

Homework 3

Jinyuan Wu

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1.1

For the free electron gas, we have

$$\mu = \frac{1}{2m} \left(\frac{3\pi^2 N}{V} \right)^{2/3}, \quad (1)$$

where we have considered the fact that electrons are spin-1/2 particles. The free electron gas Green function is

$$G_\sigma(\mathbf{k}, \omega) = \frac{1}{\omega - \frac{\mathbf{k}^2}{2m} + \mu + i \operatorname{sgn}(\omega) 0^+}. \quad (2)$$

The imaginary part is therefore

$$\begin{aligned} A_\sigma(\mathbf{k}, \omega) &= -\frac{1}{\pi} \operatorname{Im} G_\sigma(\mathbf{k}, \omega) = -\frac{1}{\pi} \cdot \operatorname{Im}(-\pi i) \operatorname{sgn}(\omega) \delta\left(\omega - \frac{\mathbf{k}^2}{2m} + \mu\right) \\ &= \operatorname{sgn}\left(\frac{\mathbf{k}^2}{2m} - \mu\right) \delta\left(\omega - \frac{\mathbf{k}^2}{2m} + \mu\right). \end{aligned} \quad (3)$$

This function is not really analytical, since its peak has infinitesimal width and infinite height; but adding a small imaginary part to the energy of a single electron, the spectral function should be analytical.

1.2

When

$$\Sigma = \lambda\omega^2 + i(\alpha\omega^2 + \beta T^2), \quad (4)$$

we have

$$\begin{aligned} G_\sigma(\mathbf{k}, \omega) &= \frac{1}{\omega - \frac{\mathbf{k}^2}{2m} - \Sigma + i \operatorname{sgn}(\omega) 0^+} \\ &= \frac{1}{(1-\lambda)\omega - \frac{\mathbf{k}^2}{2m} + \mu - i(\alpha\omega^2 + \beta T^2)} \\ &= \frac{(1-\lambda)\omega - \frac{\mathbf{k}^2}{2m} + \mu + i(\alpha\omega^2 + \beta T^2)}{\left((1-\lambda)\omega - \frac{\mathbf{k}^2}{2m} + \mu\right)^2 + (\alpha\omega^2 + \beta T^2)^2}, \end{aligned} \quad (5)$$

where the infinitesimal imaginary part can be ignored since we already have a finite imaginary part. The imaginary part gives

$$A_\sigma(\mathbf{k}, \omega) = -\frac{1}{\pi} \frac{\alpha\omega^2 + \beta T^2}{\left((1-\lambda)\omega - \frac{\mathbf{k}^2}{2m} + \mu\right)^2 + (\alpha\omega^2 + \beta T^2)^2}. \quad (6)$$

When ω is already given and is a real number, \mathbf{k} should satisfy

$$(1-\lambda)\omega - \frac{\mathbf{k}^2}{2m} + \mu = 0 \Rightarrow \omega = \frac{\mathbf{k}^2}{2m^*} - \mu^*, \quad m^* = (1-\lambda)m \quad (7)$$

so that $A_\sigma(\mathbf{k}, \omega)$ is maximized. This is the effective dispersion relation measured by ARPES; the effective mass is $(1-\lambda)m$. If ω is complex, the maximizing procedure will be much more tedious, but this isn't physical, since in ARPES the frequency of the input beam can't be complex.