

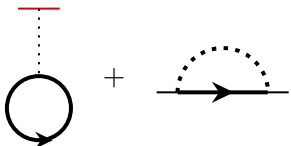
Details in GW-BSE

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Infinitesimal

We all know the word “GW” means that $\Sigma = i GW$ (of course we have Hartree term but it’s already in DFT)

$$\Sigma = \text{[Diagram 1]} + \text{[Diagram 2]}, \quad (1)$$


where W is the RPA-screened potential.

Why some say $\Sigma(1, 2) = i G(1, 2)W(1^+, 2)$?

- $G(1, 2)$ is actually $G(1, 2^+)$ (so when $1 = 2$, $G = n_{\text{occ}}$: the loop in the Hartree term above)
- $\Sigma(1, 2) = i G(1, 2^+)W(1, 2) = i G(1^-, 2)W(1, 2) = i G(1, 2)(1^+, 2)$.
- 1^+ or $2^+ \Leftrightarrow e^{\pm i\omega 0^+} \Leftrightarrow$ how to take contour

Other tricky details in diagrammatics

Time-reversal symmetry

- $W(-\mathbf{p}, -\omega) = W(\mathbf{p}, \omega)$ is always true (or otherwise we can symmetrize the Lagrangian)
- The real symmetry:

$$\begin{aligned} W(\omega, -\mathbf{k}) = W(\omega, \mathbf{k}) &\Leftrightarrow W(-\omega, \mathbf{k}) = W(\omega, \mathbf{k}) \\ &\Leftrightarrow W(\mathbf{r}, \mathbf{r}', \omega) = W(\mathbf{r}', \mathbf{r}, \omega) \Leftrightarrow W(\mathbf{r}, \mathbf{r}', \omega) = W(\mathbf{r}, \mathbf{r}', -\omega). \end{aligned} \quad (2)$$

Imaginary unit

$$iG = iG_0 + \underbrace{iG_0 \times \text{hole} \times iG}_{-i\Sigma} \Rightarrow G = \frac{1}{\omega - E^0 - \Sigma}. \quad (3)$$

“Antiparticles” You can treat holes as antiparticles (negative energy, $i\text{sgn}(\xi_{nk})$ in time-ordered Green function) but then corresponding electron modes have to be ignored.

Feynman rules I

Recall that we are working in a crystal – we need to talk about \mathbf{G} vectors
One set of rules that work:

- Propagator:

$$\begin{array}{c} n, k \\ \longrightarrow \end{array} = \frac{i}{\omega - \xi_{n\mathbf{k}} + i0^+ \text{sgn}(\omega)} =: i G_{n\mathbf{k}}^0(\omega). \quad (4)$$

- Interaction:

$$\begin{array}{c} q, \mathbf{G} \\ \cdots \cdots \cdots \end{array} = -i \frac{1}{V} v(\mathbf{q} + \mathbf{G}). \quad (5)$$

But the prefactor of the interaction Hamiltonian is still $1/2V$, and

$$v(\mathbf{q}) = \int d^3\mathbf{r} e^{-i\mathbf{q}\cdot\mathbf{r}} v(\mathbf{r}). \quad (6)$$

Feynman rules II

- For internal lines, sum over $\mathbf{k}, n, \mathbf{G}$; no additional normalization factors are needed.
- For external lines, outgoing lines are $\phi_{n\mathbf{k}}(\mathbf{r})$, ingoing lines are $\phi_{n\mathbf{k}}^*(\mathbf{r})$, as in:

$$G(\mathbf{r}, \mathbf{r}', \omega) = \sum_{n, \mathbf{k}} \frac{\phi_{n\mathbf{k}}(\mathbf{r}) \phi_{n\mathbf{k}}(\mathbf{r})^*}{\omega - \xi_{n\mathbf{k}} + i \operatorname{sgn}(\xi_{n\mathbf{k}})}, \quad (7)$$

where \mathbf{r} is the outgoing index and \mathbf{r}' is the ingoing index.

GW without G