

# Bloch equation and rate equation in two-level systems

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Consider Eq. (5.125) in Steck's lecture notes. The optical Bloch equations can be derived in a diagrammatic way. Basically, we have the following diagrams: for the  $\rho_{ee}$  we have (possibly the directions of the lines are wrong; check later)

$$\begin{array}{c} \text{e} \longrightarrow \text{e} \xrightarrow{\text{g}} \text{e} \\ \text{wavy line} \\ \times \end{array} \simeq \frac{1}{\partial_t} \frac{\Omega}{2} \rho_{eg}, \quad \begin{array}{c} \text{e} \longleftarrow \text{e} \xleftarrow{\text{g}} \text{e} \\ \text{wavy line} \\ \times \end{array} \simeq -\frac{1}{\partial_t} \frac{\Omega}{2} \rho_{ge}, \quad (1)$$

and for the  $\rho_{eg}$  part we have

$$\begin{array}{c} \text{e} \longrightarrow \text{e} \xrightarrow{\text{g}} \text{g} \\ \text{wavy line} \\ \times \end{array} \simeq \frac{1}{\partial_t} \frac{\Omega}{2} \rho_{gg}, \quad \begin{array}{c} \text{e} \longleftarrow \text{e} \xleftarrow{\text{g}} \text{g} \\ \text{wavy line} \\ \times \end{array} \simeq \frac{1}{\partial_t} \frac{\Omega}{2} \rho_{ee}, \quad (2)$$

and also

$$\text{e} \longrightarrow \text{g} \simeq \frac{1}{\partial_t} \Delta \rho_{eg}. \quad (3)$$

Here we need to note that we are dealing with lesser Green functions, and therefore the free  $G_{gg}^<$  or  $G_{ee}^<$  is just  $1/i\partial_t$ ; on the other hand, we have (3). This can be illustrated by the following diagrams:

$$\begin{array}{c} \text{e} \longrightarrow \text{g} \longrightarrow \text{g} \\ + \quad - \quad - \end{array}, \quad \begin{array}{c} \text{e} \longrightarrow \text{e} \longrightarrow \text{g} \\ + \quad + \quad - \end{array} \quad (4)$$

and we can see that the diagrams for  $G_{eg}^<$ , when the external driving field is turned off, contains “ordinary” time-ordered and anti-time-ordered Green functions, and after the summation, the difference between the energies of the e state and the g state enters the self-energy of  $G_{eg}^<$ . We can do the same thing for  $G_{ee}^<$  or  $G_{gg}^<$  but this time we just get

$$\begin{array}{c} \text{e} \longrightarrow \text{e} \longrightarrow \text{e} \\ + \quad - \quad - \end{array}, \quad \begin{array}{c} \text{e} \longrightarrow \text{e} \longrightarrow \text{e} \\ + \quad + \quad - \end{array}, \quad (5)$$

and we find that the contributions of time-ordered and anti-time-ordered Green functions just cancel each other. It should also be noted that  $\Delta$  itself contains corrections from external fields and is the energy gap minus the external phonon frequency.

Now we turn to how to obtain rate equations – equations where the transitions between e and g are modeled as “scattering” – from the aforementioned optical Bloch equation, or more generally, quantum master equations. The idea is just to first replace  $G_{eg}^<$  by diagrams in terms of  $G_{ee,gg}^<$ , and then ignore any frequency dependence in the internal components (similar to COHSEX). TODO: possible to do something like GPP?