

Dynamically, central force leads to 2D orbiting motion.

Let's try quantizing the problem in 2D first.

$$H = \frac{p_r^2}{2m} + \frac{L^2}{2mr^2} - \frac{e}{r}$$

① no  $\phi$  dependence  $\Rightarrow L$  is constant of motion with quantized value  $\ell$

② Assume  $\langle r | p_r^2 = -\frac{\hbar^2}{r} \partial_r r \partial_r \langle r | = (-\hbar^2) \left( \frac{1}{r} \partial_r + \partial_r^2 \right) \langle r |$

$$\left[ \left( -\frac{\hbar^2}{2m} \right) \left( \frac{1}{r} \partial_r + \partial_r^2 \right) + \frac{\hbar^2}{2m} \frac{\ell^2}{r^2} - \frac{e}{r} - E \right] R = 0$$

$$r \rightarrow \infty, \left( \partial_r^2 + \frac{2mE}{\hbar^2} \right) R \rightarrow 0 \Rightarrow R \rightarrow e^{-\lambda r}; \quad \lambda^2 \equiv \frac{-2mE}{\hbar^2}$$

incorporate the physical scale:  $\rho \equiv \lambda r$

$$\left( \partial_\rho^2 + \frac{1}{\rho} \partial_\rho - \frac{\ell^2}{\rho^2} + \frac{2me}{\hbar^2} \frac{1}{\lambda \rho} - 1 \right) R = 0; \quad \gamma \equiv \frac{me}{\hbar^2 \lambda}$$

$$\text{Assume } R = e^{-\rho} f(\rho), \quad \partial_\rho R = (f'(\rho) - f(\rho)) e^{-\rho}, \quad \partial_\rho^2 R = (f''(\rho) - 2f'(\rho) + f(\rho)) e^{-\rho}$$

$$(f'' + \frac{1}{\rho} f' - \frac{\ell^2}{\rho^2} f) + ((2\gamma - 1) \frac{1}{\rho} f - 2f') = 0$$

$$\text{Assume } f(\rho) = \sum_m a_m \rho^{s+m}$$

$$\sum_m a_m \left[ (s+m)(s+m-1) + (s+m) - \ell^2 \right] \rho^{s+m-2} - \sum_m a_m \left[ 2\gamma - 1 - 2(s+m) \right] \rho^{s+m-1} = 0$$

$$\text{lowest order: } s^2 - \ell^2 = 0 \Rightarrow s = \ell$$

$$\text{Assume } a_m = 0 \text{ for } m > m_\ell,$$

$$\text{highest order: } \gamma = s + m_\ell + \frac{1}{2}$$

$$E = \frac{-\hbar^2}{2m} \lambda^2 = \frac{-\hbar^2}{2m} \left( \frac{me}{\hbar^2} \frac{1}{\gamma} \right)^2 = \frac{-me^2}{2\hbar^2} \cdot \frac{1}{\gamma^2} \Rightarrow \text{fit exp. if } s = \frac{1}{2}, \frac{3}{2}, \frac{5}{2}, \dots$$