

Problem 1

- 1.1 Obtain the partition function of a distinguishable particle in a one-dimensional simple harmonic oscillator potential.

Solution:

The energy levels of a quantum harmonic oscillator is given by

$$E_n = \left(n + \frac{1}{2}\right) \hbar\omega \quad (1)$$

Since energy of this oscillator is allowed to change, we use the partition function of the canonical ensemble which is defined as

$$Z \equiv \sum_s e^{-\beta E_s} \quad (2)$$

where $\beta = 1/(k_B T)$ in which k_B is the Boltzmann constant. Substituting Eq. (1), we obtain

$$Z = \sum_n e^{-\beta(n+\frac{1}{2})\hbar\omega} = e^{-\frac{1}{2}\beta\hbar\omega} \sum_n \left(e^{-\beta\hbar\omega}\right)^n \quad (3)$$

Note that we have obtained an infinite geometric series. The sum of such a series is

$$1 + x + x^2 + x^3 \dots = (1 - x)^{-1} \quad (4)$$

Applying this relation, the Z becomes

$$Z = \frac{e^{-\frac{1}{2}\beta\hbar\omega}}{1 - e^{-\beta\hbar\omega}} = \frac{1}{e^{\frac{1}{2}\beta\hbar\omega} - e^{-\frac{1}{2}\beta\hbar\omega}} = \frac{1}{2 \sinh \frac{\beta\hbar\omega}{2}} \quad (5)$$

with the use of hyperbolic functions.

- 1.2 From your result in 1.1, obtain the system's internal energy and heat capacity.

Solution:

It is first important to find the natural logarithm of the partition function as follows:

$$\ln Z = \ln \frac{1}{2 \sinh \frac{\beta\hbar\omega}{2}} = -\ln \left(2 \sinh \frac{\beta\hbar\omega}{2} \right) \quad (6)$$

Then, using this result, we obtain the system's internal energy as

$$U = -\frac{\partial \ln Z}{\partial \beta} = \frac{\partial}{\partial \beta} \left[\ln \left(2 \sinh \frac{\beta\hbar\omega}{2} \right) \right] = \frac{2 \cosh \frac{\beta\hbar\omega}{2}}{2 \sinh \frac{\beta\hbar\omega}{2}} \cdot \frac{\hbar\omega}{2} = \frac{\hbar\omega}{2} \coth \frac{\beta\hbar\omega}{2} \quad (7)$$

Thus, we get

$$U = \frac{\hbar\omega}{2} \coth \frac{\hbar\omega}{2k_B T} \quad (8)$$

We use this U to get the heat capacity of the system:

$$C_V = \left(\frac{dU}{dT} \right)_{V,N} = \frac{\hbar\omega}{2} \frac{d}{dT} \left[\coth \frac{\hbar\omega}{2k_B T} \right] = \frac{\hbar\omega}{2} \left(-\operatorname{csch}^2 \frac{\hbar\omega}{2k_B T} \right) \left(-\frac{\hbar\omega}{2k_B T^2} \right) \quad (9)$$

Thus, we have

$$C_V = \frac{1}{k_B} \left(\frac{\hbar\omega}{2T} \operatorname{csch} \frac{\hbar\omega}{2k_B T} \right)^2 \quad (10)$$

1.3 Discuss your result in 1.2, as T approaches infinity.

Solution:

Ignoring the effect the constants in the expressions, we plot the system's internal energy and heat capacity in Fig. 1 and Fig. 2 respectively. We see in Fig. 1 that the internal energy is proportional to temperature. On the other hand, the heat capacity asymptotes to some value as T approaches infinity as shown in Fig. 2.

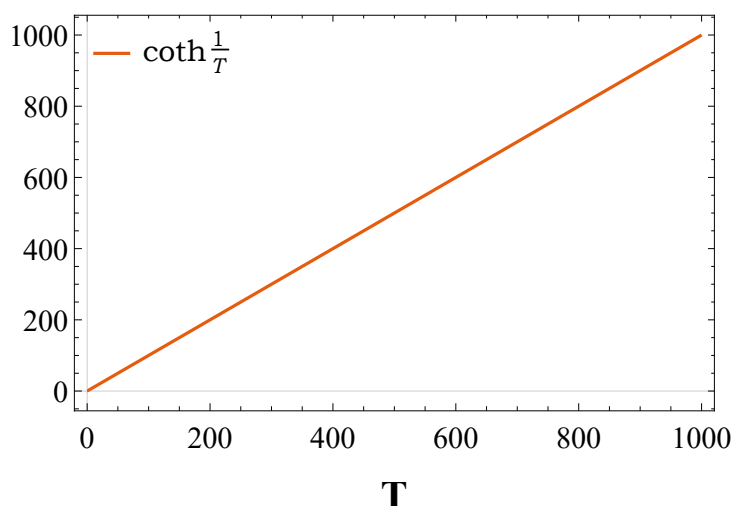


Figure 1: A plot of $\coth T^{-1}$

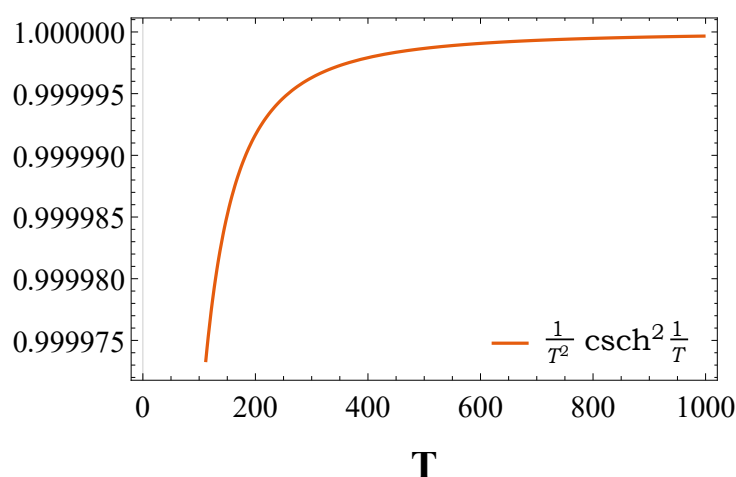


Figure 2: A plot of $T^{-2} \coth^2 T^{-1}$

Problem 2

Consider a one-dimensional array of N spins in a paramagnet with each spin having a magnetic moment μ . In the presence of a magnetic field B , the spins can align parallel (-) or antiparallel (+) to the field with energy $\pm\mu B$.

2.1 Obtain the partition function of this system.

Solution:

It is given that the energy of the system can either be

$$E = \pm \mu B \quad (11)$$

As the energy of a spin can change, we again use the partition function of the canonical ensemble as given in Eq. (2). Let us first obtain the partition function of a single spin:

$$Z_1 = e^{-\beta \mu B} + e^{\beta \mu B} = 2 \cosh \beta \mu B \quad (12)$$

For a system with N spins, the partition function is given by

$$Z_N = (Z_1)^N = (2 \cosh \beta \mu B)^N \quad (13)$$

2.2 From your result in 2.1, obtain the system's internal energy.

Solution:

In the same manner as in the previous problem, we first calculate for $\ln Z_N$:

$$\ln Z_N = \ln (2 \cosh \beta \mu B)^N = N \ln (2 \cosh \beta \mu B) \quad (14)$$

Then, the internal energy of the given system is

$$U = -\frac{\partial \ln Z_N}{\partial \beta} = -N \frac{\partial}{\partial \beta} [\ln (2 \cosh \beta \mu B)] = -N \frac{2 \sinh \beta \mu B}{2 \cosh \beta \mu B} \cdot \mu B = -N \mu B \tanh \beta \mu B \quad (15)$$