# **Cavity QED**

Quantum light-matter interaction to the extreme

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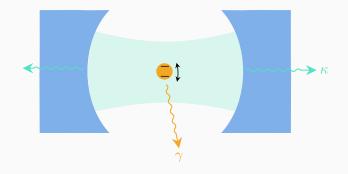
## When do all approximate theories fail?

Maxwell-Langevin equation or atomic quantum master equation high coherence quantum nature of light strong  $\omega$  dependence, strong coupling Classical light+quantum atom Effective dielectric function. scattering matrix

One scenario: in a cavity.

# Cavity and one atom

## High coherence in cavity QED



### Coupling with the environment

- ullet Cavity leaking  $\kappa$
- $\bullet$  Atomic spontaneous emission rate (outside the cavity)  $\gamma$
- (Possible non-radiative decay: phonon, etc.)

Strong coupling limit  $\kappa, \gamma \ll d \cdot E_{\text{cavity}} \Rightarrow \text{high coherence}$ 

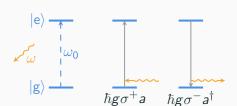
## Frequency selectivity and Jaynes-Cummings model

### Approximations already

- No atom-atom interaction
- No damping at all

#### **RWA**

- Rotating-wave approx.
- Single active photon mode



$$H^{ ext{Jaynes-Cummings}} = \hbar\omega\left(a^{\dagger}a + rac{1}{2}
ight) + rac{\hbar\omega_0}{2}\sigma^z + \hbar g(a\sigma^+ + a^{\dagger}\sigma^-)$$

Possible external field driving  $\omega_0 \to \Delta = \omega_0 - \omega_{\text{drive}}$ 

#### Quantum Rabi oscillation

#### Quantum nature of the model

 $\bullet \;\; |e\rangle \stackrel{\mathsf{Spontaneous}\;\mathsf{emission}}{\to} \; |g\rangle \; \mathsf{(but\;not\;irreversible)}$ 

**Dressed state**  $H^{\text{Jaynes-Cummings}}$  in  $\{|g, n+1\rangle, |e, n\rangle\} =$ 

$$\hbar\omega\left(n+\frac{1}{2}\right)-\frac{\hbar\omega_0}{2}+\begin{pmatrix}\hbar\omega&\hbar g\sqrt{n+1}\\\hbar g\sqrt{n+1}&\hbar\omega_0\end{pmatrix}$$

- Oscillation starting with |e>
- Markovian approx. fails
- We have experimental evidence

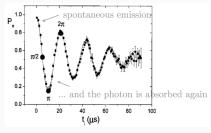


Fig. from S Haroche et al., RMP 73 565 (2001)

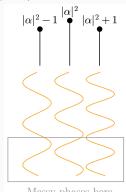
## Collapse and revival

Start with 
$$|\mathbf{e}, \alpha\rangle$$
?  $|\mathbf{e}, \alpha\rangle = e^{-\frac{|\alpha|^2}{2}} \sum_{n=0}^{\infty} \frac{\alpha^n}{\sqrt{n!}} \underbrace{|\mathbf{e}, n\rangle}_{\leftrightarrow |\mathbf{g}, n+1\rangle}$ 

$$P_{e}(t) = \frac{1}{2} \left[ 1 + e^{-|\alpha|^{2}} \sum_{n=0}^{\infty} \frac{|\alpha|^{2n}}{n!} \cos(\Omega_{n}t) \right]$$

$$\stackrel{t \ll 1/g|\alpha|}{=} \frac{1}{2} + \frac{1}{2} \cos(2g|\alpha|t) e^{-\frac{1}{2}g^{2}t^{2}}.$$

**Collapse of**  $P_{\mathbf{e}}$  **when**  $t \ll 1/g|\alpha|$  Because  $\phi^{|e,n\rangle}$  not synchronized This can be simulated by a statistical noise as well; but as  $|\psi\rangle$  is not truly incoherent...



Messy phases here

### Collapse and revival

Revival When the phases of the major components realign again:

$$2\pi = (\Omega_{|\alpha|} - \Omega_{|\alpha|^2 - 1})t$$

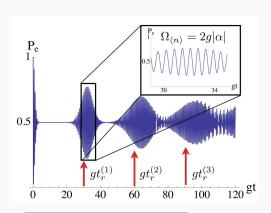


Fig. from arXiv 1111.1143.

#### Revival is a

- Coherent property: not possible with statistical noise
- "Granular" property: see  $|\alpha|^2 1$

## Creation of entangled atom pairs

#### **Protocol**

- 1. Move atom 1 (in  $|e\rangle$ ) into the cavity mode.
- $2. \hspace{0.2cm} |\mathsf{e}_1,0\rangle \overset{\frac{1}{2}\Omega_0 t = \frac{\pi}{4}}{\longrightarrow} \tfrac{1}{\sqrt{2}} (\,|\mathsf{g}_1,1\rangle +\, |\mathsf{e}_1,0\rangle).$
- 3. Move atom 1 out of the light beam. Move atom 2 (in  $|g\rangle$ ) into the light beam.
- $4. \ \ \frac{1}{\sqrt{2}} \big( \ |g_1,\underline{g_2,1}\rangle + \ |e_1,g_2,0\rangle \big) \overset{\frac{1}{2}\Omega_0t=\frac{\pi}{2}}{\longrightarrow} \ \frac{1}{\sqrt{2}} \big( \ |g_1,e_2,0\rangle + \ |e_1,g_2,0\rangle \big)$  coupling happens only here
- 5. Move all atoms out.

That's how you get an Einstein-Podolsky-Rosen pair.

#### C '. I !!

Cavity and medium

Quantum light-matter interaction to the extreme

## "Usual" medium within cavity?

The full theory of quantum optics with  $\epsilon(\omega)$  is hard!!!

In a theory only about photons:  $\epsilon(\omega) \Rightarrow \operatorname{Im} \epsilon \neq 0 \Rightarrow$  non-unitary?

- 1. Auxiliary fields needed for frequency dependence
- 2.  $\Rightarrow$  (space-resolved) Input-output formalism
- 3.  $\stackrel{\text{thermalization}}{\Rightarrow}$  quantum Langevin eq.

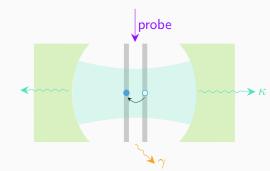
Alternatively we pick up excitations in the medium as auxiliary fields, or "atoms"

See e.g. arXiv quant-ph/0006121, PRL 110, 153602

## Cavity QED with exciton

#### Setup

- One active photon mode
- Probing by the side
- Exciton modes
   ("two-level atom": Xtwo excitons in one mode)



The effective theory: one photon mode, several atoms

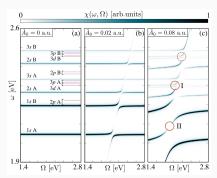
Why this matters  $\langle S|\boldsymbol{d}|0\rangle=0 \Rightarrow \langle S_1|\boldsymbol{d}|S_2\rangle=0$ : dark excitons are optically active in this setting!

### Cavity QED with excitons as atoms

### **Exciton polariton** is $c_1 | \text{exciton} \rangle + c_2 | \text{photon} \rangle$

- Easiest setup
- From the perspective of excitons: virtual photons only

**Optical spectrum**  $\Omega = \text{cavity mode freq.}, \ \omega = \text{detection freq.}$ 



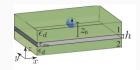
- Dark modes appear (I)
- Quench of optical activity in one band (II)

Fig. from Nano Letters 2019 19 (6), 3473-3479

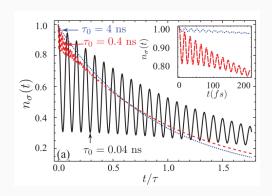
## Cavity QED with surface plasmon polaritons (SPPs) as photons

#### Setup

- Replace "simple" photons by SPPs.
- Atom = exciton on nanoparticle



#### Rabi oscillation of excited state population



## Take home message

- Cavity  $\Rightarrow$  highly selective strong light-matter interaction  $\Rightarrow$  high coherence  $\Rightarrow$  necessity of cavity QED
- Quantum Rabi oscillation, collapse and revival, entanglement
- Cavity QED happens everywhere