Details in GW-BSE

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What's GW

To be very concise... $-i\Sigma \approx GW^{RPA} =$ _______.

Assuming quasiparticles... $v(q+G) = \frac{4\pi}{V|q+G|^2}$. $\Sigma^{\int_{\text{contour}}} \Sigma^{\text{COH}} + \Sigma^{\text{SEX}}$, $\epsilon_{GG'}(q,\omega) = \delta_{GG'} - v(q+G) \chi_{GG'}(q,\omega)$, $M_{nn'}(k,q,G) = \langle nk+q | \mathrm{e}^{\mathrm{i}(q+G) \cdot r} | n'k \rangle$

$$\langle n\mathbf{k} | \Sigma^{\text{SEX}}(\omega) | n'\mathbf{k} \rangle = -\sum_{n''}^{\text{occ}} \sum_{\mathbf{q} \mathbf{G} \mathbf{G}'} M_{n''n}^*(\mathbf{k}, -\mathbf{q}, -\mathbf{G}) M_{n''n'}(\mathbf{k}, -\mathbf{q}, -\mathbf{G}')$$

$$\times \epsilon_{\mathbf{G} \mathbf{G}'}^{-1}(\mathbf{q}, \omega - E_{n'', \mathbf{k} - \mathbf{q}}) v(\mathbf{q} + \mathbf{G}'),$$
(1)

$$\langle n\mathbf{k} | \mathbf{\Sigma}^{\mathsf{COH}}(\omega) | n'\mathbf{k} \rangle = \frac{\mathsf{i}}{2\pi} \sum_{n''} \sum_{\mathbf{q}, \mathbf{G}, \mathbf{G}'} M_{n''n}^*(\mathbf{k}, -\mathbf{q}, -\mathbf{G}) M_{n''n'}(\mathbf{k}, -\mathbf{q}, -\mathbf{G}')$$

$$\times \int_0^\infty d\omega' \frac{[\epsilon_{\mathbf{G}\mathbf{G}'}^{\mathsf{c}}]^{-1}(\mathbf{q}, \omega') - [\epsilon_{\mathbf{G}\mathbf{G}'}^{\mathsf{a}}]^{-1}(\mathbf{q}, \omega')}{\omega - E_{n\mathbf{k}} - \omega' + \mathsf{i} \, 0^+ \operatorname{sgn}(E_{n\mathbf{k}})} v(\mathbf{q} + \mathbf{G}').$$
(2)

$$\chi_{\mathbf{GG'}}^{r/a}(\mathbf{q},\omega) = \sum_{\mathbf{k}} \sum_{n} \sum_{n'}^{\text{conp}} M_{nn'}(\mathbf{k},\mathbf{q},\mathbf{G}) M_{nn'}^{*}(\mathbf{k},\mathbf{q},\mathbf{G}') \times \left(\frac{1}{\omega + F_{n}} \sum_{k+n}^{n} F_{n}(\mathbf{k}+\mathbf{q},\mathbf{G}') + \frac{1}{-\omega + F_{n}} \sum_{k+n}^{n} F_{n}(\mathbf{k}+\mathbf{q},\mathbf{G}') \right).$$
(3)

What's GW

What a splendid equation system!

In practice...we always use GPP $\epsilon(\omega)$ is assumed to be plasmon model-like, so – we feed

$$\chi_{\mathbf{G}\mathbf{G}'}(\mathbf{q},\omega=0) = \sum_{\mathbf{k}} \sum_{n} \sum_{n'} \sum_{n'} M_{nn'}(\mathbf{k},\mathbf{q},\mathbf{G}) M_{nn'}^{*}(\mathbf{k},\mathbf{q},\mathbf{G}') \frac{2}{E_{n,\mathbf{k}+\mathbf{q}} - E_{n'\mathbf{k}}}$$
(4)

to analytic expressions of $\Sigma^{\text{COH, SEX}}$.

Huge simplification... Hedin
$$\overset{\text{assuming }GW}{\to}$$
 $GW \overset{\text{QP. approx.}}{\to}$ (1), (2), (3) $\overset{\text{GPP}}{\to}$ we are here

But still burdensome *Summation over empty bands – 1000-30000 bands!!!*

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The problem of summing over many empty bands

Empty states are important - but why?

One further simplification trick: pseudobands. For $E_{nk} \gg E_F$:

$$\begin{split} \{\phi_{n\pmb{k}}\}_{\text{adjacent energy block }n_b} &\to \sum_{n\in \text{block }n_b} \phi_{n\pmb{k}}, \\ \{E_{n\pmb{k}}\}_{\text{adjacent energy block }n_b} &\to \frac{1}{|n_b|} \sum_{n\in \text{block }n_b} E_{n\pmb{k}}. \end{split}$$

4000 bands $\rightarrow \sim$ 400 to \sim 1000 bands.

... but is this snake oil? It makes no sense!!!

Understanding pseudobands: energy averaging

Observation

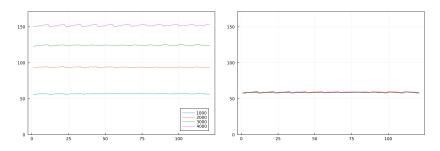
- $\chi \sim MM^* \times \text{some-function}(E_{\text{emp}} E_{\text{occ}});$
- $\Sigma^{\text{COH}} \sim MM^* \times \text{some-function}(\omega E)$;
- We are only interested in $\omega \sim E_{\rm F}$.
- **t** For $E_{n{m k}}\gg E_{\rm F}$ (ω): energy-dependent factors \sim const. for all ${m k}$.

Thus for both χ and Σ^{COH} involving summation over empty bands:

$$\sum_{n''} M_{n''n}^* M_{n''n'} \times \cdots = \sum_{\text{block } n_b} \cdots \times \sum_{n'' \in \text{block } n_b} M_{n''n}^* M_{n''n'}. \quad (5)$$

In RHS E_{emp} is replaced by the average energy.

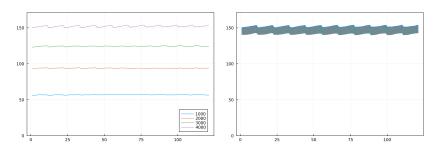
Understanding pseudobands: energy averaging



Example High-energy DFT bands of WTe₂ monolayer (ONCV-SG15, 120 electrons; 80 Ry cutoff, $20 \times 20 \times 1$ grid). x axis = k index in irreducible 1BZ; fastest varying coordinate = k_y

- **■** Bands above 1000 are generally flat (but dispersion \uparrow as $n \uparrow$)
- Low-lying pseudobands blocks are very close in energy (but $1/(E_{\rm emp}-E_{\rm occ})$ more sensitive to dispersion)

Understanding pseudobands: energy averaging



Example High-energy DFT bands of WTe₂ monolayer (ONCV-SG15, 120 electrons; 80 Ry cutoff, $20 \times 20 \times 1$ grid). x axis = k index in irreducible 1BZ; fastest varying coordinate = k_y

- **■** Bands above 1000 are generally flat (but dispersion \uparrow as $n \uparrow$)
- High-energy blocks are more dispersive (but $1/(E_{\rm emp}-E_{\rm occ})$ is smaller so no worry)

Understanding pseudobands: wave function averaging?

For pseudobands to work, we need

$$\sum_{n'' \in \text{block } n_b} M_{n''n}^* M_{n''n'} \sim M_{\text{averaged band},n}^* M_{\text{averaged band},n'}$$

$$= \sum_{n_1'', n_2'' \in \text{block } n_b} M_{n_1''n'} M_{n_2''n}^*. \tag{6}$$

The problem: it of course isn't the case in general. If $[M_{n''_1n'}M^*_{n''_2n}]_{n''_1n''_2}$ is a random matrix: LHS: RHS ≈ 0.3 .

The question: Then in which case is (6) correct in some sense?

The structure of $\phi_{n\mathbf{k}}$

Plane wave basis In BerkeleyGW WFN.h5,

$$\phi_{nk}(\mathbf{r},\sigma) = \frac{1}{\sqrt{V}} \sum_{\mathbf{G}} e^{i(\mathbf{k}+\mathbf{G})\cdot\mathbf{r}} c_{n\mathbf{k},\mathbf{G}\sigma}.$$
 (7)

Thus

$$M_{nn'}(\boldsymbol{k}, \boldsymbol{q}, \boldsymbol{G}) = \langle n\boldsymbol{k} + \boldsymbol{q} | e^{i(\boldsymbol{q} + \boldsymbol{G}) \cdot \boldsymbol{r}} | n' \boldsymbol{k} \rangle = \sum_{\boldsymbol{G}', \sigma} c_{n\boldsymbol{k} + \boldsymbol{q}, \boldsymbol{G} + \boldsymbol{G}' \sigma}^* c_{n\boldsymbol{k}, \boldsymbol{G}' \sigma}.$$
 (8)

Cutoff Each k has its own G-grid (~ 30000 vectors for 80 Ry).

Procedure

Input

- indices of **k**, **q** in **k**-grid;
- index of \boldsymbol{G} in \boldsymbol{G} -grid of \boldsymbol{k} (expect a \boldsymbol{G} in GW \boldsymbol{G} -grid, cutoff = say 30 Ry, not 80 Ry);
- \bullet n, n'.

Procedure

- find index of k
- ② find index of G + G' in G-grid of k + q, for each G' in G-grid of k
- 3 do summation $\sum_{\mathbf{G}',\sigma} c_{n\mathbf{k}+\mathbf{q},\mathbf{G}+\mathbf{G}'\sigma}^* c_{n\mathbf{k},\mathbf{G}'\sigma}$.

Performance Main bottleneck: finding ${m G} + {m G}'$. Using StaticArrays.jl helps a lot!

Under GPP:

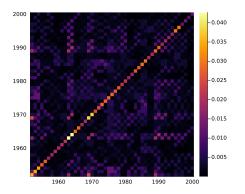
$$\chi_{\boldsymbol{G}\boldsymbol{G}'}(\boldsymbol{q},\omega)^{\text{high band terms}} \approx \sum_{\boldsymbol{k}} \sum_{\text{block } n_{b}}^{\text{emp}} \frac{2}{E_{n,\boldsymbol{k}+\boldsymbol{q}} - E_{\text{average in block } n_{b}}} \times \sum_{n' \in \text{block } n_{b}} \sum_{n}^{\text{occ}} M_{nn'}(\boldsymbol{k},\boldsymbol{q},\boldsymbol{G}) M_{nn'}^{*}(\boldsymbol{k},\boldsymbol{q},\boldsymbol{G}')$$

$$(9)$$

Our goal Finding how diagonal is

$$\sum_{n}^{\text{occ}} M_{nn'}(\boldsymbol{k}, \boldsymbol{q}, \boldsymbol{G}) M_{nn'}^*(\boldsymbol{k}, \boldsymbol{q}, \boldsymbol{G}')$$
 (10)

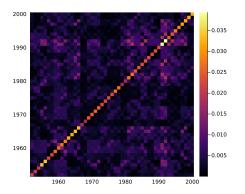
Case 1: G = G' In this case $\sum_{n}^{\text{occ}} M_{nn'}(\mathbf{k}, \mathbf{q}, G) M_{nn'}^*(\mathbf{k}, \mathbf{q}, G')$ is large and fairly diagonal



Parameters $\mathbf{G} = (0, 1, -14), \mathbf{k} = \mathbf{k}_2 = (0, 0.00, 0), \mathbf{q} = (0, 0.00, 0)$

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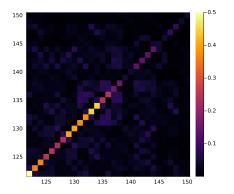


Parameters $\mathbf{G} = (0, 1, -14), \mathbf{k} = \mathbf{k}_2 = (0, 0.05, 0), \mathbf{q} = (0, 0.10, 0)$

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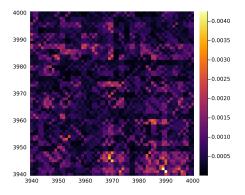
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Strikingly, $\sum_{n}^{\text{occ}} M_{nn'}(\mathbf{k}, \mathbf{q}, \mathbf{G}) M_{nn'}^*(\mathbf{k}, \mathbf{q}, \mathbf{G}')$ is still very diagonal for bands near Fermi surface!

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Case 2: $G \neq G' \sum_{n}^{\text{occ}} M_{nn'}(\mathbf{k}, \mathbf{q}, G) M_{nn'}^*(\mathbf{k}, \mathbf{q}, G')$ is very non-diagonal, but since the terms's random phases cancel each other so it's small



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Half-way generalization about pseudobands in χ

- pseudobands works when it's necessary to do so
- What prevents pseudobands from working around Fermi surface is the energy dispersion

... but do we have any theoretical explanation for this?

The secret of