

Midterm

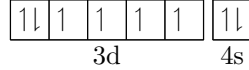
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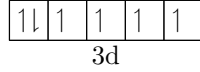
1 Problem 1

1.1

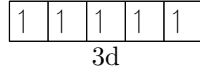
The electron configuration of Fe is $[\text{Ar}]3d^64s^2$, and according to Hund's rule, the spins are



The electron configuration of Fe^{2+} is $[\text{Ar}]3d^6$, and according to Hund's rule, the spins are



Fe^{3+} is obtained by reducing one electron and the spins are



For an iron atom, we only need to work on the 3d orbital because the 4s orbital is full. the total spin quantum number is $S = 2$, and the total orbital angular momentum quantum number is $L = 2$. Therefore

$$g_J = \frac{3}{2} + \frac{S(S+1) - L(L+1)}{2J(J+1)} = \frac{3}{2}. \quad (1)$$

Since the 3d shell is more than half filled, we have $J = L + S = 4$, and the total magnetic moment should be

$$\mu = \mu_B g_J J = 6\mu_B. \quad (2)$$

1.2

The experimentally observed atomic magnetic moment is $2.22\mu_B$, which doesn't agree with the aforementioned prediction. If somehow the orbital angular momentum is quenched, then $S = J = 2$, and $g_J = 2$. So the total angular momentum is 2, which is close to 2.22, but g_J should be multiplied to the former and the result $4\mu_B$ is no longer close to $2.22\mu_B$. The predicted $4\mu_B$ magnetic moment however agrees well with the $4.4\mu_B$ magnetic moment of Fe^{2+} [8], which should have the same magnetic moment with Fe atom because both of them only have the 3d orbital open.

2 Problem 2

2.1

A material is **metamagnetic** if when the external magnetic field passes a finite value H_c , the magnetic configuration changes all of a sudden. This is a phenomenological term and may be driven by various physical mechanisms. The material $\text{Sr}_3\text{Ru}_2\text{O}_7$ is metamagnetic, because experiments have observed that around $\mu_0 H = 7.9\text{ T}$, a sharp peak can be seen in magnetic susceptibility [3], and therefore there is indeed a sudden change in the magnetization.

The low-field phase is paramagnetic, and the high-field phase is itinerantly ferromagnetic: the material shows "a rapid change from a paramagnetic state at low fields to a more highly polarized state" [7]. (On the other hand, some other metamagnetic materials undergo an antiferromagnetism-to-ferromagnetism transition; this is not the case for $\text{Sr}_3\text{Ru}_2\text{O}_7$.)

The boundary between the two phases was once thought to be a quantum critical point: above the phase boundary between the low-field phase and the high-field phase, the resistance doesn't have typical Fermi-liquid behaviors [7]; the phase boundary between the two phases is a first-order phase transition line with a terminating end point, and this critical point is pushed to $T = 0$ when the external magnetic field is pointed towards the c direction, creating a quantum critical point [2]. Further investigations however have found that there are actually *two* peaks in susceptibility near $\mu_0 H = 7.9$ T, and this “quantum critical point” is surrounded by two first-order phase transitions [5, 3]. The exotic temperature-resistance curve likely comes from an SDW order on top of the ferromagnetic moment formed between the two aforementioned first-order phase transition, which also gives rise to anisotropic resistance (or in other words, electronic nematic) which isn't induced by the crystal structure and can't be seen away from $\mu_0 H = 7.9$ T [6, 1].

2.2

$\text{Sr}_3\text{Ru}_2\text{O}_7$ is an itinerant magnetic material and the metamagnetic transition is likely due to Fermi surface instability. This is explained in [4] with a toy model. The effect or the coupling between the electron magnetic moment and the external magnetic field is modifying the chemical potential for

This also explains why nematic electron fluid is only observed near the so-called quantum critical point $\mu_0 H = 7.9$ T: because when the external magnetic field is stronger, the

3 Problem 3

3.1

The model Hamiltonian is

$$\begin{aligned}
H = & \sum_k \varepsilon_k^c c_k^\dagger c_k + \sum_k \varepsilon_k^f f_k^\dagger f_k + \sum_q \omega_q b_q^\dagger b_q \\
& + \frac{1}{L} \sum_{k, k', q} V_q c_{k+q}^\dagger f_{k'-q}^\dagger f_{k'} c_k \\
& + \frac{1}{\sqrt{L}} \sum_{k, q} g_{kq} c_{k+q}^\dagger b_q f_k + \text{h.c.} \\
& - \mu N,
\end{aligned} \tag{3}$$

where L is the length of the sample and the last term is the chemical potential. Here the $1/2V$ factor before the Coulomb interaction term is replaced by $1/V$ because otherwise we have a $f^\dagger c^\dagger c f$ term; in the electron-phonon scattering term, we only have the $c^\dagger f b$ channel, because the conduction band is energetically higher than the valence band and therefore hopping from the latter to the former needs the additional energy from an existing phonon.

3.2

3.3

Since

3.4

TODO This can also be found by perturbation theory: it requires one $f^\dagger c$ vertex and one $c^\dagger f$ vertex to complete an energy correction to the f band or the c band, and therefore the first correction to the free energy caused by phonons is proportional to $b^2 \sim X^2$, and it's negative because of aforementioned energy saving. So this negative X^2 term is to be merged with the $(T - T_c)X^2$ term and now T_c is higher, and therefore electron-phonon coupling enhances the existing structural instability.

References

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