

Bosonic modes in Fermi liquid

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Background

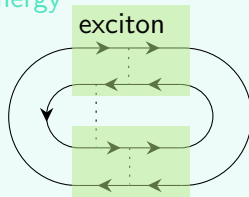
In a Fermi liquid we have ...

- Quasiparticles (electron/hole) with Σ -correction
- Any anything else?

single electron energy



exciton energy



... and more

Question

What to do

Finding modes other than the corrected single electron/hole

Why it's important

Usually not for C_V but for optical response: ϵ , $\chi^{(3)}$, etc.

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Today's topic

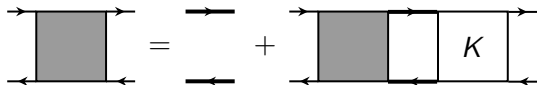
Electron-hole bosonic modes in Fermi liquid (with *some* scattering picked up back, i.e. beyond $\delta E \sim \varepsilon \delta n + f \delta n \delta n$), i.e.

$$|\text{single excitation}\rangle = \sum_{\mathbf{k}_1, \mathbf{k}_2} c_{\mathbf{k}_1 \mathbf{k}_2} \left| \begin{array}{c} \bullet \\ \text{---} \\ \text{---} \end{array} \right\rangle \quad (1)$$

No trion, higher order correlation, or even more exotic spinons, etc.
beyond Fermi liquid

Series calculation

Bethe–Salpeter equation (BSE) is for quantitative calculations.



The diagram shows a Feynman diagram equation. On the left is a single shaded rectangular box with four external lines (two incoming from the left, two outgoing to the right). This is followed by an equals sign. To the right of the equals sign is a sum of two terms. The first term is two parallel horizontal lines, one above and one below, representing a free propagator. The second term is a sum of two diagrams: the first is a shaded rectangular box with four external lines, and the second is a white rectangular box labeled 'K' with four external lines. The entire equation is labeled with a large '(2)' on the far right.

$$\text{Diagram} = \text{Diagram} + \text{Diagram} \quad (2)$$

Problem: no picture about “how the electron moves”

Series calculation

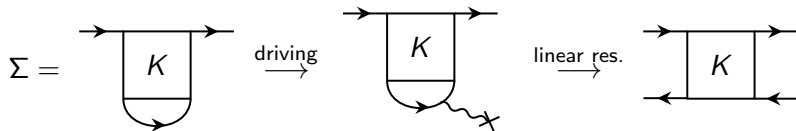
Bethe–Salpeter equation (BSE) is for quantitative calculations.


$$\text{[Shaded Box]} = \text{[Box]} + \text{[Shaded Box]} + \text{[Box K]}$$
 (2)

Problem: no picture about “how the electron moves”

Linking BSE with single-electron kinetic theory

Linear response of single-electron under external field = BSE


$$\Sigma = \text{[Shaded Box with Loop]} \xrightarrow{\text{driving}} \text{[Box with Shaded Bottom and Wavy Line]} \xrightarrow{\text{linear res.}} \text{[Box K]}$$
 (3)

simplest single-electron theory: quantum Boltzmann equation (QBE)

What to investigate

Stable oscillation modes of QBE (\Leftrightarrow infinite response to external field \Leftrightarrow bosonic mode): for $n_{\mathbf{p}\sigma\sigma'}(\mathbf{r})$,

$$\frac{\partial n_{\mathbf{p}}}{\partial t} + \underbrace{\frac{\partial \varepsilon_{\mathbf{p}}}{\partial \mathbf{p}} \cdot \frac{\partial n_{\mathbf{p}}}{\partial \mathbf{r}}}_{\text{diffusion}} - \underbrace{\frac{\partial \varepsilon_{\mathbf{p}}}{\partial \mathbf{r}} \cdot \frac{\partial n_{\mathbf{p}}}{\partial \mathbf{p}}}_{\text{force}} + \underbrace{i[\varepsilon_{\mathbf{p}}, n_{\mathbf{p}}]}_{\text{multi-band}} = \underbrace{I_{\text{Fermi golden rule}}}_{\text{collision}}. \quad (4)$$

What to expect

Three types of important bosonic modes:

- Zero sound in uncharged single-band Fermi liquid
- Plasmon in charged single-band Fermi liquid = zero sound + long range interaction
- Exciton in charged multi-band Fermi liquid

Equation governing zero sound

System Single-band Fermi liquid with spin ignored

Kinetics of uncharged Fermi liquid *Landau equation* = QBE +

$$\varepsilon_{\mathbf{p}}(\mathbf{r}) = \varepsilon_{\mathbf{p}}^0 + \frac{1}{V} \sum_{\mathbf{p}'} f_{\mathbf{p}\mathbf{p}'} \delta n_{\mathbf{p}}(\mathbf{r}) \quad (5)$$

(assumption: $\mathbf{q} \rightarrow 0$ in $c_{\mathbf{p}+\mathbf{q}}^\dagger c_{\mathbf{p}}$, i.e. $\delta n_{\mathbf{p}}(\mathbf{r})$ being smooth in \mathbf{r})

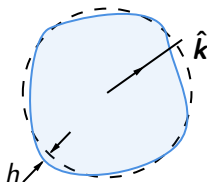
EOM governing zero sound Small disturbance, no collision, :

$$\frac{\partial \delta n_{\mathbf{p}}}{\partial t} + \frac{\partial \varepsilon_{\mathbf{p}}^{\text{static}}}{\partial \mathbf{p}} \cdot \frac{\partial \delta n_{\mathbf{p}}}{\partial \mathbf{r}} - \underbrace{\frac{\partial n_{\mathbf{p}}^{\text{static}}}{\partial \mathbf{p}} \cdot \frac{1}{V} \sum_{\mathbf{p}'} f_{\mathbf{p}\mathbf{p}'} \frac{\partial \delta n_{\mathbf{p}}}{\partial \mathbf{r}}}_{\partial \varepsilon_{\mathbf{p}} / \partial \mathbf{r}} = 0 \quad (6)$$

Fermi surface vibration

Ansatz Disturbance as small as possible ...

$$n_{\mathbf{p}}(\mathbf{r}, t) = e^{i(\mathbf{q} \cdot \mathbf{r} - i\omega t)} \theta(\mu - \varepsilon_{\mathbf{p}}^{\text{stable}} - h(\hat{\mathbf{p}})) \quad (7)$$

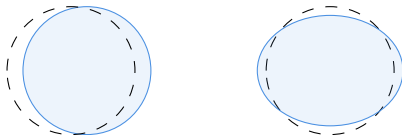


Eigenvalue problem

$$(\omega - \mathbf{q} \cdot \mathbf{v})h(\hat{\mathbf{k}}) = \mathbf{q} \cdot \mathbf{v} \int \frac{d\Omega'}{4\pi} F(\vartheta)h(\hat{\mathbf{k}}'). \quad (8)$$

where \mathbf{v} is single-electron velocity. \Rightarrow zero sound has linear dispersion;
zero sound requires $F \neq 0$

Shape of Fermi surface



Zero sound is not density wave In zero sound $V_{\text{Fermi sea}} = \text{const.} \Rightarrow$
zero sound is not ordinary sound

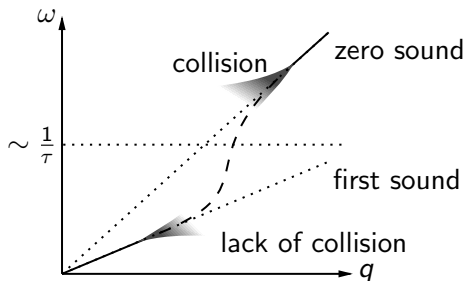
Comparison with ordinary sound

Ordinary sound Fermi liquid theory $\Rightarrow \partial\rho/\partial P \Rightarrow$ another sound mode (“first sound”, ordinary sound, density mode) from hydrodynamics

Relation with zero sound

- First sound appears when $\omega\tau \ll 1$: ordinary hydrodynamics \Leftrightarrow local equilibrium $\Leftrightarrow \tau \ll 1/\omega$
- zero sound appears when $\omega\tau \gg 1$: no collision integral $\Leftrightarrow \tau \gg 1/\omega$

The two are connected: a radical finite- T correction



Fermi liquid, uncharged: zero sound

- Linear, gapless
- From $f_{pp'}$

Fermi liquid, charged: plasmon

- Divergent Hartree term \Rightarrow self-energy correction in real space
- When $\mathbf{q} = 0$: $f_{pp'}$ not important; gapped

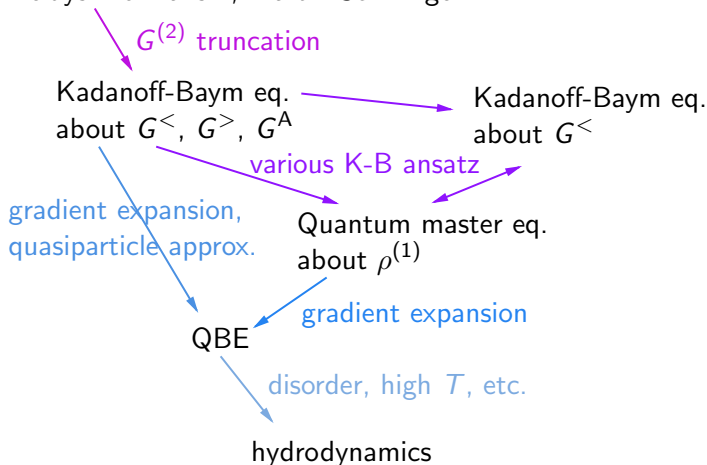
Two bands: exciton

Justifying quantum Boltzmann equation

Is QBE reliable?

Yes! When we intuitively expect it to work –

Keldysh formalism, Martin-Schwinger



Discussion