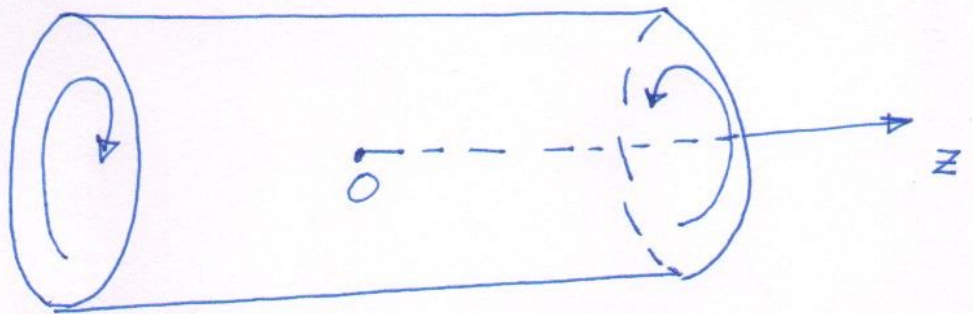


## Combined extension-torsion of Circular beams



By physical considerations:-

$$u_r(r), \quad u_\theta(r, z) = r \frac{\theta_0}{L} z, \quad u_z(z)$$

$$\Rightarrow \underline{\underline{\epsilon}} = \begin{bmatrix} u_r' & 0 & 0 \\ 0 & u_r/r & \frac{r\theta_0}{2L} \\ 0 & \frac{r\theta_0}{2L} & u_z' \end{bmatrix} \text{ in cylindrical coordinate.}$$

$$\text{So, } \sigma_{rr} = \lambda \left( u_r' + \frac{u_r}{r} + u_z' \right) + 2\mu u_r'$$

$$\sigma_{\theta\theta} = \lambda \left( \text{''} \right) + 2\mu \frac{u_r}{r}$$

$$\sigma_{zz} = \lambda \left( \text{''} \right) + 2\mu u_z'$$

$$\sigma_{r\theta} = \sigma_{zr} = 0, \quad \sigma_{\theta z} = \frac{G\theta_0}{L} r$$

Let us look at '0'-equation

$$\Rightarrow \cancel{\frac{\partial \sigma_{rr}}{\partial r}} + \frac{1}{r} \cancel{\frac{\partial \sigma_{\theta\theta}}{\partial \theta}} + \cancel{\frac{\partial \sigma_{zz}}{\partial z}} + 2 \cancel{\frac{\sigma_{r\theta}}{r}} = 0$$

Hence, it gets trivially satisfied

Z. equation.

$$\Rightarrow \cancel{\frac{\partial \sigma_{rz}}{\partial r}} + \cancel{\frac{1}{r} \frac{\partial \sigma_{\theta z}}{\partial \theta}} + \frac{\partial \sigma_{zz}}{\partial z} + \cancel{\frac{\sigma_{rz}}{r}} = 0$$

$$\Rightarrow \frac{\partial \sigma_{zz}}{\partial z} = 0 \quad \text{or, } \sigma_{zz} \text{ is independent of } z$$

$$\Rightarrow \lambda \left( u_r' + \frac{u_r}{r} + u_z' \right) + 2\mu u_z' = \text{indep. of } z$$

$$\Rightarrow \lambda \left( u_r' + \frac{u_r}{r} \right) + (\lambda + 2\mu) u_z' = \quad "$$

$$\Rightarrow \boxed{u_z' = \text{constant} = C}$$

Now, let us look at radial eq.

$$\frac{\partial \sigma_{rr}}{\partial r} + \cancel{\frac{1}{r} \frac{\partial \sigma_{\theta\theta}}{\partial \theta}} + \cancel{\frac{\partial \sigma_{rz}}{\partial z}} + \frac{\sigma_{rr} - \sigma_{\theta\theta}}{r} = 0$$

$$\Rightarrow \frac{\partial \sigma_{rr}}{\partial r} + \frac{\sigma_{rr} - \sigma_{\theta\theta}}{r} = 0$$

$$\Rightarrow \lambda \left( u_r' + \frac{u_r}{r} \right)' + 2\mu u_r'' + 2\mu \left( \frac{u_r'}{r} - \frac{u_r}{r^2} \right) = 0$$

$$\Rightarrow (\lambda + 2\mu) u_r'' + (\lambda + 2\mu) \frac{u_r'}{r} - (\lambda + 2\mu) \frac{u_r}{r^2} = 0$$

$$\Rightarrow u_r'' + \frac{u_r'}{r} - \frac{u_r}{r^2} = 0$$

$$\Rightarrow u_r'' + \left( \frac{u_r}{r} \right)' = 0 \Rightarrow \boxed{u_r' + \frac{u_r}{r} = C}$$

$$\text{or, } \boxed{\epsilon_{rr} + \epsilon_{\theta\theta} = C}$$

$$\epsilon_{rr} + \epsilon_{\theta\theta} = C$$



As,  $u_r' + \frac{u_r}{r} = C \Rightarrow \sigma_{zz} = \lambda C + (\lambda + 2\mu) \epsilon$

$$\boxed{\sigma_{zz} = \lambda (C + \epsilon) + 2\mu \epsilon} \rightarrow \text{a constant}$$

Also, note that

$$\sigma_{rr} + \sigma_{\theta\theta} = 2(\lambda + \mu) \left( u_r' + \frac{u_r}{r} \right) + 2\lambda u_z'$$

or,  $\boxed{\sigma_{rr} + \sigma_{\theta\theta} = 2(\lambda + \mu) C + 2\lambda \epsilon} \rightarrow \text{another constant}$

Let us look at radial equation again

$$\frac{d\sigma_{rr}}{dr} + \frac{\sigma_{rr} - \sigma_{\theta\theta}}{r} = 0$$

$$\Rightarrow \frac{d\sigma_{rr}}{dr} + 2\frac{\sigma_{rr}}{r} = \frac{\sigma_{rr} + \sigma_{\theta\theta}}{r} = \frac{A}{r}$$

$$\Rightarrow \frac{1}{r^2} \frac{d}{dr} (\sigma_{rr} r^2) = \frac{A}{r}$$

$$\Rightarrow \boxed{\sigma_{rr} = \frac{A}{2} + \frac{B}{r^2}}$$

Solid tube

$$\sigma_{rr} = \text{finite at } r=0$$

$$\Rightarrow B=0$$

$$\sigma_{rr} = 0 \text{ at } r=r_2$$

$$\Rightarrow A=0$$

$$\Rightarrow \sigma_{rr} = 0$$

Hollow tube

$$\sigma_{rr} = 0 \text{ at } r=r_1, r_2$$

$$\Rightarrow A=B=0$$

Hence,  $\sigma_{rr} = 0$  for both solid and hollow tubes

Plugging  $\sigma_{rr} = 0$  in radial eq

$$\Rightarrow \sigma_{\theta\theta} = 0$$

Hence,  $\sigma_{rr} = \sigma_{\theta\theta} = 0$  for both solid and hollow beams.

Again  $u_r' + \frac{u_r}{r} = C$

$$\Rightarrow \frac{1}{r} \frac{d}{dr} (r u_r) = C \Rightarrow \boxed{u_r = \frac{C}{2} r + \frac{D}{r}}$$

plugging it into expression for  $\sigma_{rr}$

$$\Rightarrow \lambda(C + \epsilon) + 2\mu \left( \frac{C}{2} - \frac{D}{r^2} \right) = 0$$

$$\Rightarrow (\lambda + \mu)C + \lambda\epsilon - \frac{2\mu D}{r^2} = 0$$

$$\Rightarrow D = 0 \quad (\text{no dependence on } r)$$

$$\Rightarrow (\lambda + \mu)C + \lambda\epsilon = 0$$

$$\Rightarrow C = -\frac{\lambda}{\lambda + \mu} \epsilon$$

$$\Rightarrow \epsilon_{rr} = \epsilon_{\theta\theta} = \frac{C}{2} = -\frac{\lambda}{2(\lambda + \mu)} \epsilon$$

or,

$$\boxed{\epsilon_{rr} = \epsilon_{\theta\theta} = -\nu \epsilon}$$

and

$$\boxed{u_r = -\nu \epsilon r}$$



Now, we have to understand how to generate  $\epsilon$  and  $\theta_z$

for this note that

$$\epsilon = \frac{1}{E} \left( \sigma_{zz} - \nu(\cancel{\sigma_{rr}} + \cancel{\sigma_{\theta\theta}}) \right)$$

$$\Rightarrow \sigma_{zz} = E \epsilon$$

or, Axial force:  $\int \sigma_{zz} dA = E A \epsilon$

So, Axial force of  $E A \epsilon$  is needed to generate strain " $\epsilon$ " and  
 $\downarrow$   
This does not change even if torsion is present

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Similarly,

$$\sigma_{\theta z} = G \frac{\theta_z}{L} r$$

$$\Rightarrow \text{Torque (T)} = \int_{r_1}^{r_2} r (2\pi r dr) \frac{G \theta_z}{L} r$$
$$= G \frac{\theta_z}{L} \int r^2 dA = G \frac{\theta_z}{L} J$$

or,  $\boxed{T = G \frac{\theta_z}{L} J}$

\* Basically torsion and axial extension become independent of each other

\* Also,  $u_r = -\nu \epsilon r$  does not arise if only torsion is present