

# Homework 4

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## Exercise 3 in lecture 12

**Solution** The normalization factor in the time domain is

$$\int_{-\infty}^{\infty} |g(t)|^2 dt = \sigma\sqrt{\pi}. \quad (1)$$

The expectation values are

$$\langle t^2 \rangle = \frac{\int_{-\infty}^{\infty} dt |g(t)|^2 t^2}{\int_{-\infty}^{\infty} dt |g(t)|^2} = t_0^2 + \frac{1}{2}\sigma^2, \quad (2)$$

and

$$\langle t \rangle = \frac{\int_{-\infty}^{\infty} dt |g(t)|^2 t}{\int_{-\infty}^{\infty} dt |g(t)|^2} = t_0. \quad (3)$$

The frequency domain version of  $g$  is

$$g[\omega] = \int_{-\infty}^{\infty} e^{i\omega t} g(t) dt = \sigma\sqrt{2\pi} e^{i(\omega-\omega_0)t - \frac{1}{2}\sigma^2(\omega-\omega_0)^2}. \quad (4)$$

The normalization factor is

$$\int_{-\infty}^{\infty} |g[\omega]|^2 d\omega = 2\pi\sigma^2 \cdot \sqrt{\frac{\pi}{\sigma^2}} = 2\pi^{3/2}\sigma. \quad (5)$$

The expectation values are

$$\langle \omega^2 \rangle = \frac{\int_{-\infty}^{\infty} d\omega |g[\omega]|^2 \omega^2}{\int_{-\infty}^{\infty} d\omega |g[\omega]|^2} = \frac{1}{2\sigma^2} + \omega_0^2, \quad (6)$$

and

$$\langle \omega \rangle = \omega_0. \quad (7)$$

So

$$\begin{aligned} \sigma_t \sigma_\omega &= \sqrt{\langle t^2 \rangle - \langle t \rangle^2} \sqrt{\langle \omega^2 \rangle - \langle \omega \rangle^2} \\ &= \sqrt{\frac{1}{2}\sigma^2 \cdot \frac{1}{2\sigma^2}} = \frac{1}{2}. \end{aligned} \quad (8)$$

So indeed for  $g(t)$ ,  $\sigma_t \sigma_\omega$  reaches its minimum value.

**Exercise 1 in lecture 13** We have the following transmission model: the sent signal  $X$  has 3 possible values  $x = 0, \pm 1$  with  $p_0 = p_1 = p_{-1} = 1/3$ . The noise  $\Xi$  has two possible values  $\zeta = \pm 1$ . We suppose the channel function  $F$  has the following expression  $Y = \text{Im} [e^{i2\pi(X+\Xi)/}]$ . What are all the possible values of  $Y$ ? Give the expression of the distributions  $w(x), w(y), w(x, y)$  and  $w(y | x)$ .

**Solution** The possible results are listed in Table 1. The possible values of  $Y$  are 0 and  $\pm\sqrt{3}/2$ . From Table 1 we have

$$w(x) = \begin{cases} \frac{1}{3}, & x = 0, \\ \frac{1}{3}, & x = 1, \\ \frac{1}{3}, & x = -1, \end{cases} \quad (9)$$

$$w(y) = \begin{cases} \frac{1}{3}, & y = 0, \\ \frac{1}{3}, & y = \sqrt{3}/2, \\ \frac{1}{3}, & y = -\sqrt{3}/2, \end{cases} \quad (10)$$

Table 1: Probabilistic distribution of  $Y$ ; the probability of each row is  $1/6$ .

$X$	$\Xi$	$Y = \text{Im} e^{i2\pi(X+\Xi)/3}$
0	1	$\frac{\sqrt{3}}{2}$
	-1	$-\frac{\sqrt{3}}{2}$
1	1	$-\frac{\sqrt{3}}{2}$
	-1	0
-1	1	0
	-1	$\frac{\sqrt{3}}{2}$

and

$$w(x, y) = \begin{cases} \frac{1}{6}, & (x, y) = (0, \sqrt{3}/2), \\ \frac{1}{6}, & (x, y) = (0, -\sqrt{3}/2), \\ \frac{1}{6}, & (x, y) = (1, -\sqrt{3}/2), \\ \frac{1}{6}, & (x, y) = (1, 0), \\ \frac{1}{6}, & (x, y) = (-1, 0), \\ \frac{1}{6}, & (x, y) = (-1, \sqrt{3}/2), \\ 0, & \text{otherwise.} \end{cases} \quad (11)$$

So

$$w(y|x) = \frac{w(x, y)}{w(x)} = \begin{cases} \frac{1}{2}, & (x, y) = (0, \sqrt{3}/2), \\ \frac{1}{2}, & (x, y) = (0, -\sqrt{3}/2), \\ \frac{1}{2}, & (x, y) = (1, -\sqrt{3}/2), \\ \frac{1}{2}, & (x, y) = (1, 0), \\ \frac{1}{2}, & (x, y) = (-1, 0), \\ \frac{1}{2}, & (x, y) = (-1, \sqrt{3}/2), \\ 0, & \text{otherwise.} \end{cases} \quad (12)$$

**Exercise 2 in lecture 14**

**Solution** Suppose  $p$  is the probability of  $X = -1$ , and we have

$$m = \langle X \rangle = -p + (1 - p) = 1 - 2p, \quad -1 \leq m \leq 1, \quad (13)$$

and therefore

$$p = \frac{1 - m}{2}, \quad (14)$$

and

$$\begin{aligned} H_\alpha(X) &= \frac{1}{1 - \alpha} \log_2(p^\alpha + (1 - p)^\alpha) \\ &= \frac{1}{1 - \alpha} \log_2 \left( \left( \frac{1 - m}{2} \right)^\alpha + \left( \frac{1 + m}{2} \right)^\alpha \right). \end{aligned} \quad (15)$$

The plots can be found in Figure 1.

**Exercise 4 in lecture 14**

**Solution** We have

$$\begin{aligned} H(Y|X) &= \sum_x H(Y, X = x)p(x) \\ &= - \sum_x p(x) \sum_y p(y|x) \log_2 p(y|x) \\ &= - \sum_{x,y} p(x, y) \log_2 p(y|x), \end{aligned} \quad (16)$$

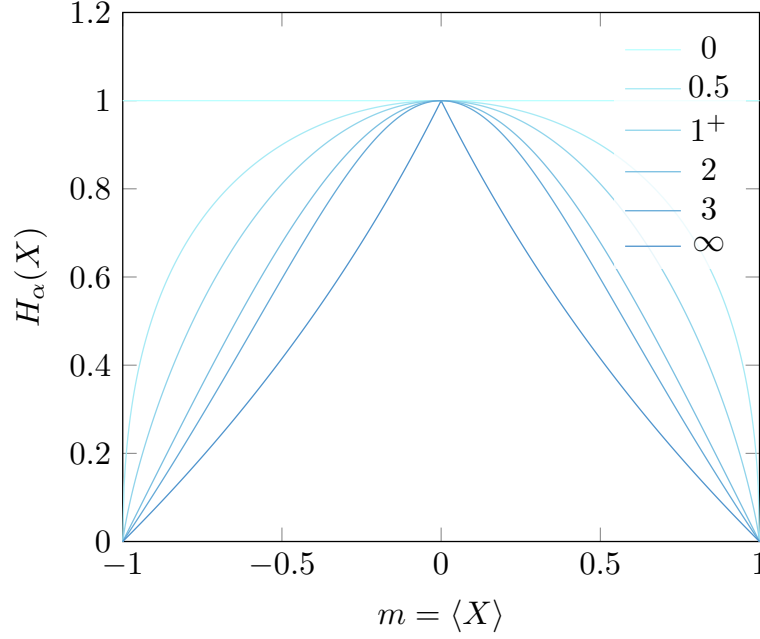


Figure 1: Plots of  $H_\alpha(X)$  with different  $\alpha$ 's

and a direct verification gives

$$\begin{aligned}
H(X, Y) &= - \sum_{x,y} p(x, y) \log_2 \underbrace{p(x, y)}_{p(x)p(y|x)} \\
&= - \sum_{x,y} p(x, y) (\log_2 p(x) + \log_2 p(y|x)) \\
&= - \sum_x p(x) \log_2 p(x) - \sum_{x,y} p(x, y) \log_2 p(y|x) \\
&= H(X) + H(Y|X).
\end{aligned} \tag{17}$$

**Problem** Differences between the classical and quantum entropies

**Solution**

(a) This can be done by Schmidt decomposition. Suppose  $\dim \mathcal{H}_A = m$ ,  $\dim \mathcal{H}_B = n$ , and without loss of generality  $m \geq n$ , then we can already find a set of basis, in which

$$|\psi_{AB}\rangle = \sum_{i=1}^m c_i |u_A^i\rangle \otimes |v_B^i\rangle. \tag{18}$$

We have

$$\begin{aligned}
\rho_A &= \sum_{j=1}^m \langle v_B^j | \psi_{AB} \rangle \langle \psi_{AB} | v_B^j \rangle \\
&= \sum_{j=1}^m c_j |u_A^j\rangle c_j^* \langle u_A^j| \\
&= \sum_{j=1}^m |c_j|^2 |u_A^j\rangle \langle u_A^j|,
\end{aligned} \tag{19}$$

and similarly (note that  $\{|v_B^j\rangle\}_{j=1}^m$  appearing in the expansion of  $|\psi_{AB}\rangle$  is over-complete, and the form of  $\rho_B$  is exactly the same as  $\rho_A$ )

$$\rho_B = \sum_{i=1}^m |c_i|^2 |v_B^i\rangle \langle v_B^i|. \tag{20}$$

So we find

$$S(\rho_A) = S(\rho_B) = - \sum_{i=1}^m |c_i|^2 \ln |c_i|^2. \quad (21)$$

Note that here

$$S(\rho_{AB}) = 0, \quad (22)$$

because we are dealing with a pure state. So here the entropy of a system is less than any of its parts, and this is not possible in the classical case.

(b) Let

$$\rho = \rho_{AB}, \quad \sigma = \rho_A \otimes \rho_B,$$

and from

$$\text{tr}(\rho \ln \rho) - \text{tr}(\rho \ln \sigma) \geq 0,$$

we have

$$\begin{aligned} 0 &\leq \text{tr}(\rho_{AB} \ln \rho_{AB}) - \text{tr}(\rho_{AB} \underbrace{\ln \rho_A \otimes \rho_B}_{\ln \rho_A + \ln \rho_B}) \\ &= -S(\rho_{AB}) - \text{tr}_A \text{tr}_B \rho_{AB} \ln \rho_A - \text{tr}_B \text{tr}_A \rho_{AB} \ln \rho_B \\ &= -S(\rho_{AB}) + S(\rho_A) + S(\rho_B), \end{aligned}$$

and thus

$$S(\rho_A) + S(\rho_B) - S(\rho_{AB}) \geq 0. \quad (23)$$

(c) What we need to prove

$$S(\rho_{AB}) \geq S(\rho_B) - S(\rho_A). \quad (24)$$

Assume the density matrix  $\rho_{AB}$  is actually the partial trace of the pure state  $|\Psi_{ABC}\rangle$ . (23) means

$$S(\rho_B) + S(\rho_C) \geq S(\rho_{BC}),$$

and from (a), we have

$$S(\rho_C) = S(\rho_{AB}), \quad S(\rho_{BC}) = S(\rho_A),$$

and therefore

$$\begin{aligned} S(\rho_B) + S(\rho_{AB}) &\geq S(\rho_A), \\ S(\rho_{AB}) &\geq S(\rho_A) - S(\rho_B). \end{aligned} \quad (25)$$

Similarly

$$S(\rho_{AB}) \geq S(\rho_B) - S(\rho_A). \quad (26)$$

(d) We have

$$p_j = \langle j | \rho_B | j \rangle, \quad (27)$$

and therefore

$$\rho_{AB|\Pi_B^j} = \frac{1}{p_j} \text{tr}_B |j\rangle\langle j| \langle j | \rho_{AB} | j \rangle = \frac{1}{p_j} \langle j | \rho_{AB} | j \rangle. \quad (28)$$

So

$$\begin{aligned} \sum_j p_j S(\rho_{AB|\Pi_B^j}) &= -\text{tr}_A \sum_j \langle j | \rho_{AB} | j \rangle \ln \frac{\langle j | \rho_{AB} | j \rangle}{p_j} \\ &= -\text{tr}_A \sum_j \langle j | \rho_{AB} | j \rangle \ln \langle j | \rho_{AB} | j \rangle + \sum_j \langle j | \rho_B | j \rangle \ln \langle j | \rho_B | j \rangle \\ &= S(\rho_{AB}^D) - S(\rho_B^D). \end{aligned} \quad (29)$$

The inequality to be proven is equivalent to

$$S(\rho_A) - S(\rho_A|\rho_B) \leq S(\rho_A) - H(\rho_A|\rho_B),$$

which further is equivalent to

$$S(\rho_A|\rho_B) = \min_{\{\Pi_B^j\}} \sum_j p_j S(\rho_{AB|\Pi_B^j}) \geq H(\rho_A|\rho_B). \quad (30)$$

(29) and

$$S(\rho_{AB}^D) - S(\rho_B^D) \geq H(\rho_A|\rho_B) \quad (31)$$

show that for any  $\{\Pi_B^j\}$ , the inequality is correct, so we have already finished the proof.

(e) The disagreements in the definition of conditional entropy intuitively arises from the entanglement between the two systems, so we expect when

$$\rho_{AB} = \rho_A \otimes \rho_B, \quad (32)$$

there is no inconsistency in the definition of conditional entropy, which is the classical case. This can be verified directly, because when (32) is correct, we have

$$\rho_{AB}^D = \rho_{AB}, \quad \rho_B^D = \rho_B,$$

and therefore according to (29), we have

$$S(\rho_A|\rho_B) = S(\rho_{AB}) - S(\rho_B), \quad (33)$$

and

$$J(\rho_A : \rho_B) = S(\rho_A) + S(\rho_B) - S(\rho_{AB}) =: I(\rho_A : \rho_B). \quad (34)$$

(f) In a pure state, we have

$$I(\rho_A|\rho_B) = S(\rho_A) + S_{\rho_B} - S_{\rho_{AB}} = 2S(\rho_A). \quad (35)$$

On the other hand, we have

$$\rho_{AB|\Pi_B^j} = \frac{1}{p_j} \langle j|\psi_{AB}\rangle \langle \psi_{AB}|j\rangle. \quad (36)$$

So now we can use Schmidt decomposition again: We have (18), and therefore (here  $|j\rangle$  means a basis vector for B, and  $|j\rangle_A$  is used to refer to a basis vector for A)

$$\langle j|\psi_{AB}\rangle = c_i |j\rangle_A,$$

so  $\rho_{AB|\Pi_B^j}$  is actually a pure state, and therefore we have  $S(\rho_A|\rho_B) = 0$ , because each term in it is zero. So

$$J(\rho_A : \rho_B) = S(\rho_A), \quad (37)$$

and thus

$$\delta(\rho_A : \rho_B) = 2S(\rho_A) - S(\rho_A) = S(\rho_A) = S(\rho_B). \quad (38)$$