J.R.Schrieffer, 1964, Theory of Superconductivity

Koji Higasa

2020 年吉日

Preface

これは [1] の計算を追ったものである.

目次

| 第1章 | Introduction | 7 |
|-----|--|---|
| 1.1 | Simple Experimental Facts | 7 |
| 1.2 | Phenomenological Theories | 7 |
| 第2章 | The Pairing Theory of Superconductivity | 9 |
| 2.1 | Physical Nature of the Superconducting State | 9 |
| 2.2 | The One-Pair Problem | 9 |
| 2.3 | Landau's Theory of a Fermi Liquid | 9 |
| 2.4 | The Pairing Approximation | 9 |
| 2.5 | Quasi-Particle Excitations | 9 |
| 2.6 | Linearized Equations of Motion | 9 |

第1章

Introduction

- 1.1 Simple Experimental Facts
- 1.2 Phenomenological Theories
- 1.2.1 Gorter-Casimir Model

(1-6).

$$\begin{split} \frac{\partial F}{\partial x} &= \frac{1}{2} \frac{1}{\sqrt{x}} \left(-\frac{1}{2} \gamma T^2 \right) - 1 (-\beta) = 0 \\ &\frac{\gamma T^2}{4} \frac{1}{\sqrt{x}} = \beta \\ &\therefore \sqrt{x} = \frac{\gamma T^2}{4\beta} \\ &\therefore x = \left(\frac{\gamma}{4\beta} \right)^2 T^4 \end{split}$$

(1-8).

$$H_c^2(T) = 8\pi (F_n(T) - F_s(T)) = 8\pi (F(1,T) - F(0,T))$$

$$= 8\pi (f_n(T) - f_s(T)) = 8\pi \left(-\frac{1}{2}\gamma T^2 + \beta\right)$$

$$= 8\pi \beta \left(1 - \frac{\gamma}{2\beta}T^2\right) \equiv H_0\left(1 - \left(\frac{T}{T_c}\right)\right)$$

(1-9).

$$C_{es}(T) = -T\left(\frac{\partial^2 F}{\partial T^2}\right) = -T\sqrt{x}(-\gamma) = \gamma T\sqrt{x}$$
$$= \gamma T_c \left(\frac{T}{T_c}\right)^3 :: \text{Eq.}(1\text{-}6)$$

- 1.2.2 The London Theory
- 1.2.3 F.London's Justification of the London Theory
- 1.2.4 Pippard's Nonlocal Generalization of the London Theory
- 1.2.5 Ginsburg-Landau Theory

(1-43).

$$f(T) = a\left(-\frac{a}{b}\right) + \frac{1}{2}b\left(-\frac{a}{b}\right)^2$$
$$= -\frac{a^2}{b} + \frac{a^2}{2b} = -\frac{a^2(T)}{2b(T)}$$

(1-44).

$$\frac{\lambda^2(0)}{\lambda^2(T)} = \frac{n_s^2(T)}{n_s^2(0)} = \frac{|\Psi_e(T)|^2}{|\Psi_e(0)|^2} :: \text{Eq.}(1\text{-}38)$$
$$= |\Psi_e(T)|^2 :: n = n_s \quad \text{for} \quad T = 0$$
$$: : |\Psi_e(0)| = 1$$

(1-48).

$$\begin{split} \frac{\partial f}{\partial \Psi(\mathbf{r})} &= 0 \\ \Rightarrow 0 &= -\frac{\hbar^2}{2m^*} \bigg(\nabla + \frac{\mathrm{i} e^*}{\hbar c} A(\mathbf{r}) \bigg) 2 \Psi(\mathbf{r}) \\ 0 &= -\frac{H_c^2(T)}{4\pi m^*} \frac{\lambda^2(T)}{\lambda^2(0)} \bigg(1 - \frac{\lambda^2(T)}{\lambda^2(0)} |\Psi(\mathbf{r})|^2 \bigg) 2 \Psi(\mathbf{r}) \\ \therefore & \text{Eq.} (1\text{-}48) \end{split}$$

- (1-51).
- (1-52).

第2章

The Pairing Theory of Superconductivity

- 2.1 Physical Nature of the Superconducting State
- 2.2 The One-Pair Problem
- 2.3 Landau's Theory of a Fermi Liquid
- 2.4 The Pairing Approximation
- 2.5 Quasi-Particle Excitations
- 2.6 Linearized Equations of Motion

参考文献

[1] J.R.Schrieffer. 1964. Theory of Superconductivity [Revised Printing]. http://zimp.zju.edu.cn/~qchen/Teaching/AdvStat/Schrieffer.pdf