Why `pseudobands` works?

Numerical experiments, analysis, and possible enhancement of the method

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A review of GW

GW in momentum space

$-\mathrm{i}\,\Sigma \approx GW^{\mathrm{RPA}} =$  . In $\mathbf{k}$ space, $M_{nn'}(\mathbf{k},\mathbf{q},\mathbf{G}) = \underbrace{\langle n\mathbf{k}+\mathbf{q}|\mathrm{e}^{\mathrm{i}(\mathbf{q}+\mathbf{G})\cdot\mathbf{r}}|n'\mathbf{k}\rangle}_{\text{quasiparticle approx.}}$; $v(\mathbf{q}+\mathbf{G}) = \frac{4\pi}{V|\mathbf{q}+\mathbf{G}|^2}$.

In epsilon $\epsilon_{\mathbf{GG}'}(\mathbf{q},\omega) = \delta_{\mathbf{GG}'} - v(q+\mathbf{G})\chi_{\mathbf{GG}'}(\mathbf{q},\omega)$,

$$\chi^{\mathrm{r/a}}_{\mathbf{GG}'}(\mathbf{q},\omega) = \sum_{\mathbf{k}} \sum_{n}^{\mathrm{occ}} \sum_{n'}^{\mathrm{emp}} M_{nn'}(\mathbf{kqG}) M^*_{nn'}(\mathbf{k},\mathbf{q},\mathbf{G}') \left(\frac{1}{\omega + E_{n,\mathbf{k}+\mathbf{q}} - E_{n'\mathbf{k}} \pm \mathrm{i}\,0^+} + \frac{1}{-\omega + E_{n,\mathbf{k}+\mathbf{q}} - E_{n'\mathbf{k}} \mp \mathrm{i}\,0^+} \right).$$

In sigma $\Sigma \overset{\int \text{contour}}{=} \Sigma^{\mathrm{CH}} + \Sigma^{\mathrm{SX}}$,

$$\langle n\mathbf{k}|\Sigma^{\mathrm{SX}}(\omega)|n'\mathbf{k}\rangle = -\sum_{n''}^{\mathrm{occ}} \sum_{\mathbf{qGG}'} M^*_{n''n}(\mathbf{k},-\mathbf{q},-\mathbf{G}) M_{n''n'}(\mathbf{k},-\mathbf{q},-\mathbf{G}') \epsilon^{-1}_{\mathbf{GG}'}(\mathbf{q},\omega - E_{n'',\mathbf{k}-\mathbf{q}}) v(\mathbf{q}+\mathbf{G}')$$

$$\langle n\mathbf{k}|\Sigma^{\mathrm{CH}}(\omega)|n'\mathbf{k}\rangle = \frac{\mathrm{i}}{2\pi} \sum_{n''} \sum_{\mathbf{qGG}'} M^*_{n''n}(\mathbf{k},-\mathbf{q},-\mathbf{G}) M_{n''n'}(\mathbf{k},-\mathbf{q},-\mathbf{G}') \int_0^\infty \frac{[\epsilon^{\mathrm{r}}(\mathbf{q},\omega') - \epsilon^{\mathrm{a}}(\mathbf{q},\omega')]^{-1}_{\mathbf{GG}'}\,\mathrm{d}\omega'}{\omega - E_{n\mathbf{k}} - \omega' + \mathrm{i}\,0^+ \operatorname{sgn}(E_{n\mathbf{k}})} v(\mathbf{q}+\mathbf{G}').$$

Infinite $\sum^{\text{emp.}}$ in χ and $\Sigma^{\mathrm{CH}}_{nn'}$: **A major bottleneck of *GW* in momentum space**

The problem of summing over many empty bands

One further simplification trick: `pseudobands`. For $E_{n\mathbf{k}} \gg E_{\mathrm{F}}$:

$$\{\phi_{n\mathbf{k}}\}_{\text{adjacent energy block } n_{\mathrm{b}}} \to \sum_{n \in \text{block } n_{\mathrm{b}}} \phi_{n\mathbf{k}},$$

$$\{E_{n\mathbf{k}}\}_{\text{adjacent energy block } n_{\mathrm{b}}} \to \frac{1}{|n_{\mathrm{b}}|} \sum_{n \in \text{block } n_{\mathrm{b}}} E_{n\mathbf{k}}.$$

4000 bands $\to \sim 400$ to ~ 1000 bands.

...but is this snake oil? It makes no sense!!!

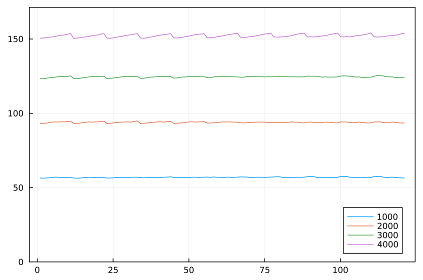
`pseudobands`, and when it works

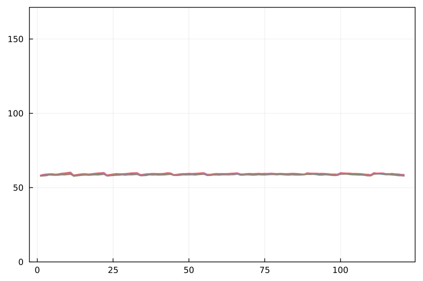
Understanding `pseudobands`: energy averaging

Flatness of high energy bands

High-energy DFT bands of WTe_2 monolayer (ONCV-SG15, 120 electrons; 80 Ry cutoff, $20 \times 20 \times 1$ grid). x axis $=$ $\mathbf{k}$ index in irreducible 1BZ; fastest varying coordinate $= k_y$

☛ $n \sim 1000$: bands are generally flat





Understanding `pseudobands`: energy averaging

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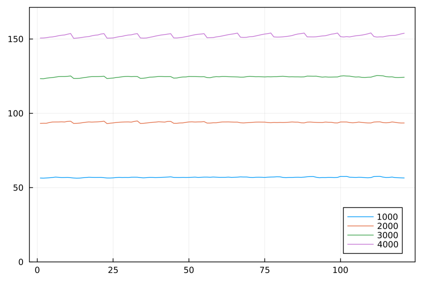
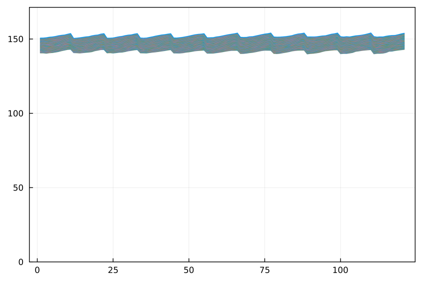
👉High-lying bands are more dispersive but since they are high, $1/(E_{\text{emp}} - E_{\text{occ}})$ is still flat



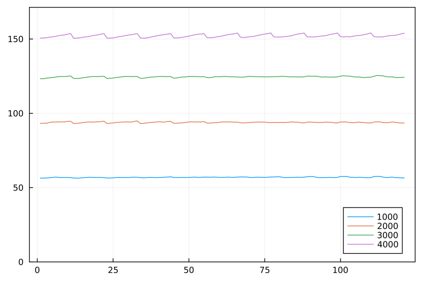


Understanding pseudobands: energy averaging

Flatness of high energy bands

- $\chi \sim MM^* \times f_1(E_{\text{emp}} - E_{\text{occ}})$;
- $\Sigma^{\text{CH}} \sim MM^* \times f_2(\omega - E), \quad \omega \sim E_{\text{F}}$.

For $E_{n\mathbf{k}} \gg E_{\text{F}}$ (ω): $f_{1,2} \sim$ const. for all $\mathbf{k}$.



Thus in both χ and Σ^{CH}:

$$\sum_{n''}^{\text{high emp.}} M^*_{n''n} M_{n''n'} \times f(E_{n''\mathbf{k}}) = \sum_{\text{block } n_{\text{b}}}^{\text{high emp.}} f(\bar{E}_{n_{\text{b}}}) \sum_{n''}^{\text{block } n_{\text{b}}} M^*_{n''n} M_{n''n'}.$$

Understanding `pseudobands`: wave function averaging?

For `pseudobands` to work, we need

$$\begin{aligned}\underbrace{\sum_{n'' \in \text{block } n_\mathrm{b}} M^*_{n'' n} M_{n'' n'}}_{\text{normal}} &\sim M^*_{\text{averaged band},n} M_{\text{averaged band},n'} \\ &= \underbrace{\sum_{n''_1, n''_2 \in \text{block } n_\mathrm{b}} M_{n''_1 n'} M^*_{n''_2 n}}_{\text{pseudobands}} .\end{aligned} \tag{1}$$

The question: is $[M_{n''_1 n'} M^*_{n''_2 n}]_{n''_1 n''_2}$ **diagonal?** Easy to verify: it's not.

Then in which case is (1) **correct in some sense?**

Technical issues in calculating $M_{nn'}(\mathbf{k}, \mathbf{q}, \mathbf{G})$

The structure of $\phi_{n\mathbf{k}}$

Plane wave basis In BerkeleyGW `WFN.h5`,

$$\phi_{n\mathbf{k}}(\mathbf{r}, \sigma) = \frac{1}{\sqrt{V}} \sum_{\mathbf{G}} \mathrm{e}^{\mathrm{i}(\mathbf{k}+\mathbf{G})\cdot\mathbf{r}} \, c_{n\mathbf{k},\mathbf{G}\sigma}.$$

Thus

$$M_{nn'}(\mathbf{k}, \mathbf{q}, \mathbf{G}) = \langle n\mathbf{k}+\mathbf{q} | \mathrm{e}^{\mathrm{i}(\mathbf{q}+\mathbf{G})\cdot\mathbf{r}} | n'\mathbf{k} \rangle = \sum_{\mathbf{G}', \sigma} c^*_{n\mathbf{k}+\mathbf{q},\mathbf{G}+\mathbf{G}'\sigma} c_{n\mathbf{k},\mathbf{G}'\sigma}.$$

Cutoff Each $\mathbf{k}$ has its own $\mathbf{G}$-grid ($\sim$ 30000 vectors for 80 Ry).

Procedure

Input

- indices of $\mathbf{k}, \mathbf{q}$ in $\mathbf{k}$-grid;
- index of $\mathbf{G}$ in $\mathbf{G}$-grid of $\mathbf{k}$ (expect a $\mathbf{G}$ in GW $\mathbf{G}$-grid, cutoff = say 30 Ry, not 80 Ry);
- n, n'.

Procedure

1. find index of $\mathbf{k}$
2. find index of $\mathbf{G} + \mathbf{G}'$ in $\mathbf{G}$-grid of $\mathbf{k} + \mathbf{q}$, for each $\mathbf{G}'$ in $\mathbf{G}$-grid of $\mathbf{k}$
3. do summation $\sum_{\mathbf{G}', \sigma} c^*_{n\mathbf{k}+\mathbf{q}, \mathbf{G}+\mathbf{G}'\sigma} c_{n\mathbf{k}, \mathbf{G}'\sigma}$.

Performance Main bottleneck: finding $\mathbf{G} + \mathbf{G}'$. Using `StaticArrays.jl` helps a lot!

`pseudobands` for χ

χ revisited

Under GPP:

$$
\begin{aligned}
&\chi_{\mathbf{G}\mathbf{G}'}(\mathbf{q}, \omega = 0)^{\text{pseudobands}} \\
\approx &\sum_{\mathbf{k}} \sum_{\text{block } n_{\text{b}}}^{\text{emp}} \frac{2}{E_{n,\mathbf{k}+\mathbf{q}} - E_{\text{average in block } n_{\text{b}}}} \\
&\times \sum_{n'_{1,2} \in \text{block } n_{\text{b}}} \sum_{n}^{\text{occ}} M_{nn'_1}(\mathbf{k}, \mathbf{q}, \mathbf{G}) M^*_{nn'_2}(\mathbf{k}, \mathbf{q}, \mathbf{G}')
\end{aligned}
$$

Our goal Show that

$$
\sum_{n}^{\text{occ}} M_{nn'_1}(\mathbf{k}, \mathbf{q}, \mathbf{G}) M^*_{nn'_2}(\mathbf{k}, \mathbf{q}, \mathbf{G}')
$$

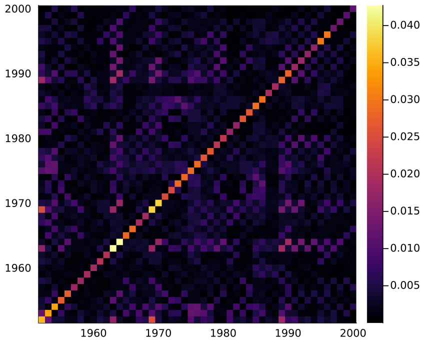
is diagonal for `pseudobands` to work.

χ revisited

Under GPP:

$$\chi_{\mathbf{GG}'}(\mathbf{q}, \omega = 0)^{\text{pseudobands}}$$
$$\approx \sum_{\mathbf{k}} \sum_{\text{block } n_b}^{\text{emp}} \frac{2}{E_{n,\mathbf{k}+\mathbf{q}} - E_{\text{average in block } n_b}}$$
$$\times \sum_{n'_{1,2} \in \text{block } n_b} \sum_{n}^{\text{occ}} M_{nn'_1}(\mathbf{k}, \mathbf{q}, \mathbf{G}) M^*_{nn'_2}(\mathbf{k}, \mathbf{q}, \mathbf{G}')$$

For $\mathbf{G} = \mathbf{G}'$, $\sum_{n}^{\text{occ}} M_{nn'_1}(\mathbf{k}, \mathbf{q}, \mathbf{G}) M^*_{nn'_2}(\mathbf{k}, \mathbf{q}, \mathbf{G}')$ is diagonal.

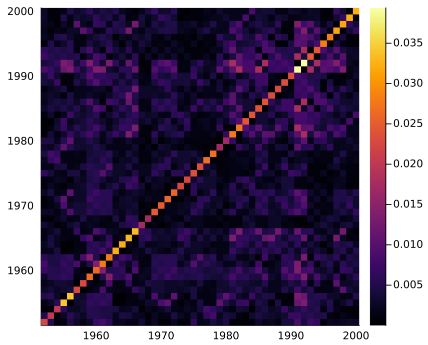


The same WTe_2 run, $\mathbf{G} = (0, 1, -14), \mathbf{k} = \mathbf{k}_2 = (0, 0.00, 0), \mathbf{q} = (0, 0.00, 0)$

Under GPP:

$$\chi_{\mathbf{GG}'}(\mathbf{q},\omega=0)^{\text{pseudobands}}$$
$$\approx \sum_{\mathbf{k}} \sum_{\text{block } n_b}^{\text{emp}} \frac{2}{E_{n,\mathbf{k}+\mathbf{q}} - E_{\text{average in block } n_b}}$$
$$\times \sum_{n'_{1,2}\in\text{block } n_b} \sum_{n}^{\text{occ}} M_{nn'_1}(\mathbf{k},\mathbf{q},\mathbf{G}) M^*_{nn'_2}(\mathbf{k},\mathbf{q},\mathbf{G}')$$

For $\mathbf{G}=\mathbf{G}'$, $\sum_{n}^{\text{occ}} M_{nn'_1}(\mathbf{k},\mathbf{q},\mathbf{G}) M^*_{nn'_2}(\mathbf{k},\mathbf{q},\mathbf{G}')$ is diagonal.



The same WTe_2 run, $\mathbf{G}=(0,1,-14)$, $\mathbf{k}=\mathbf{k}_2=(0,0.05,0)$, $\mathbf{q}=(0,0.10,0)$

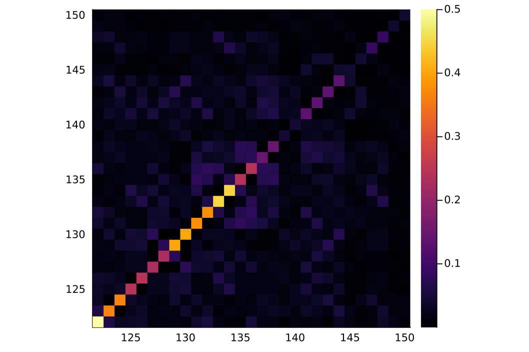
χ revisited

Under GPP:

$$\chi_{\mathbf{GG}'}(\mathbf{q}, \omega = 0)^{\text{pseudobands}}$$
$$\approx \sum_{\mathbf{k}} \sum_{\text{block } n_b}^{\text{emp}} \frac{2}{E_{n,\mathbf{k}+\mathbf{q}} - E_{\text{average in block } n_b}}$$
$$\times \sum_{n'_{1,2} \in \text{block } n_b} \sum_{n}^{\text{occ}} M_{nn'_1}(\mathbf{k}, \mathbf{q}, \mathbf{G}) M^*_{nn'_2}(\mathbf{k}, \mathbf{q}, \mathbf{G}')$$

For $\mathbf{G} = \mathbf{G}'$, $\sum_{n}^{\text{occ}} M_{nn'_1}(\mathbf{k}, \mathbf{q}, \mathbf{G}) M^*_{nn'_2}(\mathbf{k}, \mathbf{q}, \mathbf{G}')$ is diagonal.

Strikingly, $\sum_n^{\text{occ}} M_{nn'_1}(\mathbf{k},\mathbf{q},\mathbf{G}) M^*_{nn'_2}(\mathbf{k},\mathbf{q},\mathbf{G}')$ is still very diagonal for bands near Fermi surface!

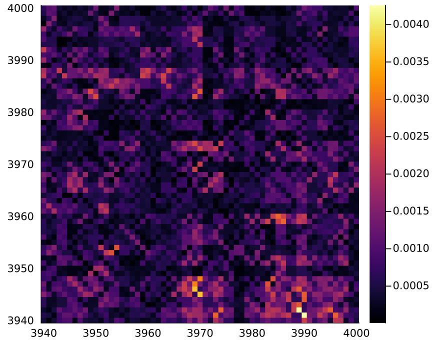


χ revisited

Under GPP:

$$
\begin{aligned}
&\chi_{\mathbf{G}\mathbf{G}'}(\mathbf{q}, \omega = 0)^{\text{pseudobands}} \\
\approx &\sum_{\mathbf{k}} \sum_{\text{block } n_\text{b}}^{\text{emp}} \frac{2}{E_{n,\mathbf{k}+\mathbf{q}} - E_{\text{average in block } n_\text{b}}} \\
&\times \sum_{n'_{1,2} \in \text{block } n_\text{b}} \sum_{n}^{\text{occ}} M_{nn'_1}(\mathbf{k}, \mathbf{q}, \mathbf{G}) M^*_{nn'_2}(\mathbf{k}, \mathbf{q}, \mathbf{G}')
\end{aligned}
$$

When $\mathbf{G} \neq \mathbf{G}'$, $\sum_{n}^{\text{occ}} M_{nn'_1}(\mathbf{k}, \mathbf{q}, \mathbf{G}) M^*_{nn'_2}(\mathbf{k}, \mathbf{q}, \mathbf{G}')$ is not diagonal at all.



χ revisited

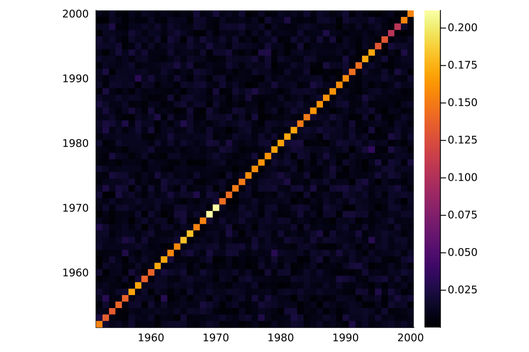
Under GPP:

$$\chi_{\mathbf{G}\mathbf{G}'}(\mathbf{q}, \omega = 0)^{\text{pseudobands}}$$

$$\approx \sum_{\mathbf{k}} \sum_{\text{block } n_b}^{\text{emp}} \frac{2}{E_{n,\mathbf{k}+\mathbf{q}} - E_{\text{average in block } n_b}}$$

$$\times \sum_{n'_{1,2} \in \text{block } n_b} \sum_{n}^{\text{occ}} M_{nn'_1}(\mathbf{k}, \mathbf{q}, \mathbf{G}) M^*_{nn'_2}(\mathbf{k}, \mathbf{q}, \mathbf{G}')$$

But since $E_{n\mathbf{k}+\mathbf{q}} \ll E_{\text{average in block } n_b}$, the $1/\Delta E$ factor is a constant w.r.t. $\mathbf{k}$, and after summing over $\mathbf{k}$, MM^* becomes diagonal.



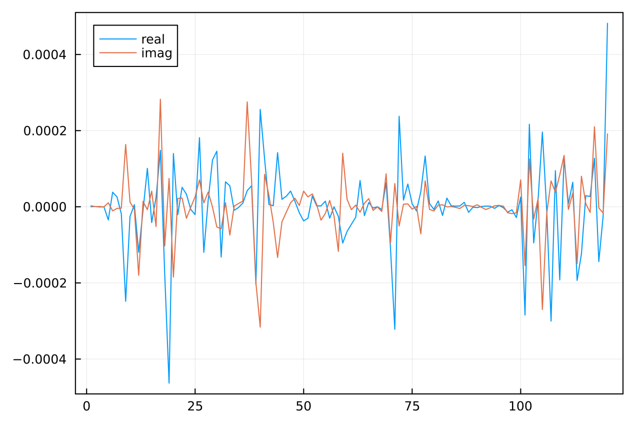
Half-way generalization about `pseudobands` in χ

- `pseudobands` works when it's necessary to do so
- What prevents `pseudobands` from working around Fermi surface is the energy dispersion

. . . but do we have any theoretical explanation for this?

Cancellation of cross-terms in χ

y axis: $M_{nn'_1}(\mathbf{k},\mathbf{q},\mathbf{G})M^*_{nn'_2}(\mathbf{k},\mathbf{q},\mathbf{G}')$, $n'_1 \neq n'_2$ x axis: n (occupied band index)



👉For a single n, MM^* can be large; but they terms cancel each other $\Rightarrow$ diagonal $\sum_n^{\text{occ}} M_{nn'_1}(\mathbf{k},\mathbf{q},\mathbf{G})M^*_{nn'_2}(\mathbf{k},\mathbf{q},\mathbf{G}')$
Why is this the case?

👉**The main reason: random phase factor in DFT diagonalization**

Cancellation of cross-terms in χ

$$
\begin{aligned}
&\sum_{\mathbf{k}} \sum_{n}^{\text{occ}} M_{nn_1'}(\mathbf{k}, \mathbf{q}, \mathbf{G}) M^*_{nn_2'}(\mathbf{k}, \mathbf{q}, \mathbf{G}') \\
=& \sum_{\mathbf{k}} \sum_{n}^{\text{occ}} \sum_{\mathbf{G}_1 \mathbf{G}_2} c^*_{n\mathbf{k}+\mathbf{q}}(\mathbf{G} + \mathbf{G}_1) c_{n_1'\mathbf{k}}(\mathbf{G}_1) c_{n\mathbf{k}+\mathbf{q}}(\mathbf{G}' + \mathbf{G}_2) c^*_{n_2'\mathbf{k}}(\mathbf{G}_2)
\end{aligned}
$$

Cancellation of cross-terms in χ

$$\sum_{\mathbf{k}} \sum_{n}^{\mathrm{occ}} M_{nn_1'}(\mathbf{k},\mathbf{q},\mathbf{G}) M^*_{nn_2'}(\mathbf{k},\mathbf{q},\mathbf{G}')$$

$$= \sum_{\mathbf{k}} \sum_{n}^{\mathrm{occ}} \sum_{\mathbf{G}_1\mathbf{G}_2} c^*_{n\mathbf{k}+\mathbf{q}}(\mathbf{G}+\mathbf{G}_1) c_{n_1'\mathbf{k}}(\mathbf{G}_1) c_{n\mathbf{k}+\mathbf{q}}(\mathbf{G}'+\mathbf{G}_2) c^*_{n_2'\mathbf{k}}(\mathbf{G}_2)$$

$$= \sum_{\mathbf{k}} \sum_{\mathbf{G}_1\mathbf{G}_2} c_{n_1'\mathbf{k}}(\mathbf{G}_1) c^*_{n_2'\mathbf{k}}(\mathbf{G}_2) \underbrace{\sum_{n}^{\mathrm{occ}} c^*_{n\mathbf{k}+\mathbf{q}}(\mathbf{G}+\mathbf{G}_1) c_{n\mathbf{k}+\mathbf{q}}(\mathbf{G}'+\mathbf{G}_2)}_{\propto \delta_{\mathbf{G}_1+\mathbf{G},\mathbf{G}_2+\mathbf{G}'} \text{ from approx. completeness}}$$

In $\sum_{n}^{\mathrm{occ}}$: the number of $\mathbf{G}$ vectors making $c_{n\mathbf{k}}(\mathbf{G})$ non-zero isn't much larger than N_{occ}: we have an approximate, unnormalized completeness relation

$$
\begin{aligned}
&\sum_{\mathbf{k}} \sum_{n}^{\text{occ}} M_{nn_1'}(\mathbf{k},\mathbf{q},\mathbf{G}) M^*_{nn_2'}(\mathbf{k},\mathbf{q},\mathbf{G}') \\
&= \sum_{\mathbf{k}} \sum_{n}^{\text{occ}} \sum_{\mathbf{G}_1\mathbf{G}_2} c^*_{n\mathbf{k}+\mathbf{q}}(\mathbf{G}+\mathbf{G}_1) c_{n_1'\mathbf{k}}(\mathbf{G}_1) c_{n\mathbf{k}+\mathbf{q}}(\mathbf{G}'+\mathbf{G}_2) c^*_{n_2'\mathbf{k}}(\mathbf{G}_2) \\
&= \sum_{\mathbf{k}} \sum_{\mathbf{G}_1\mathbf{G}_2} c_{n_1'\mathbf{k}}(\mathbf{G}_1) c^*_{n_2'\mathbf{k}}(\mathbf{G}_2) \underbrace{\sum_{n}^{\text{occ}} c^*_{n\mathbf{k}+\mathbf{q}}(\mathbf{G}+\mathbf{G}_1) c_{n\mathbf{k}+\mathbf{q}}(\mathbf{G}'+\mathbf{G}_2)}_{\propto \delta_{\mathbf{G}_1+\mathbf{G},\mathbf{G}_2+\mathbf{G}'} \text{ from approx. completeness}} \\
&\propto \sum_{\mathbf{G}_1} \sum_{\mathbf{k}} \underbrace{c_{n_1'\mathbf{k}}(\mathbf{G}_1) c^*_{n_2'\mathbf{k}}(\mathbf{G}_1+\mathbf{G}-\mathbf{G}')}_{\text{random phase cancellation}} \propto \delta_{n_1'n_2'}.
\end{aligned}
$$

In $\sum_{n}^{\text{occ}}$: the number of $\mathbf{G}$ vectors making $c_{n\mathbf{k}}(\mathbf{G})$ non-zero isn't much larger than N_{occ}: we have an approximate, unnormalized completeness relation

In $\sum_{\mathbf{k}}$: suppose $\theta_{n\mathbf{k}}$ is the random DFT diagonalization phase of $|\phi_{n\mathbf{k}}\rangle$, $\sum_{\mathbf{k}} \mathrm{e}^{\mathrm{i}(\theta_{n_1'\mathbf{k}} - \theta_{n_2'\mathbf{k}})} \propto \delta_{n_1'n_2'}$

`pseudobands` for Σ

Σ^{CH} revisited

Only Σ^{CH} involves $\sum^{\mathrm{emp}}$:

$$\langle n\mathbf{k}|\Sigma^{\mathrm{CH}}(\omega)|n'\mathbf{k}\rangle = \frac{\mathrm{i}}{2\pi}\sum_{n''}\sum_{\mathbf{qGG'}} M^*_{n''n}(\mathbf{k},-\mathbf{q},-\mathbf{G})M_{n''n'}(\mathbf{k},-\mathbf{q},-\mathbf{G'})\int_0^\infty \frac{[\epsilon^{\mathrm{r}}(\mathbf{q},\omega') - \epsilon^{\mathrm{a}}(\mathbf{q},\omega')]^{-1}_{\mathbf{GG'}}\,\mathrm{d}\omega'}{\omega - E_{n\mathbf{k}} - \omega' + \mathrm{i}\,0^+ \,\mathrm{sgn}(E_{n\mathbf{k}})} v(\mathbf{q}+\mathbf{G'})$$

Problems

- Summation over $\mathbf{G}$ very complicated $\Rightarrow$ analysis based on $\mathbf{G}$ very hard.
- But at high empty bands, $\mathbf{G'}$ dominates $v(\mathbf{q}+\mathbf{G'})$ so the part besides MM^* should be smooth w.r.t. $\mathbf{q}$
- and the summation over $\mathbf{q}$ still leads to the desirable diagonality

– But I don't have my own implementation of `sigma` so that's pure speculation.

Enhancement of pseudobands: Felipe H. Jornada's recent paper

"Stochastic" pseudobands

Energy averaging is also done in their paper:

$$G = \sum_{n \text{ near F.S.}} \frac{|\phi_{n\mathbf{k}}\rangle\langle\phi_{n\mathbf{k}}|}{\omega - E_{n\mathbf{k}}} + \sum_{\text{pseudobands block}} \frac{1}{\omega - \bar{E}} \sum_{n \in \text{block}} |\phi_{n\mathbf{k}}\rangle\langle\phi_{n\mathbf{k}}| \,.$$

Their procedure Replace each pseudobands block by *several bands*:

$$\{\phi_{n\mathbf{k}}\}_{n \text{ in block}} \longrightarrow \left\{ |\phi_{\xi\mathbf{k}}\rangle = \sum_{n \text{ in block}} \mathrm{e}^{\mathrm{i}\,\theta_{n\xi}} \, |\phi_{n\mathbf{k}}\rangle \right\}_{\xi}$$

Justification When $N_\xi \to \infty$, $\left\langle \mathrm{e}^{-\mathrm{i}\,\theta_{n'\xi}} \, \mathrm{e}^{\mathrm{i}\,\theta_{n\xi}} \right\rangle_\xi = \delta_{nn'}$

Comparison with the current script When $N_\xi = 1$, $\mathrm{e}^{\mathrm{i}\,\theta_{n\xi}}$ comes from DFT diagonalization

FHJ's justification of pseudobands

Convergence conditions When $N_{\text{pseudobands blocks}} \to \infty$ and in each block $N_\xi \to \infty$,

- the expectations of χ and Σ goes to the true value;
- the standard error goes to zero.

This convergence condition however is equivalent to not doing pseudobands at all. . . and their recommended values are not large enough anyway.

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This convergence condition however is equivalent to not doing pseudobands at all. . . and their recommended values are not large enough anyway.

They also realized this "the polarizability tends to converge much faster than G, partially due to the *rapidly oscillating nature* of the matrix elements involving Kohn-Sham states used in the evaluation of the polarizability"

As is shown above.

Compressing the bands even harder: size of blocks

Recall that band dispersion is the main problem...

$$\Delta\chi \sim \frac{\Delta E}{E^2}, \quad \chi \sim \frac{1}{E} \Rightarrow \frac{\Delta\chi}{\chi} \simeq \frac{\Delta E}{E}$$
$$\Rightarrow \frac{\Delta\chi}{\chi} \lesssim \text{const} \Leftrightarrow \boxed{\frac{\Delta E}{E} \lesssim \text{const}}$$

Exponential growth of the energy spread of blocks!

Compressing the bands even harder: pseudobands near E_F?

They claim that

- for `epsilon`, bands near Fermi surface can also be pseudo-ized;
- for `sigma`, only the bands near Fermi surface shouldn't be pseudo-ized.

Important claim: protection window is not a convergence parameter: just protect a handful of bands and you're fine.

My own numerical experiments show that this is NOT the case when the band gap is very small so maybe a handful of bands still need to be protected.

Discussion

Discussion

Implications to *GW* acceleration There are lots of garbage in the giant input files we feed to *GW*

Implications to machine learning There exists an *analytic* relation between $\epsilon^{\text{RPA}}, \Sigma^{GW}$ and $\{\bar{E}, |\phi_{\mathbf{k}}\rangle^{\text{pseudo}}\}_{\text{blocks}}$

- Pseudobands (with further compression of autoencoders) should be the starting point of all ML tasks
- Starting from pseudobands is *not* feature engineering