Problem A

Consider an isolated system in equilibrium with two parts, A and B, having a mixture of particles where $U_{\text{total}} = U_A + U_B$, $V_{\text{total}} = V_A + V_B$, and $N_{\text{total}} = N_A + N_B$ (for each type of particle). Since it is isolated, the change in these parameters is $dU_{\text{total}} = dV_{\text{total}} = dN_{\text{total}} = 0$. The change in entropy can be expressed as

$$\Delta S_{T} = \sum_{\alpha=A}^{B} \left[\left(\frac{\partial S_{\alpha}}{\partial U_{\alpha}} \right)_{V_{\alpha} N_{\alpha}}^{0} \Delta U_{\alpha} + \left(\frac{\partial S_{\alpha}}{\partial V_{\alpha}} \right)_{U_{\alpha} N_{\alpha}}^{0} \Delta V_{\alpha} + \left(\frac{\partial S_{\alpha}}{\partial N_{\alpha}} \right)_{V_{\alpha} U_{\alpha}}^{0} \Delta N_{\alpha} \right] + \frac{1}{2} \left[\Delta \left(\frac{\partial S_{\alpha}}{\partial U_{\alpha}} \right)_{V_{\alpha} N_{\alpha}}^{0} \Delta U_{\alpha} + \Delta \left(\frac{\partial S_{\alpha}}{\partial V_{\alpha}} \right)_{U_{\alpha} N_{\alpha}}^{0} \Delta V_{\alpha} + \sum_{j=1}^{l} \Delta \left(\frac{\partial S_{\alpha}}{\partial N_{\alpha,j}} \right)_{V_{\alpha} U_{\alpha}}^{0} \Delta N_{\alpha,j} \right] + \dots$$

$$(1)$$

Show that the two terms within the red box will yield zero. What conditions between A and B can be derived from these two terms?

Solution:

Using the Born Square, we get the following relation for the change in internal energy:

$$dU = TdS - PdV + \mu dN \tag{2}$$

Rearranging this, we obtain a relation for the change in entropy:

$$dS = \frac{1}{T}dU + \frac{P}{T}dV - \frac{\mu}{T}dN$$
(3)

Comparing the obtained relation with Eq. (1), we can infer that

$$\frac{1}{T} = \left(\frac{\partial S}{\partial U}\right)_{V,N}, \qquad \frac{P}{T} = \left(\frac{\partial S}{\partial V}\right)_{U,N}, \qquad \frac{\mu}{T} = -\left(\frac{\partial S}{\partial N}\right)_{V,U} \tag{4}$$

We assume that the interface between parts A and B can transmit heat, mechanical energy and particles of all types. In this case, the thermodynamic variables in both parts of the system (*i.e.* temperature, pressure, and chemical potential for each type of particle) must be the same for the system to be in equilibrium. Thus, the first term in the red box becomes

$$\sum_{\alpha=A}^{B} \left[\left(\frac{\partial S_{\alpha}}{\partial V_{\alpha}} \right)_{U_{\alpha}^{0} N_{\alpha}} \Delta V_{\alpha} \right] = \frac{P_{A}}{T_{A}} \Delta V_{A} + \frac{P_{B}}{T_{B}} \Delta V_{B} = \left(\frac{P_{A}}{T_{A}} \frac{P_{B}}{T_{B}} \right)^{0} \Delta V_{A} = 0$$
 (5)

and the second term becomes

$$\sum_{\alpha=A}^{B} \left[\left(\frac{\partial S_{\alpha}}{\partial N_{\alpha}} \right)_{U_{\alpha}V_{\alpha}}^{0} \Delta N_{\alpha} \right] = -\frac{\mu_{A}}{T_{A}} \Delta N_{A} - \frac{\mu_{B}}{T_{B}} \Delta N_{B} = \left(\frac{\mu_{B}}{T_{B}} \right)_{A}^{0} \Delta N_{A} = 0$$
 (6)

Problem B

Derive the Maxwell's relations from the Gibb's free energy assuming constant number of particles.

Solution:

The total differential of G can be written as

$$dG = \left(\frac{\partial G}{\partial P}\right)_{T,N} dP + \left(\frac{\partial G}{\partial T}\right)_{P,N} dT + \left(\frac{\partial G}{\partial N}\right)_{P,T} dN^{\bullet 0}$$
(7)

where dN = 0 since there is a constant number of particles. Using the Born Square, we get

$$dG = VdP - SdT + \mu dN^{-0}$$
(8)

Comparing these two equations, we have the following relations

$$V = \left(\frac{\partial G}{\partial P}\right)_{T,N}, \qquad S = -\left(\frac{\partial G}{\partial T}\right)_{P,N} \tag{9}$$

Making use of the symmetry of the second derivatives and applying the relations in Eq. (9), we obtain

$$\left(\frac{\partial}{\partial T}\frac{\partial G}{\partial P}\right)_{P,N} = \left(\frac{\partial}{\partial P}\frac{\partial G}{\partial T}\right)_{T,N} \longrightarrow \left(\frac{\partial V}{\partial T}\right)_{P,N} = -\left(\frac{\partial S}{\partial P}\right)_{T,N} \tag{10}$$

which is the Maxwell's relation derived from Gibb's free energy with fixed N.