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## Literature Review

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# Particle Identification in Modern Cherenkov Detectors

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### Abstract

With major detector assemblies lining up for construction at high-energy facilities such as EIC and FAIR, technologies for particle identification are becoming ever more relevant. This review motivates the theory of the field, introduces the technologies behind modern Cherenkov detectors, and discusses the implementation of these detectors into upcoming experiments. Emphasis is placed on ..... Persistent challenges include ..... We conclude by outlining promising future developments for .....

*Keywords:* Particle Identification, Cherenkov Detectors, DIRC, RICH, PANDA, EIC, GSI

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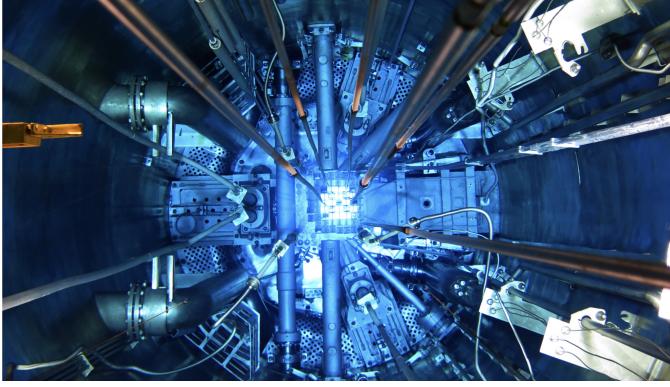


Figure 1: Cherenkov radiation in a nuclear reactor pool. The characteristic glow is indicative of relativistic electrons moving through the water. (Photo: IAEA)

## 1. Introduction

### 2. Theoretical Foundations

#### 2.1. Cherenkov Radiation

Cherenkov radiation is a phenomenon that occurs when a charged particle travels through a dielectric medium with refractive index  $n(\lambda)$  at a speed greater than the phase velocity in that medium. This effect was first observed by Pavel Čherenkov in 1934 [1] and later explained theoretically by Ilya Frank and Igor Tamm [2].

From Maxwell's equations arises that electromagnetic propagation with wavelength  $\lambda$  will have its phase velocity modified by the medium it is travelling through. This leads to the idea of refractive index  $n(\lambda) = c/v$ , defined as the ratio between the speed of light in vacuum  $c$  and the phase velocity  $v$  in the medium. The condition for Cherenkov radiation to occur is then given by  $v > c/n$ .

The Cherenkov radiation itself propagates at the phase velocity of the medium. As such, the radiation forms a conical wavefront with characteristic angle  $\theta$ , similar to how shock cones are formed by supersonic aircraft. The angle  $\theta$  can be derived from simple geometric considerations shown in Fig 2, yielding the following:

$$\cos \theta = \frac{v_p}{v} = \frac{c}{nv} \quad (1)$$

The main principle behind a Cherenkov detector is to measure the angle  $\theta$  of the emitted photons, which can then be used to determine the velocity  $v$  of the charged particle. With an associated momentum measurement from a forward tracking detector, the mass of the particle can be inferred, allowing for spectroscopy. This velocity spectroscopy for various particle species is shown in Fig 3. It can be seen that pions, for example, will produce a larger Cherenkov angle than kaons or protons at the same momentum, due to them being less massive.

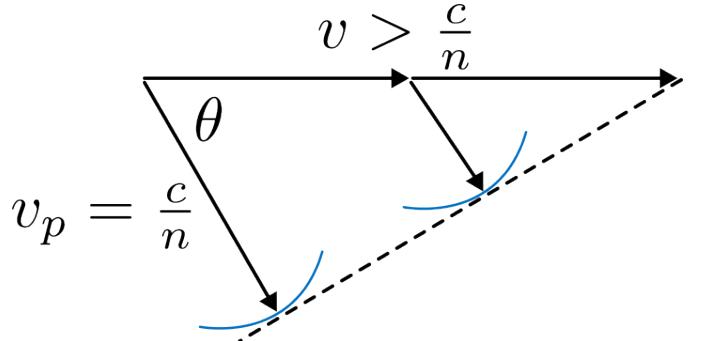


Figure 2: Schematic of a Cherenkov cone. A particle travels through a medium with velocity  $v$ . Photons are emitted at the phase velocity  $v_p$  and angle  $\theta$  to the path of travel of the particle, giving the characteristic blue shock cone.

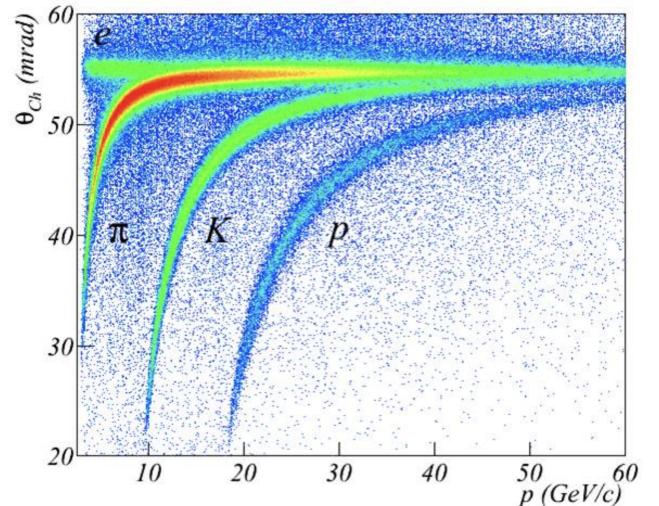


Figure 3: Cherenkov spectroscopy of the COMPASS RICH-1 detector at CERN. There are clear curves for different particle species. These curves can be determined mathematically using (1) and relativistic dynamics. (Photo: [3])

The energy  $dE$  emitted per unit length is given by the Frank-Tamm formula, which can be expressed as

$$\frac{dE}{dx} = \frac{q^2}{4\pi} \int_{v > \frac{c}{n(\omega)}} \mu(\omega) \omega \left(1 - \frac{c^2}{v^2 n^2(\omega)}\right) d\omega \quad (2)$$

for a particle with charge  $q$  and frequency  $\omega$ , moving through a medium with permeability  $\mu$  and refractive index  $n(\omega)$ . The integral is taken over all frequencies where the Cherenkov condition  $v = c/n(\omega)$  is met. By changing variables to wavelength  $\lambda$  and using that  $E = Nh\omega$ , one can show that the number of photons  $N$  emitted per unit length is inversely proportional to the square of the wavelength. This explains the characteristic blue of Cherenkov radiation, as shorter wavelength (blue light) photons are emitted more.

## 2.2. Relativistic Kinematics

Relativistic kinematics is the study of the motion of particles at near-light speeds. The purposes of this review do not require an exhaustive treatment of the subject, but a few key concepts pertaining to Cherenkov radiation and detector physics are summarised.

The famous energy-mass relation is written as:

$$E^2 = p^2 c^2 + m_0^2 c^4 \quad (3)$$

for a particle with energy  $E$ , momentum  $p$  and rest mass  $m_0$ . The Lorentz factor arises from considering the four-velocity, and defines a mapping between the rest frame of the particle (the frame in which the particle's velocity is zero) and the observer's frame. It is defined as:

$$\gamma = \frac{1}{\sqrt{1 - \beta^2}} \quad (4)$$

where  $\beta = v/c$  is the velocity of the particle as a fraction of the speed of light. We can use the Lorentz factor to transform Newtonian quantities into their relativistic equivalents using the equations

$$p = \gamma m_0 v \quad (5)$$

$$E = \gamma m_0 c^2 \quad (6)$$

Using (5) and (6) we can rewrite the Cherenkov angle condition from (1) as

$$\cos \theta = \frac{c}{nv} = \frac{E}{np c} = \frac{\sqrt{p^2 + m^2 c^2}}{np} \quad (7)$$

where one can clearer see how a combination of the Cherenkov angle and the momentum measurements can be used to infer the mass of the particle. This equation can be used to plot the theoretical curves for the experimental data in Fig 3.

Often in detector assemblies, the Cherenkov detectors will receive momentum measurements from forward detectors. Then, the statistics from the Cherenkov detectors

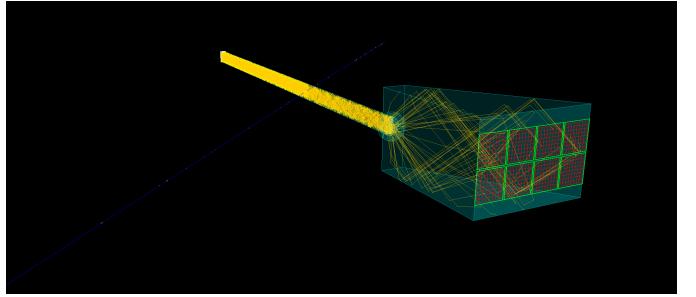


Figure 4: Simulated event display of PANDA Barrel DIRC. The blue track is a charged particle incident on the detector, and the yellow tracks are all Cherenkov photons. There is a mirror at the far end of the radiator bar, causing all photons to eventually reach the detector array.

are so impressive, that the detectors are usually able to decisively improve on the forward detector's measurements. This is especially true for DIRC detectors, which have an excellent angular resolution.

## 2.3. Optical Detector Theory

Cherenkov detectors rely heavily on the principles of optics to guide and detect emitted photons. Key concepts include Snell's law, total internal reflection, and the behaviour of light in various media. Fig 4 shows a simulated event display of PANDA's Barrel DIRC detector [4], illustrating how photons are internally reflected within the radiator bars and subsequently focused going into the trapezoidal expansion volume before being detected by the photodetector array.

One can imagine how important the optical properties and the geometry of the detector are to optimise things such as photon yield and angular resolution. All things told, the PANDA Barrel DIRC is actually a relatively rudimentary design compared to more recent simulations for EIC DIRC detectors, which employ more complex focusing optics to improve performance. This will be discussed further in § [sec:photon-optics](#) and § [sec:applications-at-future-facilities](#).

## 2.4. Particle Physics Motivation

Detector assemblies are becoming ever more complex out of necessity, as the continual push for higher precision measurements dominates the investigation into quantum chromodynamics. Many types of particle physics experiments require particle identification (PID) of hadrons such as pions, kaons and protons to observe rare processes or other heavy-ion collisions. These PID requirements often span across extensive momentum and solid angle ranges; Cherenkov detectors are excellent at fulfilling these requirements.

For example, at the upcoming ePIC experiment at the Electron-Ion Collider (EIC), Cherenkov methods will be

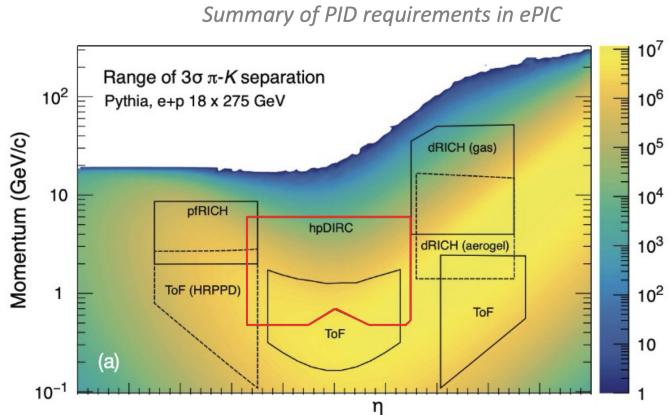


Figure 5: **The ePIC PID requirements.** The pfRICH, hpDIRC and dRICH are all Cherenkov based detectors that are expected to perform  $3\sigma$  pion and kaon separation at momenta between around 1 and 10 GeV. Here,  $\eta$  represents the pseudorapidity. (Photo: Greg Kalicy, ePIC Collaboration [5])

used to provide hadron PID across essentially all high momentum ranges, as shown in Fig 5. These detector performance metrics will be crucial to the experiment's physics goals towards investigating things such as nucleon spin, sea quarks, and high energy gluon behaviour.

To summarise, hadron detection is extremely important for flavour tagging and thus event reconstruction in these QCD processes. That is why Cherenkov detectors are imperative components of cutting edge detector assemblies.

## 2.5. Statistical Event Reconstruction Methods

## 2.6. Radiation-Matter Interactions

## 3. Photomultipliers

### 3.1. MCP-PMTs

### 3.2. SiPMs

### 3.3. LAPPDs

### 3.4. HRPPDs *might get cut*

## 4. Photon Optics *might get cut*

### 4.1. Lenses and Reflections

### 4.2. Laser Calibration

## 5. Ring Imaging Cherenkov Detectors (RICH)

## 6. Detection of Internally Reflected Cherenkov Light (DIRC)

## 7. Neutrino Identification with Cherenkov Detectors *might get cut*

### 7.1. Atmospheric Muons

### 7.2. Water Cherenkov Detectors

## 8. Applications at Future Facilities

### 8.1. PANDA DIRC at FAIR

### 8.2. hpDIRC at ePIC Brookhaven

### 8.3. Other Proposed Far Future Detectors *might get cut*

## 9. Current Challenges and Limitations

### 9.1. ML PID

### 9.2. High Rate Environments

### 9.3. Radiation Damage

### 9.4. Budget Constraints

## 10. Emerging Technologies

### 10.1. HRPPDs

### 10.2. Advanced Optics

### 10.3. New Materials

## 11. Summary

## References

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