

Trion in time-resolved ARPES

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The most generic theory requires Keldysh formalism but let's spare ourselves the burden. . .

Ingredients of our model of ARPES

- Electric dipole interaction only
- Sudden approximation
- *Separation between pump and probe*: system *not* driven when probed; pump prepares an initial state and nothing more
- Fermi's golden rule in probing

Main result Output intensity $I_{\mathbf{k}}(t) \propto \sum_c P_{\mathbf{k}}^c(t)$,

$$P_{\mathbf{k}}^c(t) = \sum_{n, \mathbf{k}'} \rho_n |M_{\mathbf{k}\mathbf{k}'}^{fc}|^2 \int_{t_0}^t dt_1 \int_{t_0}^{t_1} dt_2 e^{-i(E_n - \omega)(t_1 - t_2)} \langle \Psi_n(t_0) | c_{\mathbf{k}'}^\dagger U(t_2, t_1) c_{\mathbf{k}} | \Psi_n(t_0) \rangle s(t_1) s(t_2).$$

1. ρ_n : distribution of the final state of pumping (initial state of probing).
2. Probe field is $\mathbf{E} = s(t) \mathbf{E}_0 e^{-i\omega_0 t} + \text{c.c.}$, and the transition matrix is $M_{\mathbf{k}\mathbf{k}'}^{fc} = \langle f\mathbf{k} | -\mathbf{d} \cdot \mathbf{E}_0 | c\mathbf{k}' \rangle$; c is the probed electron, f is the out-going state.
3. The $e^{-i(E_n - \omega)(t_1 - t_2)}$ factor gives half of energy conservation condition; ω is driving frequency ω_0 shifted by work function.
4. $\langle \Psi_n | \cdots | \Psi_n \rangle$: electron Green function with excited state background $|\Psi_n\rangle$; it gives the structure of $|\Psi_n\rangle$ in electron basis, and the second half of energy conservation condition (energy after one electron being kicked out);
5. $s(t_1)s(t_2)$: shape of probe pulse; broadening $\delta(E_n - \omega - E_{\text{after}})$

Successful example: exciton

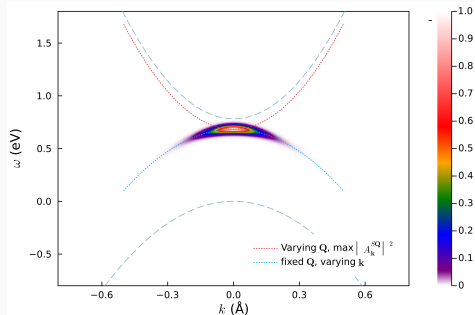
- Below we work with 2D material $\Rightarrow \mathbf{k}_{\parallel}^{\text{out}} = \mathbf{k}'$; we refer to $\mathbf{k}_{\parallel}^{\text{out}}$ as \mathbf{k}
- Only consider valence band top and conduction band valley: parabolic bands

ARPES response of an exciton (besides valence band electron)

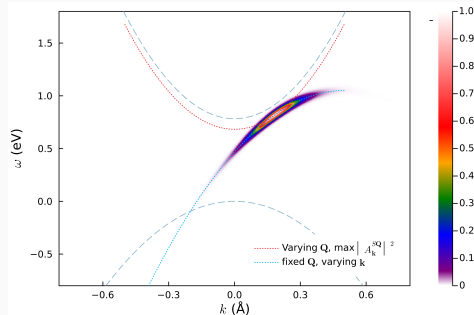
$$P(t) = \sum_{S, Q} \rho_{SQ} |A_{\mathbf{k}'}^{SQ}|^2 |M_{\mathbf{k}\mathbf{k}'}^{fc}|^2 \int_{t_0}^t dt_1 \int_{t_0}^t dt_2 e^{-i(E_{SQ} + \epsilon_{v\mathbf{k}'} - Q - \omega)(t_1 - t_2)} s(t_1) s(t_2).$$

- Detection of $|A_{\mathbf{k}}^{SQ}|^2$
- $Q = 0 \Rightarrow$ dispersion relation $\omega = \epsilon_{v\mathbf{k}}$ (energy conservation)
- Q obeys $e^{-\beta E_Q}$ distribution: dispersion relation is $\omega = E_g + E_B + \frac{k^2}{2m_e}$ (energy conservation; $|A_{\mathbf{k}}^{SQ}|^2$ reaches peak when $\mathbf{v}_e = \mathbf{v}_h$)

Successful example: exciton



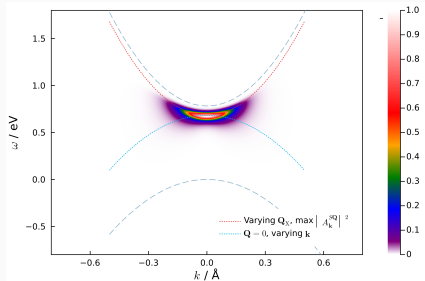
Single exciton, $Q = 0.0 \text{ \AA}^{-1}$



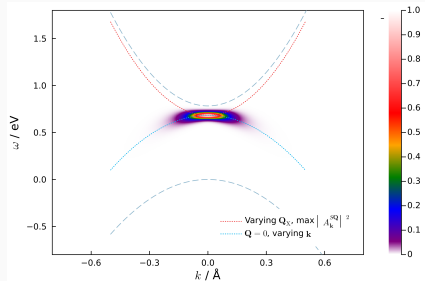
Single exciton, $Q = 0.5 \text{ \AA}^{-1}$

- The center of the signature is on a replica of the conduction band
- The shape of the signature is a replica of the valence band

Successful example: exciton



Thermalized exciton, $\beta = 10$ eV



Thermalized exciton, $\beta = 100$ eV

- The center of the signature is on a replica of the conduction band
- The shape of the signature is a replica of the valence band
- When excitons are hot, the dispersion relation is the conduction band; when they are cold, the dispersion relation is the valence band

Trion in two-band model

Configuration of trion structure

- Two holes, one electron (what likely happens in ARPES)
- *Two holes on one band* – to simplify analysis
- *No hole scattering after the electron is kicked out* – realistic or not?

$$H = \frac{(\mathbf{k}_e - \mathbf{w})^2}{2m_e} + E_g + \frac{\mathbf{k}_{h1}^2}{2m_h} + \frac{\mathbf{k}_{h2}^2}{2m_h} + V(\mathbf{r}_{h1} - \mathbf{r}_{h2}) - V(\mathbf{r}_e - \mathbf{r}_{h1}) - V(\mathbf{r}_e - \mathbf{r}_{h2}). \quad (1)$$

Trion in two-band model

Redefining coordinates What we want: $\mathbf{k}_{e,h1,h2} \rightarrow \mathbf{P}, \mathbf{k}_1, \mathbf{k}_2$, following these constraints:

$$\mathbf{k}_{1,2} = -i\hbar\partial_{\mathbf{r}_{1,2}} = \text{const}, \quad \mathbf{r}_{1,2} = \mathbf{r}_{h1,2} - \mathbf{r}_e, \quad \mathbf{R} = i\partial_{\mathbf{P}} = \frac{m_h\mathbf{r}_{h1} + m_h\mathbf{r}_{h2} + m_e\mathbf{r}_e}{2m_h + m_e}.$$

It can be verified that with

$$\mathbf{P} = \mathbf{k}_e + \mathbf{k}_{h1} + \mathbf{k}_{h2}, \quad \mathbf{k}_{1,2} = -i\hbar\partial_{\mathbf{r}_{1,2}} + \frac{m_h}{2m_h + m_e}\mathbf{w}, \quad (2)$$

$$M = 2m_h + m_e, \quad \mu = \frac{m_e m_h}{m_e + m_h}, \quad H_X = -\frac{\nabla^2}{2\mu} + V(\mathbf{r}), \quad (3)$$

$$H = \underbrace{H_X(\mathbf{r}_1) + H_X(\mathbf{r}_2) - \frac{\hbar^2 \nabla^2}{2m_e} + V(\mathbf{r}_1 - \mathbf{r}_2)}_{\text{causing } E_B} + \frac{\mathbf{P}^2}{2m_T} + E_g \quad (4)$$

Trion in two-band model

That's to say: a trion, when the two holes of it are of the same species, can be seen as the bound state of two excitons.

Trion wave function Thus in this specific case:

$$\phi_T(\mathbf{r}_e, \mathbf{r}_{h1}, \mathbf{r}_{h2}) = \sum_{S_1, S_2} \phi_X^S(\mathbf{r}_1) \phi_X^S(\mathbf{r}_2) e^{i\mathbf{P} \cdot \mathbf{R}}, \quad (5)$$

where $\mathbf{r}_{1,2}$ are defined above.

The below ansatz is frequently used:

$$\phi_T(\mathbf{r}_1, \mathbf{r}_2) \propto \phi_{1s}(\mathbf{r}_1; a) \phi_{1s}(\mathbf{r}_2; b) + \phi_{1s}(\mathbf{r}_1; b) \phi_{1s}(\mathbf{r}_2; a) \quad (6)$$

where a, b are exciton radii.

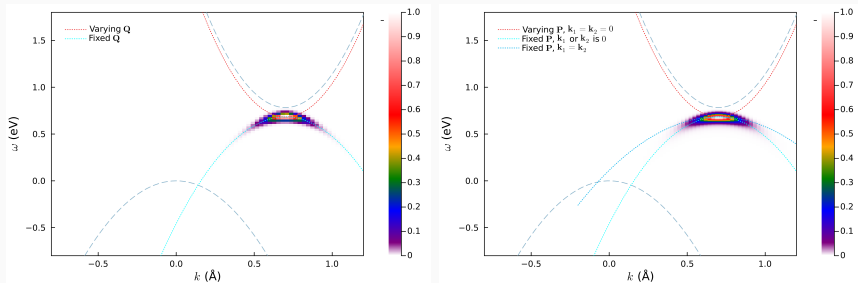
Trion in two-band model

In the tentative plots below We ignore spin and choose $S_1 = S_2 = 1s$. The ARPES response (here ϵ_v are *negative*):

$$P_{\mathbf{k}}(t) = \sum_{\mathbf{P}} \rho_{\mathbf{P}} \sum_{\mathbf{k}_{h1}, \mathbf{k}_{h2}} \sum_{S_1, S_2} |M_{\mathbf{k}\mathbf{k}_e}^{fc}|^2 |A_{S_1 S_2}|^2 |\phi_{S_1}(\mathbf{k}_1)|^2 |\phi_{S_2}(\mathbf{k}_2)|^2 \times \int_{t_0}^t dt_1 \int_{t_0}^t dt_2 e^{-i(E_{SP} + \epsilon_{v\mathbf{k}_{h1}} + \epsilon_{v\mathbf{k}_{h2}} - \omega)(t_1 - t_2)} s(t_1) s(t_2) \quad (7)$$

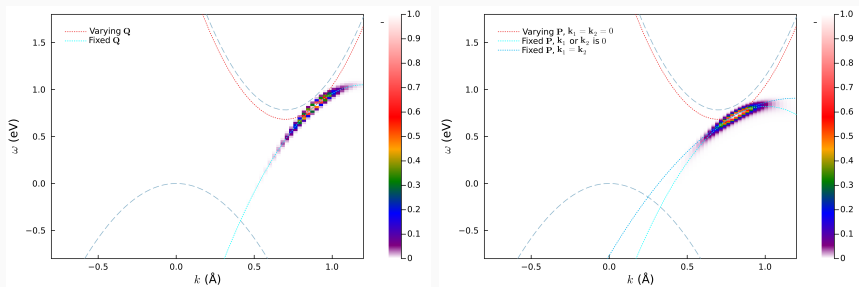
Signature maximum: $\mathbf{k}_1 = 0$ or $\mathbf{k}_2 = 0$, or $\mathbf{k}_1 = \mathbf{k}_2$?

Trion and exciton: $Q = w$



Main difference: curvature. In trion the signature partially has the shape of $2\epsilon_{v,k/2}$

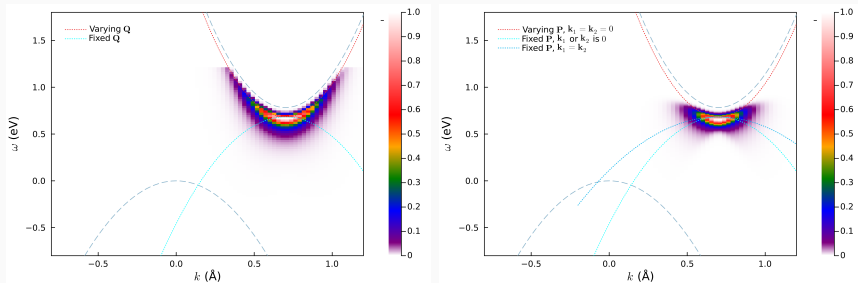
Exciton and trion: finite $Q - w$



- The signature curvature difference is more obvious.
- The positions of the signatures are different!

$$\mathbf{k}_{\text{exciton, max}} = \frac{m_e \mathbf{Q} + m_h \mathbf{w}}{m_e + m_h}, \quad \mathbf{k}_{\text{trion, max}} = \frac{m_e}{2m_h + m_e} \mathbf{P} + \frac{2m_h}{2m_h + m_e} \mathbf{w}.$$

Exciton and trion: thermalized, $\beta = 10$



The trion calculation is not well-conserved: the sampling grid of \mathbf{P} is too small; I don't expect to see any large difference because the shape of the two signatures should both follow ϵ_c .

Conclusion

- Trion signature on ARPES spectrum looks similar to that of exciton
- The finite- T induced “upper” dispersion of the two is the same
- The trion and exciton binding energies are different
- With finite momenta, exciton and trion signatures are different in shape and position.