

Homework 11

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Problem 1 Consider the band structure of InP as having only the electron band and the heavy hole band. Use numerical values from Kittel for the material. a) Where in the band gap is the chemical potential located at 300 K for an ultrapure sample of InP? b) Explain briefly the physical reason why the chemical potential has the value you found in part (a) (i.e., why it changes in the direction it does with $T > 0$ compared to $T = 0$).

Solution

(a) In Kittel Section 8.1, we find $E_g = 1.27 \text{ eV}$, and in Section 8.2.5, we find $m_e^* = 0.073m$, $m_h = 0.4m$, so according to

$$\mu = E_v + \frac{1}{2}E_g + \frac{3}{4}k_B T \ln\left(\frac{m_h^*}{m_e^*}\right), \quad (1)$$

we find $\mu - E_v = 0.68 \text{ eV}$, so it's slightly above the middle point between E_v and E_c .

(b) We know the Fermi-Dirac distribution function for electrons above μ and the Fermi-Dirac distribution function for holes below μ are symmetric to each other. Since m_h^* is larger, the valence band is flatter, and near the top of the valence band, the number of states is larger than that near the bottom of the conduction band. So if $\mu - E_v = E_g/2$, the Fermi-Dirac factor in n and p is the same, while the density of state factors are different, and we find $p > n$, which is wrong. So we have to push μ higher to reduce the $f(E)$ factor in p to ensure $n = p$.

Problem 2 At very low temperatures, donor electrons in n-type semiconductors are bound to the impurity (donor) ions (only at higher temperatures will they be ionized and be mobile). Staying at very low temperatures, consider increasing the doping concentration to the point when the bound electron orbits on adjacent impurity atoms being to "touch" (overlap). When this occurs, electrons can become mobile by tunneling from one impurity orbit to the next and an "impurity band" forms (very close in philosophy to the tight-binding picture). Consider n-type InAs and use any needed values from Kittel. a) What is the donor ionization energy in InAs? Assume the hydrogenic model holds so you can simply scale from the Hydrogen atom. b) What is the radius of the electron orbit in the bound state on the impurity? Use the same assumptions as (a) and then check your assumptions are selfconsistently correct. c) Estimate the donor concentration above which the adjacent electron orbits overlap significantly and we start to get an impurity band. (You can define a term like "significantly" in your own sensible way but make sure you explain your definition.)

Solution

(a) When the hydrogenic model works, the inner details of the donor atom don't matter: the doped material can be seen as a positive charge placed in an undoped material, and the Coulomb field introduced by the former is then screened by the latter. So the spectrum of an electron near the donor atom can be obtained by replacing m by m_e^* and ϵ_0 by $\epsilon\epsilon_0$ in the Bohr theory of hydrogen, and we get

$$E_n = \frac{1}{n^2} \frac{e^4 m_e^*}{2(4\pi\epsilon\epsilon_0\hbar)^2} = \frac{1}{\epsilon^2} \frac{m_e^*}{m} \frac{1}{n^2} \underbrace{E_0}_{13.6 \text{ eV}}. \quad (2)$$

In Section 8.4.1 we find $\epsilon = 14.55$, and in Section 8.2.5 we find $m_e^* = 0.026m$, so finally we find $E_1 = -1.67 \times 10^{-3} \text{ eV}$, and therefore the ionization energy is $1.67 \times 10^{-3} \text{ eV}$.

(b) Similarly we have

$$a = \frac{4\pi\epsilon\epsilon_0\hbar^2}{m_e^* e^2} = \frac{\epsilon m}{m_e^*} \cdot 0.53 \text{ \AA}, \quad (3)$$

and for InAs, $a = 29.6 \text{ nm}$. Wikipedia says the lattice constant of InAs is $\sim 6.05 \text{ \AA}$, so the radius of the electron donated is much larger than the characteristic length of the crystal structure,

which means our approximation is reasonable, or otherwise at some momentum values involved, the hyperbolic approximation of the band is no longer correct, and the details of the crystal structure is visible to the motion of the impurity-bound electron, and the approximation breaks down.

(c) When the distance between impurities is comparable to a , hopping between nearest impurities becomes frequent. So when

$$n \gtrsim \frac{1}{a^3} = 3.86 \times 10^{22} \text{ m}^{-3},$$

the impurity band becomes visible.

Problem 3 Kittel problem 2 of chapter 8. The temperature here is so low, that the formulae we derived in lecture are not applicable (the chemical potential is now quite close to the donor bound state energy and not all the donors are ionized). You will have to read on the right relations to use for such low temperatures (A&M Chapter 28 problem 6 is a good place to look; Kittel also has a brief discussion, but you will have to see if the two of them agree...).

In a particular semiconductor there are $1 \times 10^{13} \text{ donors/cm}^{-3}$ with an ionization energy E_d of 1meV and an effective mass 0.01 m. (a) Estimate the concentration of conduction electrons at 4 K. (b) What is the value of the Hall coefficient? Assume no acceptor atoms are present and that $E_g \gg k_B T$.

Solution

(a) Thermal hopping from the valence band to the conduction band can be ignored and approximately

$$n = N_d - n_d = N_d \left(1 - \frac{1}{1 + \frac{1}{2} e^{(E_d - \mu)/k_B T}} \right), \quad (4)$$

and

Problem 4

Solution

(a) The intrinsic chemical potential is

$$\mu_i = \frac{1}{2} E_g + \frac{3}{4} k_B T \ln \left(\frac{m_h^*}{m_e^*} \right) = 0.74 \text{ eV}. \quad (5)$$

So

$$n_i = 2 \left(\frac{m_e^* k_B T}{2\pi \hbar^2} \right)^{3/2} e^{-(E_c - \mu_i)/k_B T} = 3.48 \times 10^{12} \text{ m}^{-3}. \quad (6)$$

Using Eq. (28.39) in A&M (further assuming $E_d - \mu \gg k_B T$), which is

$$\frac{N_d - N_a}{n_i} = 2 \sinh \frac{\mu - \mu_i}{k_B T}, \quad (7)$$

we get $\mu = 0.74 \text{ eV} + 0.54 \text{ eV} = 1.28 \text{ eV}$.

(b) We have (here E_d is the donor binding energy)

$$n_d = \frac{N_d}{1/2 + e^{(E_c - E_d - \mu)/k_B T}} = 1.8 \times 10^{14} \text{ cm}^{-3}. \quad (8)$$

It can be found that most of the donor electrons indeed enter the conduction band, and $N_d - n_d \gg n_i$. The system is therefore in the extrinsic limit.

(c) We have

$$\frac{E_c - \mu}{k_B T} = 4.6, \quad \frac{E_c - E_d - \mu}{k_B T} = 3.3,$$

so the non-degenerate assumption that the above ratios should be very large doesn't hold – but since $\Delta E/k_B T$ appears in the exponential function, in practice it's still good enough.

(d) We have

$$n = \frac{\Delta n + \sqrt{4n_i^2 + \Delta n^2}}{2}, \quad \Delta n = N_d - n_d, \quad (9)$$

and the result is $n = 4.82 \times 10^{15} \text{ cm}^{-3}$.

(e) It can be found that the system is indeed in the extrinsic limit: we have $n = \Delta n$, and the contribution of n_i can be ignored.

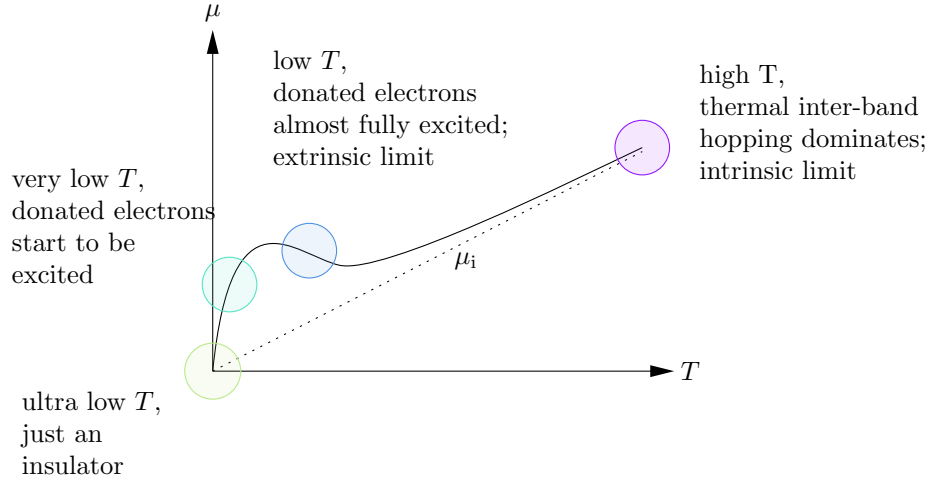


Figure 1: The temperature dependence of μ . The energy zero point is $E_v + E_g/2$. When T is very, very close to 0, doped electron

Problem 5

Solution

(a) When $T = 700$ K, we have

$$n_i = 2 \left(\frac{m_e^* k_B T}{2\pi \hbar^2} \right)^{3/2} e^{-(E_c - \mu_i)/k_B T} = 7.54 \times 10^{18} \text{ m}^{-3}, \quad (10)$$

and (here we assume the middle of the gap has zero energy)

$$\mu_i = \frac{3}{4} k_B T \ln \left(\frac{m_h^*}{e_e^*} \right) = 0.92 \text{ eV}. \quad (11)$$

Again using (7), we have $\mu = 0.92 \text{ eV} - 0.07 \text{ eV} = 0.85 \text{ eV}$. This is close to μ_i : the relative error is 0.08. The reason is now the system is in the intrinsic limit, because doping concentration is much smaller than n_i .

(b) See Figure 1.