

# Dzyaloshinskii-Moriya interaction

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We use the wannier function as the bases to expand the Hamiltonian.

$$\{\omega_{n\uparrow}(\mathbf{r} - \mathbf{R}), \omega_{n\downarrow}(\mathbf{r} - \mathbf{R})\} \quad (1)$$

And the  $\hat{\alpha}_{n\uparrow}(\mathbf{R})$  and  $\hat{\alpha}_{n\downarrow}^\dagger(\mathbf{R})$  are the annihilation and the creation operators of the electrons in the state  $\omega_{n\uparrow}(\mathbf{r} - \mathbf{R})$ , etc.

Then the Hamiltonian of a system after considering spin-orbit coupling (SOC) can be written as,

$$\hat{H}_1 = \frac{\hat{\mathbf{p}}^2}{2m} + V(\hat{\mathbf{r}}) + \frac{\hbar}{2m^2c^2} \hat{\mathbf{S}} \cdot [\nabla V(\hat{\mathbf{r}}) \times \hat{\mathbf{p}}] \quad (2)$$

The last term is SOC term ( $\hat{H}_{\text{SOC}}$ ) drive from the Dirac equation<sup>a</sup>.

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<sup>a</sup>Spin-orbit coupling: Dirac equation

Suppose  $V(\mathbf{r}) = V(|\mathbf{r}|) = V(r)$ ,

$$\begin{aligned} H_{\text{SOC}} &= \frac{\hbar}{2m^2c^2} \hat{\mathbf{S}} \cdot \left[ \frac{dV(r)}{dr} \frac{\hat{\mathbf{r}}}{r} \times \hat{\mathbf{p}} \right] \\ &= \lambda \hat{\mathbf{L}} \cdot \hat{\mathbf{S}} \\ &= \lambda \left( \hat{L}_z \hat{S}_z + \frac{1}{2} (\hat{L}_+ \hat{S}_- + \hat{L}_- \hat{S}_+) \right) \end{aligned} \quad (3)$$

# Mechanisms of DMI $\triangleright$ Basis-set expansion

$$\begin{aligned}
 \hat{H} = \hat{H}_0^{\text{all}} + \hat{T}^{\text{all}} = & \sum_{\mathbf{R}} \sum_n \epsilon_n(\mathbf{R}) \left[ \hat{\alpha}_{n\uparrow}^\dagger(\mathbf{R}) \hat{\alpha}_{n\uparrow}(\mathbf{R}) + \hat{\alpha}_{n\downarrow}^\dagger(\mathbf{R}) \hat{\alpha}_{n\downarrow}(\mathbf{R}) \right] \\
 & + \sum_{\mathbf{R} \neq \mathbf{R}'} \sum_{n, n'} \left\{ b_{n'n}(\mathbf{R}' - \mathbf{R}) \left[ \hat{\alpha}_{n'\uparrow}^\dagger(\mathbf{R}') \hat{\alpha}_{n\uparrow}(\mathbf{R}) + \hat{\alpha}_{n'\downarrow}^\dagger(\mathbf{R}') \hat{\alpha}_{n\downarrow}(\mathbf{R}) \right] \right. \\
 & \quad + C_{n'n}^z(\mathbf{R}' - \mathbf{R}) \left[ \hat{\alpha}_{n'\uparrow}^\dagger(\mathbf{R}') \hat{\alpha}_{n\uparrow}(\mathbf{R}) - \hat{\alpha}_{n'\downarrow}^\dagger(\mathbf{R}') \hat{\alpha}_{n\downarrow}(\mathbf{R}) \right] \\
 & \quad + C_{n'n}^-(\mathbf{R}' - \mathbf{R}) \hat{\alpha}_{n'\uparrow}^\dagger(\mathbf{R}') \hat{\alpha}_{n\downarrow}(\mathbf{R}) \\
 & \quad \left. + C_{n'n}^+(\mathbf{R}' - \mathbf{R}) \hat{\alpha}_{n'\downarrow}^\dagger(\mathbf{R}') \hat{\alpha}_{n\uparrow}(\mathbf{R}) \right\} \quad (4)
 \end{aligned}$$

$$b_{n'n}(\mathbf{R}' - \mathbf{R}) + C_{n'n}^z(\mathbf{R}' - \mathbf{R}) = \int \omega_{n'\uparrow}^*(\mathbf{r} - \mathbf{R}') H_1 \omega_{n\uparrow}(\mathbf{r} - \mathbf{R}) d\mathbf{r} \quad (5a)$$

$$b_{n'n}(\mathbf{R}' - \mathbf{R}) - C_{n'n}^z(\mathbf{R}' - \mathbf{R}) = \int \omega_{n'\downarrow}^*(\mathbf{r} - \mathbf{R}') H_1 \omega_{n\downarrow}(\mathbf{r} - \mathbf{R}) d\mathbf{r} \quad (5b)$$

$$C_{n'n}^x(\mathbf{R}' - \mathbf{R}) - iC_{n'n}^y(\mathbf{R}' - \mathbf{R}) = C_{n'n}^-(\mathbf{R}' - \mathbf{R}) = \int \omega_{n'\uparrow}^*(\mathbf{r} - \mathbf{R}') H_1 \omega_{n\downarrow}(\mathbf{r} - \mathbf{R}) d\mathbf{r} \quad (5c)$$

$$C_{n'n}^x(\mathbf{R}' - \mathbf{R}) + iC_{n'n}^y(\mathbf{R}' - \mathbf{R}) = C_{n'n}^+(\mathbf{R}' - \mathbf{R}) = \int \omega_{n'\downarrow}^*(\mathbf{r} - \mathbf{R}') H_1 \omega_{n\uparrow}(\mathbf{r} - \mathbf{R}) d\mathbf{r} \quad (5d)$$

# Mechanisms of DMI ▷ Perturbation

Basically, there are three techniques to achieve our aims, that perturbed the low energy subspace with the high energy excitation space.

- Downfolding<sup>1</sup>
- Lowdin Perturbation<sup>2</sup>
- Green's Function

All of those three methods gives the same result:

$$\hat{H}_{\text{eff}} = \hat{H}_0 + \frac{\hat{T}^\dagger \hat{T}}{U} = \hat{H}_0 + \hat{H}_M \quad (6)$$

Where  $U$  is the Hubbard  $U$  refer to the energy cost when move two electrons to the same site. And  $\hat{T}$  only content the hoping terms contribute to this process.

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<sup>1</sup> Introduction to the Magnetic Exchange Mechanisms, Yang Li, 2020

<sup>2</sup> <https://link.springer.com/content/pdf/bbm%3A978-1-4615-5673-2%2F1.pdf>

# Mechanisms of DMI $\triangleright$ Simplification

$$\hat{H}_M = \frac{1}{U} \hat{T}^\dagger \hat{T}$$

$$\begin{aligned} \hat{T} = & b_{n'n}(\mathbf{R}' - \mathbf{R}) \left[ \hat{\alpha}_{n'\uparrow}^\dagger(\mathbf{R}') \hat{\alpha}_{n\uparrow}(\mathbf{R}) + \hat{\alpha}_{n'\downarrow}^\dagger(\mathbf{R}') \hat{\alpha}_{n\downarrow}(\mathbf{R}) \right] \\ & + C_{n'n}^z(\mathbf{R}' - \mathbf{R}) \left[ \hat{\alpha}_{n'\uparrow}^\dagger(\mathbf{R}') \hat{\alpha}_{n\uparrow}(\mathbf{R}) - \hat{\alpha}_{n'\downarrow}^\dagger(\mathbf{R}') \hat{\alpha}_{n\downarrow}(\mathbf{R}) \right] \\ & + C_{n'n}^-(\mathbf{R}' - \mathbf{R}) \hat{\alpha}_{n'\uparrow}^\dagger(\mathbf{R}') \hat{\alpha}_{n\downarrow}(\mathbf{R}) \\ & + C_{n'n}^+(\mathbf{R}' - \mathbf{R}) \hat{\alpha}_{n'\downarrow}^\dagger(\mathbf{R}') \hat{\alpha}_{n\uparrow}(\mathbf{R}) \end{aligned} \quad (7)$$

The relation between fermi annihilation/creation operators and spin operators is,

$$\hat{S}_{z,n}(\mathbf{R}) = \frac{1}{2} \left[ \hat{\alpha}_{n\uparrow}^\dagger(\mathbf{R}) \hat{\alpha}_{n\uparrow}(\mathbf{R}) - \hat{\alpha}_{n\downarrow}^\dagger(\mathbf{R}) \hat{\alpha}_{n\downarrow}(\mathbf{R}) \right] \quad (8a)$$

$$\hat{S}_{+,n}(\mathbf{R}) = \hat{\alpha}_{n\uparrow}^\dagger(\mathbf{R}) \hat{\alpha}_{n\downarrow}(\mathbf{R}) \quad (8b)$$

$$\hat{S}_{-,n}(\mathbf{R}) = \hat{\alpha}_{n\downarrow}^\dagger(\mathbf{R}) \hat{\alpha}_{n\uparrow}(\mathbf{R}) \quad (8c)$$

$$\begin{aligned}
 \hat{H}_M = & J_{\mathbf{R},\mathbf{R}'} \hat{\mathbf{S}}(\mathbf{R}) \cdot \hat{\mathbf{S}}(\mathbf{R}') \\
 & + \mathbf{D}_{\mathbf{R},\mathbf{R}'} \cdot \left[ \hat{\mathbf{S}}(\mathbf{R}) \times \hat{\mathbf{S}}(\mathbf{R}') \right] \\
 & + \hat{\mathbf{S}}(\mathbf{R}) \cdot \overleftrightarrow{\Gamma}_{\mathbf{R},\mathbf{R}'} \cdot \hat{\mathbf{S}}(\mathbf{R}')
 \end{aligned} \tag{9}$$

Where,

$$J_{\mathbf{R},\mathbf{R}'} = 2|b_{nn'}(\mathbf{R} - \mathbf{R}')|^2/U \tag{10a}$$

$$\mathbf{D}_{\mathbf{R},\mathbf{R}'} = (4i/U) [b_{nn'}(\mathbf{R} - \mathbf{R}')\mathbf{C}_{n'n}(\mathbf{R}' - \mathbf{R}) - b_{n'n}(\mathbf{R}' - \mathbf{R})\mathbf{C}_{nn'}(\mathbf{R} - \mathbf{R}')] \tag{10b}$$

$$\begin{aligned}
 \overleftrightarrow{\Gamma}_{\mathbf{R},\mathbf{R}'} = & 4/U [\mathbf{C}_{n'n}(\mathbf{R}' - \mathbf{R}) \otimes \mathbf{C}_{nn'}(\mathbf{R} - \mathbf{R}') + \mathbf{C}_{nn'}(\mathbf{R} - \mathbf{R}') \otimes \mathbf{C}_{n'n}(\mathbf{R}' - \mathbf{R}) \\
 & - (\mathbf{C}_{n'n}(\mathbf{R}' - \mathbf{R}) \cdot \mathbf{C}_{nn'}(\mathbf{R} - \mathbf{R}')) \mathbb{1}]
 \end{aligned} \tag{10c}$$

And,

$$\begin{aligned}
 b_{nn'}(\mathbf{R} - \mathbf{R}') &= b_{n'n}^*(\mathbf{R}' - \mathbf{R}) \\
 \mathbf{C}_{nn'}(\mathbf{R} - \mathbf{R}') &= \mathbf{C}_{n'n}^*(\mathbf{R}' - \mathbf{R}) \\
 \mathbf{C} &= (C_x, C_y, C_z)
 \end{aligned} \tag{11}$$

$$\hat{H}_{\text{DM}} = \mathbf{D} \cdot (\hat{\mathbf{S}}_1 \times \hat{\mathbf{S}}_2) \quad (12)$$

Suppose the two ions contribute to the electron transfer is located at point  $A$  and  $B$ , the point bisecting  $AB$  is denoted by  $C$ .

- When a center of inversion is located at  $C$ ,  $\mathbf{D} = 0$
- When a mirror plane perpendicular to  $AB$  passes through  $C$ ,  $\mathbf{D} \parallel \text{mirror plane}$
- When a mirror plane including  $AB$ ,  $\mathbf{D} \perp \text{mirror plane}$
- When a  $C_2$  axis perpendicular to  $AB$  passes  $C$ ,  $\mathbf{D} \perp C_2 \text{ axis}$
- When there is a  $C_n$  ( $n \geq 2$ ) along  $AB$ ,  $\mathbf{D} \parallel AB$