GW* **G**-grid, cutoff = say 30 Ry, not 80 Ry);
- $\bullet$  n, n'.

#### **Procedure**

- 1. find index of k
- 2. find index of  $\mathbf{G} + \mathbf{G}'$  in  $\mathbf{G}$ -grid of  $\mathbf{k} + \mathbf{q}$ , for each  $\mathbf{G}'$  in  $\mathbf{G}$ -grid of  $\mathbf{k}$
- 3. do summation  $\sum_{\mathbf{G}',\sigma} c_{n\mathbf{k}+\mathbf{q},\mathbf{G}+\mathbf{G}'\sigma}^* c_{n\mathbf{k},\mathbf{G}'\sigma}$ .

**Performance** Main bottleneck: finding G + G'. Using StaticArrays.jl helps a lot!

## pseudobands for $\chi$

$$\begin{split} &\chi_{\mathbf{G}\mathbf{G}'}(\mathbf{q},\omega=0)^{\text{high band terms}} \\ &\approx \sum_{\mathbf{k}} \sum_{\text{block } n_{\text{b}}}^{\text{emp}} \frac{2}{E_{n,\mathbf{k}+\mathbf{q}} - E_{\text{average in block } n_{\text{b}}}} \\ &\times \sum_{n' \in \text{block } n_{\text{b}}} \sum_{n}^{\text{occ}} M_{nn'}(\mathbf{k},\mathbf{q},\mathbf{G}) M_{nn'}^{*}(\mathbf{k},\mathbf{q},\mathbf{G}') \end{split}$$

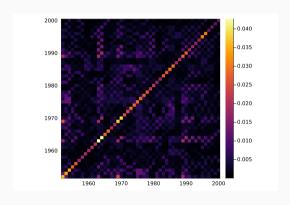
#### Our goal Show that

$$\sum_{n}^{\text{occ}} M_{nn'_1}(\mathbf{k}, \mathbf{q}, \mathbf{G}) M_{nn'_2}^*(\mathbf{k}, \mathbf{q}, \mathbf{G}')$$

is diagonal for pseudobands to work.

$$\begin{split} &\chi_{\mathbf{G}\mathbf{G}'}(\mathbf{q},\omega=0)^{\mathsf{high band terms}} \\ &\approx \sum_{\mathbf{k}} \sum_{\mathsf{block}}^{\mathsf{emp}} \frac{2}{E_{n,\mathbf{k}+\mathbf{q}} - E_{\mathsf{average in block } n_{\mathsf{b}}}} \\ &\times \sum_{n' \in \mathsf{block } n_{\mathsf{b}}} \sum_{n}^{\mathsf{occ}} M_{nn'}(\mathbf{k},\mathbf{q},\mathbf{G}) M_{nn'}^*(\mathbf{k},\mathbf{q},\mathbf{G}') \end{split}$$

For  $\mathbf{G}=\mathbf{G}'$ ,  $\sum\limits_{n}^{\mathrm{occ}}M_{nn_1'}(\mathbf{k},\mathbf{q},\mathbf{G})M_{nn_2'}^*(\mathbf{k},\mathbf{q},\mathbf{G}')$  is diagonal.

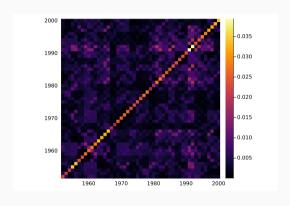


The same WTe<sub>2</sub> run,  $\mathbf{G} = (0, 1, -14), \mathbf{k} = \mathbf{k}_2 = (0, 0.00, 0), \mathbf{q} = (0, 0.00, 0)$ 

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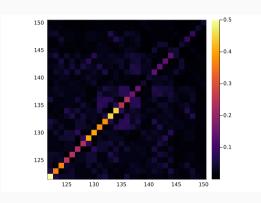
The same WTe<sub>2</sub> run,  $\mathbf{G} = (0, 1, -14), \mathbf{k} = \mathbf{k}_2 = (0, 0.05, 0), \mathbf{q} = (0, 0.10, 0)$ 

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$$\begin{split} &\chi_{\mathbf{G}\mathbf{G}'}(\mathbf{q},\omega=0)^{\text{high band terms}} \\ &\approx \sum_{\mathbf{k}} \sum_{\text{block } n_{\text{b}}}^{\text{emp}} \frac{2}{E_{n,\mathbf{k}+\mathbf{q}} - E_{\text{average in block } n_{\text{b}}}} \\ &\times \sum_{n' \in \text{block } n_{\text{b}}} \sum_{n}^{\text{occ}} M_{nn'}(\mathbf{k},\mathbf{q},\mathbf{G}) M_{nn'}^{*}(\mathbf{k},\mathbf{q},\mathbf{G}') \end{split}$$

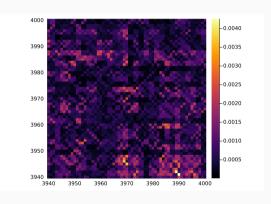
For 
$$\mathbf{G}=\mathbf{G}'$$
,  $\sum\limits_{n}^{\mathrm{occ}}M_{nn_1'}(\mathbf{k},\mathbf{q},\mathbf{G})M_{nn_2'}^*(\mathbf{k},\mathbf{q},\mathbf{G}')$  is diagonal.

Strikingly,  $\sum_{n}^{\text{occ}} M_{nn'_{1}}(\mathbf{k},\mathbf{q},\mathbf{G}) M_{nn'_{2}}^{*}(\mathbf{k},\mathbf{q},\mathbf{G}')$  is still very diagonal for bands near Fermi surface!



$$\begin{split} &\chi_{\mathbf{G}\mathbf{G}'}(\mathbf{q},\omega=0)^{\mathsf{high band terms}} \\ &\approx \sum_{\mathbf{k}} \sum_{\mathsf{block}\ n_\mathsf{b}}^{\mathsf{emp}} \frac{2}{E_{n,\mathbf{k}+\mathbf{q}} - E_{\mathsf{average in block}\ n_\mathsf{b}}} \\ &\times \sum_{n'\in\mathsf{block}} \sum_{n_\mathsf{b}}^{\mathsf{occ}} M_{nn'}(\mathbf{k},\mathbf{q},\mathbf{G}) M_{nn'}^*(\mathbf{k},\mathbf{q},\mathbf{G}') \end{split}$$

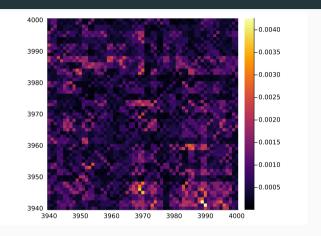
When 
$$\mathbf{G} \neq \mathbf{G}'$$
, 
$$\sum_{n}^{\text{occ}} M_{nn'_1}(\mathbf{k}, \mathbf{q}, \mathbf{G}) M_{nn'_2}^*(\mathbf{k}, \mathbf{q}, \mathbf{G}') \text{ is not diagonal at all.}$$



#### Numerical charaterization of

Case 2:  $\mathbf{G} \neq \mathbf{G}' \sum_{n=0}^{\infty} M_{nn'_1}(\mathbf{k},\mathbf{q},\mathbf{G}) M_{nn'_2}^*(\mathbf{k},\mathbf{q},\mathbf{G}')$  is very non-diagonal, but since the terms's random phases cancel each other so the overall sum after  $\sum_{n'}^{\text{emp}}$  is small

#### Numerical charaterization of

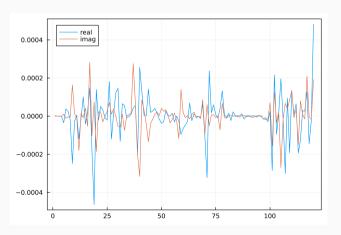


## Half-way generalization about pseudobands in $\chi$

- pseudobands works when it's necessary to do so
- What prevents pseudobands from working around Fermi surface is the energy dispersion

... but do we have any theoretical explanation for this?

 $M_{nn_1'}({f k},{f q},{f G})M_{nn_2'}^*({f k},{f q},{f G}')$  and n (occupied band index), when  $n_1' 
eq n_2'$ :



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■Although for a single n, the value can be large, as we sum over n the terms cancel each other  $\Rightarrow$  diagonal  $\sum_{n}^{\text{occ}} M_{nn'_1}(\mathbf{k},\mathbf{q},\mathbf{G}) M_{nn'_2}^*(\mathbf{k},\mathbf{q},\mathbf{G}')$ 

Why is this the case?

Case 1:  $|n\mathbf{k}\rangle \propto e^{i(\mathbf{G}_n+\mathbf{k})\cdot\mathbf{r}}$  For high-energy bands primarily containing only one **G** components, unwanted terms are:

$$\sum_{n}^{\text{occ}} M_{nn_1}(\mathbf{k}, \mathbf{q}, \mathbf{G}) M_{nn_2}^*(\mathbf{k}, \mathbf{q}, \mathbf{G}) = \sum_{n}^{\text{occ}} c_{n\mathbf{k}+\mathbf{q}} (\mathbf{G} + \mathbf{G}_1)^* c_{n\mathbf{k}+\mathbf{q}} (\mathbf{G} + \mathbf{G}_2)$$

The terms rapid fast when  $\mathbf{G}_1 \neq \mathbf{G}_2$ :

Case 2:  $|n\mathbf{k}\rangle$  contains several **G** components suppose high-energy bands  $n_{1,2}$  primarily contain two **G** components:

$$|n_i \mathbf{k}\rangle = (c_{i1} |\mathbf{G}_1\rangle + c_{i2} |\mathbf{G}_2\rangle) e^{i \mathbf{k} \cdot \mathbf{r}}, \quad i = 1, 2.$$
 (2)

In this way

$$\begin{split} &\sum_{n}^{\text{occ}} M_{nn_1}(\mathbf{k}, \mathbf{q}, \mathbf{G}) M_{nn_2}^*(\mathbf{k}, \mathbf{q}, \mathbf{G}) \\ &= \sum_{n}^{\text{occ}} (c_{11} c_{n\mathbf{k}+\mathbf{q}}^* (\mathbf{G} + \mathbf{G}_1) + c_{12} c_{n\mathbf{k}+\mathbf{q}}^* (\mathbf{G} + \mathbf{G}_2)) \times \\ &\quad (c_{21}^* c_{n\mathbf{k}+\mathbf{q}} (\mathbf{G} + \mathbf{G}_1) + c_{22}^* c_{n\mathbf{k}+\mathbf{q}} (\mathbf{G} + \mathbf{G}_2)). \end{split}$$

The following terms vanish:

$$\sum_{n=0}^{\infty} (c_{11}c_{21}^*|c_{n\mathbf{k}+\mathbf{q}}(\mathbf{G}+\mathbf{G}_1)|^2 + c_{12}c_{22}^*|c_{n\mathbf{k}+\mathbf{q}}(\mathbf{G}+\mathbf{G}_2)|^2)$$

- 1. After summation over n,  $\sum_n |c_{n\mathbf{k}+\mathbf{q}}(\mathbf{G}+\mathbf{G}_1)|^2 = \sum_n |c_{n\mathbf{k}+\mathbf{q}}(\mathbf{G}+\mathbf{G}_2)|^2$
- 2. Orthogonal relation:  $c_{11}c_{21}^* + c_{12}c_{22}^* = 0$ .

The other terms

$$\sum_{n=0}^{\infty} (c_{11}c_{22}^*c_{nk+q}^*(\mathbf{G}+\mathbf{G}_1)c_{nk+q}(\mathbf{G}+\mathbf{G}_2)+c_{12}c_{21}^*c_{nk+q}^*(\mathbf{G}+\mathbf{G}_2)c_{nk+q}(\mathbf{G}+\mathbf{G}_1))$$

oscillate in the same way in the first case.

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# pseudobands for $\Sigma$

Only  $\Sigma^{CH}$  involves  $\Sigma^{emp}$ :

$$\langle n\mathbf{k}|\Sigma^{\mathsf{CH, GPP}}|n'\mathbf{k}\rangle = \frac{1}{2}\sum_{n''}\sum_{\mathbf{q}\mathbf{G}\mathbf{G}'}M^*_{n''n}(\mathbf{k},-\mathbf{q},-\mathbf{G})M_{n''n'}(\mathbf{k},-\mathbf{q},-\mathbf{G}')\times$$
something about  $\mathbf{q},\mathbf{k},\mathbf{G},\mathbf{G}'$ 

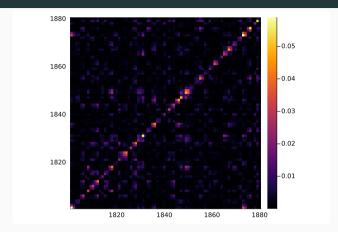
#### **Problems**

- No summation over occupied states ⇒ Xthe aforementioned cancellation mechanism
- Summation over G very complicated  $\Rightarrow$  analysis based on G very hard.

#### Some tentative ideas

- It has been verified that a single  $M_{n_i''n}^* M_{n_2''n}$  term is never diagonal
- But it seems naively summing over  $\mathbf{G} = \mathbf{G}'$  (without considering the weight factor) makes it diagonal...

#### Some tentative ideas



•  $M^*_{n_1''n}(\mathbf{G}_1)M_{n_2''n}(\mathbf{G}_2)$  is not diagonal at all when  $\mathbf{G}_1 \neq \mathbf{G}_2$ ;

#### Some tentative ideas

• but probably only the  $G_1 = G_2$  terms matter?

Comments from Felipe H. Jornada's

recent paper

## "Stochastic" pseudobands

Energy averaging is also done in their paper:

$$G = \sum_{n \text{ near F.S.}} \frac{|\phi_{n\mathbf{k}}\rangle\!\langle\phi_{n\mathbf{k}}|}{\omega - E_{n\mathbf{k}}} + \sum_{\text{pseudobands block}} \frac{1}{\omega - \bar{E}} \sum_{n \in \text{block}} |\phi_{n\mathbf{k}}\rangle\!\langle\phi_{n\mathbf{k}}| \,.$$

Their procedure Replace each pseudobands block by several bands:

$$\{\phi_{n\mathbf{k}}\}_{n \text{ in block}} \longrightarrow \left\{ |\phi_{\xi\mathbf{k}}\rangle = \sum_{n \text{ in block}} \mathrm{e}^{\mathrm{i}\,\theta_{n\xi}} \; |\phi_{n\mathbf{k}}\rangle 
ight\}_{\xi}$$

**Justification** When  $N_{\xi} \to \infty$ ,  $\left\langle e^{-i\theta_{n'\xi}} e^{i\theta_{n\xi}} \right\rangle_{\xi} = \delta_{nn'}$ 

Comparison with the current script When  $N_{\xi}=1$ ,  $\mathrm{e}^{\mathrm{i}\,\theta_{n\xi}}$  comes from DFT diagonalization

## FHJ's justification of pseudobands

**Convergence conditions** When  $N_{\text{pseudobands blocks}} \to \infty$  and in each block  $N_{\xi} \to \infty$ ,

- the expectations of  $\chi$  and  $\Sigma$  goes to the true value;
- the standard error goes to zero.

**Is this result enough as a justification?** My answer: no; because the limits = not doing pseudobands at all. (And their recommended values are not large enough anyway.)

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**Additional justification** "the polarizability tends to converge much faster than G, partially due to the *rapidly oscillating nature* of the matrix elements involving Kohn-Sham states used in the evaluation of the polarizability"

## Compressing the bands even harder: size of blocks

Recall that band dispersion is the main problem...

$$\Delta \chi \sim \frac{\Delta E}{E^2}, \quad \chi \sim \frac{1}{E} \Rightarrow \frac{\Delta \chi}{\chi} \simeq \frac{\Delta E}{E}$$

$$\Rightarrow \frac{\Delta \chi}{\chi} \lesssim \text{const} \Leftrightarrow \boxed{\frac{\Delta E}{E} \lesssim \text{const}}$$

Exponential growth of the energy spread of blocks!

## Compressing the bands even harder: pseudobands near $E_F$ ?

They claim that for epsilon, bands near Fermi surface can also be pseudo-ized

It somehow goes against my own numerical experiments; but they are quite sure about that hmm...

They also claim that for sigma, only the bands near Fermi surface shouldn't be pseudo-ized.

Important claim: protection window is not a convergence parameter.

## **Discussion**

Implications to GW acceleration There are lots of garbage in the giant input files we feed to GW

Implications to machine learning There exists a analytic relation between  $\Sigma^{GW}$  and  $\{\bar{E},\,|\phi_{\mathbf{k}}\rangle^{\mathrm{pseudo}}\}_{\mathrm{blocks}}$ 

- Starting from pseudobands is *not* feature engineering
- Pseudobands (with further compression of autoencoders) should be the starting point of all ML tasks