

Time-dependent adiabatic GW

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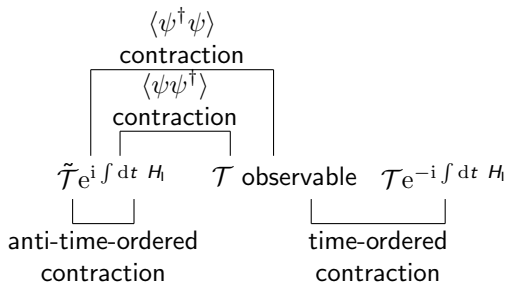
Non-equilibrium Green function

Motivation

$$\langle A \rangle = \langle S^{-1} \mathcal{T}_t(S A_I(t)) \rangle, \quad S = U(\infty, -\infty) \quad (1)$$

Non-equilibrium state: not pure; contains excited state components;
 $|\Psi_n\rangle$ is excited state $\Rightarrow S |\Psi_n\rangle \neq e^{i\alpha} |\Psi_n\rangle \Rightarrow$ we can't peel the S^{-1} off!!

Solution Four (instead of one) types of propagators: (note S^{-1} is *anti*-time ordered)

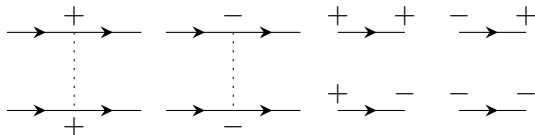


Keldysh formalism

Four types of (fermionic) propagators

$$\begin{aligned} iG^{--} = iG^c &= \langle \mathcal{T} \psi_1 \psi_2^\dagger \rangle, & iG^{++} = iG^a &= \langle \tilde{\mathcal{T}} \psi_1 \psi_2^\dagger \rangle, \\ iG^{+-} = iG^> &= \langle \psi_1 \psi_2^\dagger \rangle, & iG^{-+} = iG^< &= -\langle \psi_2^\dagger \psi_1 \rangle. \end{aligned} \quad (2)$$

Diagrams

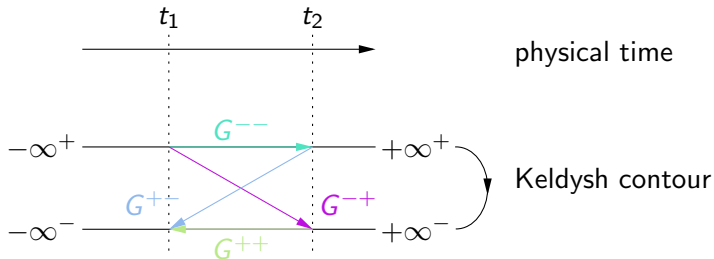


Self-energy

$$G = \begin{pmatrix} G^{--} & G^{-+} \\ G^{+-} & G^{++} \end{pmatrix}, \quad \Sigma = \begin{pmatrix} \Sigma^{--} & \Sigma^{-+} \\ \Sigma^{+-} & \Sigma^{++} \end{pmatrix}, \quad G = G_0 + G_0 \Sigma G. \quad (3)$$

Alternative formulation: Keldysh contour

Keldysh contour The information in the G matrix can be alternatively stored in a time-ordered Green function on *Keldysh contour*



From Keldysh contour to physical contour Lengreth theorem:

$$\begin{aligned}(AB)^{<} &= A^R B^{<} + A^{<} B^A, & (AB)^{>} &= A^R B^{>} + A^{>} B^A, \\ (AB)^R &= A^R B^R, & (AB)^A &= A^A B^A,\end{aligned}\tag{4}$$

where

$$\begin{aligned}A^{>}(t_1, t_2) &= A(t_1^+, t_2^-), & A^{<}(t_1, t_2) &= A(t_1^-, t_2^+), \\ A^R(t_1, t_2) &= \theta(t_1 - t_2)(A^{>} - A^{<}).\end{aligned}\tag{5}$$

Mapping an equation on Keldysh contour to its counterpart on the physical time axis!

Derivation of EOM of $G^{<, >}$ and G^A I

Recommended references The following series:

- Václav Špička, Bedřich Velický, and Anděla Kalvová. “Long and short time quantum dynamics: I. Between Green’s functions and transport equations”. In: *Physica E: Low-dimensional Systems and Nanostructures* 29.1-2 (2005), pp. 154–174
- Jørgen Rammer and H Smith. “Quantum field-theoretical methods in transport theory of metals”. In: *Reviews of modern physics* 58.2 (1986), p. 323

Derivation of EOM of $G^{<,>}$ and G^A II

From self-energy correction to EOM From Lengreth theorem:

$$G = G_0 + G_0 \Sigma G \Rightarrow G^{<} = G_0^{<} + G_0^{<} \Sigma^A G^A + G_0^R \Sigma^R G^{<} + G_0^R \Sigma^{<} G^A, \quad (6)$$

$$G = G_0 + G \Sigma G_0 \Rightarrow G^{<} = G_0^{<} + G_0^R \Sigma^R G_0^{<} + G^R \Sigma^{<} G_0^A + G^{<} \Sigma^A G^A, \quad (7)$$

$$G^A = G_0^A + G_0^A \Sigma^A G^A, \quad G^R = G_0^R + G_0^R \Sigma^R G^R. \quad (8)$$

Getting rid of G_0 We define

$$G_0^{-1} := i \partial_t - H_0, \quad (9)$$

and

$$G_0^{-1} G_0^{A,R} = I, \quad G_0^{-1} G_0^{<,>} = 0. \quad (10)$$

Taking complex conjugate of the def. of $G_0^{<,>}$ we find (left arrow = apply ∂_t and H_0 to the second index of $G_0^{<,>}$)

$$G_0^{<,>} (-i \overleftarrow{\partial}_{t_2} - H_0) = 0. \quad (11)$$

Derivation of EOM of $G^{<,>}$ and G^A III

The Schrödinger-like EOM Applying G_0^{-1} to the left of (6) and to the right of (7):

$$(i \partial_{t_1} - H_0) G^{<}(1, 2) = \Sigma^R G^{<} + \Sigma^{<} G^A, \quad (12)$$

$$-i \partial_{t_2} G^{<}(1, 2) - G^{<} H_0 = G^R \Sigma^{<} + G^{<} \Sigma^A, \quad (13)$$

$$\Rightarrow i(\partial_{t_1} + \partial_{t_2}) G^{<} - [H_0, G^{<}] = \Sigma^R G^{<} + \Sigma^{<} G^A - G^R \Sigma^{<} - G^{<} \Sigma^A. \quad (14)$$

Mixed coordinates We define “average time” and “relative time”:

$$T = \frac{t_1 + t_2}{2}, \quad t = t_1 - t_2, \quad (15)$$

$$\Rightarrow \frac{\partial}{\partial T} = \frac{\partial}{\partial t_1} + \frac{\partial}{\partial t_2}. \quad (16)$$

We then do Fourier transform over t : similar to the equilibrium case. ($T \simeq$ driving, $t \simeq$ internal time evolution)

Towards a single-time formalism

Summary up to now

- Accurate EOMs about $G^{A,R}$, and EOM of $G^<$:

$$i \partial_T G^< - [H_0, G^<] = \Sigma^R G^< + \Sigma^< G^A - G^R \Sigma^< - G^< \Sigma^A. \quad (17)$$

The RHS contains t (or ω) and $G^<$.

- Note: we can actually put the $t = 0$ part of Σ into H_0 ! \Rightarrow Example: COHSEX TD-aGW

Goal Obtaining quantum kinetics:

- Quantum master equation (QME), i.e. EOM of $\rho(\mathbf{r}_1, \mathbf{r}_2, t)$,
- and its long wave length limit, the quantum Boltzmann equation (QBE)

Problem Both LHS and RHS contain ω : problem too large.

What we want Obtaining a close form EOM about $G^<(T, t = 0)$

Reduced density matrix Single-electron density matrix:

$$i\rho(T) = G^<(T, t=0) = \int \frac{d\omega}{2\pi} G^<(T, \omega) \quad (18)$$

What we want Two types of reduction:

- Reducing Σ to an easy function of G , ideally $G^<$
- Reducing $G^<$ to $\rho(T)$

Reducing Σ

Gradient expansion: from QME to QBE

A radical move: quantum Boltzmann equation

Approximations leading to QBE

- Gradient expansion: smooth U_{ext} :

$$[H_0, \rho] \quad (19)$$

- Quasiparticle approx.

$$G^<(\mathbf{x}, \mathbf{p}, T, \omega) = 2\pi\delta(\omega -) \quad (20)$$

Immediate problem:

TODO: how to get the Fermi golden rule???

Example: TODO